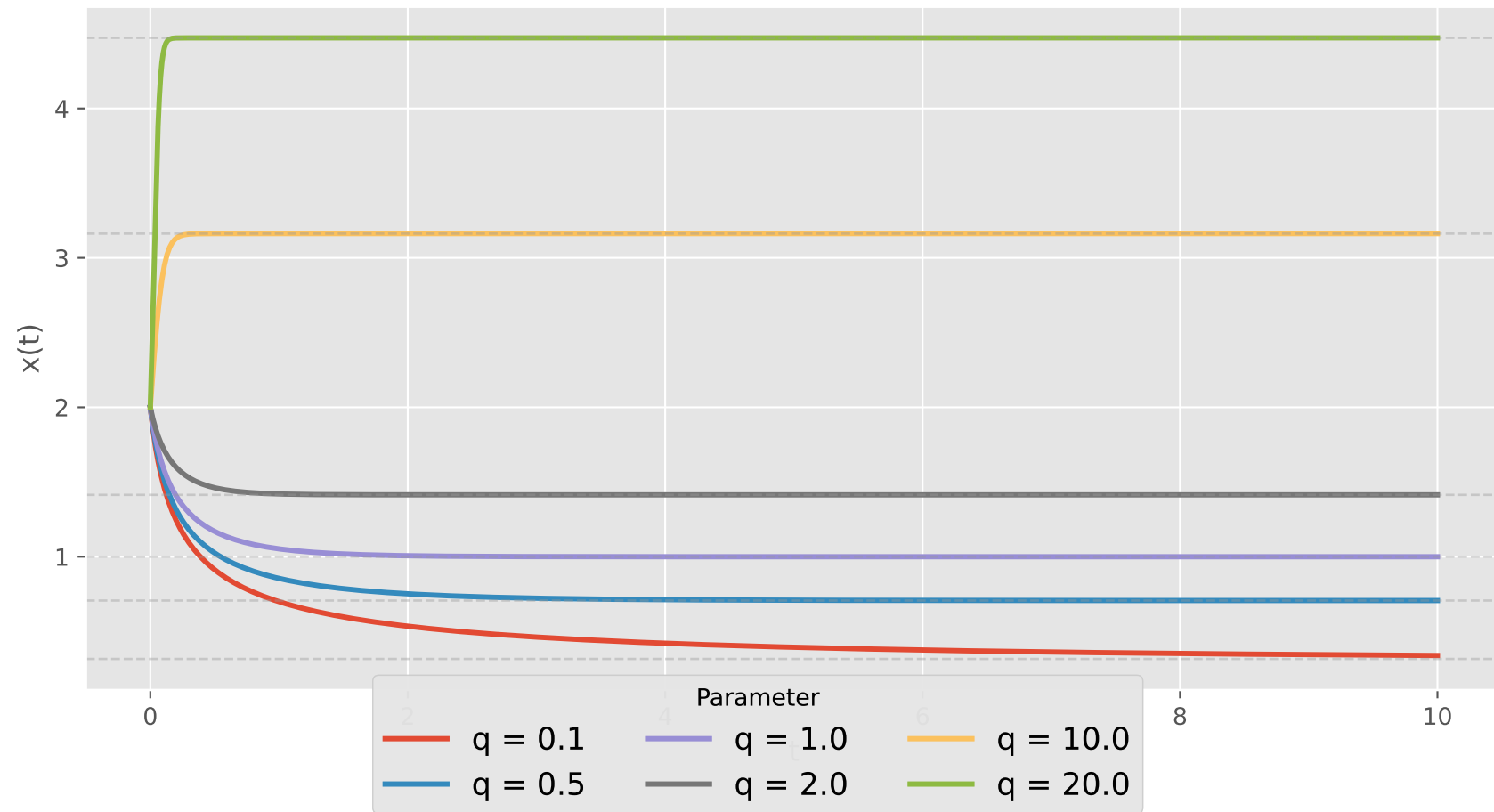
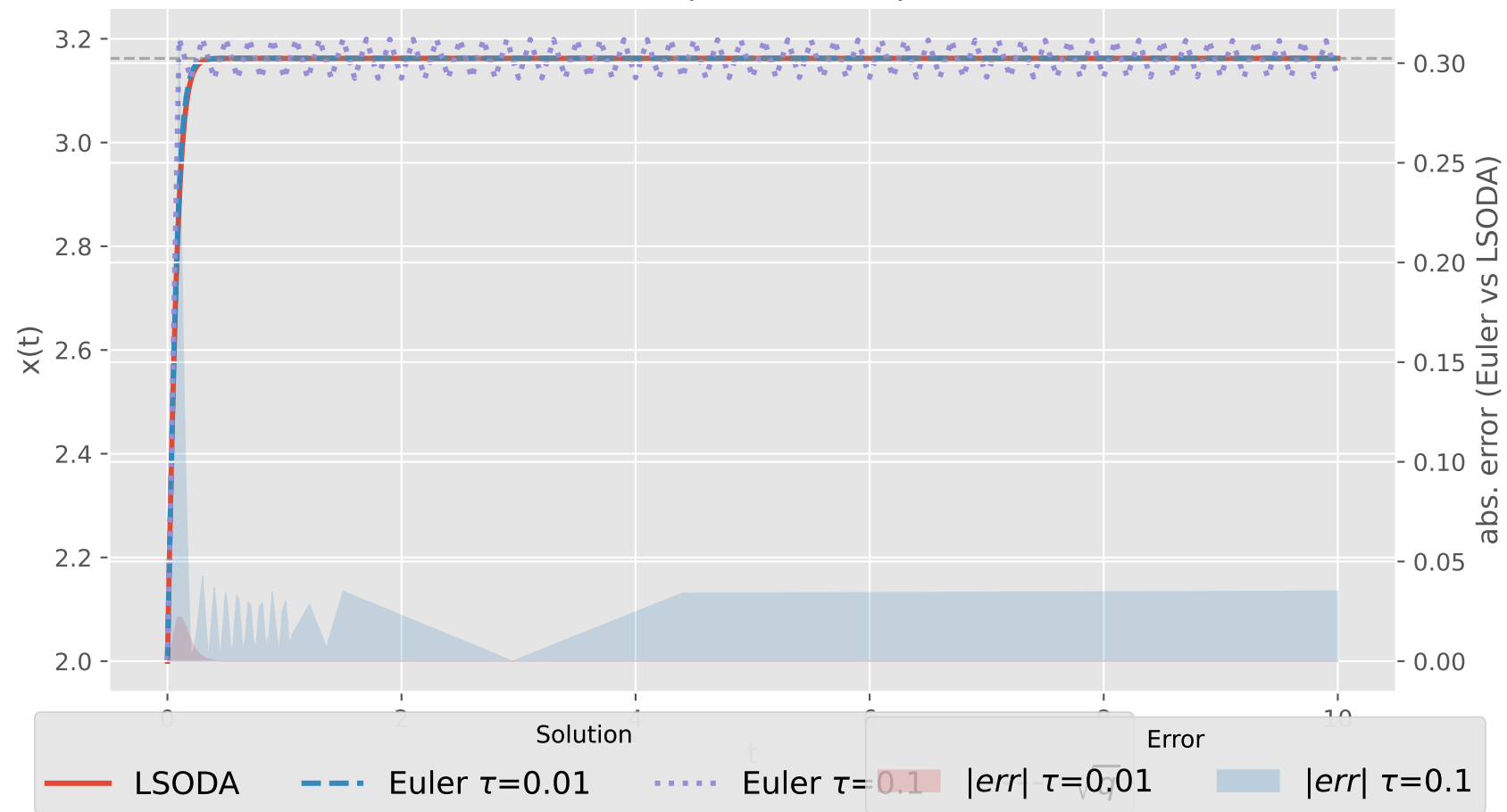


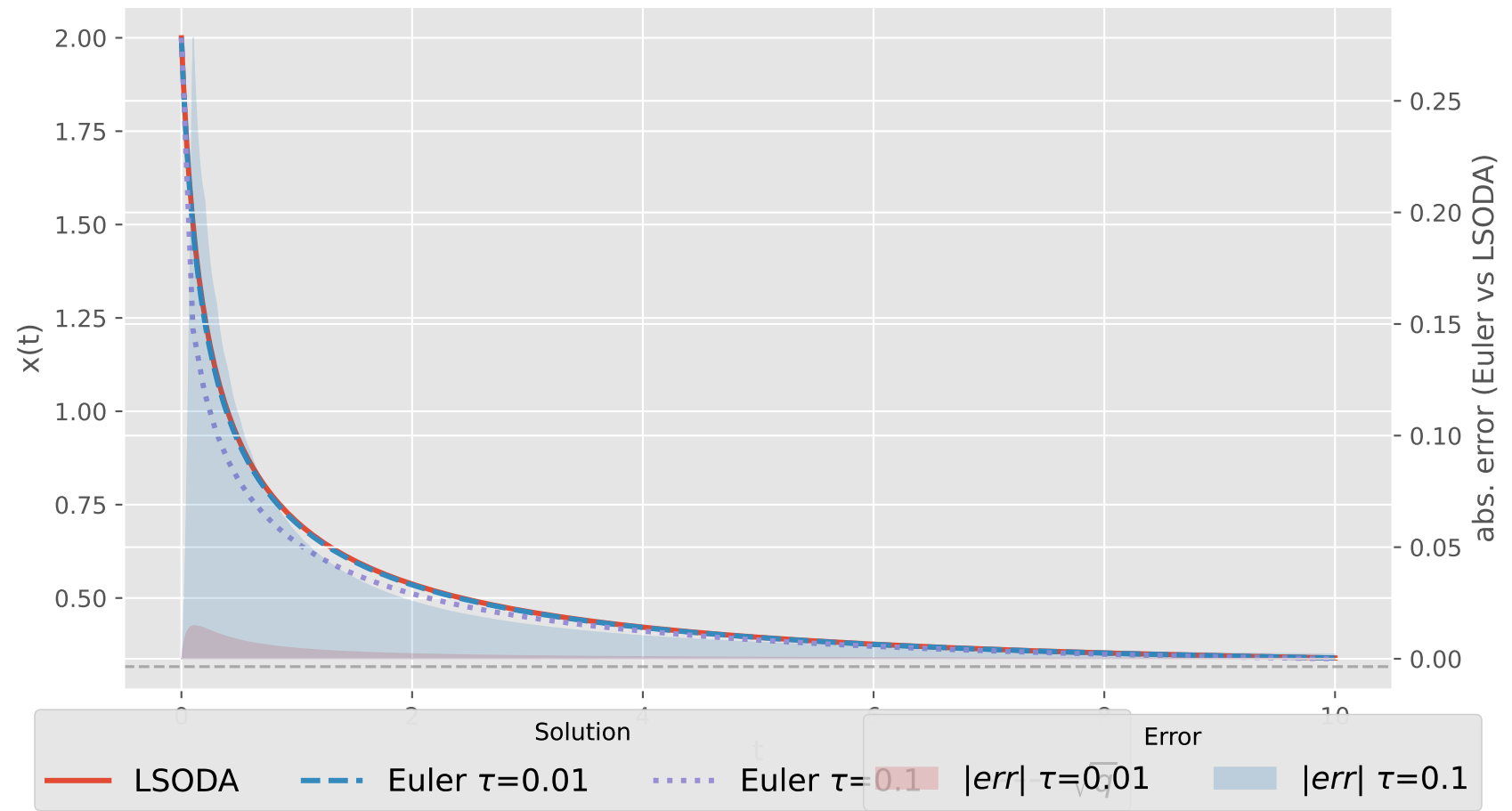
Cubic ODE parameter sweep (equilibria shown)



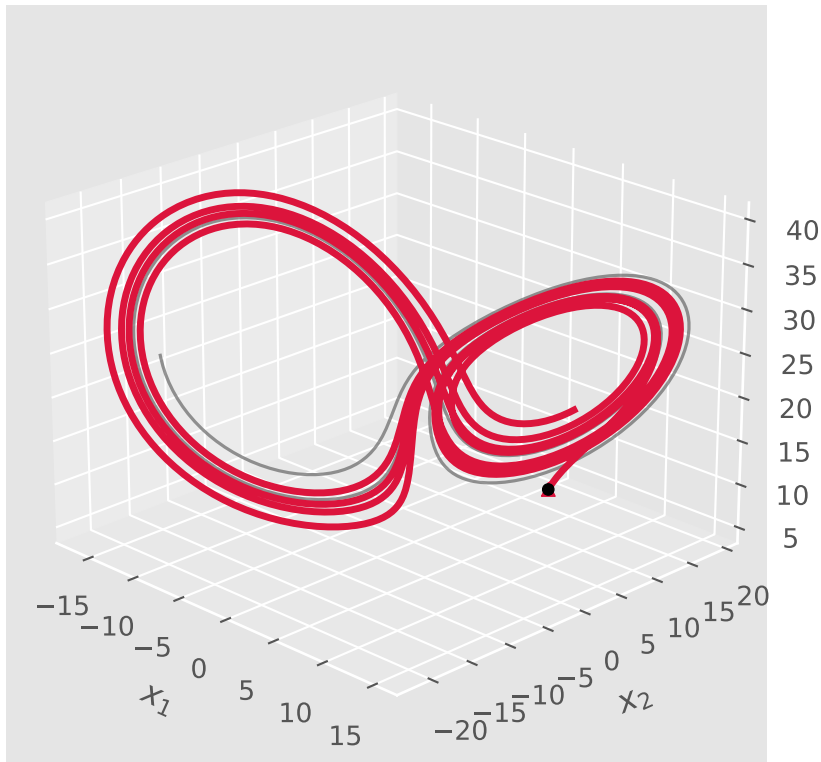
# Method comparison for $q = 10.0$



Method comparison for  $q = 0.1$

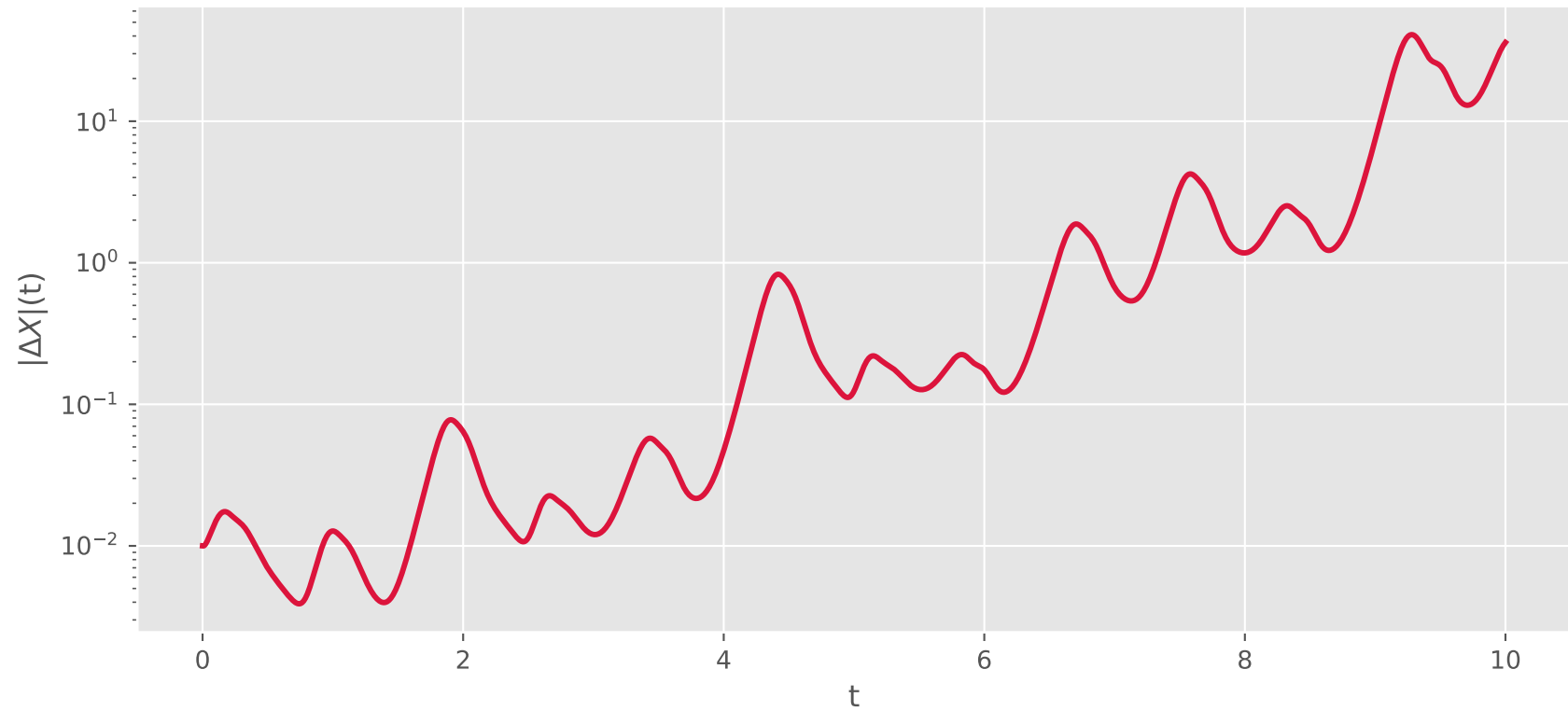


# Lorenz sensitivity to perturbing $x_2(0)$

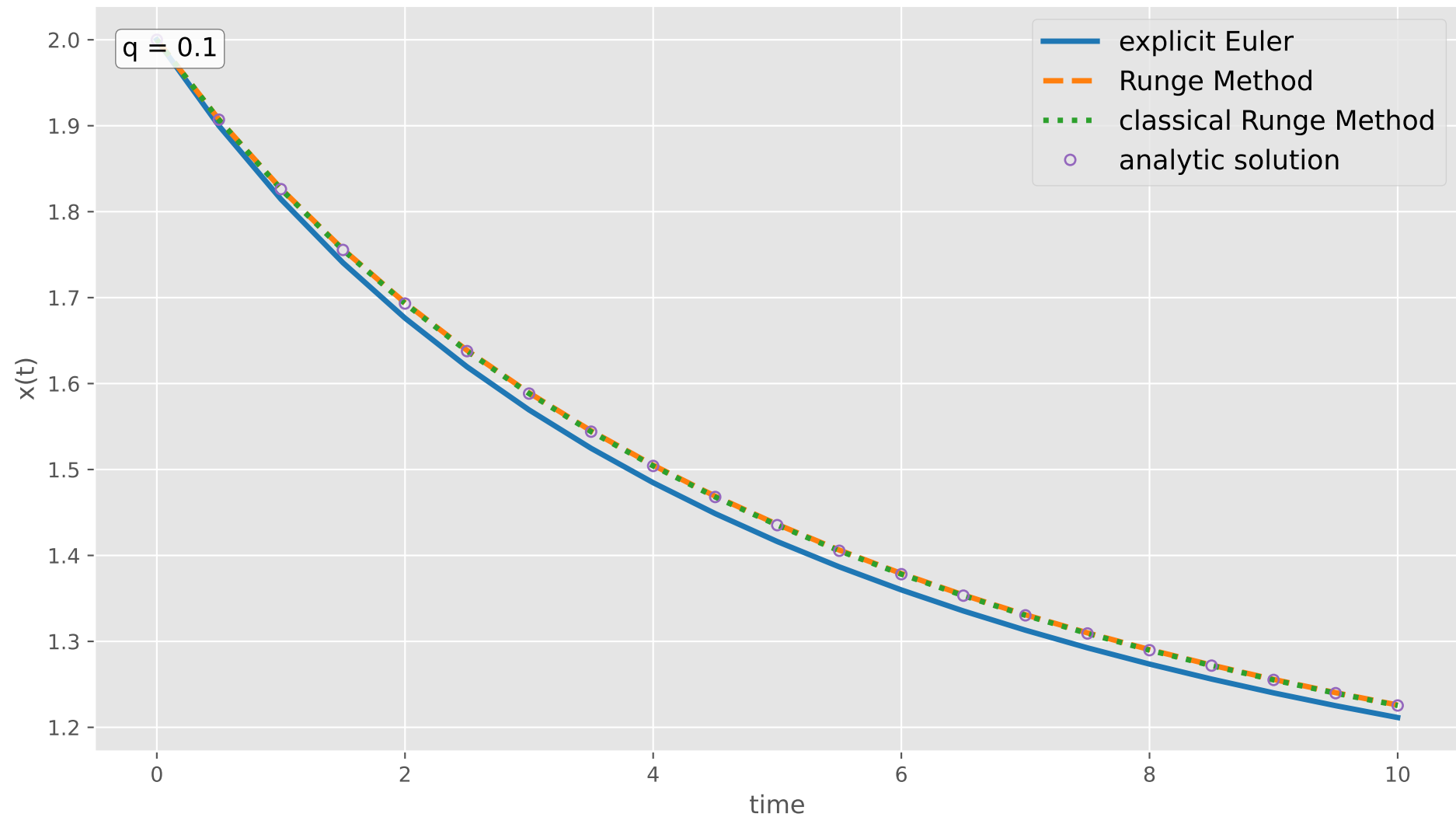


— Baseline      —  $x_2(0) = 5.01$  perturbed

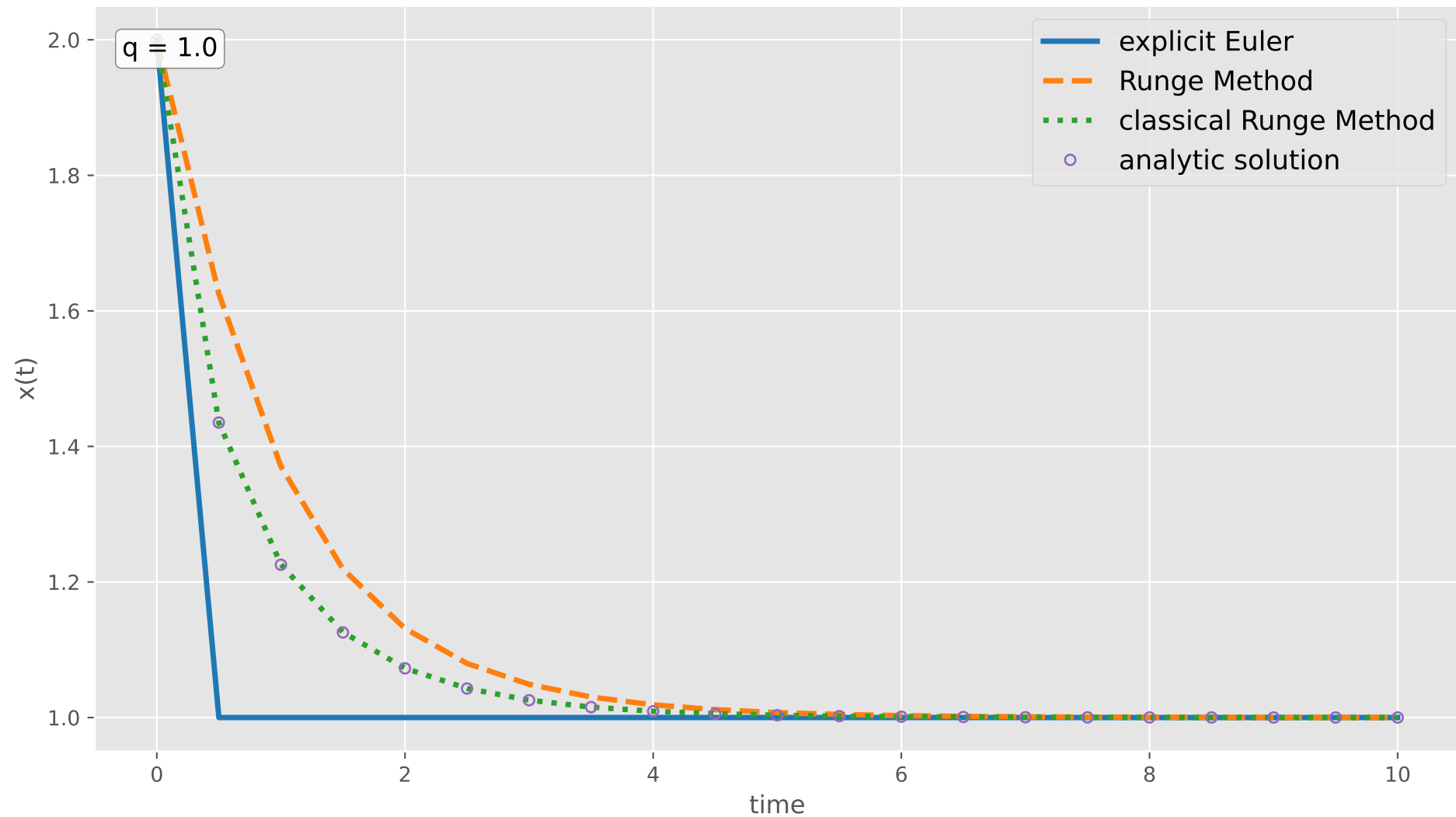
Lorenz trajectory separation



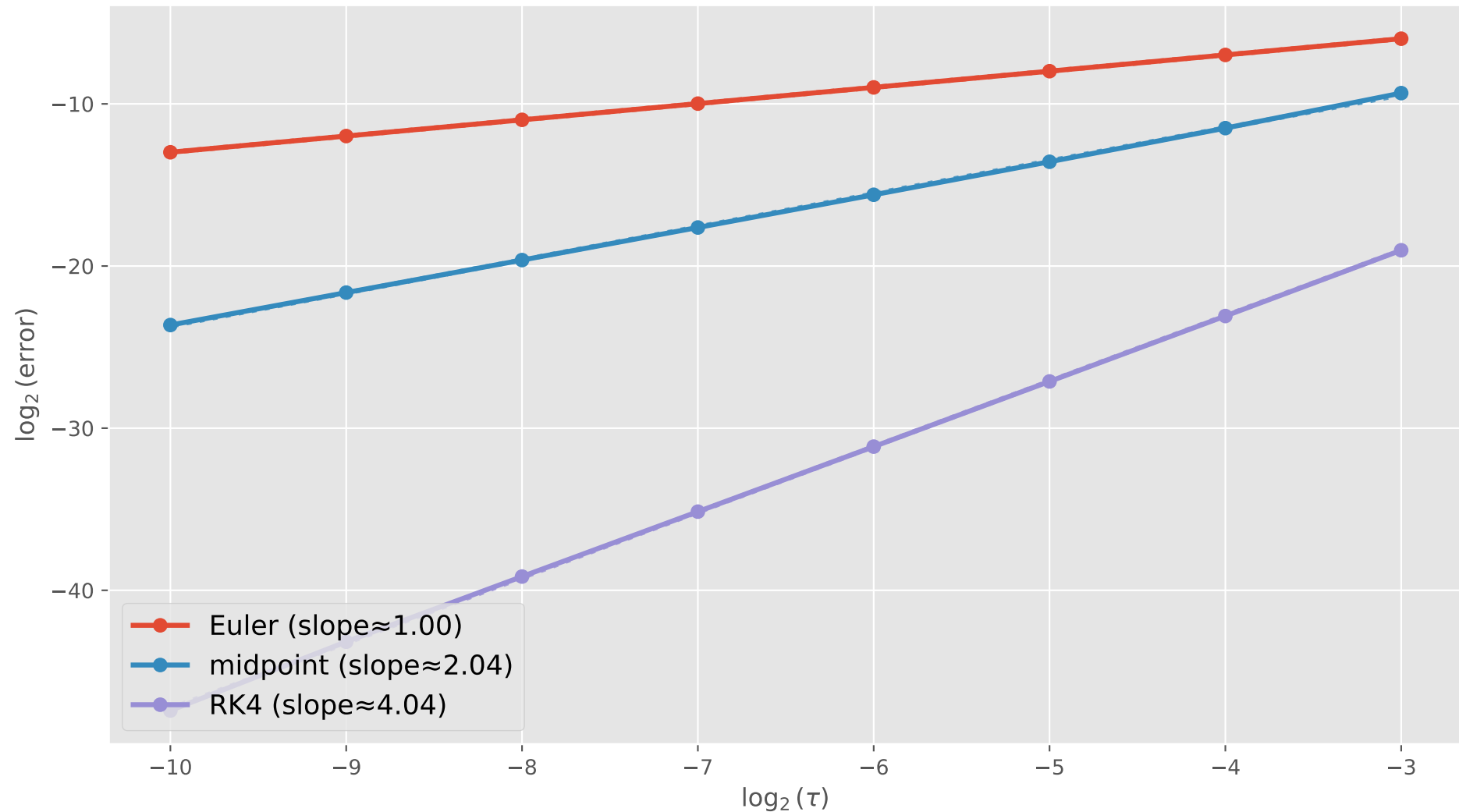
Approximation of the solution for different RK methods



Approximation of the solution for different RK methods

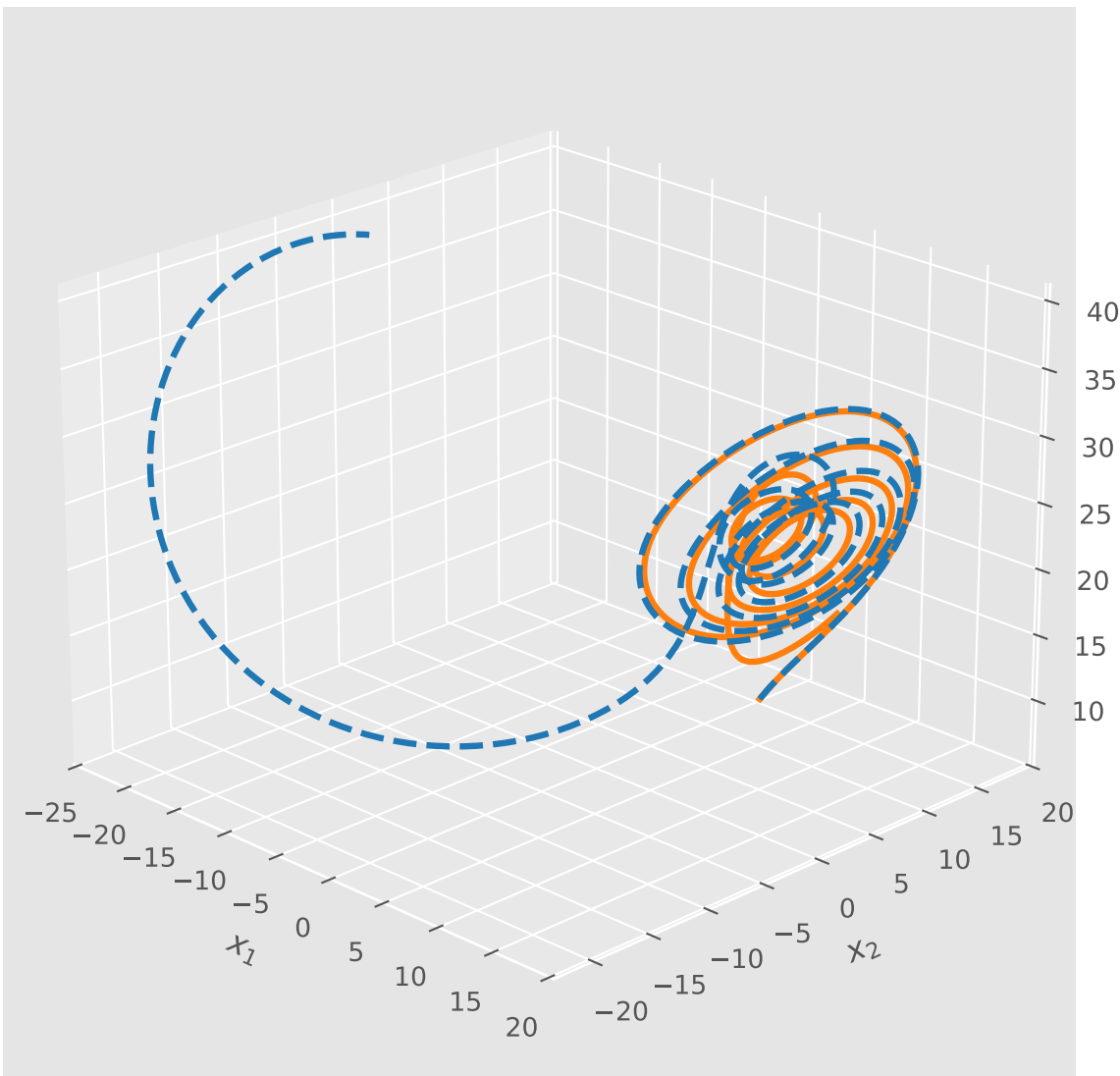


$\log_2, \log_2$  plot of the error for the three different RK methods



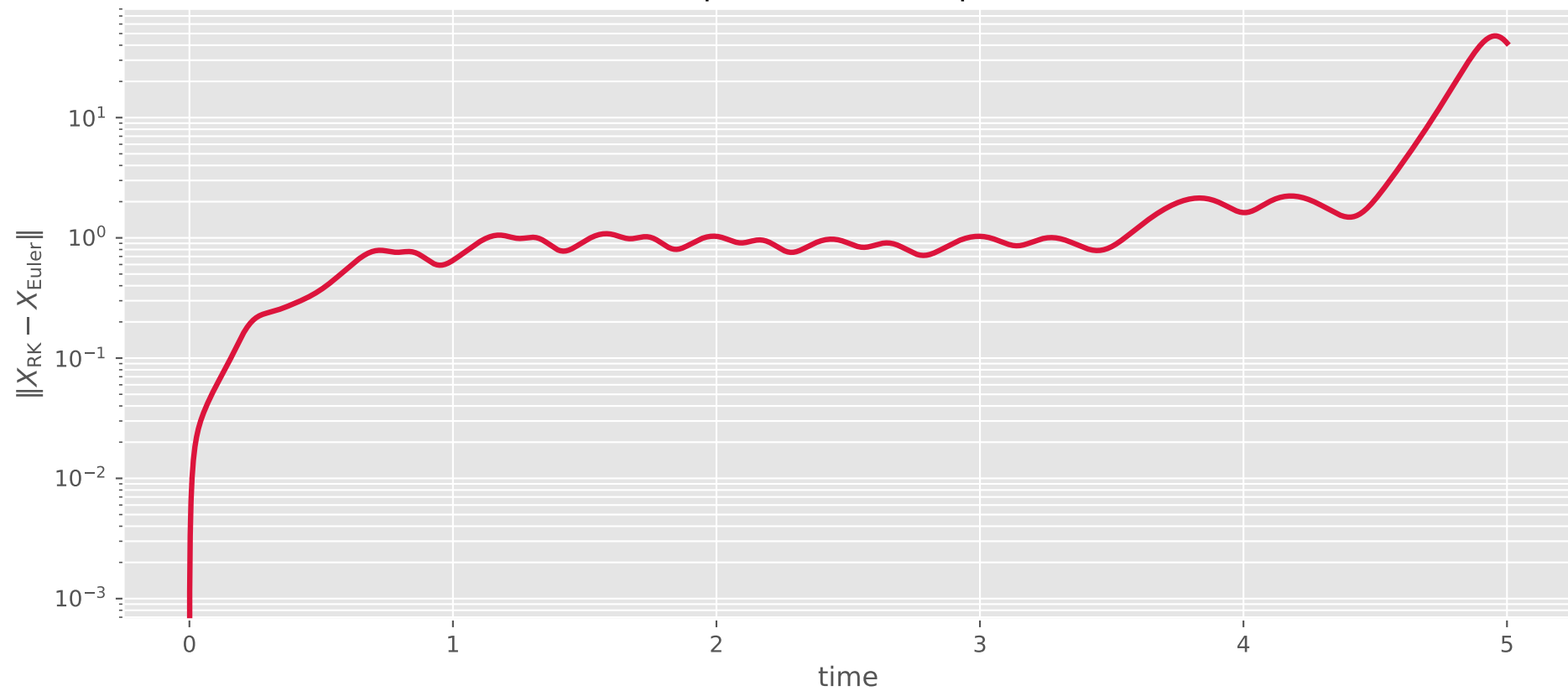


Forced Lorenz trajectories

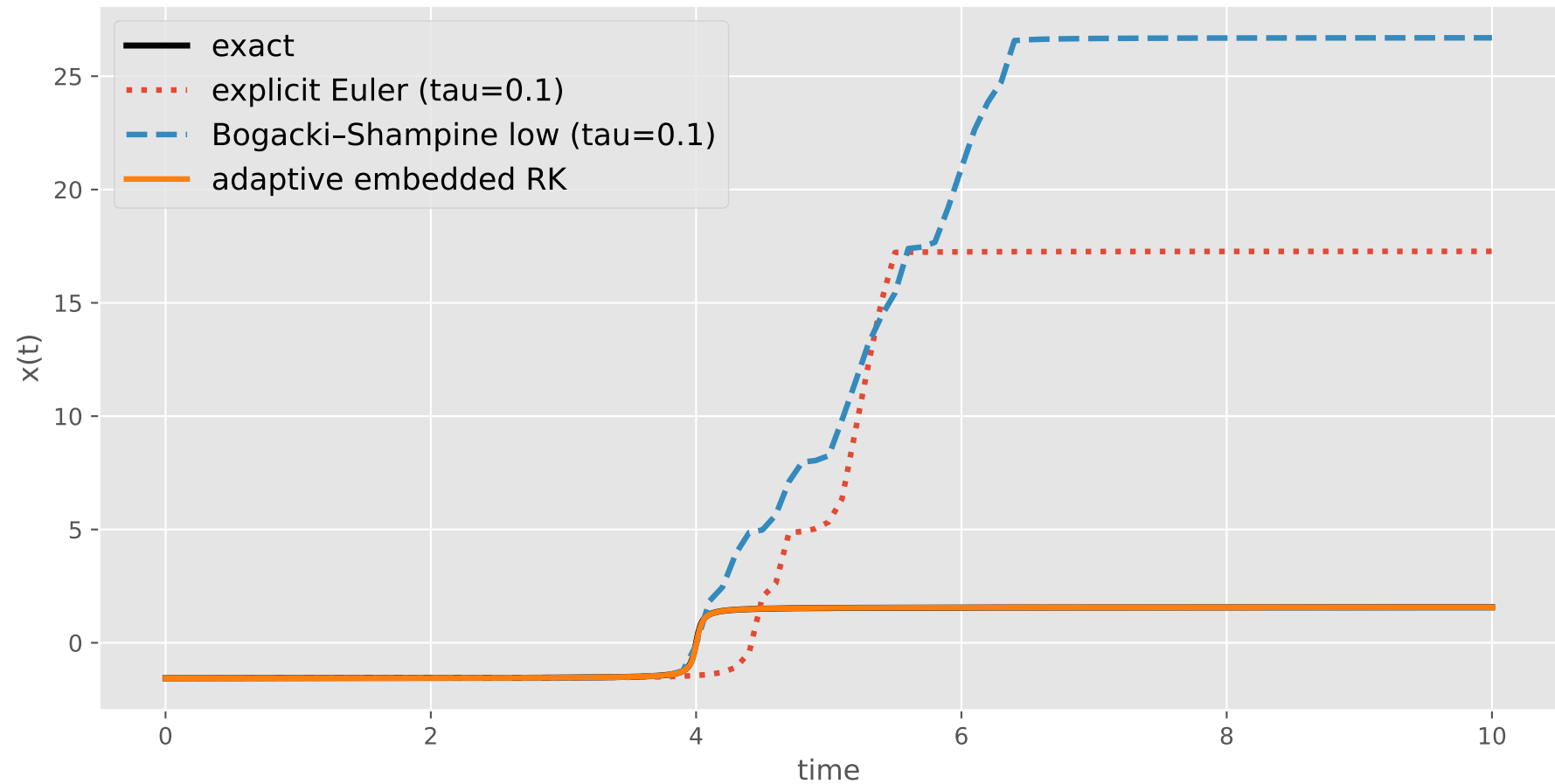


midpoint RK      explicit Euler

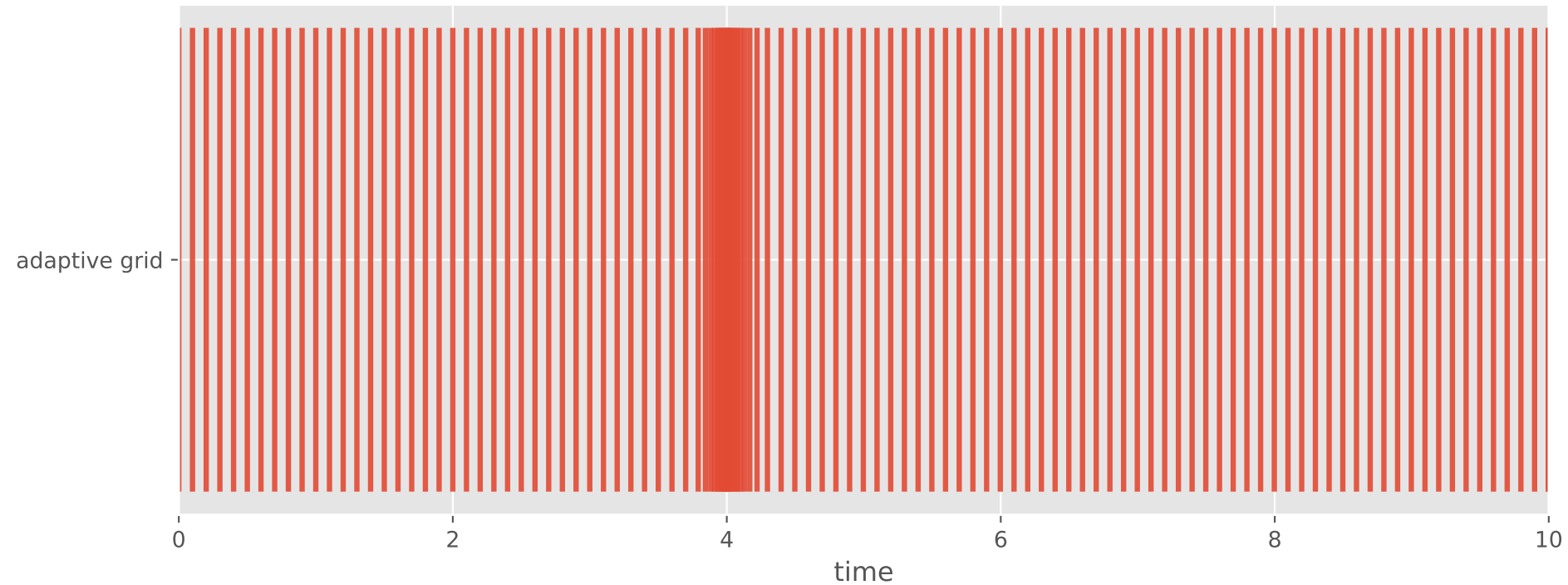
Difference between midpoint RK and explicit Euler (forced Lorenz)



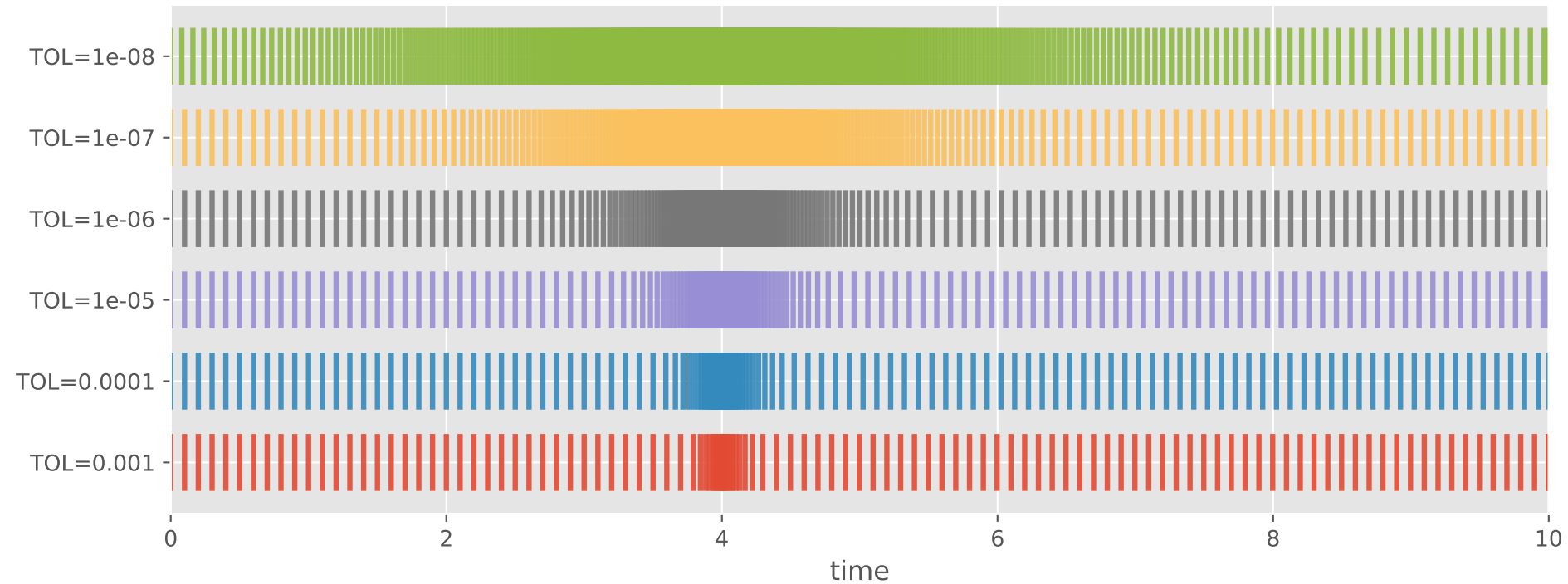
Arctan problem: adaptive BS23 vs references



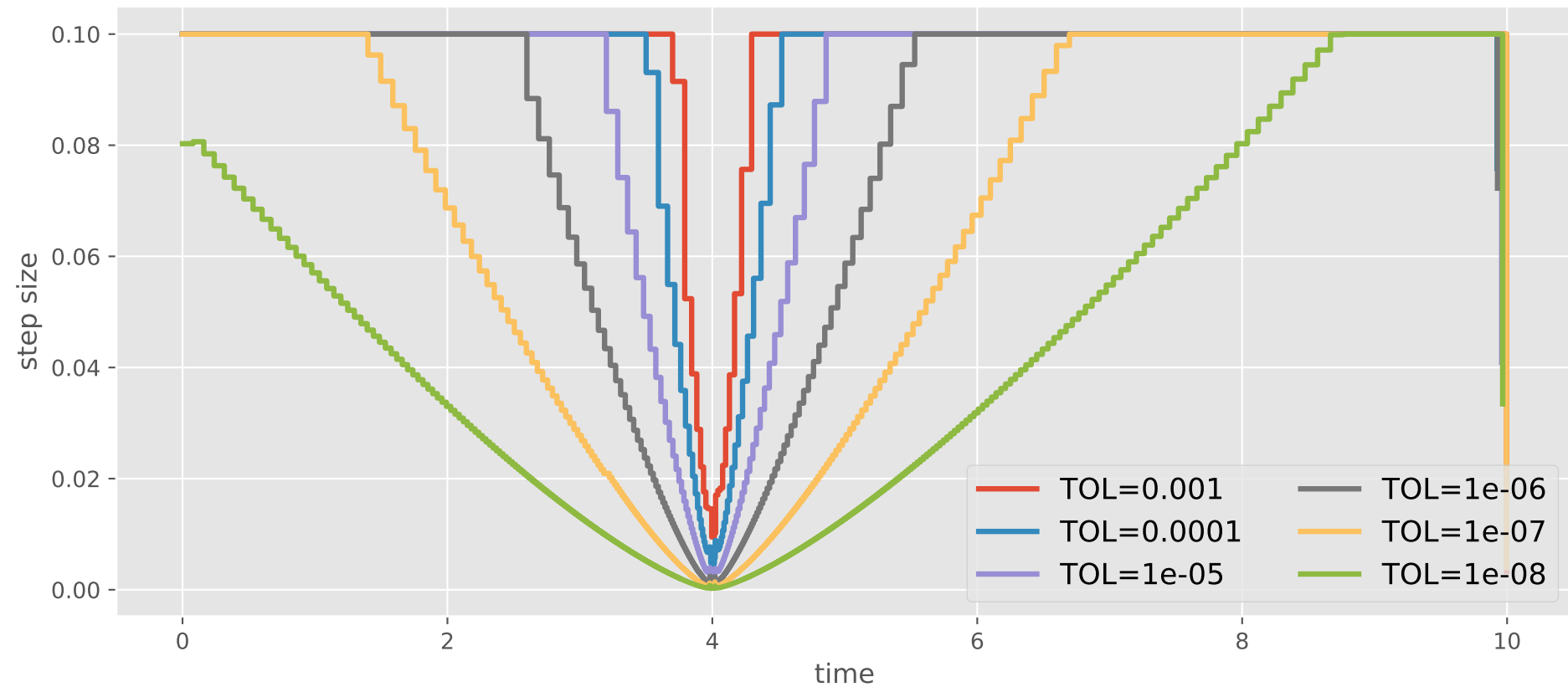
Adaptive discretization points (BS23)



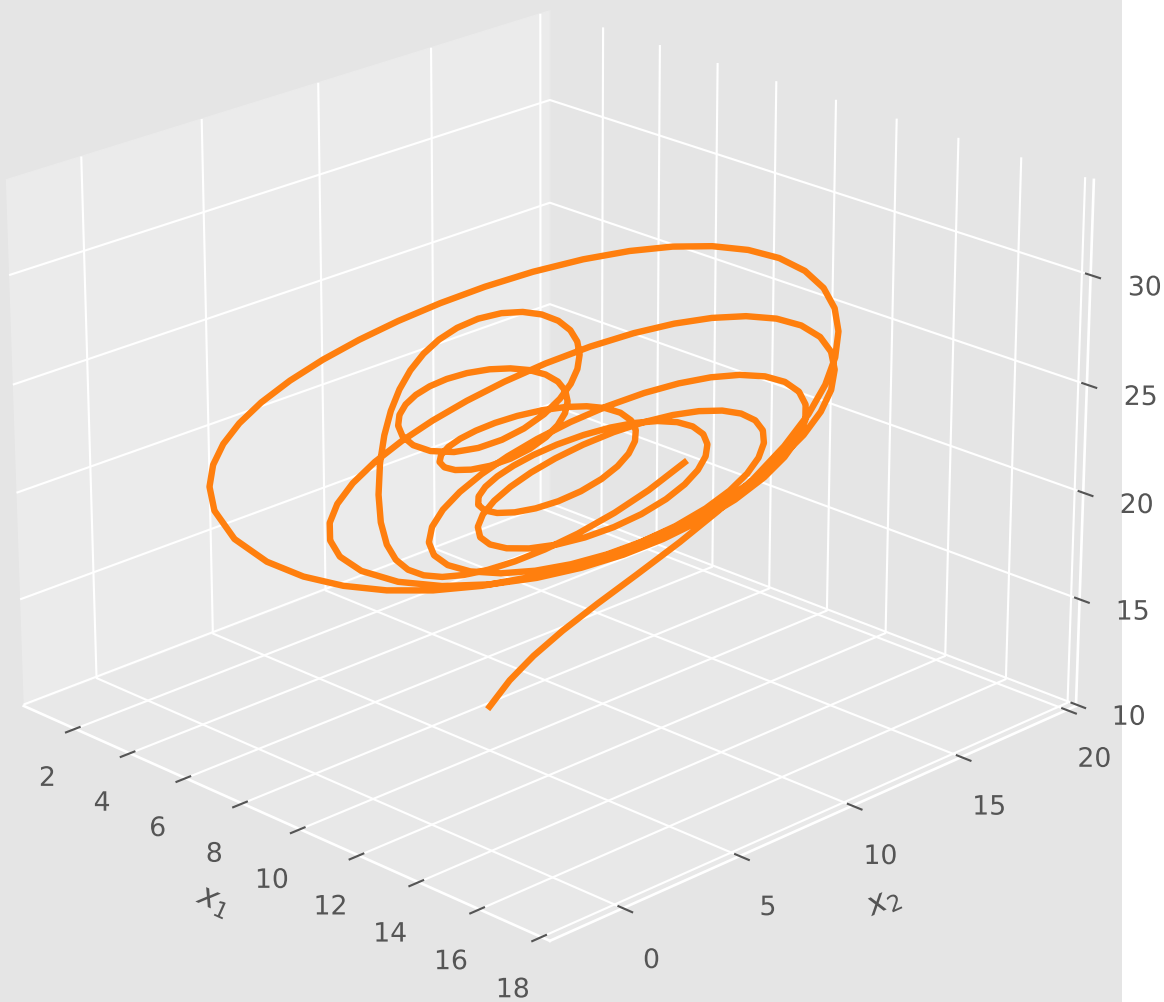
# Adaptive grids for different tolerances



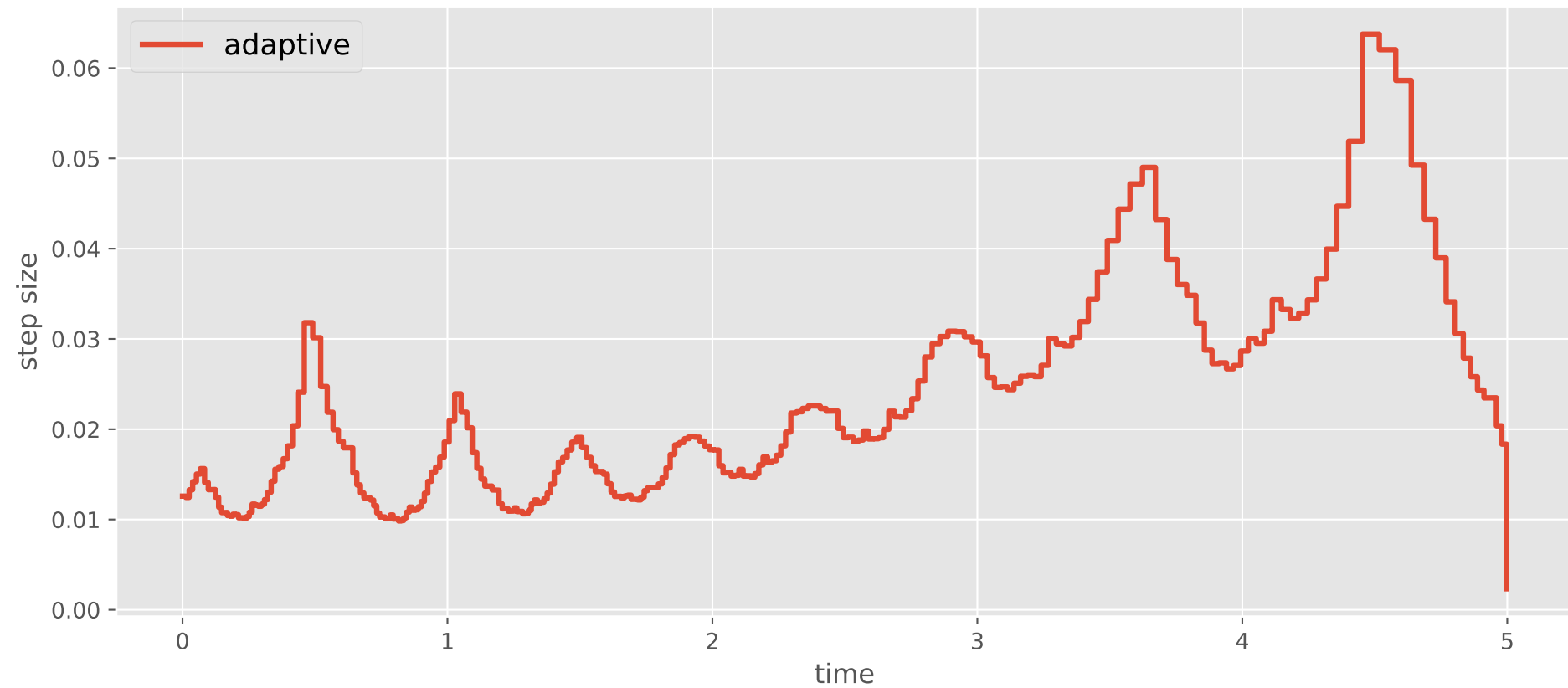
Step sizes over time for different TOL



Forced Lorenz (adaptive BS23)



Adaptive step sizes (Lorenz)





# Answers

(b) Long-term behaviour as a function of  $q$

\textbf{(b) Long-term behaviour as a function of  $q$ .}

We solve  $x'(t)=qx-x^3$  with  $x(0)=2$  and  $q>0$ . The equilibria are  $0$  and  $\pm\sqrt{q}$ . For  $q>0$ ,  $x=0$  is unstable and  $x=\pm\sqrt{q}$  are asymptotically stable since  $f'(x)=q-3x^2$  implies  $f'(\pm\sqrt{q})=-2q<0$ .

With  $x(0)=2>0$ , the trajectory approaches the stable equilibrium  $+\sqrt{q}$ :

```
\[
\begin{cases}
q<4\!: & \sqrt{q}<2 \rightarrow x(t)\ \text{decreases monotonically to}\ \sqrt{q},\ \ [2pt]
q=4\!: & x(t)\equiv 2\ \text{(equilibrium; adding this case gives a flat line at }2),\ \ [2pt]
q>4\!: & \sqrt{q}>2 \rightarrow x(t)\ \text{increases monotonically to}\ \sqrt{q}.
\end{cases}
\]
```

This matches the parameter-sweep plot: for small  $q$  the approach to  $\sqrt{q}$  is slower, so at  $T=10$  the solution can still be slightly above the limiting value.

(c) Method comparison (Euler vs. LSODA) and effect of  $q$

\textbf{(c) Method comparison (Euler vs. LSODA) and effect of  $q$ .}

We compare explicit Euler with step sizes  $\tau=0.1$  and  $\tau=0.01$  against an LSODA reference on  $[0,10]$ .

\emph{Accuracy order.} Explicit Euler is first order: the global error scales as  $\mathcal{O}(\tau)$  for smooth problems on a fixed time horizon. Hence, reducing  $\tau$  from 0.1 to 0.01 should reduce the error by about a factor of 10 (modulo transients).

```
\emph{Linear stability near the attractor.} Linearizing at the stable equilibrium
 $x^*=\sqrt{q}$  gives  $y'=f'(x^*)y=-2qy$ . For the test equation
 $y'=\lambda y$  with  $\lambda=-2q$ , explicit Euler is stable iff
\[
|1-\tau\lambda|<1 \quad \Longleftrightarrow \quad 0<\tau<\frac{1}{q}.
\]
```

\emph{Case  $q=10$ .} The stability bound is  $\tau<0.1$ , so  $\tau=0.1$  lies on the boundary and yields visible phase/amplitude error and mild oscillation around the equilibrium;  $\tau=0.01$  is well inside the stable region and closely tracks LSODA. Empirically, the absolute error curve for  $\tau=0.1$  sits roughly an order of magnitude above that for  $\tau=0.01$  over most of  $[0,10]$ , consistent with first-order convergence \emph{and} the stability-edge effect at  $\tau=0.1$ .

\emph{Case  $q=0.1$ .} The bound is  $\tau<10$ , so both  $\tau=0.1$  and 0.01 are deep inside the stability region and the dynamics are slow. Both Euler solutions lie very close to LSODA; the  $\tau=0.01$  error is still smaller (by about the expected  $\sim 10\times$  factor), but the difference is barely visible in the solution plot because all errors are small.

(d) Sensitivity for the Lorenz system

\textbf{(d) Sensitivity for the Lorenz system.}

With standard parameters  $(a,b,c)=(10,25,8/3)$  the Lorenz system exhibits sensitive dependence on initial conditions (positive largest Lyapunov exponent). We integrate on  $[0,10]$  with explicit Euler ( $\tau=0.001$ ) from  $(x_1(0),x_2(0),x_3(0))=(10,5,12)$  and from the perturbed  $(10,5.01,12)$ .

The two trajectories coincide initially but separate clearly after a short time, ultimately exploring different parts of the attractor. This is the expected behaviour for a chaotic system: for a small perturbation  $\|\delta x(0)\|$  the separation typically grows like  $\|\delta x(t)\| \approx \|\delta x(0)\| e^{\lambda t}$  with  $\lambda>0$ .

\emph{Conclusion.} Yes, the solution changes significantly when  $x_2(0)$  is perturbed to 5.01; the 3D plot makes this divergence clearly visible.