

2 – Discontinuous Galerkin Methods for Flow Problems

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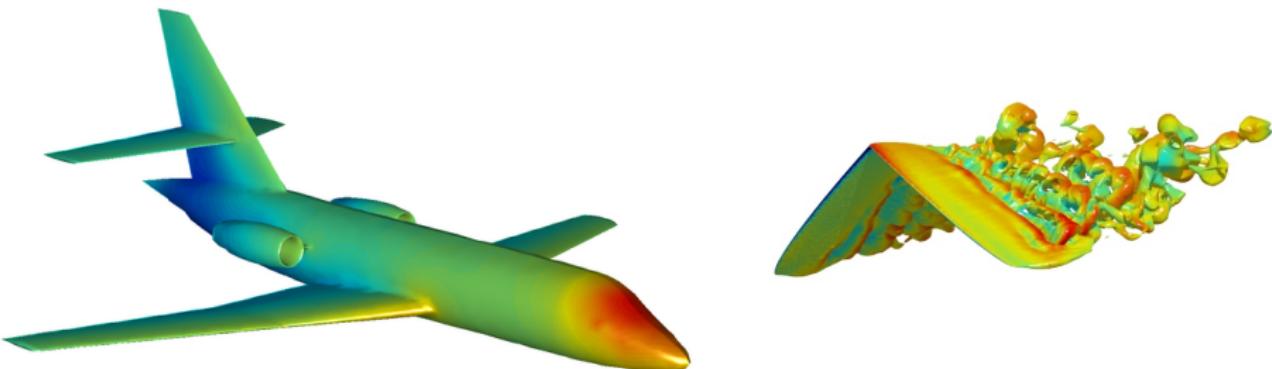


Outline

- 1 Introduction
- 2 The DG Method for Navier-Stokes
 - Application: Implicit Large Eddy Simulation
- 3 Curved Mesh Generation
- 4 Artificial Viscosity and Shock Capturing
 - Application: RAE2822
- 5 ALE for Deforming Domains
 - Application: Vertical Axis Wind Turbines
 - Application: Flapping Flight of Bat

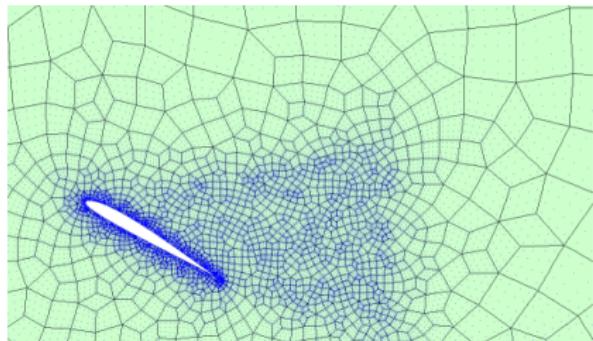
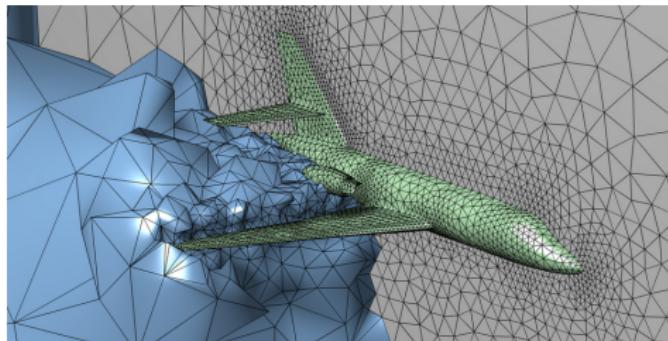
Motivation

- Need for higher fidelity predictions in computational fluid dynamics
 - Turbulent flows, fluid/structure interaction, flapping flight
 - Wave propagation, multiscale phenomena, non-linear interactions
- Widely believed that high-order methods will become the standard in future generations of simulation software
- The DG method has emerged as one of the most promising schemes, since it provides high-order accuracy on fully unstructured meshes



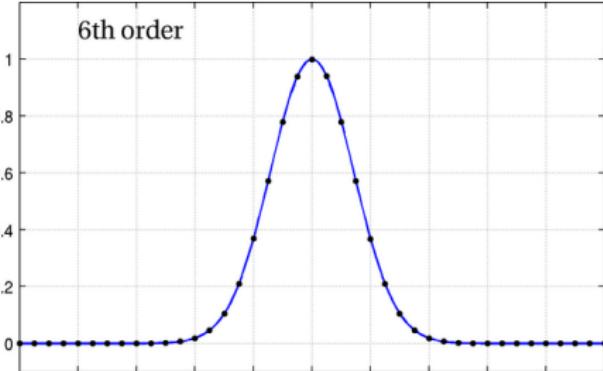
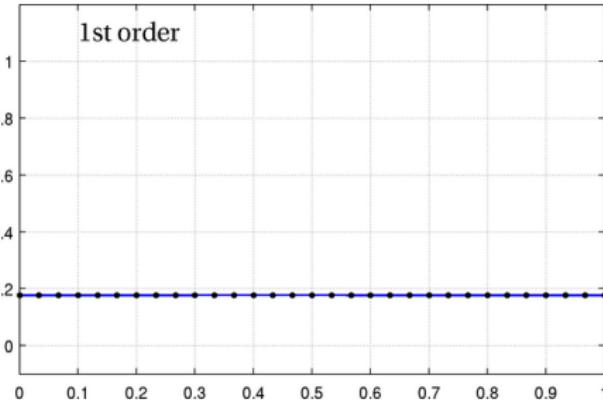
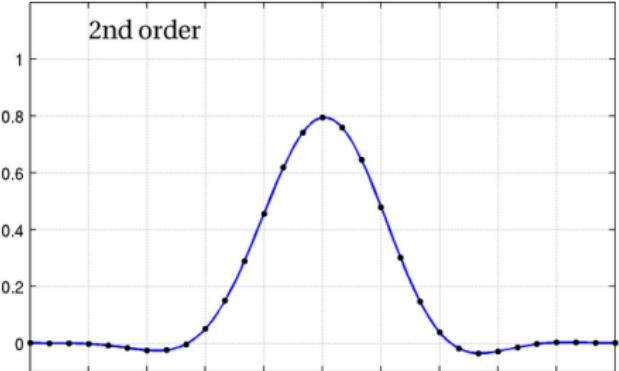
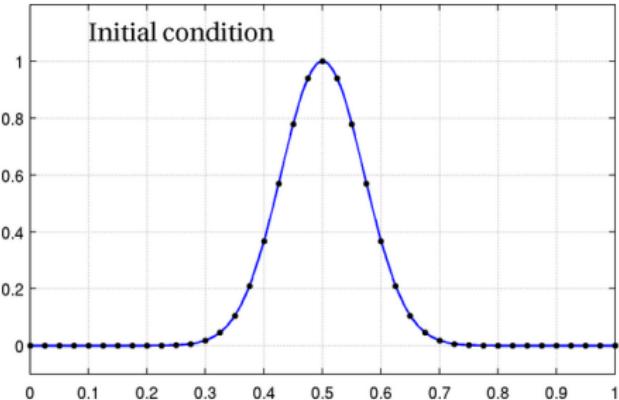
Why Unstructured Meshes?

- Complex *geometries* need flexible element topologies
- Complex *solution fields* need spatially variable resolution
- Fully automated mesh generators for CAD geometries are based on unstructured simplex elements
- Real-world simulation software dominated by unstructured mesh discretization schemes



Why high-order accurate methods?

- Scalar convection equation $u_t + u_x = 0$
- High-order gives *superior performance for equal resolution*



Numerical Schemes for Flow Problems

- Complex flow features → High-order accuracy required
- Complex geometries, adaptation → Unstructured grids required
- Discontinuous Galerkin (DG) methods have these properties:

	FVM	FDM	FEM	DG
1) High-order/Low dispersion	✗	✓	✓	✓
2) Unstructured meshes	✓	✗	✓	✓
3) Stability for conservation laws	✓	✓	✗	✓

- However, still several problems to resolve:
 - High CPU/memory requirements (compared to FVM or H-O FDM)
 - Low tolerance to under-resolved features
 - High-order geometry representation and mesh generation

Governing Equations

- The (compressible) *Navier-Stokes equations*:

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_i}(\rho u_i) = 0,$$

$$\frac{\partial}{\partial t}(\rho u_i) + \frac{\partial}{\partial x_i}(\rho u_i u_j + p) = + \frac{\partial \tau_{ij}}{\partial x_j} \quad \text{for } i = 1, 2, 3,$$

$$\frac{\partial}{\partial t}(\rho E) + \frac{\partial}{\partial x_i}(u_j(\rho E + p)) = - \frac{\partial q_j}{\partial x_j} + \frac{\partial}{\partial x_j}(u_j \tau_{ij}),$$

with

$$\tau_{ij} = \mu \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} - \frac{2}{3} \frac{\partial u_k}{\partial x_j} \delta_{ij} \right), \quad q_j = - \frac{\mu}{\Pr} \frac{\partial}{\partial x_j} \left(E + \frac{p}{\rho} - \frac{1}{2} u_k u_k \right)$$

$$p = (\gamma - 1) \rho \left(E - \frac{1}{2} u_k u_k \right)$$

- When $\mu = 0$, reduces to first-order system – *Euler's equations of gas dynamics*

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The Discontinuous Galerkin Method

- (Reed/Hill 1973, Lesaint/Raviart 1974, Cockburn/Shu 1989-, etc)
- Write the first-order equations as a system of conservation laws:

$$\mathbf{u}_t + \nabla \cdot \mathbf{F}_i(\mathbf{u}) = 0$$

- Triangulate domain Ω into elements $\kappa \in \mathcal{T}_h$
- Seek approximate solution \mathbf{u}_h in space of element-wise polynomials:

$$V_h^p = \{\mathbf{v} \in L^2(\Omega) : \mathbf{v}|_{\kappa} \in P^p(\kappa) \ \forall \kappa \in \mathcal{T}_h\}$$

- Multiply by test function $\mathbf{v}_h \in V_h^p$ and integrate over element κ :

$$\int_{\kappa} [(\mathbf{u}_h)_t + \nabla \cdot \mathbf{F}_i(\mathbf{u}_h)] \mathbf{v}_h \, d\mathbf{x} = 0$$

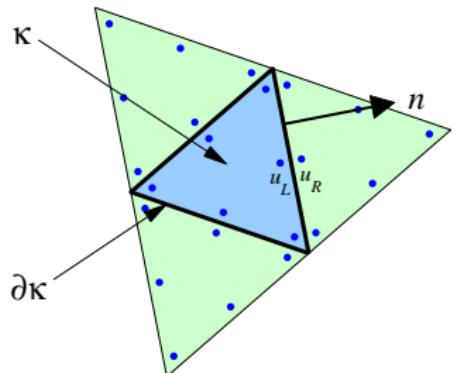
The Discontinuous Galerkin Method

- Integrate by parts:

$$\int_{\kappa} [(\mathbf{u}_h)_t] \mathbf{v}_h \, d\mathbf{x} - \int_{\kappa} \mathbf{F}_i(\mathbf{u}_h) \nabla \mathbf{v}_h \, d\mathbf{x} + \int_{\partial\kappa} \hat{\mathbf{F}}_i(\mathbf{u}_h^+, \mathbf{u}_h^-, \hat{\mathbf{n}}) \mathbf{v}_h^+ \, ds = 0$$

with numerical flux function $\hat{\mathbf{F}}_i(\mathbf{u}_L, \mathbf{u}_R, \hat{\mathbf{n}})$ for left/right states $\mathbf{u}_L, \mathbf{u}_R$ in direction $\hat{\mathbf{n}}$ (Godunov, Roe, Osher, Van Leer, Lax-Friedrichs, etc)

- Global problem: Find $\mathbf{u}_h \in V_h^p$ such that this weighted residual is zero for all $\mathbf{v}_h \in V_h^p$
- Error = $\mathcal{O}(h^{p+1})$ for smooth solutions



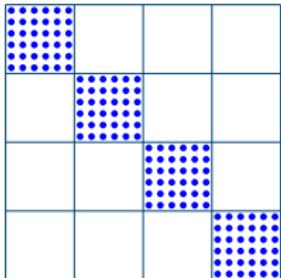
The DG Method – Observations

- Reduces to the finite volume method for $p = 0$:

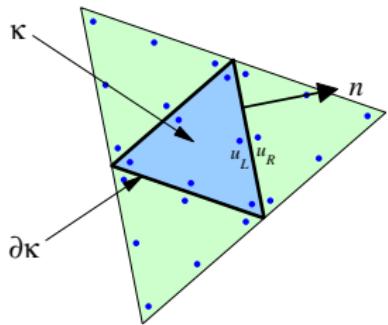
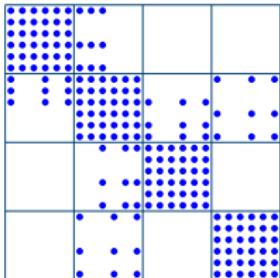
$$(\mathbf{u}_h)_t A_\kappa + \int_{\partial\kappa} \hat{\mathbf{F}}_i(\mathbf{u}_h^+, \mathbf{u}_h^-, \hat{\mathbf{n}}) ds = 0$$

- Boundary conditions enforced naturally for any degree p
- Block-diagonal mass matrix (no overlap between basis functions)
- Block-wise compact stencil – neighboring elements connected

Mass Matrix



Jacobian



Viscous Discretization

- Write equations as system of first order equations:

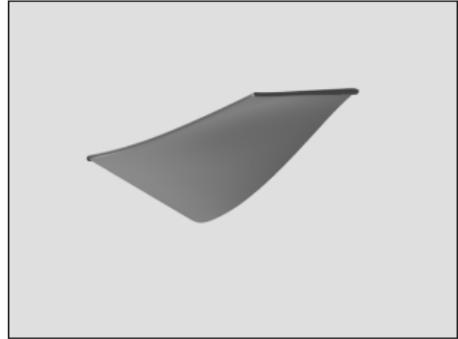
$$\mathbf{u}_t + \nabla \cdot \mathbf{F}_i(\mathbf{u}) - \nabla \cdot \mathbf{F}_v(\mathbf{u}, \boldsymbol{\sigma}) = 0$$

$$\boldsymbol{\sigma} - \nabla \mathbf{u} = 0$$

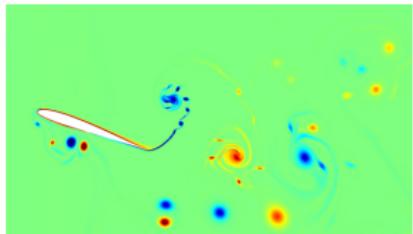
- Discretize using DG, choose appropriate numerical fluxes:
 - For inviscid component \mathbf{F}_i , use approximate Riemann solvers (Godunov, Roe, Osher, Van Leer, Lax-Friedrichs, etc)
 - For viscous component \mathbf{F}_v , use numerical fluxes $\hat{\boldsymbol{\sigma}}$, $\hat{\mathbf{u}}$ according to methods such as IP/BR2/LDG/CDG

Implementation: The 3DG Software Package

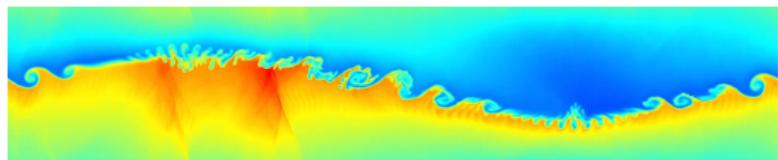
- High-order discretizations on unstructured meshes
- Optimized C++ code with MATLAB and Python interfaces
- Capable of simulating challenging problems:
 - complex real-world geometries
 - transitional flows, multiple scales
 - moving and deforming domains
 - fluid-structure interactions
- General multiphysics framework applicable to a wide range of challenging problems



Thin Structures



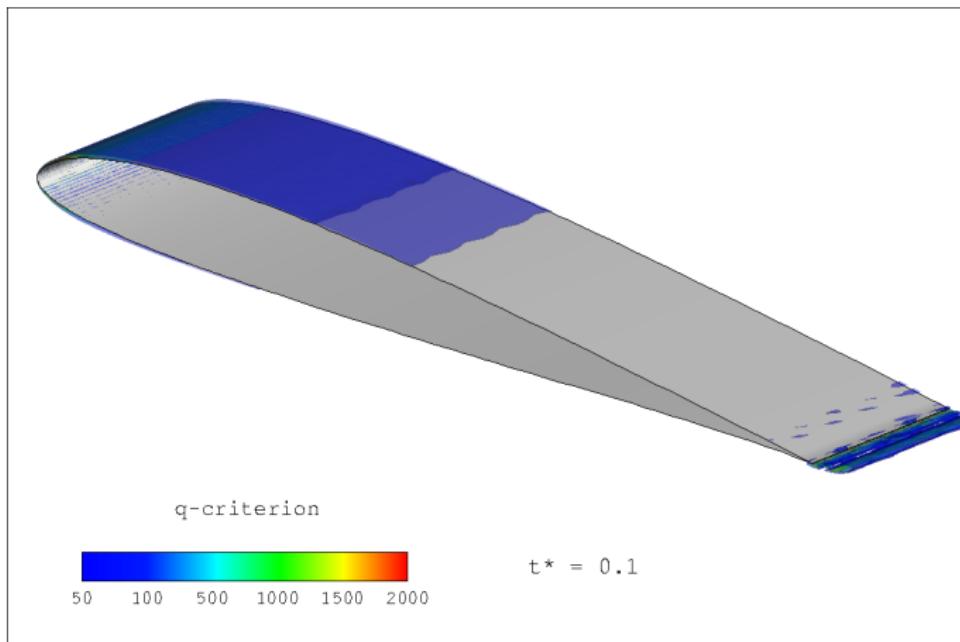
Unsteady Flows



Aeroacoustics

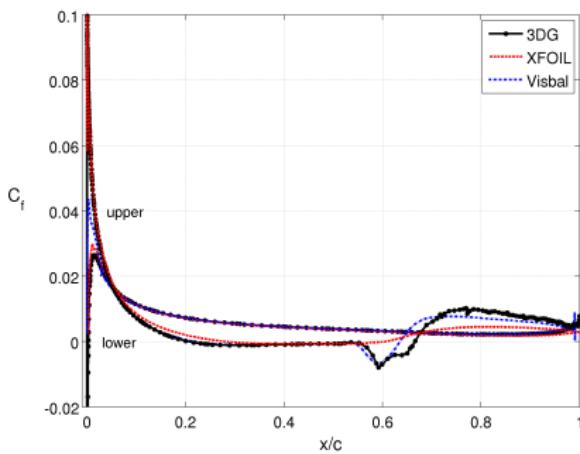
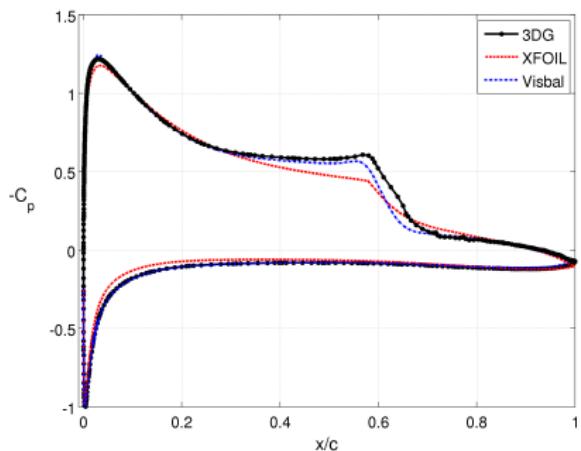
Example: ILES at $Re = 60,000$

- Implicit Large Eddy Simulations for flow past airfoil
- Separation and transition well captured
- Vortical structures: iso-surfaces of q-criterion ($\nabla^2 p / 2\rho$)



Example: ILES at $Re = 60,000$

- Good agreement with XFOIL and previously published ILES
[Uranga/Persson/Drela/Peraire '11]



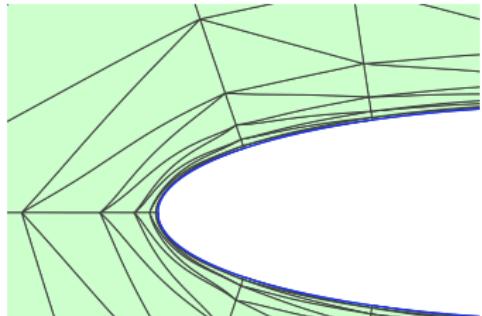
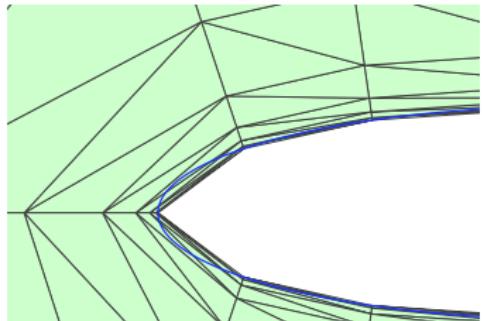
Average pressure and skin friction coefficients

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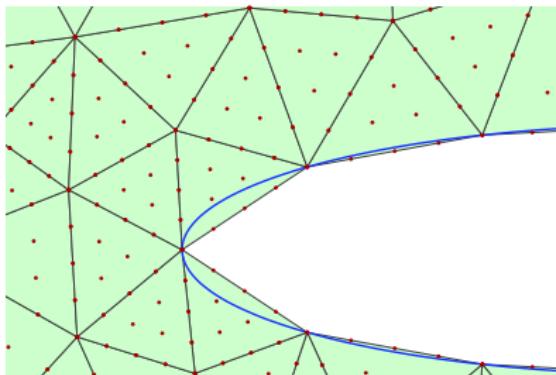
Curved Mesh Generation

- Automatic generation of non-inverted curved elements largely an unresolved problem
- In general this is a global problem, affecting many elements except for simple isotropic 2-D meshes
- In [Persson/Peraire '09], we proposed a *non-linear solid mechanics* approach, where the mesh is considered an elastic deformable solid

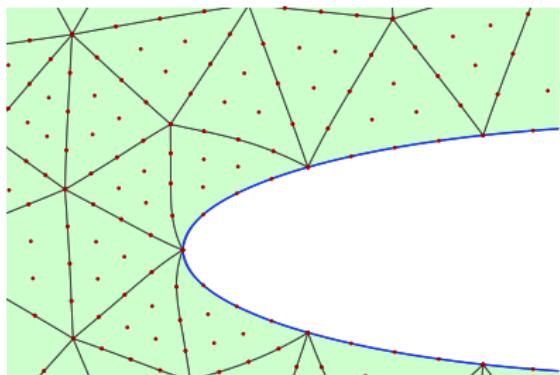


Curved Mesh Generation using Solid Mechanics

- The initial, straight-sided mesh corresponds to undeformed solid
- External forces come from the true boundary data
- Solving for a force equilibrium gives the deformed, curved, boundary conforming mesh
- Bottom-up approach can be used to obtain the boundary data



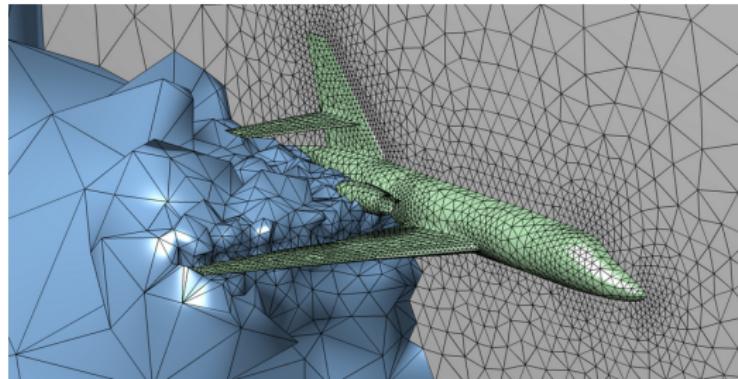
Reference domain, initial configuration



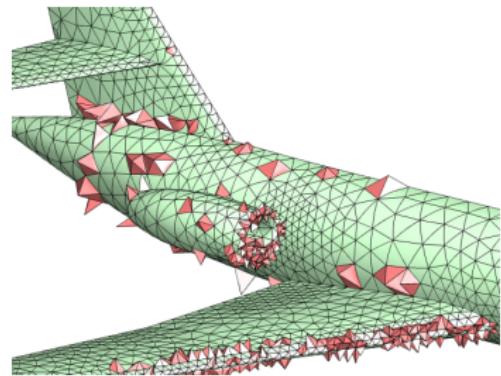
Equilibrium solution, final curved mesh

Tetrahedral Mesh of Falcon Aircraft

- Real-world mesh with coarse but realistic elements
- Unstructured Delaunay refinement mesh, with highly curved boundary segments
- Many elements would invert with a local element-wise approach



Tetrahedral mesh



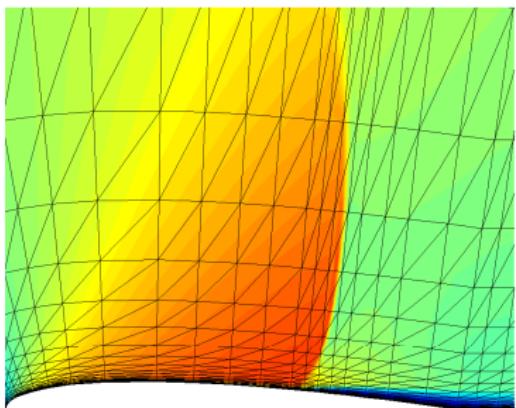
Elements with $I < 0.5$

Outline

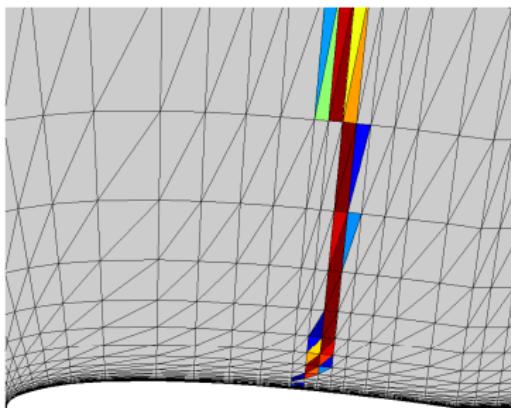
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Artificial Viscosity for Underresolved Features

- Cannot resolve all solution features (shocks, RANS, singularities)
- Low dissipation makes DG sensitive to underresolution
- Detect by sensors and add viscosity [Persson/Peraire 06,07]
- Enables shock capturing with sub-cell resolution and robust solution of Spalart-Alamaras RANS model



Mach



Sensor

Shock Sensor

- Regularity of solution determined from the decay rate of expansion coefficients in orthogonal basis
- Example: Periodic Fourier case: $f(x) = \sum_{k=-\infty}^{\infty} g_k e^{ikx}$
If $f(x)$ has m continuous derivatives $\rightarrow |g_k| \sim k^{-(m+1)}$
- For simplices: Expand solution in orthonormal Koornwinder basis:

$$u = \sum_{i=1}^{N(p)} u_i \psi_i, \quad \hat{u} = \sum_{i=1}^{N(p-1)} u_i \psi_i, \quad s_e = \log_{10} \left(\frac{(u - \hat{u}, u - \hat{u})_e}{(u, u)_e} \right)$$

- Determine elemental piecewise constant ε_e

$$\varepsilon_e = \begin{cases} 0 & \text{if } s_e < s_0 - \kappa \\ \frac{\varepsilon_0}{2} \left(1 + \sin \frac{\pi(s_e - s_0)}{2\kappa} \right) & \text{if } s_0 - \kappa \leq s_e \leq s_0 + \kappa \\ \varepsilon_0 & \text{if } s_e > s_0 + \kappa \end{cases} .$$

where $\varepsilon_0 \sim h/p$, $s_0 \sim 1/p^4$ and κ empirical

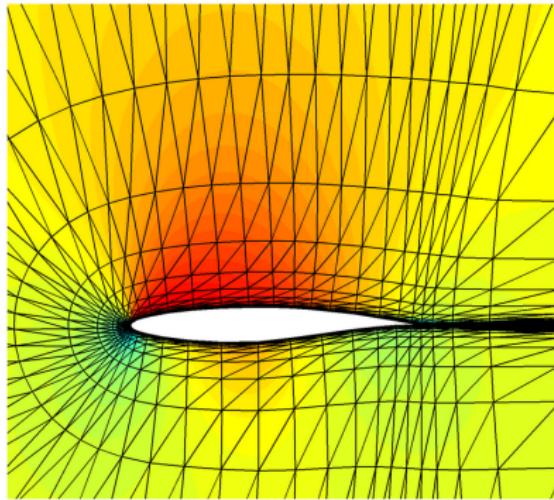
Example: RAE2822

- Turbulent RANS flow ($M = 0.675, \alpha = 2.31^\circ, \text{Re} = 6.5 \cdot 10^6$)
- p -converged solution, fixed resolution h/p

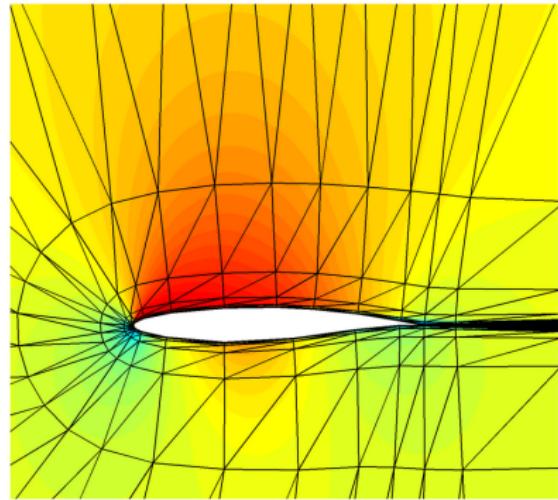
$p = 2$

(constant h/p)

$p = 4$



$$C_L = 0.6144 \quad C_D = 0.0104$$



$$C_L = 0.6131 \quad C_D = 0.0103$$

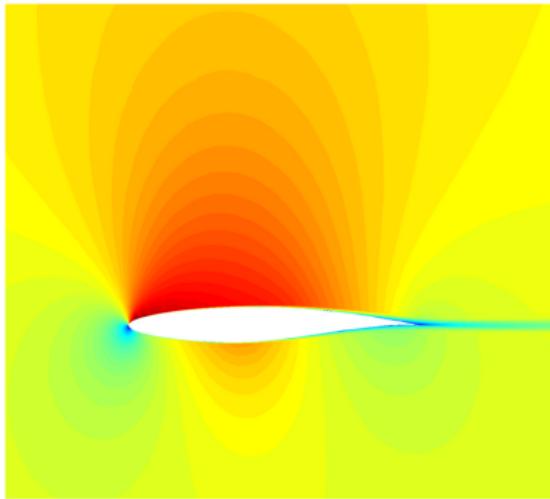
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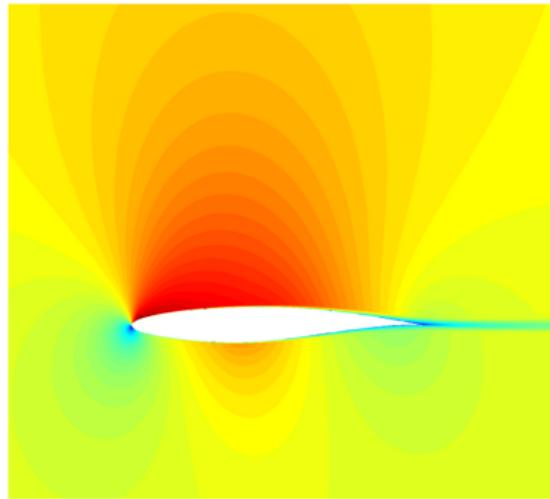
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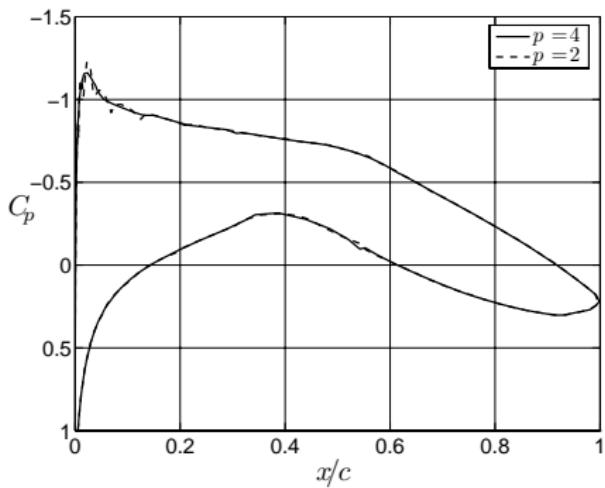


$$C_L = 0.6131 \quad C_D = 0.0103$$

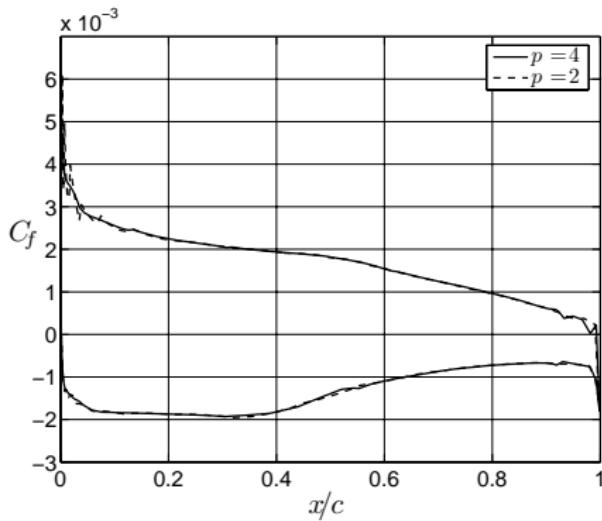
Example: RAE2822

- Highly accurate boundary forces even with coarse meshes

C_p

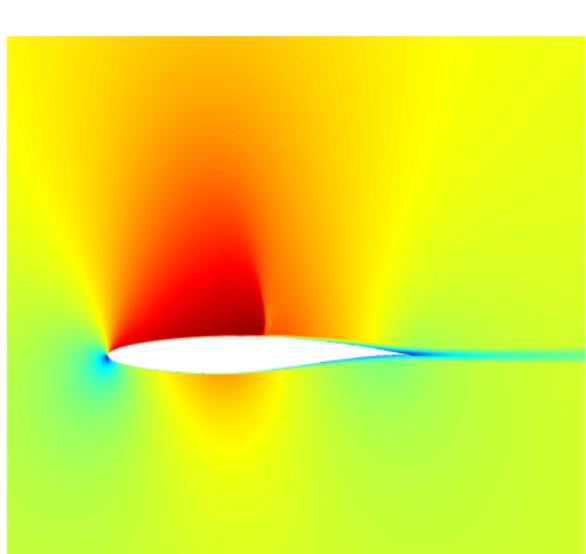


C_f

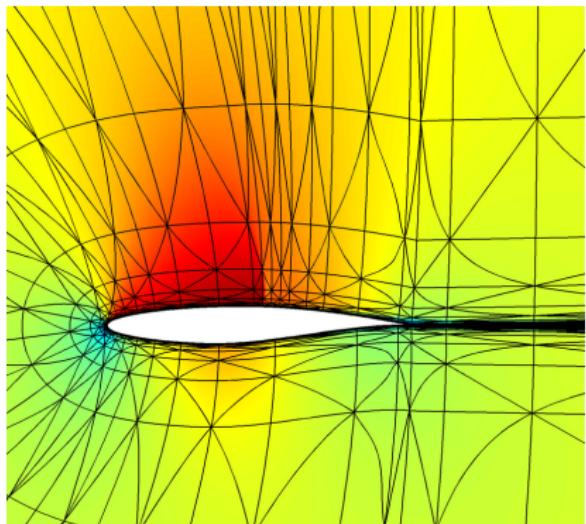


Example: RAE2822, Transonic

- Transonic flow ($M = 0.729$, $\text{Re} = 6.5 \cdot 10^6$)
- Sub-cell resolution of shocks

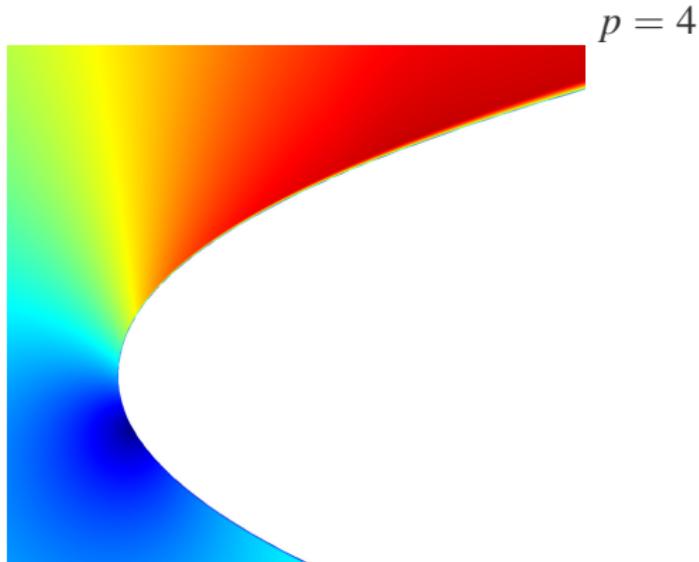


$p = 4$

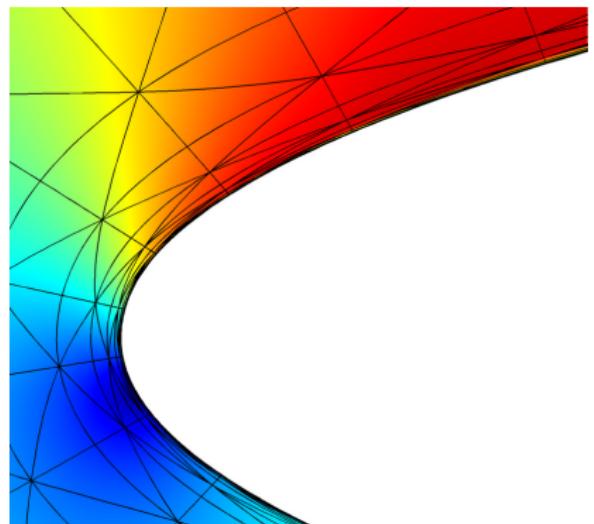


Example: RAE2822, Transonic

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$p = 4$



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ALE Formulation for Deforming Domains

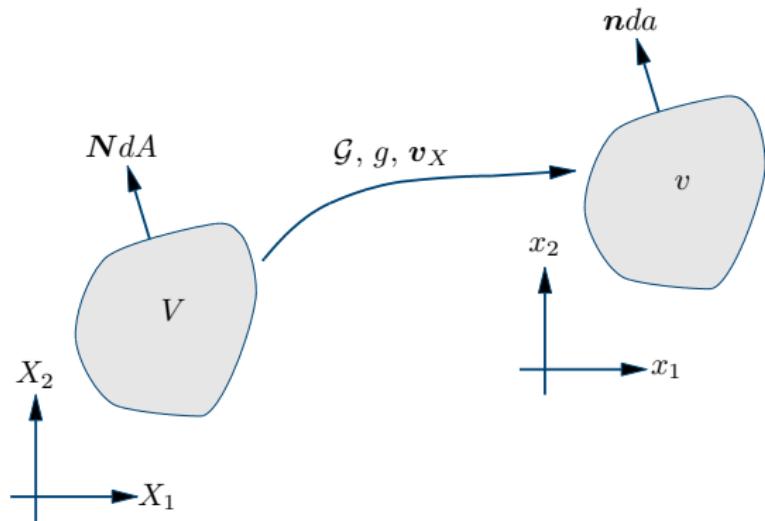
- Use mapping-based ALE formulation for moving domains
[Visbal/Gaitonde '02], [Persson/Bonet/Peraire '09]
- Map from reference domain V to physical deformable domain $v(t)$
- Introduce the *mapping deformation gradient* \mathcal{G} and the *mapping velocity* \mathbf{v}_X as

$$\mathcal{G} = \nabla_X \mathcal{G}$$

$$\mathbf{v}_X = \frac{\partial \mathcal{G}}{\partial t} \Big|_X$$

and set $g = \det(\mathcal{G})$

- Transform equations to account for the motion



Transformed Equations

- The system of conservation laws in the physical domain $v(t)$

$$\frac{\partial \mathbf{U}_x}{\partial t} \Big|_x + \nabla_x \cdot \mathbf{F}_x(\mathbf{U}_x, \nabla_x \mathbf{U}_x) = 0$$

can be written in the reference configuration V as

$$\frac{\partial \mathbf{U}_X}{\partial t} \Big|_X + \nabla_X \cdot \mathbf{F}_X(\mathbf{U}_X, \nabla_X \mathbf{U}_X) = 0$$

where

$$\mathbf{U}_X = g \mathbf{U}_x, \quad \mathbf{F}_X = g \mathbf{G}^{-1} \mathbf{F}_x - \mathbf{U}_X \mathbf{G}^{-1} \mathbf{v}_X$$

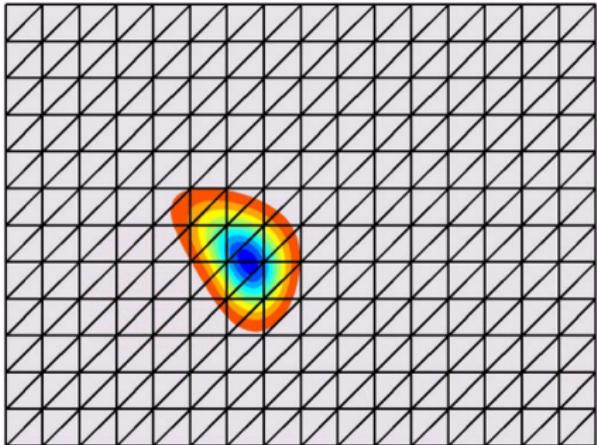
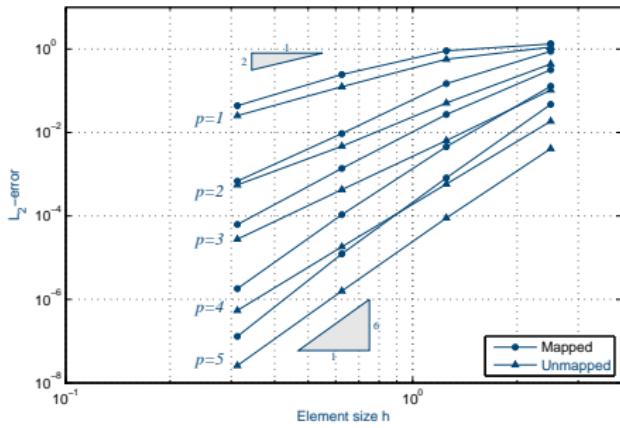
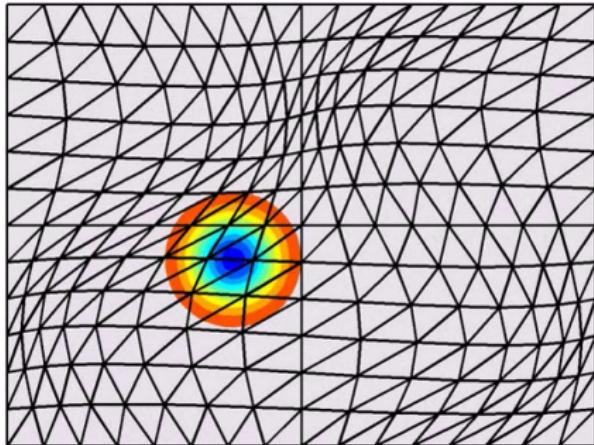
and

$$\nabla_x \mathbf{U}_x = \nabla_X(g^{-1} \mathbf{U}_X) \mathbf{G}^{-T} = (g^{-1} \nabla_X \mathbf{U}_X - \mathbf{U}_X \nabla_X(g^{-1})) \mathbf{G}^{-T}$$

- Details in [Persson/Bonet/Peraire '09], including how to satisfy the so-called Geometric Conservation Law (GCL)

ALE Formulation for Deforming Domains

- Mapping-based formulation gives arbitrarily high-order accuracy in space and time



Vertical Axis Wind Turbines

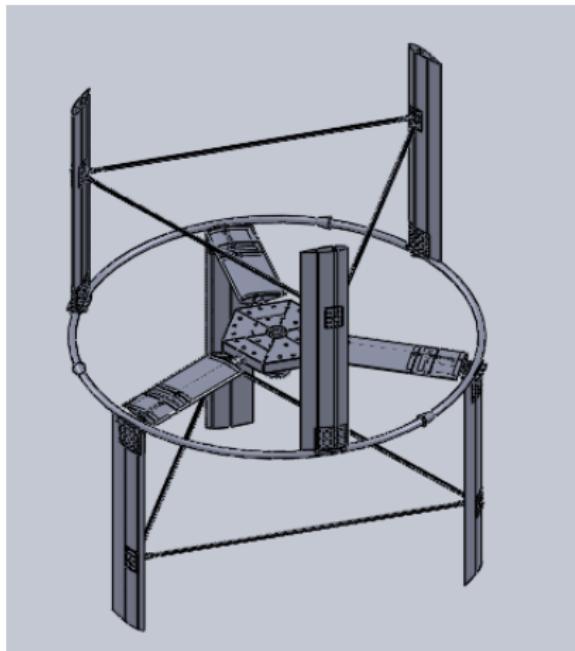
- Recent interest in vertical axis wind turbines (VAWT):
 - 2D airfoils, easy to manufacture, supportable at both ends
 - Omnidirectional (good in gusty, low wind, e.g. close to ground)
 - Lower blade speeds – lower noise and impact
 - Can be packed close together in wind farms
- Numerical simulations can help overcome remaining challenges:
 - Lower theoretical (and practical) efficiency than HAWTs
 - Sensitive to design conditions
 - Structural problems, fatigue and catastrophic failure



Windterra ECO 1200 1Kw VAWT

Vertical Axis Wind Turbines

- Experimental design by G. Dahlbacka (LBNL) and collaborators
- 3kW unit, CAD design (left) and assembled unit (right)

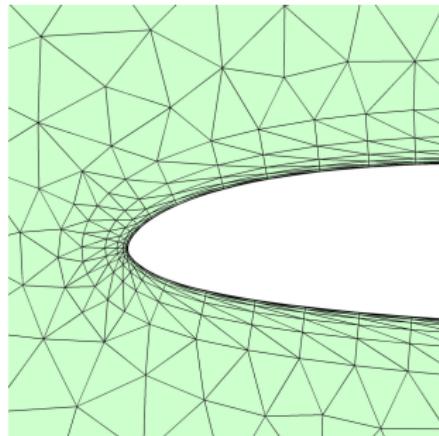
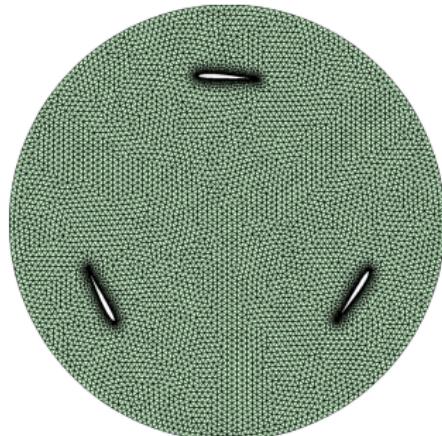


VAWT – Mathematical Model and Discretization

- Preliminary 2D simulation, using vertical symmetry
- Solve the Navier-Stokes equations in a rotating frame:

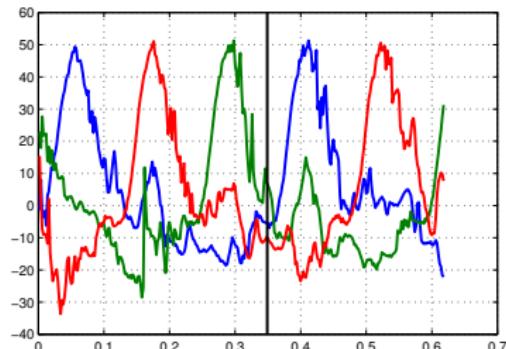
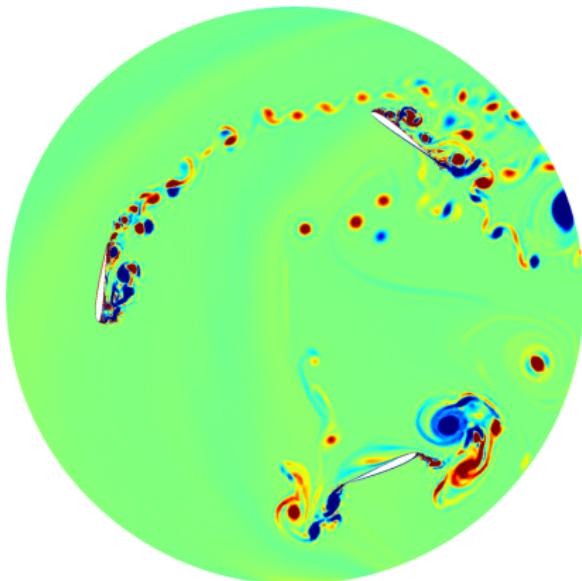
$$\mathcal{G}(X, Y, t) = \begin{bmatrix} \cos \omega t & -\sin \omega t \\ \sin \omega t & \cos \omega t \end{bmatrix} \begin{bmatrix} X \\ Y \end{bmatrix}$$

- Hybrid boundary layer/unstructured mesh, element degree $p = 3$



VAWT – Numerical Results

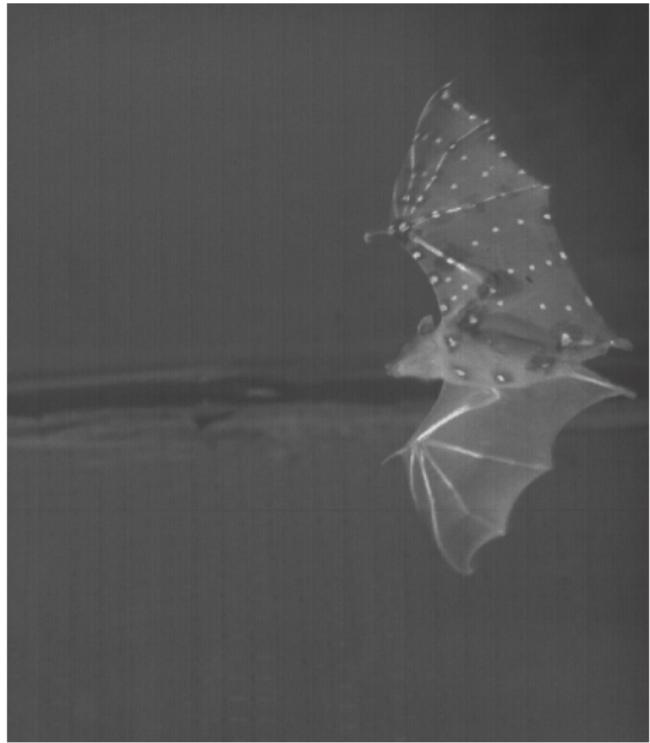
- Simulation with freestream wind speed 12 m/s (horizontally, from the left) and wing tip speed ratio 1.5
- Visualization by z -vorticity in rotating frame



Moment on each blade vs. time

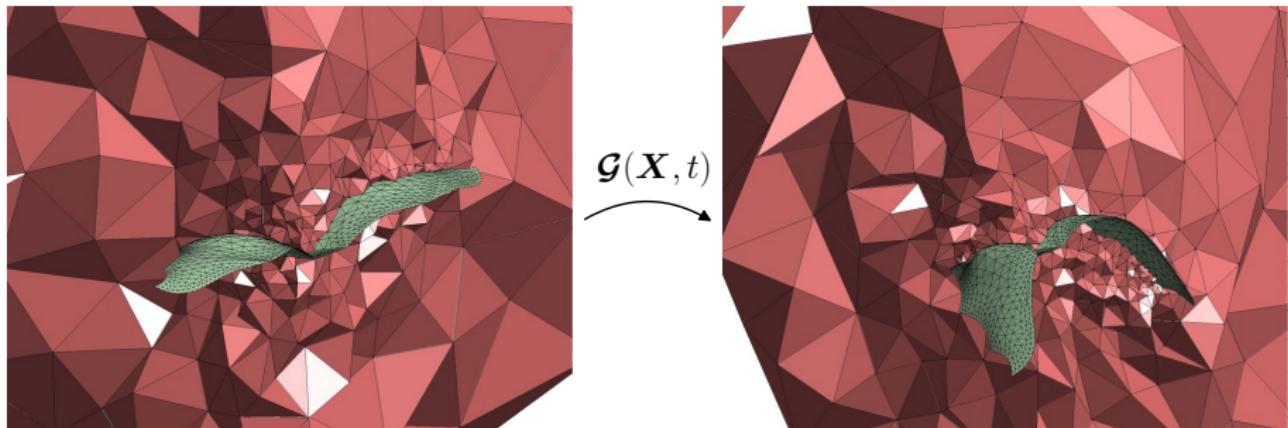
Bio-Inspiration for Flapping Wing MAVs

- Develop high-order accurate simulation capabilities that capture the complex physics in flapping flight
- Use the computational tools for increased understanding and to design optimized flapping kinematics



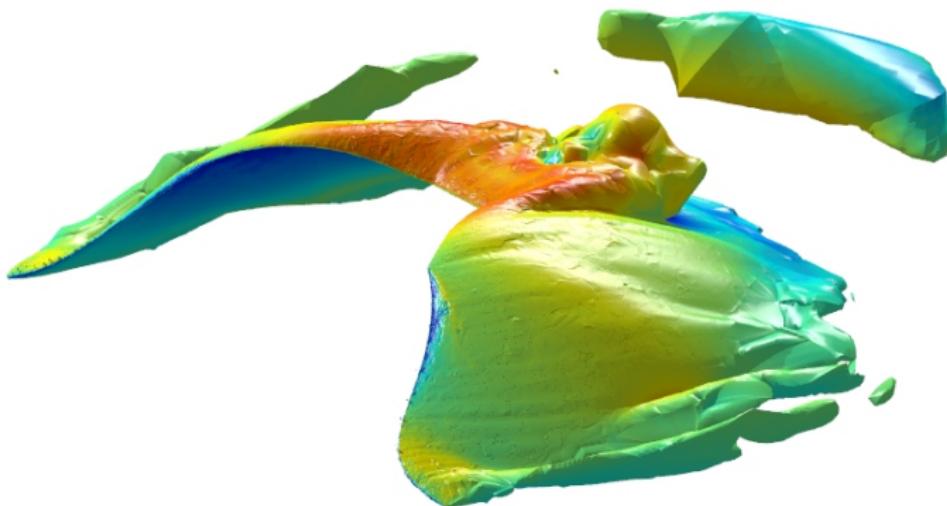
Domain Mapping

- Highly complex wing motion from measured data
- Construct mapping $\mathcal{G}(X, t)$ *numerically* by nonlinear solid mechanics approach [Persson '09]
- A reference mesh (left) is deformed elastically to smoothly align with the prescribed wing motion (right)
- Grid velocity $v_X = \frac{\partial \mathcal{G}}{\partial t} |_X$ defined consistently with DIRK scheme



Flapping Bat Flight Simulation

- Visualization of Mach number on isosurface of entropy
- Unphysical separation around simplified animal “body”



Optimal Design of Flapping Wings

- Goal: Automatically generate optimized flapping wing kinematics [Persson/Willis '11]
- A multifidelity approach, with wake-only, panel, and high-order DG methods
- Example: Flapping wing pair, prescribed camber, solve for optimal wing twist distribution

