

How do yield stress materials start to flow?



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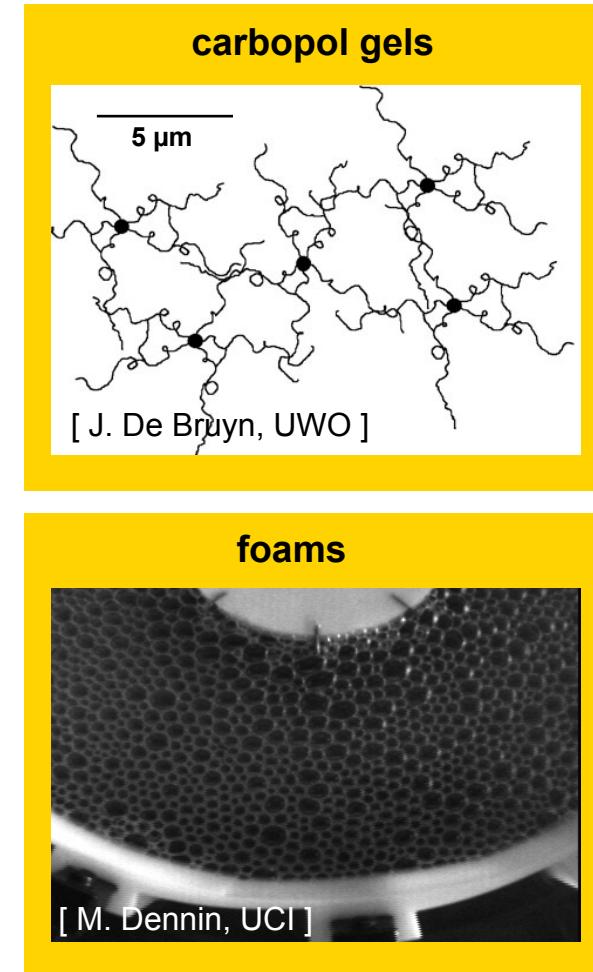
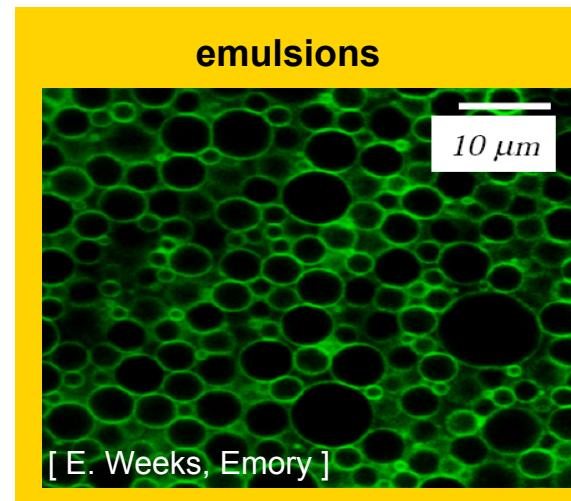
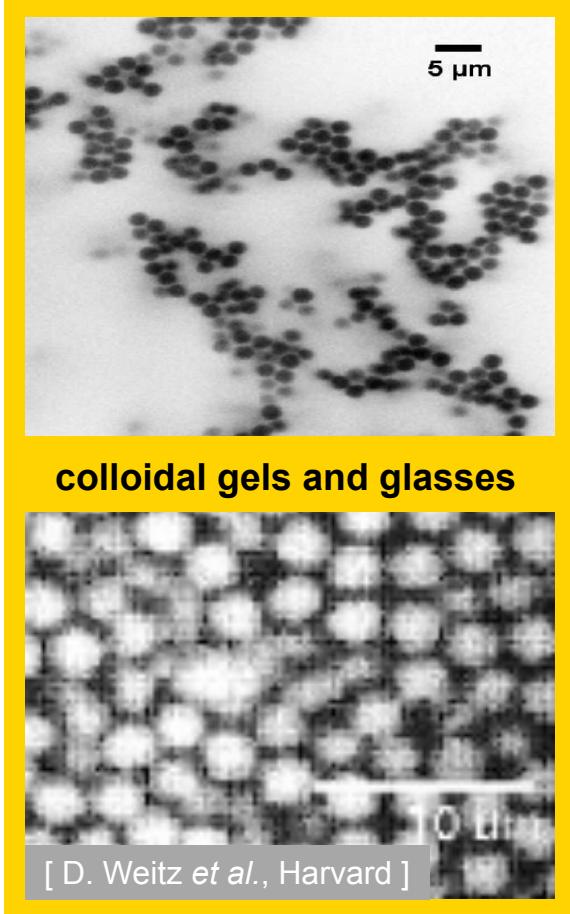
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(res. eng. 2009)

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(PhD 2009-2012)

Soft jammed materials are yield stress fluids



solid-like below σ_c
liquid-like above σ_c

⇒ how does this yielding transition work?

Outline

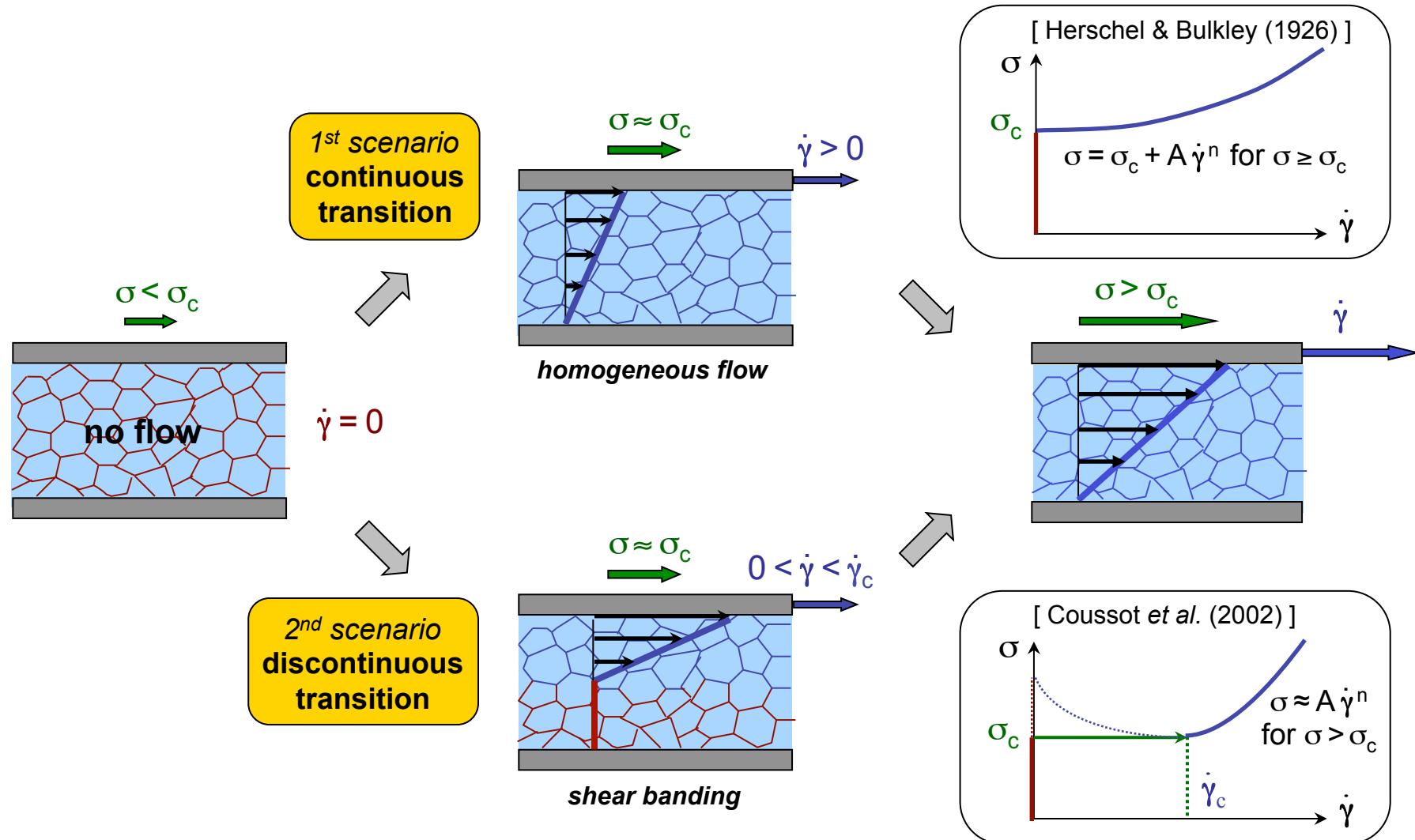
What is the steady-state reached after yielding?

- continuous vs discontinuous yielding transition
- a very short review of experimental results
- towards a classification of yield stress materials

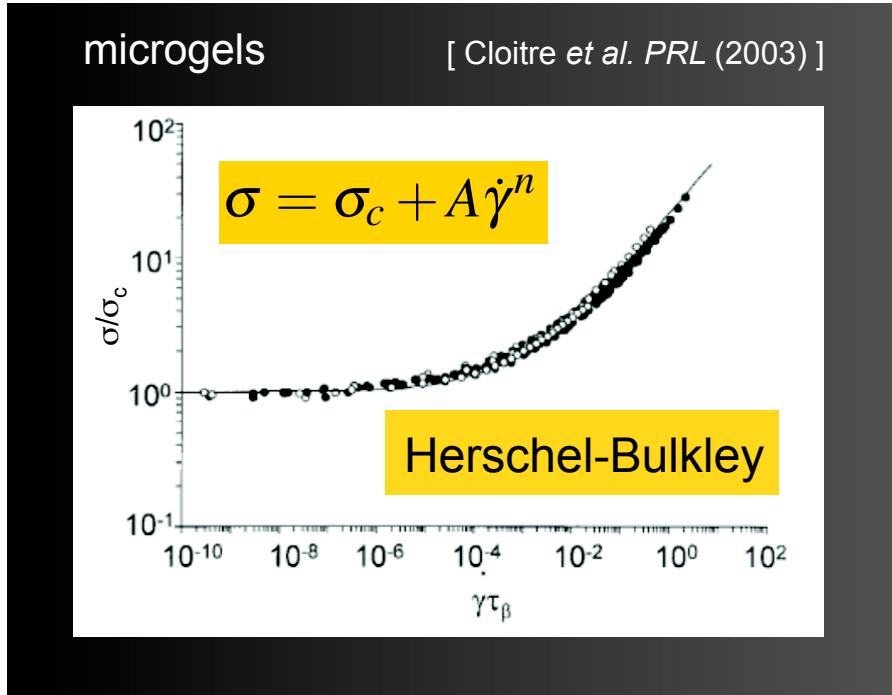
How do yield stress materials reach steady state?

- ultrasonic velocimetry in carbopol microgels
- shear rate-imposed experiments
- stress-imposed experiments
- comparison with other materials

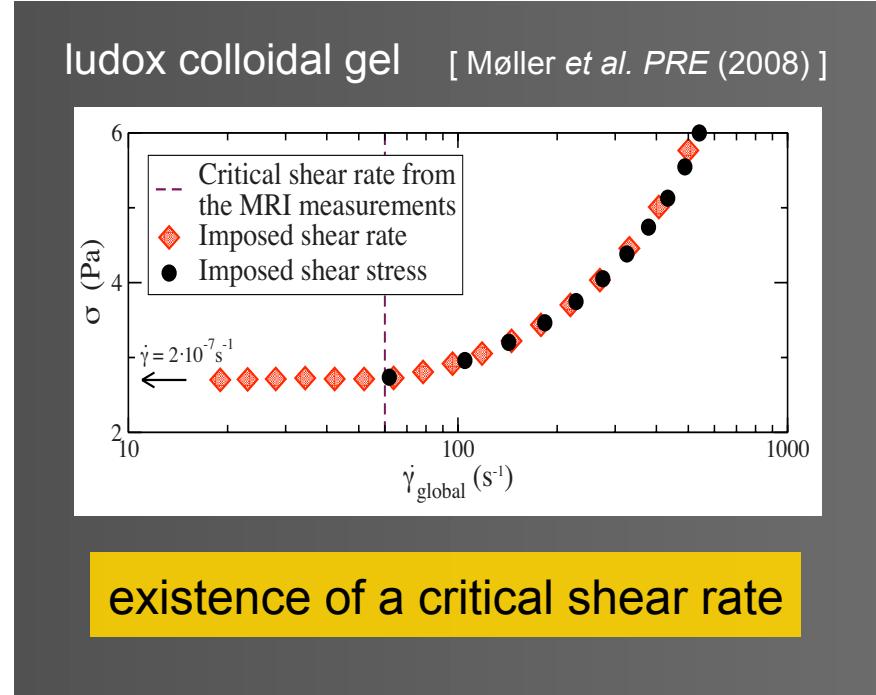
The rheological response of soft jammed materials



What is the steady-state flow curve?



continuous transition



discontinuous transition

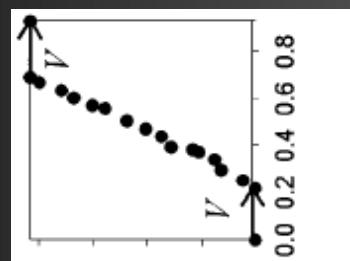
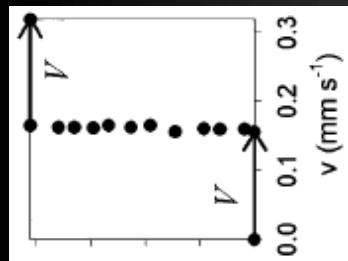
- ⇒ in practice, it may be hard to decide from rheology alone
- ⇒ look at the steady-state local flow field at the macroscale ($50 \mu\text{m}$)

What is the steady-state local flow field?

continuous transition

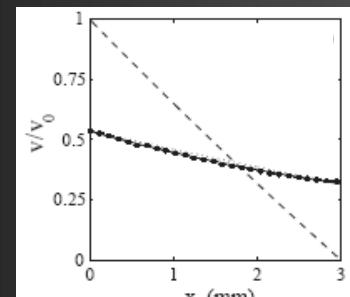
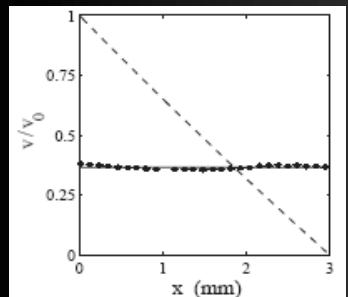
microgels

[Meeker et al. *PRL* (2004)]



emulsions

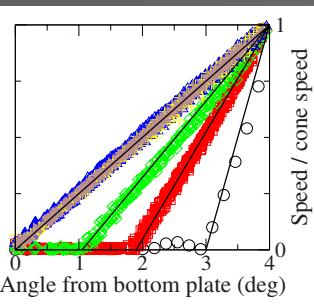
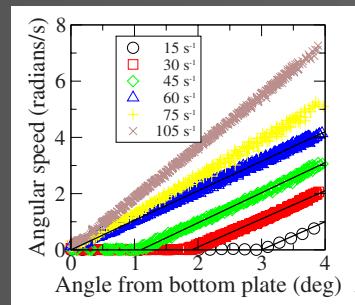
[Salmon et al. *EPJE* (2003)]



discontinuous transition

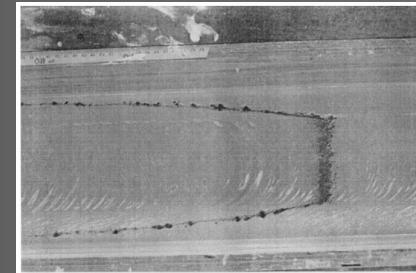
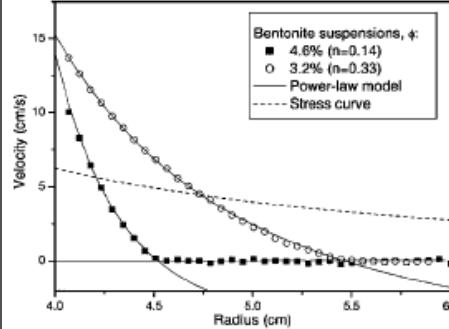
ludox colloidal gel

[Møller et al. *PRE* (2008)]



clays (bentonite, laponite) [Coussot et al. *PRL* (2002)]

Bentonite suspensions, ϕ :
■ 4.6% ($n=0.14$)
○ 3.2% ($n=0.33$)
— Power-law model
- - Stress curve



homogeneous steady-state flow

vs

shear-banded flow

Towards a classification of yield stress materials

“Simple” yield stress materials

microgels, emulsions, wet foams

- continuous yielding transition
- Herschel-Bulkley flow curve
- homogeneous flow in steady state



“Other” yield stress materials

colloidal, granular & clay suspensions

- discontinuous yielding transition
- existence of a critical shear rate $\dot{\gamma}_c$
- shear-banded flow in steady state

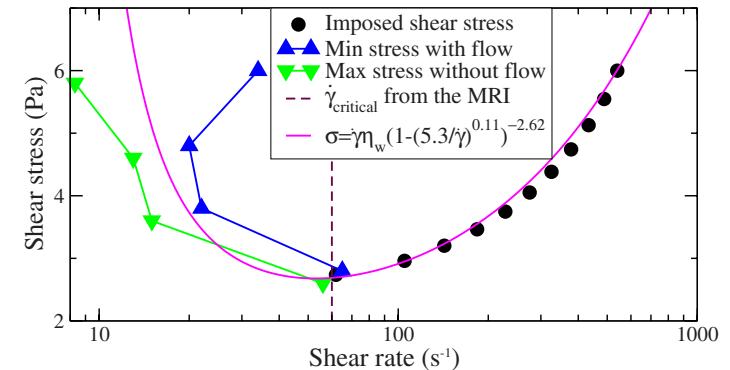


What exactly causes steady-state shear banding?

- aging vs shear rejuvenation ➤ link with thixotropy

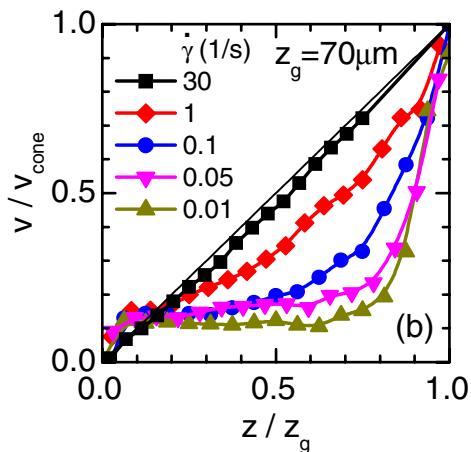
- structural parameter(s) : λ model, fluidity model
[Coussot *et al.* (2002), Picard *et al.* (2002), Rogers *et al.* (2008)]
- aggregation & breakup of fractal clusters [Møller *et al.* (2008)]

⇒ see P. Coussot's talk

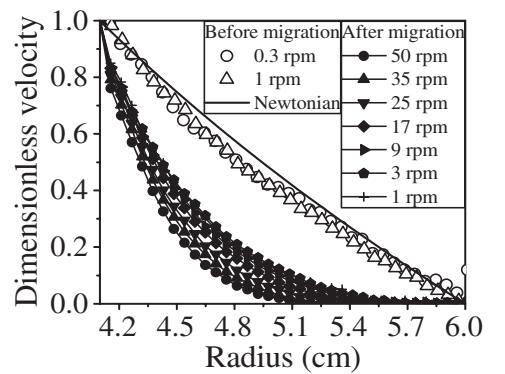
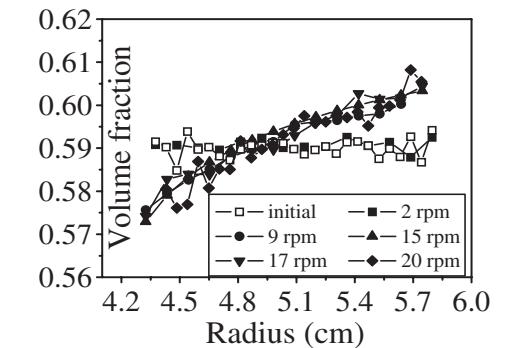
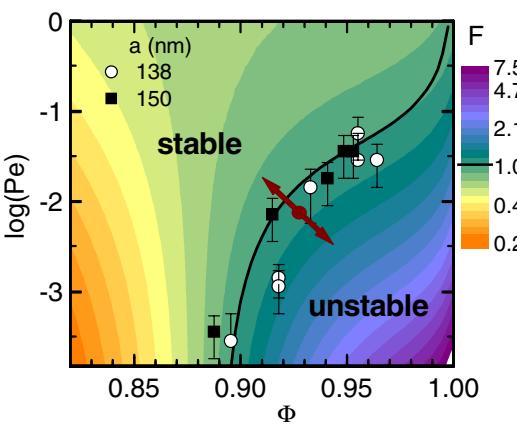


- flow-concentration coupling / particle migration

↓
colloidal glasses [Besseling *et al.* *PRL* (2010)]



→
granular suspensions [Fall *et al.* *PRL* (2010)]



Outline

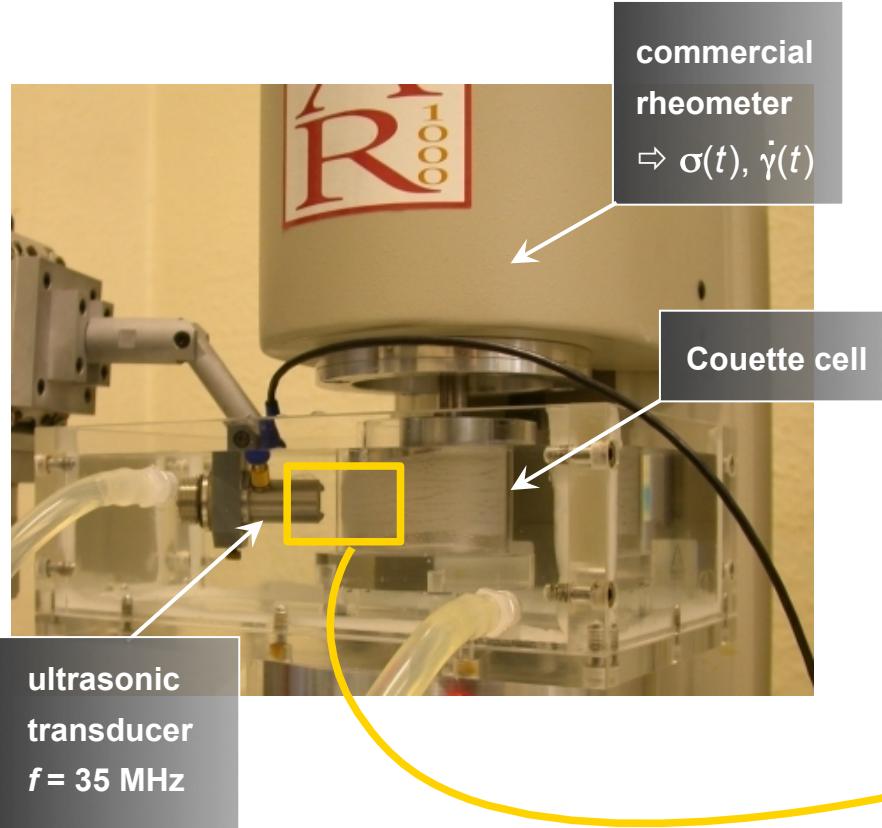
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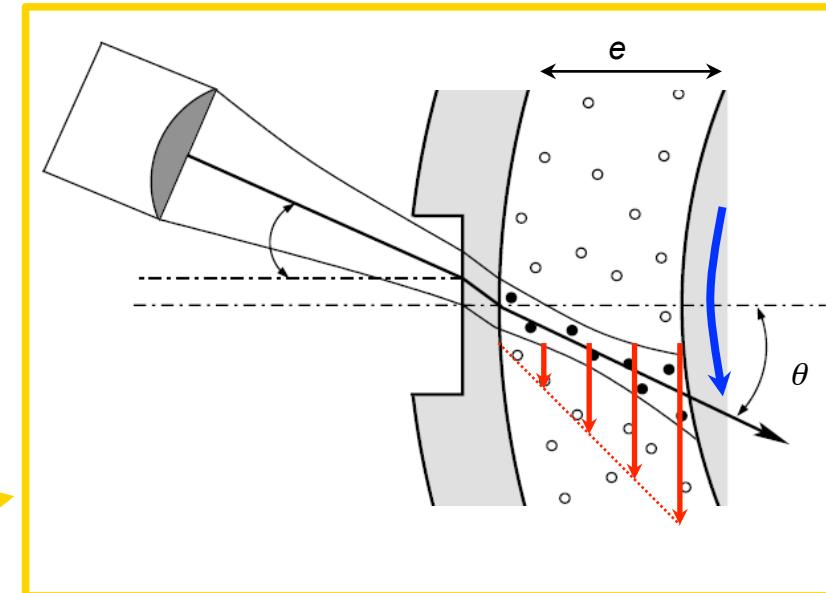
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Ultrasonic velocimetry



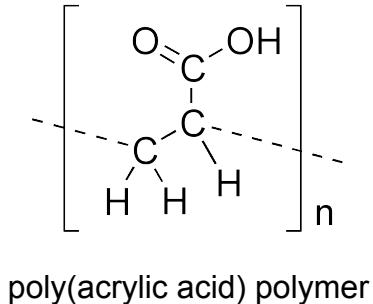
rheometer Anton Paar MCR 301
gap $e = 0.45 - 3 \text{ mm}$
surface polished Plexiglas or sandpaper



- speckle tracking algorithm $\Rightarrow v(r,t) \sin \theta$
- spatial resolution $\sim 40 \mu\text{m}$
- temporal resolution $\sim 0.1 \text{ s}$ per velocity profile

[Manneville et al., Eur. Phys. J. AP 28, 361 (2004)]

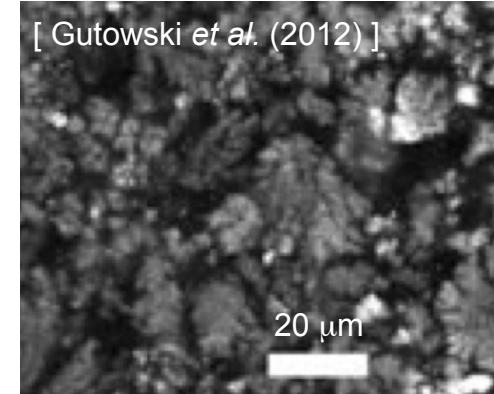
Carbopol microgels, a “simple” yield stress material



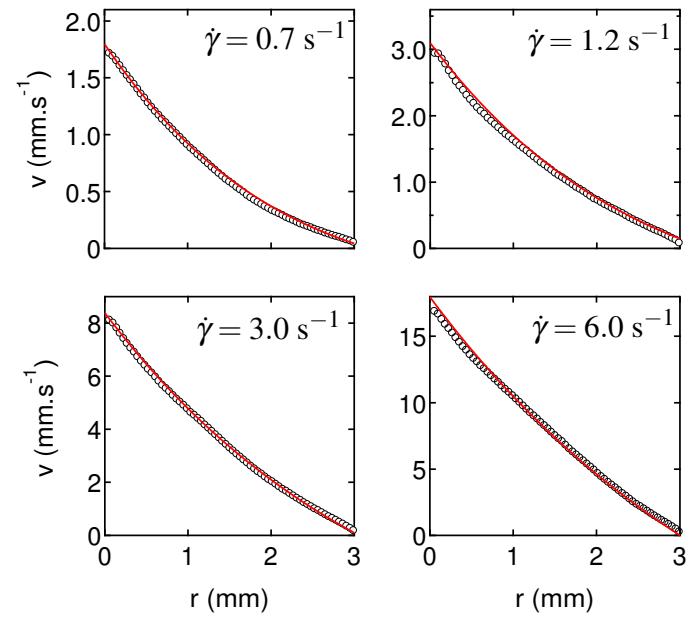
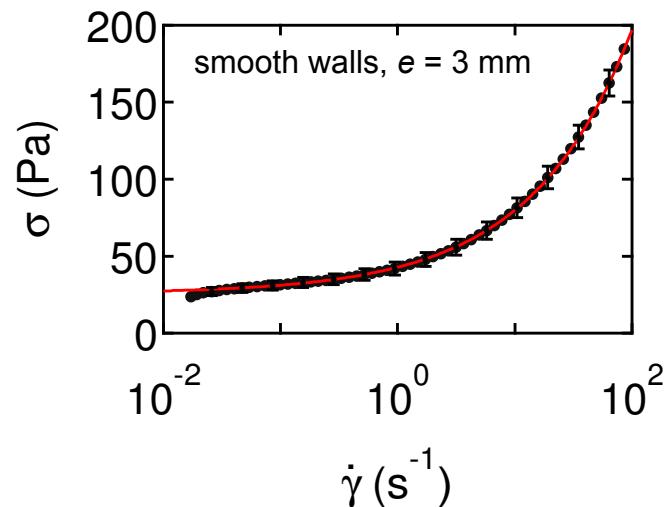
- 1% wt carbopol ETD 2050
- 0.5% wt hollow glass spheres

$$\begin{aligned} G' &= 100 \pm 10 \text{ Pa} \\ G'' &= 12 \pm 2 \text{ Pa} \end{aligned}$$

[Ketz *et al.* (1988), Baudonnet *et al.* (2004), Oppong *et al.* (2006), Lee (2011)]



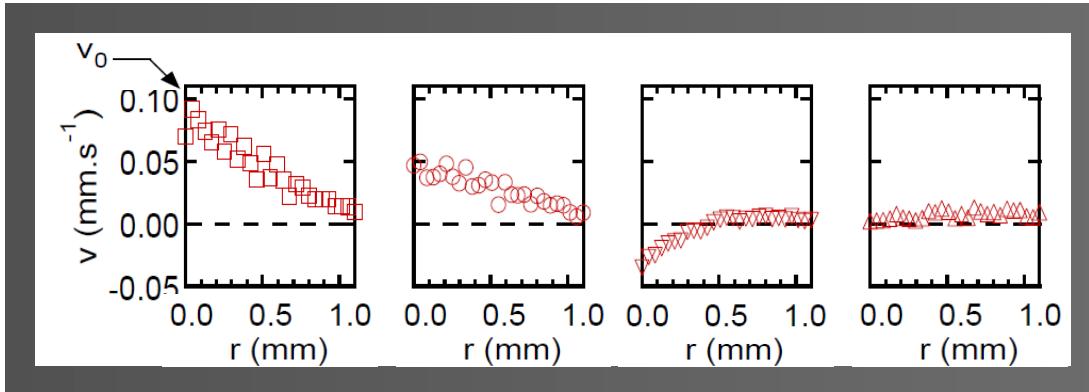
⇒ a jammed assembly of soft microgel particles with highly crosslinked cores



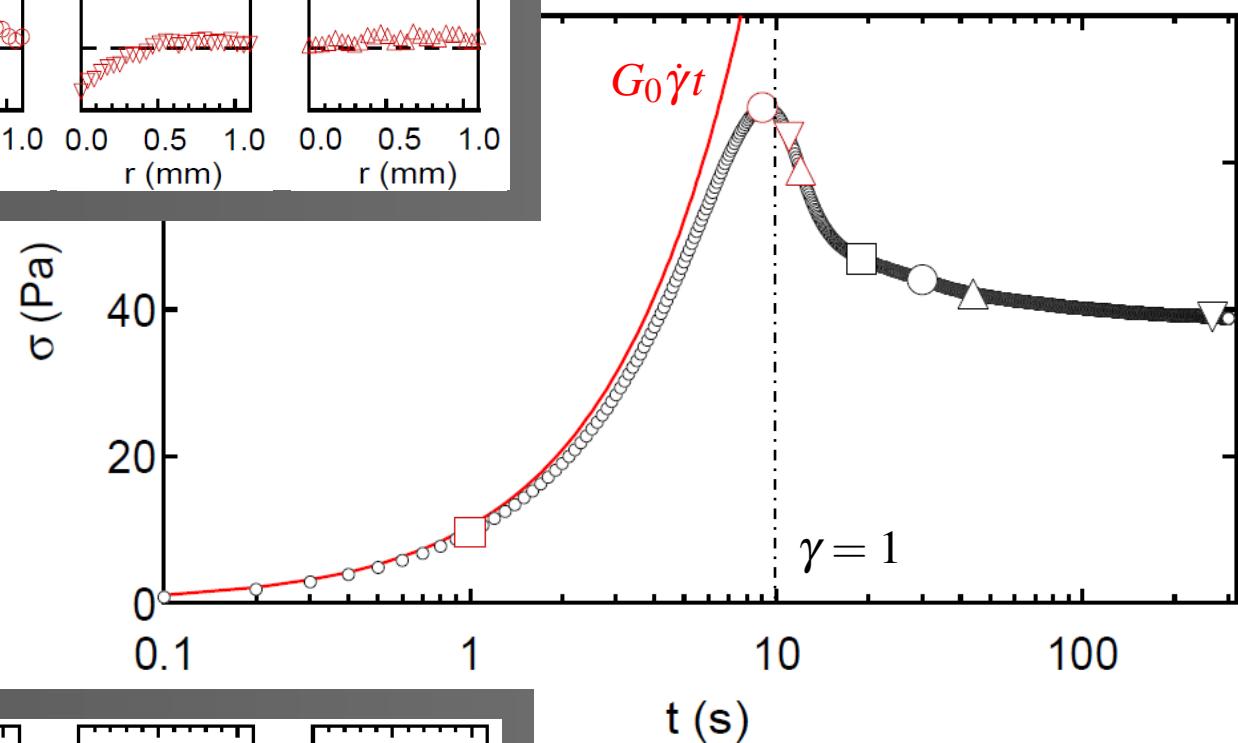
⇒ perfect Herschel-Bulkley & “simple” yielding

Shear rate-imposed experiment

[$\dot{\gamma} = 0.1 \text{ s}^{-1}$, rough walls, $e = 1.1 \text{ mm}$]



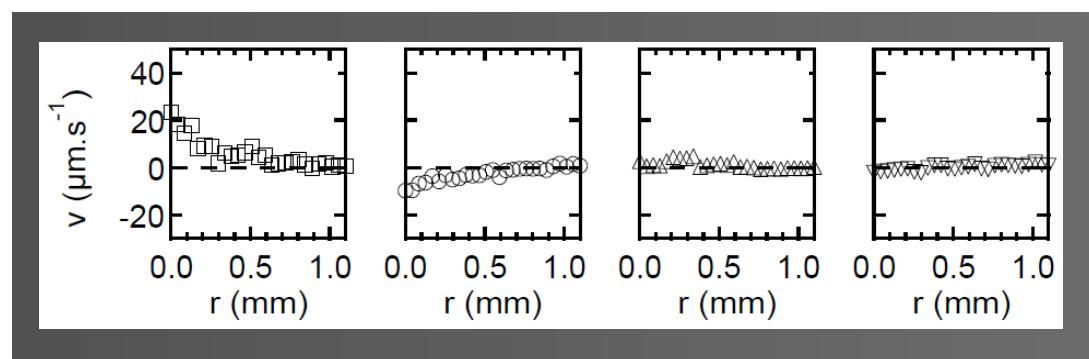
- presence of a stress overshoot at $\gamma \approx 1$
- plasticity before the overshoot



⇒ elastic deformation

⇒ failure at the rotor

⇒ elastic recoil



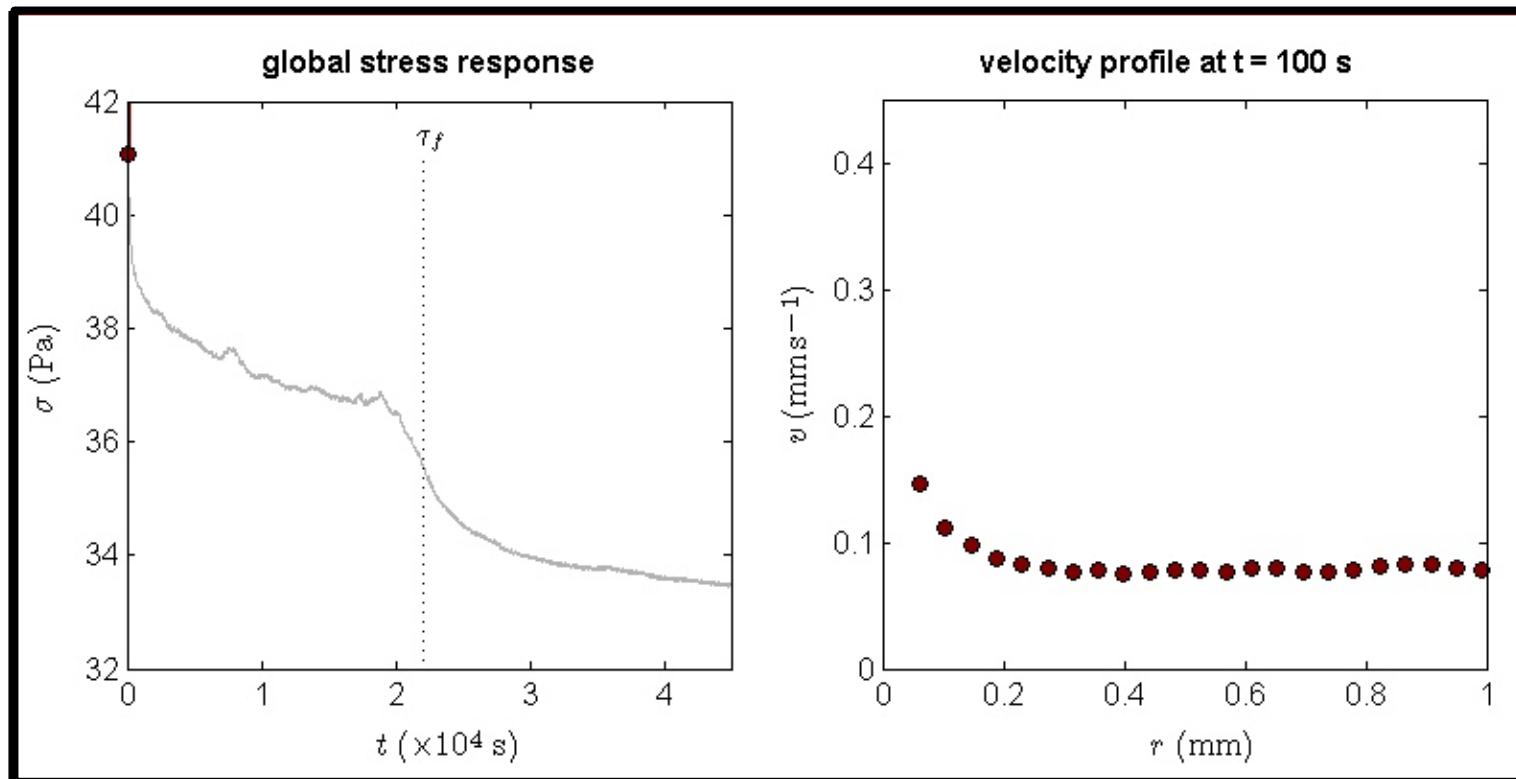
⇒ damped oscillations

⇒ total slippage at the rotor

Shear rate-imposed experiment

[$\dot{\gamma} = 0.5 \text{ s}^{-1}$, smooth walls, $e = 1 \text{ mm}$]

- long-time stress response



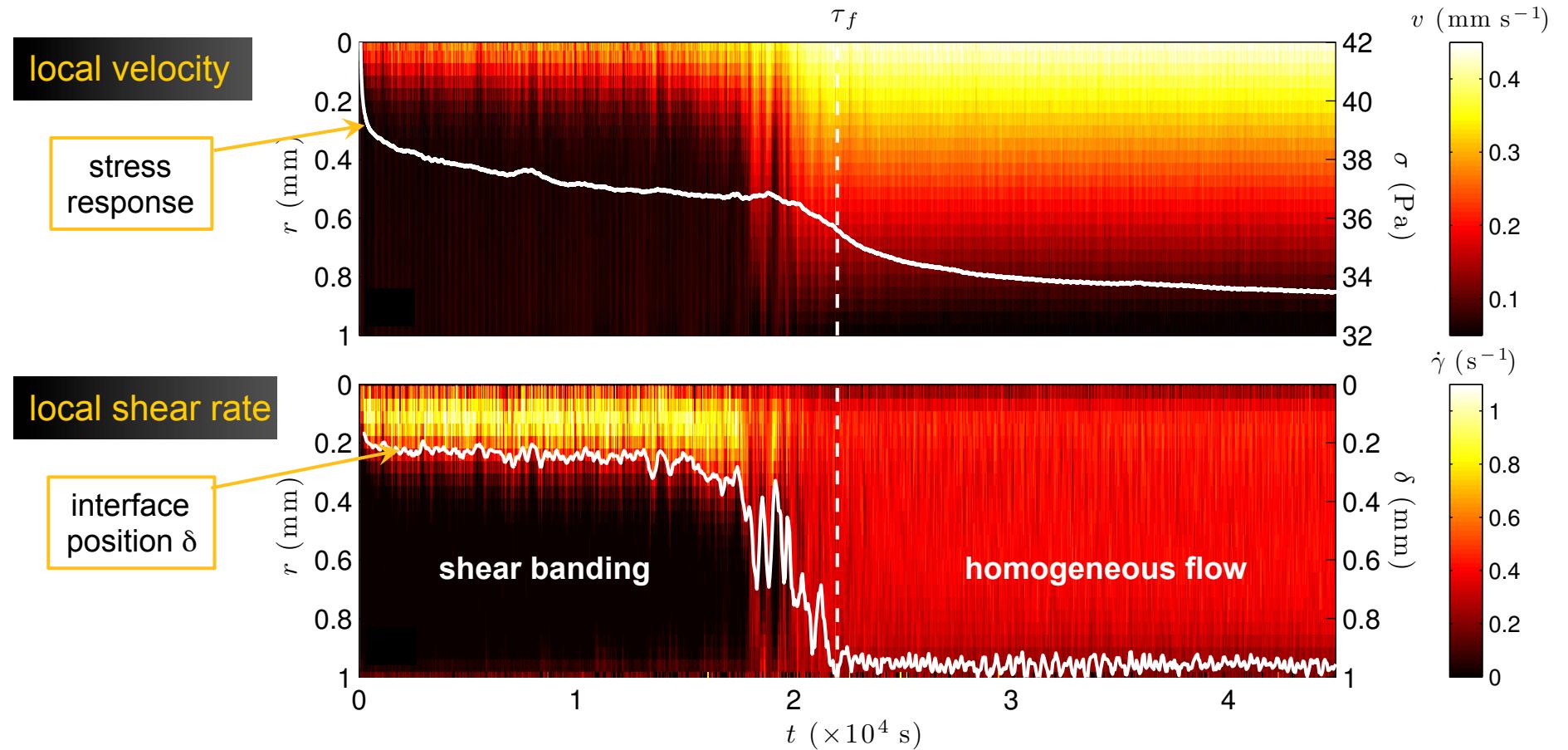
⇒ growth of a shear band & long-lasting shear-banding regime

⇒ strong fluctuations & “sudden” fluidization at τ_f

⇒ rheological signature: “kink” around τ_f

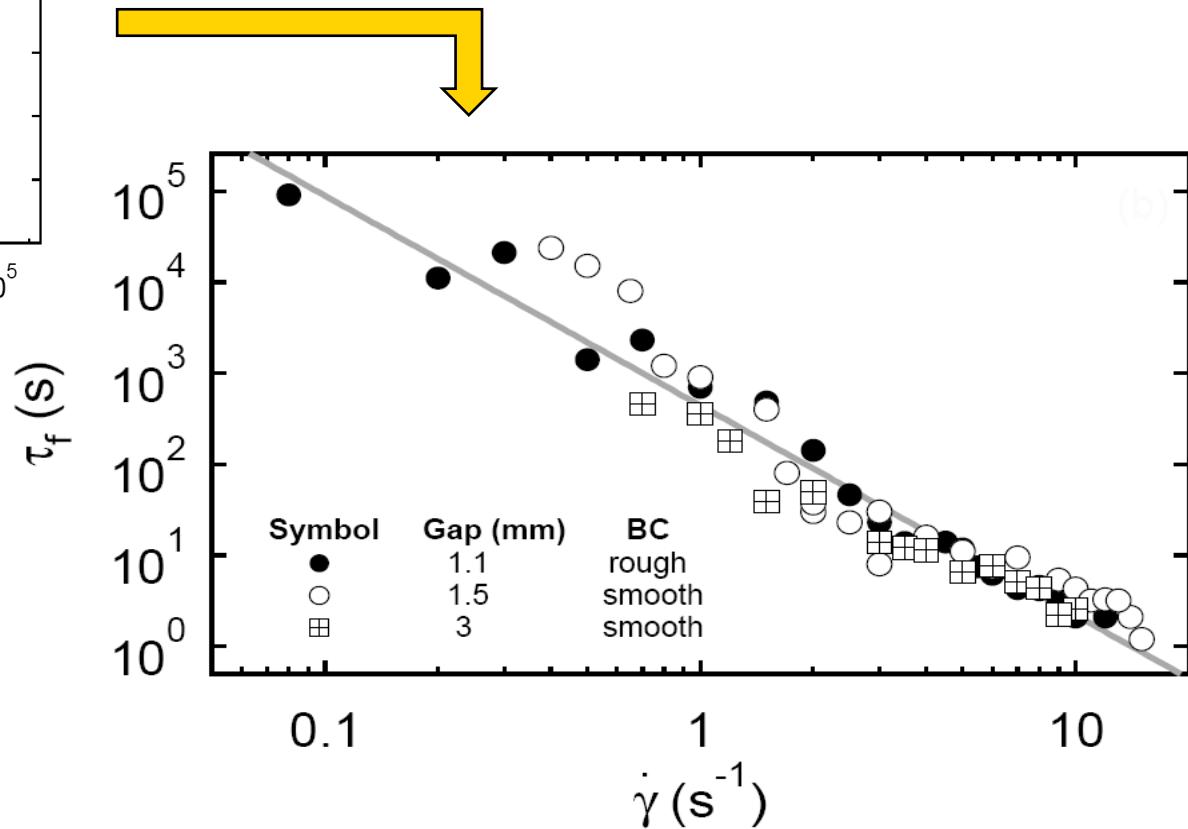
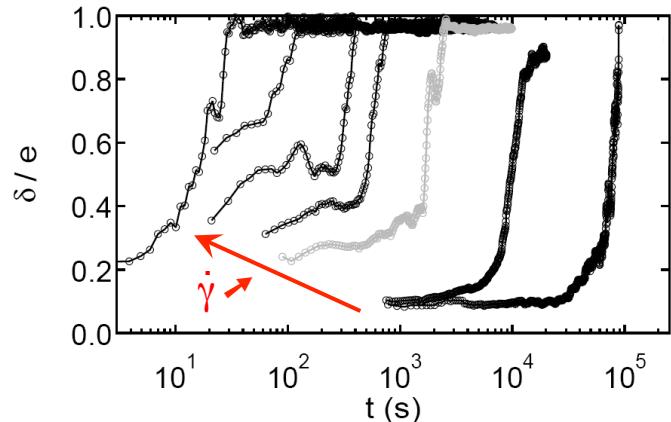
Data analysis

[$\dot{\gamma} = 0.5 \text{ s}^{-1}$, smooth walls, $e = 1 \text{ mm}$]



⇒ well-defined fluidization time τ_f

Shear rate-imposed experiments



$$\tau_f \sim \dot{\gamma}^{-\alpha}$$

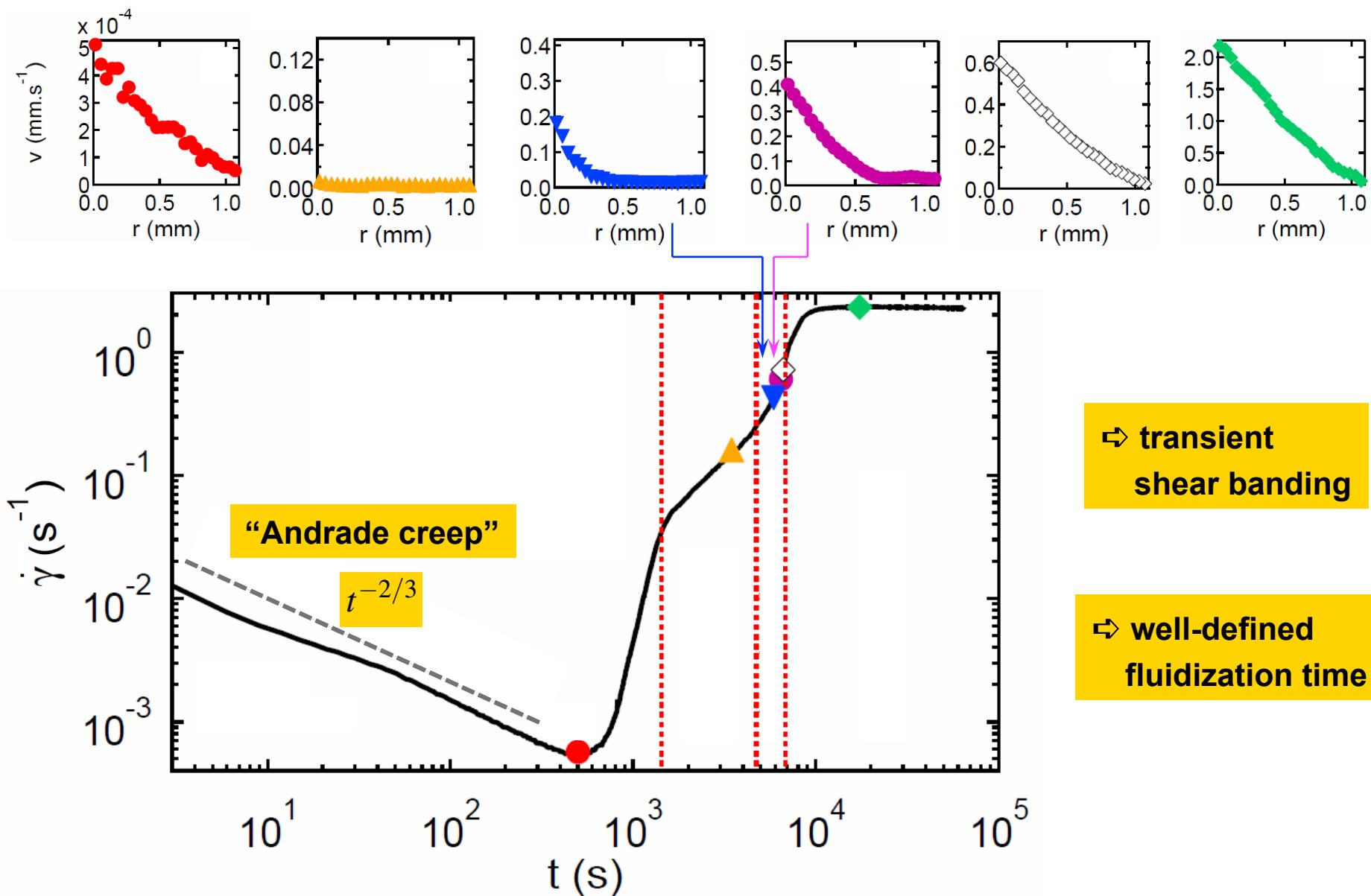
with

$$\alpha = 2.3 \pm 0.1$$

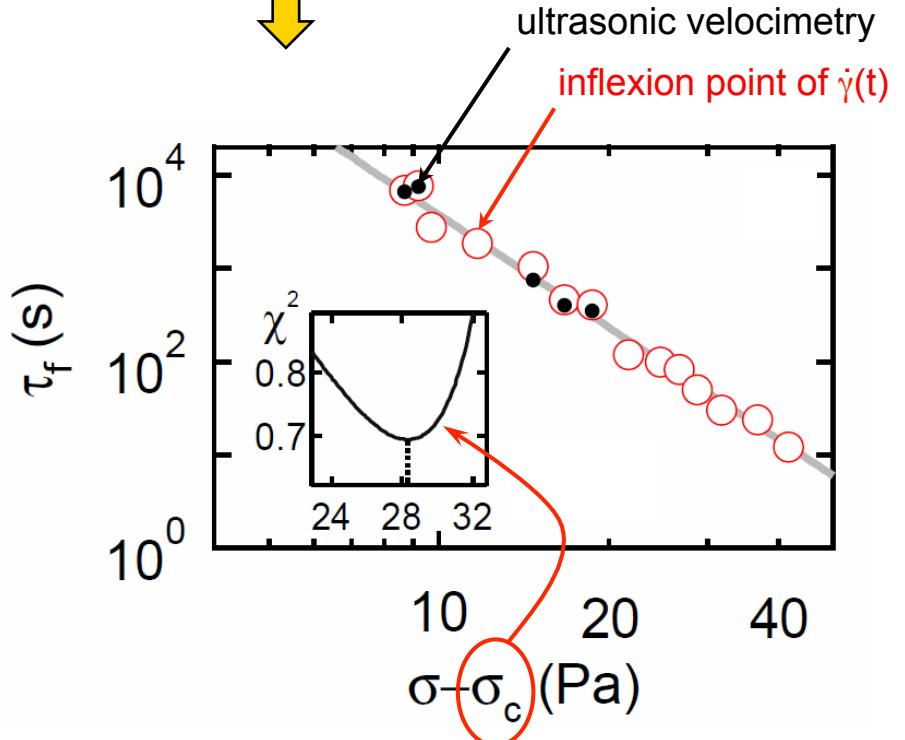
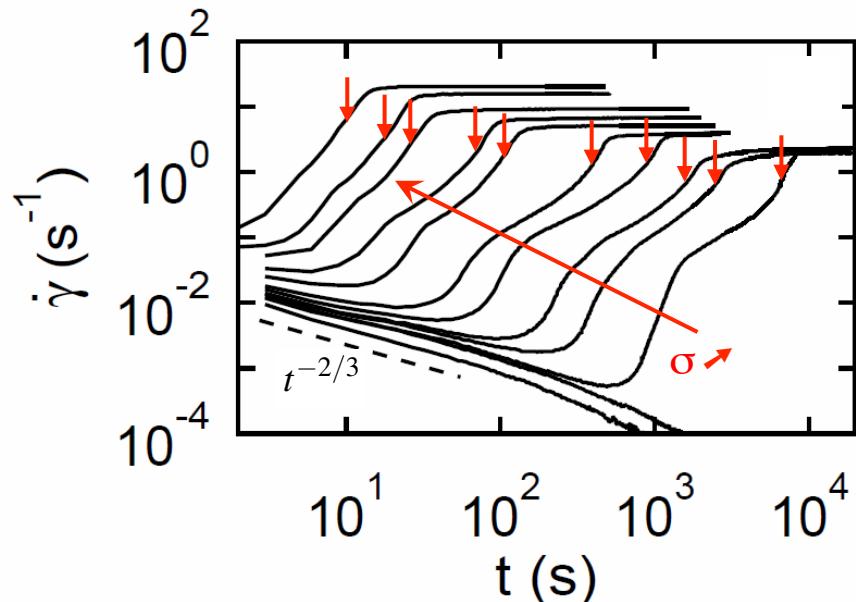
- independent of gap width and wall roughness
- dependent on carbopol batch (preparation, concentration)

Stress-imposed experiment

[$\sigma = 37 \text{ Pa}$, rough walls, $e = 1.1 \text{ mm}$]



Stress-imposed experiments



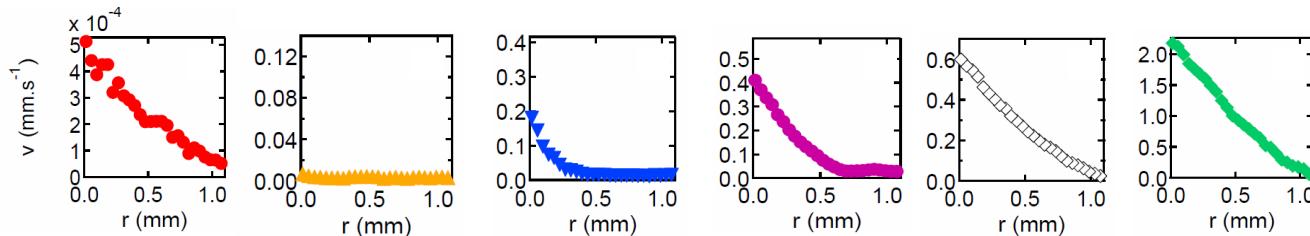
$$\tau_f \sim (\sigma - \sigma_c)^{-\beta}$$

with

$$\beta = 4.0 \pm 0.1$$

- independent of gap width and wall roughness
- dependent on carbopol batch (preparation, concentration)
- estimate of σ_c consistent with steady-state flow curve

Link between transient regime and steady state



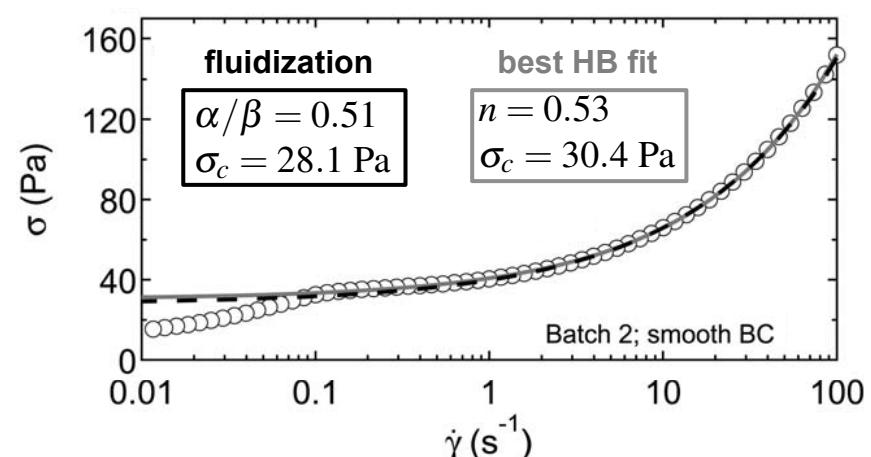
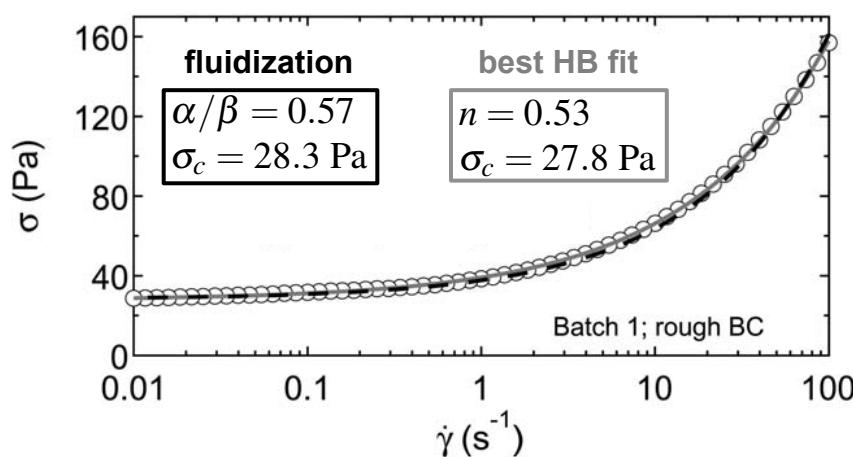
- a similar fluidization scenario under imposed $\dot{\gamma}$ and σ

$$\left. \begin{array}{l} \tau_f^{(\dot{\gamma})} = a/\dot{\gamma}^\alpha \\ \tau_f^{(\sigma)} = b/(\sigma - \sigma_c)^\beta \end{array} \right\} \text{let us assume} \quad \tau_f^{(\sigma)} = \lambda \tau_f^{(\dot{\gamma})} \quad \Rightarrow \quad \sigma = \sigma_c + A \dot{\gamma}^n$$

$n = \alpha/\beta$
 $A = (b/\lambda a)^{1/\beta}$

- steady-state Herschel-Bulkley rheology is recovered

one free parameter: $\lambda \approx 10$



What about other soft jammed materials?

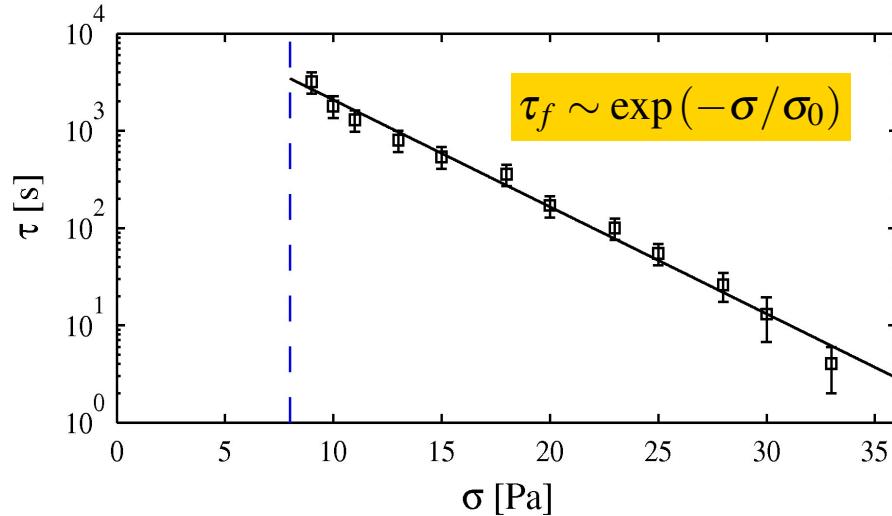
- fractures in yield stress fluids?

- Taylor model for fracture in elastic solids

- but no local measurements

[Caton & Baravian, *Rheol. Acta* **47**, 601-607 (2008)]

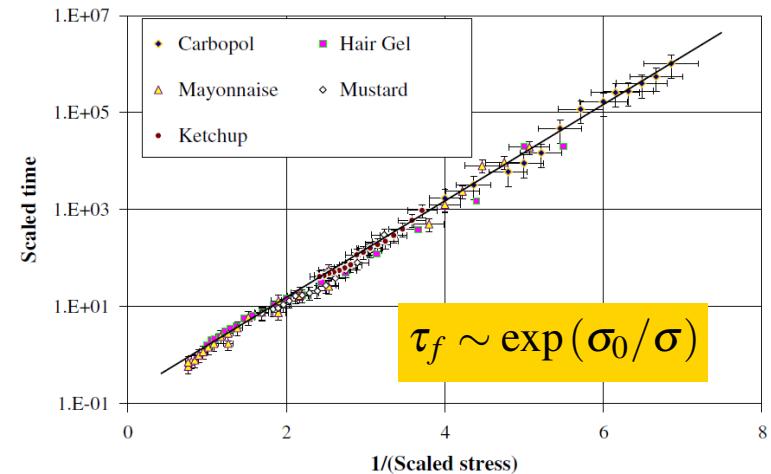
- activated processes in colloidal gels?



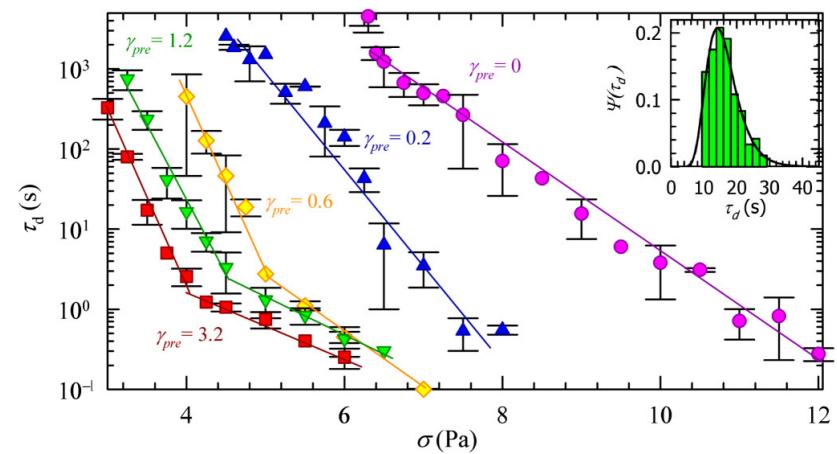
- gels of stearylated silica
- bond-breaking models

[Sprakel et al., *PRL* **106**, 248303 (2011)]

[Mora, *Soft Matter* **7**, 4908-4917 (2011)]



- carbon black particles in mineral oil
 - Kramer's law with $E(\sigma) = E_a - \sigma v$
- [Gibaud et al., *Soft Matter* **4**, 3482-3486 (2010)]



Conclusions and outlook

Critical-like dynamics of a “simple” yield stress material

- long-lasting transient shear banding
- link with steady-state rheology

$$\tau_f^{(\dot{\gamma})} = a/\dot{\gamma}^\alpha \quad \& \quad \tau_f^{(\sigma)} = b/(\sigma - \sigma_c)^\beta$$

$$\leftrightarrow \quad \sigma = \sigma_c + A\dot{\gamma}^n \quad \text{with} \quad n = \alpha/\beta$$

- ⇒ close to the yield point, check carefully that steady-state is achieved
 - ⇒ “critical” exponents that remain to be fully understood
 - ⇒ a behaviour that is certainly not general

- need for space-resolved structural measurements in the two bands
 - what time-resolved models to describe fluidization dynamics?
 - how far does the analogy with creep in solids go?

The cement slide

T. Gibaud *et al.*

PRL **101**, 258302 (2008)

Soft Matter **5**, 3026-3037 (2009)

Soft Matter **6**, 3482-3488 (2010)

T. Divoux *et al.*

PRL **104**, 208301 (2010)

Soft Matter **7**, 8409-8418 (2011)

Soft Matter **7**, 9335-9349 (2011)

Soft Matter **8**, 4151-4164 (2012)

