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clustering of particles in inhomogeneous flows:

1. Introduction

The problem we want to address is that of particle clustering for the case where the flow is NOT homogeneous.

Let us start with an analogy. If I release particles (inertial, heavy) in a homogeneous and isotropic turbulent flow, the speed of the particles follow a Maxwellian distribution:

$$\boxed{\begin{aligned} f(v) &= \frac{\exp\left(-\frac{v^2}{2\beta}\right)}{4\pi\beta^{3/2}} \end{aligned}} \quad (1)$$

If we draw an analogy with ideal gas particles, then β plays the role of inverse (scaled) temperature

$$\beta = \frac{m}{2k_B T} \quad (2)$$

For the heavy inertial particles, β depends on the Stokes number; $\beta(St)$.

Furthermore, β should also be a function of the fluid velocity; dimensionally speaking

(7)

$$(3) \quad \beta(st) = u_{rms}^2 g(st)$$

If we ~~push~~ push the analogy with temperature further we should observe the following:

If we bring into contact two bodies with different temperatures then heat flows from higher to lower temperature.

Continuing the analogy for heavy inertial particles: if we bring into contact two part of a simulation, with the same number density but different β there should be a flux of energy; carried by the particles, from lower β to higher β .

Interpreted this way turbophoresis should be determined by the flux carried by particles from lower u_{rms}^2 to higher u_{rms}^2 .

(3)

There are two fluxes that one should consider:

$J \equiv$ flux of particles

(4)

$$= n v$$

(3)

$J_q \equiv$ flux of energy carried by the particles

(5)

$$= n v^2 v$$

(4)

By the thermal analogy

(6)

$$J_q = -\kappa \nabla \left(\frac{1}{\beta} \right)$$

(5)

(7) \Rightarrow

$$J_q(st) = -g(st) \kappa \nabla \left(\frac{1}{u_{rms}^2} \right)$$

(6)

Q1. Is this true? Can we measure this effect in a simulation? And also study the fluctuation of this flux?

The ~~main~~ main question is not really about this flux; but the flux of particles. The phenomenon of thermophoresis implies that due to the temperature gradient; there is not only a flux of energy but also

a flux of particles. Phenomenologically the flux of particles may look like the following

$$(8) \quad \bar{J} = - \rho D \left(\nabla \ln \rho + \kappa_T \nabla \ln T + \kappa_p \nabla \ln p_h \right)$$

$$\left[\text{Goldhirsch and Ronis, PRA } \underline{27} \text{ 1616 1983} \right]$$

where ρ is the density, and D , a diffusion coefficient.

By analogy let us write the flux of particles in a turbophoretic situation to be the following

$$(9) \quad \bar{J} = -\chi_n \nabla n - \chi_u \nabla \left(\frac{1}{\beta} \right)$$

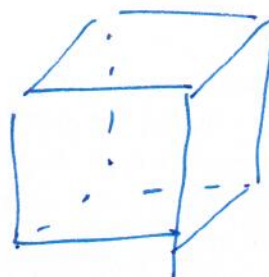
where both χ_n and χ_u are functions of Stokes number.

2. Model

Let us start with a very simple setup.

- periodic box.

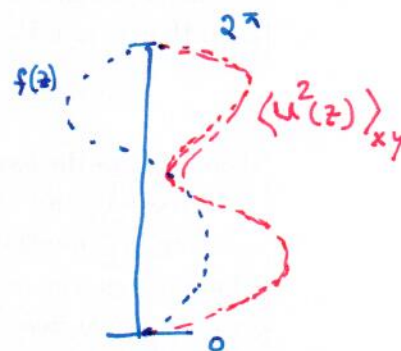
- forced turbulence (non-helical)



- The amplitude of the force varies as a periodic function of the z coordinate

$$f(x, y, z) = f_{NH}(x, y, k) f(z)$$

↑
usual forcing
function in
pencil-code



with $f(z) = \sin z$

Then we expect $\langle u^2 \rangle_{xy}(z) \sim \sin^2 z$

The net flux of inertial particles will be:

$$(10) \quad J = -\alpha_n \frac{\partial n}{\partial z} - \alpha_u \frac{\partial}{\partial z} \left(\frac{1}{\beta} \right)$$

If the system reaches a stationary state then these two fluxes are going to cancel each other.

(6)

$$\Rightarrow \frac{\partial n}{\partial z} = - \frac{x_u}{x_n} \frac{\partial}{\partial z} \left(\frac{1}{\beta} \right)$$

$$(10) \quad = - \frac{x_u(st)}{x_n} \frac{1}{g(st)} \frac{\partial}{\partial z} \left(\frac{1}{u_{rms}^2} \right)$$

$$= - \frac{x_u(st)}{x_n} \frac{1}{g(st)} u_{rms}^2 \frac{\partial}{\partial z} \left[\ln(u_{rms}^2) \right]$$

$$(11) \quad \boxed{\frac{\partial n}{\partial z} = - \frac{x_u(st)}{g(st)} \left(\frac{u_{rms}^2}{x_n} \right) \frac{\partial}{\partial z} \left[\ln(u_{rms}^2) \right]}$$

\Rightarrow Given the dependence of u_{rms}^2 and n as a function of z , ω and looking at the st dependence one should be able to figure out the ~~thermo~~ turbo-phoretic coefficient $x_u(st)$.