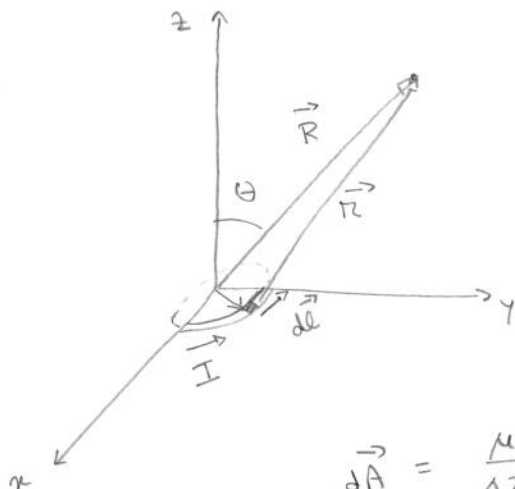


5.5

The magnetic dipole



$$d\vec{A} = \frac{\mu_0 I}{4\pi} \frac{d\vec{l} \times \vec{r}}{r^3}$$

$$\vec{A} = \frac{\mu_0 I}{4\pi} \int_0^{2\pi} \frac{d\phi r (-\sin\phi \hat{x} + \cos\phi \hat{y})}{(r^2 + R^2 - 2Rr \sin\theta \sin\phi)^{3/2}}$$

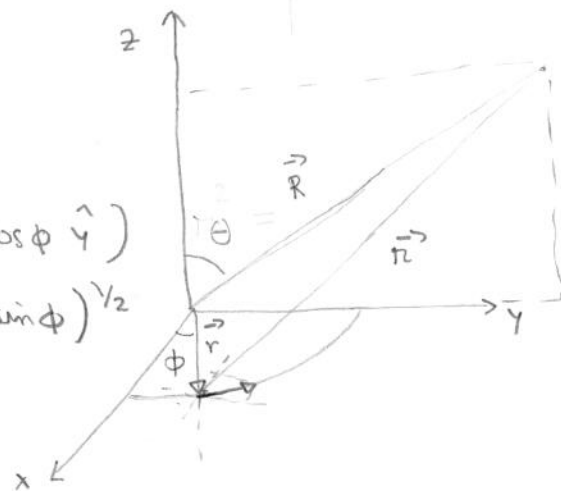
$$\int_0^{2\pi} \frac{\cos\phi d\phi}{(P - Q \sin\phi)^{3/2}} = 0$$

So the vector potential has only the x component; which is

$$-\frac{\mu_0 I}{4\pi} \int_0^{2\pi} \frac{\sin\phi d\phi}{(r^2 + R^2 - 2Rr \sin\theta \sin\phi)^{3/2}}$$

which in its full glory can look quite complex.

Let us calculate the vector potential of a current loop at a point \vec{R} from its center.



$$\vec{R} = (\hat{z} \cos\theta + \hat{y} \sin\theta) R$$

$$\vec{r} = (\hat{x} \cos\phi + \hat{y} \sin\phi) r$$

$$\vec{R} = \vec{R} - \vec{r}$$

$$= -\hat{x} \cos\phi r + \hat{y} (R \sin\theta - r \sin\phi) + \hat{z} R \cos\theta$$

$$R^2 = r^2 \cos^2\phi + R^2 \cos^2\theta$$

$$+ R^2 \sin^2\theta + r^2 \sin^2\phi$$

$$- 2Rr \sin\theta \sin\phi$$

$$= r^2 + R^2 - 2Rr \sin\theta \sin\phi$$

$$d\vec{l} = r d\phi \hat{e}_\phi$$

$$= r d\phi (-\sin\phi \hat{x} + \cos\phi \hat{y})$$

But the main purpose of this exercise is to look at the potential for large R

In that approximation:

$$\frac{1}{(R^2 + r^2 - 2rR \sin\theta \sin\phi)^{1/2}} = \frac{1}{R} (1 + \xi^2 - 2\xi \sin\theta \sin\phi)^{-1/2}, \quad \xi = \frac{r}{R}$$

$$= \frac{1}{R} \left[1 + \xi \sin\theta \sin\phi + O(\xi^2) \right]$$

↑ terms of order ξ^2 and higher which will be ignored.

$$\Rightarrow \vec{A}(0, R \sin\theta, R \cos\theta)$$

$$= -\hat{x} \frac{\mu_0 I}{4\pi} \frac{r}{R} \int_0^{2\pi} d\phi \sin\phi \left(1 + \frac{r}{R} \sin\theta \sin\phi \right)$$

$$= -\hat{x} \frac{\mu_0 I}{4\pi} \frac{r}{R} \left(0 + \frac{r}{R} \sin\theta \pi^2 \right) \quad \left| \quad \int_0^{2\pi} \sin^2\phi d\phi = \frac{1}{2} 2\pi \right.$$

$$= -\hat{x} \frac{\mu_0}{4\pi} \frac{I \pi r^2}{R^2} \sin\theta$$

$$= -\hat{x} \frac{\mu_0}{4\pi} \frac{m \sin\theta}{R^2}$$

$m \equiv$ magnetic dipole moment
 $= \pi r^2 I$

There is a nicer way to write this result,

define $\vec{m} = I \pi r^2 \hat{z}$ (by giving the area of the loop a direction)
 $= m \hat{z}$

Evaluate $\vec{m} \times \hat{R}$

$$= \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 0 & 0 & m \\ 0 & \sin\theta & R \cos\theta \end{vmatrix}$$

$$= -\hat{x} m R \sin\theta$$

$$\Rightarrow \boxed{\vec{A} = \frac{\mu_0}{4\pi} \frac{\vec{m} \times \hat{R}}{R^2}}$$

$$m = I \pi r^2 \hat{n}$$

vector potential of a magnetic dipole.

The field of a dipole can be calculated from this expression

$$\vec{B} = \nabla \times \vec{A}$$

$$\nabla \times \vec{A} = \frac{\mu_0}{4\pi} (\vec{m} \cdot \nabla) \left(\frac{\hat{R}}{R^2} \right) \cdot \left(\frac{\hat{R}}{R^2} \right)$$

5.6

The Feynman trick to calculate

$$\vec{\nabla} \times (\vec{F} \times \vec{G})$$

$\vec{\nabla}$ as a differential operation acts on both \vec{A} and \vec{B} using the chain rule. For example

$$\frac{d}{dx} (fg) = \left(\frac{df}{dx} \right) g + f \frac{dg}{dx}$$

The same expression can be symbolically written as

$$\left(\frac{d}{dx} \right) (fg) = \left(\frac{d}{dx} \right)_f (fg) + \left(\frac{d}{dx} \right)_g (fg)$$

where, by definition,

$$\left(\frac{d}{dx} \right)_f \text{ operates only on } f$$

$$\text{and } \left(\frac{d}{dx} \right)_g \text{ operates only on } g$$

In the same way

$$\vec{\nabla} \times (\vec{F} \times \vec{G}) = \vec{\nabla}_F \times (\vec{F} \times \vec{G}) + \vec{\nabla}_G \times (\vec{F} \times \vec{G})$$

Now in each of these consider $\vec{\nabla}$ to be a vector:

$$\begin{aligned} & \vec{\nabla}_F \times (\vec{F} \times \vec{G}) \\ &= \vec{F} (\vec{\nabla}_F \cdot \vec{G}) - \vec{G} (\vec{\nabla}_F \cdot \vec{F}) \\ &= (\vec{G} \cdot \vec{\nabla}) \vec{F} - \vec{G} (\vec{\nabla} \cdot \vec{F}) \end{aligned}$$

From the second term

$$\begin{aligned}
 & \vec{\nabla}_G \times (\vec{F} \times \vec{G}) \\
 = & \vec{F} (\vec{\nabla}_G \cdot \vec{G}) - \vec{G} (\vec{\nabla}_G \cdot \vec{F}) \\
 = & (\vec{\nabla} \cdot \vec{G}) \vec{F} - (\vec{F} \cdot \vec{\nabla}) \vec{G}
 \end{aligned}$$

Putting together

$$\begin{aligned}
 & \nabla \times (\vec{F} \times \vec{G}) \\
 = & (\vec{G} \cdot \vec{\nabla}) \vec{F} - (\vec{F} \cdot \vec{\nabla}) \vec{G} \\
 + & (\vec{\nabla} \cdot \vec{G}) \vec{F} - (\vec{\nabla} \cdot \vec{F}) \vec{G}
 \end{aligned}$$