

Exact solution for quantum-classical hybrid mode of a classical harmonic oscillator quadratically coupled to a degenerate two-level quantum system

The code in green should be uncommented (and the code in orange needs to be commented out) if one wants to study the general solution.

Settings

restart;
with (Physics) :

Hamiltonian
alias (H = H(q, p)) :

alias (f₁ = f₁(q), f₂ = f₂(q), f₃ = f₃(q)) :

Wavefunction
alias (Y₁ = Y₁(q, p, t), Y₂ = Y₂(q, p, t)) :

Y := Vector([Y₁, Y₂]) :

Pauli matrices

*alias (σ₁ = Library:-RewriteInMatrixForm(Psigma[1]), σ₂ = Library:-
RewriteInMatrixForm(Psigma[2]), σ₃ = Library:-RewriteInMatrixForm(Psigma[3])) :*

Quantum-classical hybrid equation of motion for the Koopman wavefunction Y

Hybrid equation of motion

$$\begin{aligned} \text{HybridEq} := & -I \cdot \hbar \cdot \frac{\partial}{\partial t} Y - I \cdot \hbar \cdot \left(\frac{\partial}{\partial p} H \cdot \frac{\partial}{\partial q} Y - \frac{\partial}{\partial q} H \cdot \frac{\partial}{\partial p} Y - \left(\frac{\partial}{\partial q} f_1 \cdot \sigma_1 + \frac{\partial}{\partial q} f_2 \cdot \sigma_2 \right. \right. \\ & \left. \left. + \frac{\partial}{\partial q} f_3 \cdot \sigma_3 \right) \cdot \frac{\partial}{\partial p} Y \right) \\ & - \frac{1}{2} \cdot \left(q \cdot \frac{\partial}{\partial q} H \cdot Y + q \cdot \left(\frac{\partial}{\partial q} f_1 \cdot \sigma_1 + \frac{\partial}{\partial q} f_2 \cdot \sigma_2 + \frac{\partial}{\partial q} f_3 \cdot \sigma_3 \right) \cdot Y + p \cdot \frac{\partial}{\partial p} H \cdot Y \right) \\ & + H \cdot Y + (f_1 \cdot \sigma_1 + f_2 \cdot \sigma_2 + f_3 \cdot \sigma_3) \cdot Y : \end{aligned}$$

Hybrid density matrix expressed through the Koopman wave function

Define the matrix components of Hybrid density matrix

$$D_{1,1} := 2 \cdot |Y_1|^2 + \Re \left(q \cdot Y_1 \cdot \frac{\partial}{\partial q} \text{conjugate}(Y_1) + p \cdot Y_1 \cdot \frac{\partial}{\partial p} \text{conjugate}(Y_1) + 2 \cdot I \cdot \hbar \cdot \frac{\partial}{\partial q} Y_1 \cdot \frac{\partial}{\partial p} \text{conjugate}(Y_1) \right) :$$

$$D_{1,2} := 2 \cdot Y_1 \cdot \text{conjugate}(Y_2) + I \cdot \hbar \cdot \left(\frac{\partial}{\partial q} Y_1 \cdot \frac{\partial}{\partial p} \text{conjugate}(Y_2) - \frac{\partial}{\partial q} \text{conjugate}(Y_2) \cdot \frac{\partial}{\partial p} Y_1 \right) + \frac{Y_1}{2} \cdot \left(q \cdot \frac{\partial}{\partial q} \text{conjugate}(Y_2) + p \cdot \frac{\partial}{\partial p} \text{conjugate}(Y_2) \right) + \frac{\text{conjugate}(Y_2)}{2} \cdot \left(q \cdot \frac{\partial}{\partial q} Y_1 + p \cdot \frac{\partial}{\partial p} Y_1 \right) :$$

$$D_{2,2} := 2 \cdot |Y_2|^2 + \Re \left(q \cdot Y_2 \cdot \frac{\partial}{\partial q} \text{conjugate}(Y_2) + p \cdot Y_2 \cdot \frac{\partial}{\partial p} \text{conjugate}(Y_2) + 2 \cdot I \cdot \hbar \cdot \frac{\partial}{\partial q} Y_2 \cdot \frac{\partial}{\partial p} \text{conjugate}(Y_2) \right) :$$

Classical density

$$\rho_{\text{classical}} := D_{1,1} + D_{2,2} :$$

Density matrix for the quantum subsystem

$$\rho_{\text{quant}} := \text{map} \left(x \rightarrow \text{Int}(x, [q = -\infty .. +\infty, p = -\infty .. +\infty]), \begin{bmatrix} D_{1,1} & D_{1,2} \\ \text{conjugate}(D_{1,2}) & D_{2,2} \end{bmatrix} \right) :$$

Example 1: free particle classical Boltzmann state

$$\text{ExampleFreeParticle} := \left\{ Y_1 = \frac{1}{\sqrt{2} \cdot \sqrt[4]{\frac{2 \cdot \text{Pi}}{\beta}}} \cdot \exp \left(-\beta \cdot \frac{p^2}{4} + \frac{I \cdot p \cdot q}{2 \cdot \hbar} \right), Y_2 = 0 \right\}$$

$$\text{ExampleFreeParticle} := \left\{ \text{Upsi}_1 = \frac{2^{1/4} \cdot e^{-\frac{\beta p^2}{4} + \frac{I p q}{2 \hbar}}}{2 \left(\frac{\pi}{\beta} \right)^{1/4}}, \text{Upsi}_2 = 0 \right\} \quad (3.1.1)$$

Get the hybrid density matrix

$$\text{simplify} \left(\text{eval} \left(\begin{bmatrix} D_{1,1} & D_{1,2} \\ \text{conjugate}(D_{1,2}) & D_{2,2} \end{bmatrix}, \text{ExampleFreeParticle} \right) \right) \text{ assuming } p :: \mathbb{R}, q :: \mathbb{R}, \hbar > 0, \beta > 0$$

$$\begin{bmatrix} \frac{\sqrt{2} \sqrt{\beta} e^{-\frac{\beta p^2}{2}}}{2 \sqrt{\pi}} & 0 \\ 0 & 0 \end{bmatrix} \quad (3.1.2)$$

Example 2: harmonic oscillator classical Boltzmann state

$$H_0 := \frac{p^2}{2 \cdot m} + \frac{m \cdot \omega^2 \cdot q^2}{2} :$$

$$\text{ExampleHarmonicOscillator} := \left\{ \Upsilon_1 = \frac{\sqrt{\frac{\omega}{2 \cdot \text{Pi} \cdot \beta}}}{H_0} \cdot \sqrt{1 - (1 + \beta \cdot H_0) \cdot \exp(-\beta \cdot H_0)}, \Upsilon_2 \right.$$

$$\left. = 0 \right\} :$$

Solving the special case of a classical harmonic oscillator quadratically coupled to a degenerate two-level quantum system

Analytically derived exact solution

Take the random unitary matrix U

#RandM := Matrix((i,j) → rand(), 2, 2) :

$$\#U := \text{simplify} \left(\text{LinearAlgebra:-MatrixExponential} \left(\frac{I \cdot (\text{RandM} + \text{RandM}^*)}{\text{rand}(\)} \right) \right) :$$

Example from the paper

$$U := \begin{bmatrix} \frac{1}{2} \sqrt{2} & \frac{1}{2} \sqrt{2} \\ -\frac{1}{2} \sqrt{2} & \frac{1}{2} \sqrt{2} \end{bmatrix} :$$

Coupling term

$$\boldsymbol{\alpha\sigma} := \text{simplify}(\lambda \cdot U^* \cdot \sigma_3 \cdot U) :$$

Construct the Hybrid equation to be solved

$$\text{HarmonicOscEq} := \text{eval} \left(\text{HybridEq}, \left\{ H = H_0, f_1 = \frac{q^2}{2} \cdot \frac{\text{Trace}(\boldsymbol{\alpha\sigma} \cdot \sigma_1)}{2}, f_2 = \frac{q^2}{2} \cdot \frac{\text{Trace}(\boldsymbol{\alpha\sigma} \cdot \sigma_2)}{2}, f_3 = \frac{q^2}{2} \cdot \frac{\text{Trace}(\boldsymbol{\alpha\sigma} \cdot \sigma_3)}{2} \right\} \right) :$$

Introducing notations

$$\omega_+ := \sqrt{\omega^2 + \frac{\lambda}{m}} : \omega_- := \sqrt{\omega^2 - \frac{\lambda}{m}} :$$

In the following exact solution, $F_1 :$ and $F_2 :$ denotes for the initial condition for $Y :$

$$\left(\text{subs}(t = 0, Y_1) = F_1(q, p), \text{subs}(t = 0, Y_2) = F_2(q, p) \right) :$$

$$\#y := (\Omega) \rightarrow U. \begin{bmatrix} F_1(q, p) \\ F_2(q, p) \end{bmatrix}$$

$$\left| \begin{array}{l} q = q \cdot \cos(\Omega \cdot t) - \frac{p \cdot \sin(\Omega \cdot t)}{m \cdot \Omega}, p = p \cdot \cos(\Omega \cdot t) + m \cdot \Omega \cdot q \cdot \sin(\Omega \cdot t) \end{array} \right. :$$

Using initial condition from the paper

$$y := (\Omega) \rightarrow U. \left[\begin{array}{c} \frac{\sqrt{\frac{\omega}{2 \cdot \text{Pi} \cdot \beta}}}{H_0} \cdot \sqrt{1 - (1 + \beta \cdot H_0) \cdot \exp(-\beta \cdot H_0)} \\ 0 \end{array} \right]$$

$$q = q \cdot \cos(\Omega \cdot t) - \frac{p \cdot \sin(\Omega \cdot t)}{m \cdot \Omega}, p = p \cdot \cos(\Omega \cdot t) + m \cdot \Omega \cdot q \cdot \sin(\Omega \cdot t)$$

Here is the sought exact solution

$$Y := U^* \cdot \begin{bmatrix} y(\omega_+) _1 \\ y(\omega_-) _2 \end{bmatrix} :$$

$$ExSolution := \{Y_1 = Y_1, Y_2 = Y_2\} :$$

Verifying the initial conditions

$$\#simplify(eval(eval(Y_1, ExSolution), t = 0) - F_1(q, p))$$

$$\#simplify(eval(eval(Y_2, ExSolution), t = 0) - F_2(q, p))$$

Verifying the exact solution

$$\#simplify(combine(pdetest(ExSolution, HarmonicOscEq)))$$

Exact solution of Aleksandrov–Gerasimenko equation

Defyning equation

$$PoissonBracket := (f, g) \rightarrow \frac{\partial}{\partial q} (f) \cdot \frac{\partial}{\partial p} (g) - \frac{\partial}{\partial p} (f) \cdot \frac{\partial}{\partial q} (g) :$$

$$AleksandrovGerasimenkoEq := (DM, H) \rightarrow -\frac{\partial}{\partial t} (DM) - \frac{I}{\hbar} \cdot (H \cdot DM - DM \cdot H) + \frac{1}{2} \cdot PoissonBracket(H, DM) - \frac{1}{2} \cdot PoissonBracket(DM, H) :$$

Generic Initial condition

$$\#InitDM := \begin{bmatrix} c_{1,1}(q, p) & c_{1,2}(q, p) \\ c_{2,1}(q, p) & c_{2,2}(q, p) \end{bmatrix} :$$

Initial condition from the paper

$$InitDM := \begin{bmatrix} \frac{\omega \cdot \beta}{2 \cdot \pi} \cdot \exp(-\beta \cdot H_0) & 0 \\ 0 & 0 \end{bmatrix} :$$

$$d := (\Omega) \rightarrow (U \cdot InitDM \cdot U^*)$$

$$\begin{aligned} & \left| \begin{aligned} q &= q \cdot \cos(\Omega \cdot t) - \frac{p \cdot \sin(\Omega \cdot t)}{m \cdot \Omega}, p = p \cdot \cos(\Omega \cdot t) + m \cdot \Omega \cdot q \cdot \sin(\Omega \cdot t) \end{aligned} \right. : \\ \varphi &:= \frac{\lambda}{2 \cdot m \cdot \hbar \cdot \omega^3} \cdot \left(\frac{p^2 - (m \cdot \omega \cdot q)^2}{2 \cdot m} \cdot \sin(2 \cdot \omega \cdot t) - \omega \cdot (2 \cdot t \cdot H_0 + p \cdot q \cdot (\cos(2 \cdot \omega \cdot t) - 1)) \right) : \\ DM &:= U^* \cdot \begin{bmatrix} d(\omega_+)_{1,1} & e^{I \cdot \varphi} \cdot d(\omega)_{1,2} \\ e^{-I \cdot \varphi} \cdot d(\omega)_{2,1} & d(\omega_-)_{2,2} \end{bmatrix} \cdot U : \end{aligned}$$

Testing initial condition

$$simplify(eval(DM, t = 0) - InitDM)$$

Testing the full solution

$$\#simplify\left(\text{combine}\left(AleksandrovGerasimenkoEq\left(DM, H_0 + \frac{q^2}{2} \cdot \alpha \sigma\right)\right)\right)$$

Solving Pauli equation and generating code

The exact solution to the Paul equation for an oscillator with frequency Ω and the initial quantum state is the ground state with frequency ω

$$PD := (\Omega) \rightarrow \frac{\sqrt{\Omega} \sqrt{-I} m^{1/4} \omega^{1/4} e^{-\frac{\Omega m q^2 (\cos(\Omega t) \omega I - \sin(\Omega t) \Omega)}{2 \hbar (\Omega \cos(\Omega t) I - \omega \sin(\Omega t))}}}{\hbar^{1/4} \pi^{1/4} \sqrt{-I \Omega \cos(\Omega t) + \omega \sin(\Omega t)}} :$$

$$\Psi := U^* \cdot \begin{bmatrix} U_{1,1} \cdot PD(\omega_+) \\ U_{2,1} \cdot PD(\omega_-) \end{bmatrix} :$$

Perform the Wigner transform

$$PauliParams := \left\{ \hbar = 1, m = 1, \omega = 1, \lambda = \frac{95}{100} \right\} :$$

$$WignerPD := (\Omega) \rightarrow \frac{1}{\pi} \cdot \int_{-\infty}^{+\infty} eval(eval(subs(I=-I, PD(\Omega)), PauliParams), q = q + s) \\ \cdot eval(eval(PD(\Omega), PauliParams), q = q - s) \cdot \exp(2 \cdot I \cdot p \cdot s) ds :$$

$$TotalWigner := simplify(|U_{1,1}|^2 \cdot WignerPD(\omega_+) + |U_{2,1}|^2 \cdot WignerPD(\omega_-)) \text{ assuming } t > 0 :$$

Get code for the Wigner function

CodeGeneration:-Python(TotalWigner, optimize, resultname='W')

$$\left[\begin{array}{cc} \int_{-\infty}^{+\infty} |U_{1,1} \cdot PD(\omega_+)|^2 \Big|_{PauliParams} dq & 0 \\ 0 & \int_{-\infty}^{+\infty} |U_{2,1} \cdot PD(\omega_-)|^2 \Big|_{PauliParams} dq \end{array} \right] \text{ assuming } q \\ :: \mathbb{R}, t > 0$$

$$rho_12 := simplify \left(\int_{-\infty}^{+\infty} simplify \left(eval(U_{1,1} \cdot PD(\omega_+) \cdot conjugate(U_{2,1} \cdot PD(\omega_-)), \right. \right. \\ \left. \left. PauliParams \right) \right) dq \Bigg) \text{ assuming } q :: \mathbb{R}, t > 0 :$$

CodeGeneration:-Python(rho_12, optimize, resultname='rotated_rho_12')

$$simplify \left(\int_{-\infty}^{+\infty} |U_{1,1} \cdot PD(\omega_+)|^2 \Big|_{PauliParams} dq \right) \text{ assuming } q :: \mathbb{R}, t > 0$$

$$simplify \left(\int_{-\infty}^{+\infty} |U_{2,1} \cdot PD(\omega_-)|^2 \Big|_{PauliParams} dq \right) \text{ assuming } q :: \mathbb{R}, t > 0$$

$$plot(\text{Re}(rho_12), t = 0..4)$$

$$plot(\text{Im}(rho_12), t = 0..4)$$

▼ **Code generation for illustration**

Specify the initial condition - the same as in Example 2: harmonic oscillator classical Boltzmann state

$$\begin{aligned} \#F_1 &:= (q, p) \rightarrow \frac{\sqrt{\frac{\omega}{2 \cdot \pi \cdot \beta}}}{H_0} \cdot \sqrt{1 - (1 + \beta \cdot H_0) \cdot \exp(-\beta \cdot H_0)} : \\ \#F_2 &:= (q, p) \rightarrow 0 : \end{aligned}$$

$ParamsToPlot := \{ \hbar = 1, m = 1, \}$:

$SolutionToPlot := eval(ExSolution, ParamsToPlot)$:

One more (redundant) check that the obtain solution satisfies the equation

$pdetest(SolutionToPlot, eval(HarmonicOscEq, ParamsToPlot))$ assuming $\beta > 0, \omega > 0$
CodeGeneration:-Python($eval(eval(D_{1,1}, SolutionToPlot), ParamsToPlot)$, *optimize, resultname*
 = "_D11") assuming $p :: \mathbb{R}, q :: \mathbb{R}, \beta > 0, \omega > 0$

CodeGeneration:-Python($eval(eval(D_{1,2}, SolutionToPlot), ParamsToPlot)$, *optimize, resultname*
 = "_D12") assuming $p :: \mathbb{R}, q :: \mathbb{R}, \beta > 0, \omega > 0$

CodeGeneration:-Python($eval(eval(D_{2,2}, SolutionToPlot), ParamsToPlot)$, *optimize, resultname*
 = "_D22") assuming $p :: \mathbb{R}, q :: \mathbb{R}, \beta > 0, \omega > 0$

Code for plotting the solution for Aleksandrov–Gerasimenko equation

Specify the initial condition - the same as in Example 2: harmonic oscillator classical Boltzmann state

$$\begin{aligned} \#c_{1,1}(q, p) &:= \frac{\omega \cdot \beta}{2 \cdot \pi} \cdot \exp(-\beta \cdot H_0) : \\ \#c_{1,2}(q, p) &:= 0 : \\ \#c_{2,1}(q, p) &:= 0 : \\ \#c_{2,2}(q, p) &:= 0 : \end{aligned}$$

Check the normalization

$int(c_{1,1}(q, p), [q = -\infty .. +\infty, p = -\infty .. +\infty])$ assuming $\beta > 0, m > 0, \omega > 0$

*CodeGeneration:-Python(simplify(eval(DM[1, 1], ParamsToPlot)), optimize, resultname
= "_D11") assuming $p :: \mathbb{R}, q :: \mathbb{R}, \beta > 0, \omega > 0$*

*CodeGeneration:-Python(simplify(eval(DM[1, 2], ParamsToPlot)), optimize, resultname
= "_D12") assuming $p :: \mathbb{R}, q :: \mathbb{R}, \beta > 0, \omega > 0$*

*CodeGeneration:-Python(simplify(eval(DM[2, 2], ParamsToPlot)), optimize, resultname
= "_D22") assuming $p :: \mathbb{R}, q :: \mathbb{R}, \beta > 0, \omega > 0$*