

## IV. Trajectory without drift

*restart* : *with*(*Physics*) : *with*(*Physics*[*Vectors*]) : *with*(*LinearAlgebra*) : *Setup*(*mathematicalnotation* = *true*) :

Making some assumptions

*Setup*(*realobjects* = {*e*, *m*, *c*}) :  
*Setup*(*realobjects* = {*t*, *x*(*t*),  $\Omega(t)$ });  
 $[realobjects = \{\widehat{i}, \widehat{j}, \widehat{k}, \widehat{\phi}, \widehat{r}, \widehat{\rho}, \widehat{\theta}, c, e, m, \phi, r, \rho, t, \theta, x, y, z, \Omega(t), x(t)\}]$  **(1)**

Cliford conjugation

*ClifordConj* := (*X*) →  $\langle \langle X[2, 2] \parallel -X[1, 2] \rangle, \langle -X[2, 1] \parallel X[1, 1] \rangle \rangle$  :

Define a contravariant vector *v* from a paravector *P*

*DefContrVect* := (*v*, *P*) → *Define* $\left( v_{\sim\mu} = \text{Simplify}\left( \text{Array}\left( 1..4, (l) \rightarrow \frac{1}{2} \cdot \text{Trace}(P.Psigma[l]) \right) \right) \right)$  :

Extract the electric field vector from the Faraday matrix

*GetElectricField* := (*F*) → *Simplify* $\left( \text{Vector}\left( 3, (l) \rightarrow \frac{1}{2} \text{Re}(\text{Trace}(Psigma[l].F)) \right) . \text{Vector}([_i, _j, _k]) \right)$  :

Extract the magnetic field vector from the Faraday matrix

*GetMagneticField* := (*F*) → *Simplify* $\left( \text{Vector}\left( 3, (l) \rightarrow \frac{1}{2 \cdot c} \text{Im}(\text{Trace}(Psigma[l].F)) \right) . \text{Vector}([_i, _j, _k]) \right)$  :

The radiated energy per unit of time is given by Larmor's formula

*PowerEmitted* := (*U*) →  $\frac{e^2}{6 \cdot c \cdot \text{Pi} \cdot \epsilon_0} \cdot \frac{w}{2} \cdot \text{Trace}(\text{diff}(U, t).ClifordConj(\text{diff}(U, t)))$  :

The Lorentz factor

local  $\gamma$ :

$$\gamma := \left( 1 - \left( \frac{\text{diff}(x(t), t)}{c} \right)^2 \right)^{-\frac{1}{2}} ;$$

$$\gamma := \frac{1}{\sqrt{1 - \frac{\dot{x}(t)^2}{c^2}}} \quad (2)$$

Set the weight

$w := \gamma$ :

Specifying the eigenspinor

$$R := c \cdot t \cdot \text{Matrix}(\text{Psigma}[0]) + x(t) \cdot \text{Matrix}(\text{Psigma}[1]);$$

$$R := \begin{bmatrix} c t & x(t) \\ x(t) & c t \end{bmatrix} \quad (3)$$

$$U := \frac{\gamma}{c} \cdot \text{diff}(R, t) :$$

$$\Lambda := \text{MatrixFunction}(U, \text{sqrt}(z), z) \cdot \text{MatrixExponential} \left( - \frac{I \cdot \Omega(t) \cdot \text{Matrix}(\text{Psigma}[2])}{2} \right) :$$

The electromagnetic field (w is a parametric measure)

$$F := w \cdot \frac{2 \cdot m \cdot c}{e} \cdot \text{diff}(\Lambda, t) \cdot \text{CliffordConj}(\Lambda) :$$

Extract the electromagnetic fields from the matrix F

$$E_- := \text{GetElectricField}(F) :$$

$$'E_-[1]' = \text{simplify}(\text{Component}(E_-, 1)) \text{ assuming } c > 0 \text{ and } 0 \leq \dot{x}(t) < c;$$

$$\vec{E}_1 = - \frac{c^3 m \ddot{x}(t)}{\sqrt{\dot{x}(t) + c} (\dot{x}(t)^2 - c^2) \sqrt{c - \dot{x}(t)} e} \quad (4)$$

$$'E_-[1]' = \text{simplify}(\text{Component}(E_-, 1)) \text{ assuming } c > 0 \text{ and } -c < \dot{x}(t) \leq 0;$$

$$\vec{E}_1 = - \frac{c^3 m \ddot{x}(t)}{\sqrt{\dot{x}(t) + c} \left( \dot{x}(t)^2 - c^2 \right) \sqrt{c - \dot{x}(t)} e} \quad (5)$$

$$'E\_ [2]' = \text{simplify}(\text{Component}(E\_ , 2));$$

$$\vec{E}_2 = 0 \quad (6)$$

$$'E\_ [3]' = \text{simplify}(\text{Component}(E\_ , 3));$$

$$\vec{E}_3 = - \frac{\dot{x}(t) \dot{\Omega}(t) m |c|^2}{\left( \dot{x}(t)^2 - c^2 \right) e} \quad (7)$$

$$B\_ := \text{GetMagneticField}(F) :$$

$$'B\_ [1]' = \text{simplify}(\text{Component}(B\_ , 1)) \text{ assuming } c > 0 \text{ and } 0 \leq \dot{x}(t) < c;$$

$$\vec{B}_1 = 0 \quad (8)$$

$$'B\_ [1]' = \text{simplify}(\text{Component}(B\_ , 1)) \text{ assuming } c > 0 \text{ and } -c < \dot{x}(t) \leq 0;$$

$$\vec{B}_1 = 0 \quad (9)$$

$$'B\_ [2]' = \text{simplify}(\text{Component}(B\_ , 2));$$

$$\vec{B}_2 = \frac{\dot{\Omega}(t) c^2 m}{\left( \dot{x}(t)^2 - c^2 \right) e} \quad (10)$$

$$'B\_ [3]' = \text{simplify}(\text{Component}(B\_ , 3));$$

$$\vec{B}_3 = 0 \quad (11)$$