### **Dykhne Adiabatic Approximation for Pair-creation**

## Making some assumptions

restart: with (Physics): with (Physics [Vectors]): with (Linear Algebra): Setup (mathematical notation = true):

$$Setup \left( real objects = \left\{ e, m, c, \varepsilon_0 \right\} \right) :$$

$$Setup \left( real objects = \left\{ t, V_x(t), V_y(t), \eta(t), \Omega(t) \right\} \right);$$

$$\left[ real objects = \left\{ \hat{i}, \hat{j}, \hat{k}, \hat{\phi}, \hat{r}, \hat{\rho}, \hat{\theta}, c, e, m, \phi, r, \rho, t, \theta, x, y, z, \varepsilon_0, \Omega(t), V_x(t), V_y(t), \eta(t) \right\} \right] \qquad \textbf{(1.1)}$$

$$alias \left( V_x = V_x(t), V_y = V_y(t), \eta = \eta(t), \Omega = \Omega(t) \right)$$

$$V_x, V_y, \eta, \Omega \qquad \textbf{(1.2)}$$

#### ▼ cRDI mathematical formalism

Cliford conjugation

$$ClifordConj := (X) \rightarrow \langle \langle X[2,2] | -X[1,2] \rangle, \langle -X[2,1] | X[1,1] \rangle \rangle$$
:

Define a contravariant vector v from a paravector P

$$DefContrVect := (v, P) \rightarrow Define \left( v_{\text{mu}} = Simplify \left( Array \left( 1 ...4, (l) \rightarrow \frac{1}{2} \cdot Trace(P ...Psigma[l]) \right) \right) \right) :$$

Extract the 3D vector from a paravector P

$$\textit{ExtractVector} := (P) \rightarrow \frac{1}{2} \cdot \textit{Simplify}(\textit{Vector}(3, (l) \rightarrow \textit{Trace}(\textit{Psigma}[l].P)).\textit{Vector}([\_i, \_j, \_k])) :$$

Get the Faraday matrix from the eigenspinor

$$GetFarayday := (\Lambda, W) \rightarrow w \cdot \frac{2 \cdot m \cdot c}{e} \cdot (diff(\Lambda, t).ClifordConj(\Lambda) + \Lambda.diff(W, t).ClifordConj(W)$$

$$.ClifordConj(\Lambda)):$$

Extract the electric field vector from the Faraday matrix

$$GetElectricField := (F) \rightarrow \frac{1}{2} \cdot Simplify(Vector(3, (l) \rightarrow Re(Trace(Psigma[l].F)))).Vector([\_i, \_j, -i])$$

Extract the magnetic field vector from the Faraday matrix

$$GetMagneticField := (F) \rightarrow \frac{1}{2 \cdot c} \cdot Simplify(Vector(3, (l)) \rightarrow Im(Trace(Psigma[l].F))) \cdot Vector([\_i, \_j, \_k])) :$$

Extracting the scalar and vector parts

$$ScalarPart := (L) \rightarrow \frac{(L + ClifordConj(L))}{2}$$
:

$$VectorPart := (L) \rightarrow \frac{(L-ClifordConj(L))}{2}$$
:

Extract the Poynting vector

PoyntingCliford := 
$$(F) \rightarrow \frac{c \cdot \varepsilon_0}{2} \cdot VectorPart(F.Dagger(F))$$
:

$$PoyntingVector := (F) \rightarrow ExtractVector(PoyntingCliford(F))$$
:

The radiated energy per unit of time is given by Larmor's formula

$$\textit{PowerEmitted} := (U) \rightarrow \frac{-e^2}{6 \cdot c \cdot \text{Pi} \cdot \epsilon_0} \cdot \frac{w^2}{2} \cdot \textit{Trace}(\textit{diff}(U, t). \textit{ClifordConj}(\textit{diff}(U, t))) :$$

# Particle in electromagnetic plane wave combined with homogeneous axial electric fields

The eigenspinor employed

$$\Lambda := MatrixExponential \left( \frac{\eta(t) \cdot Matrix(Psigma[3])}{2} \right) \cdot \begin{bmatrix} 1 & \frac{V_x(t)}{c} - \frac{I \cdot V_y(t)}{c} \\ 0 & 1 \end{bmatrix}$$

$$\Lambda := \begin{bmatrix} \frac{\eta(t)}{2} & \frac{\eta(t)}{2} & \left(\frac{V_x(t)}{c} - \frac{IV_y(t)}{c}\right) \\ & & -\frac{\eta(t)}{2} \end{bmatrix}$$

$$(3.1)$$

Proper velocity

 $U := simplify(\Lambda.Dagger(\Lambda))$ 

$$U := \begin{bmatrix} \frac{e^{\eta(t)} \left(V_{y}(t)^{2} + V_{x}(t)^{2} + c^{2}\right)}{c^{2}} & \frac{-I V_{y}(t) + V_{x}(t)}{c} \\ \frac{I V_{y}(t) + V_{x}(t)}{c} & e^{-\eta(t)} \end{bmatrix}$$
(3.2)

DefContrVect(u, U)

Defined objects with tensor properties

$$\left\{ \mathbf{\gamma}_{\mu}, \mathbf{\sigma}_{\mu}, \partial_{\mu}, g_{\mu, \nu}, u^{\mu}, \delta_{\mu, \nu}, \epsilon_{\alpha, \beta, \mu, \nu} \right\} \tag{3.3}$$

Print proper velocity components

 $|u| \sim 0$   $= expand(u \sim 0)$ 

$$u^{0} = \frac{1}{2 e^{\eta}} + \frac{e^{\eta} V_{x}^{2}}{2 c^{2}} + \frac{e^{\eta} V_{y}^{2}}{2 c^{2}} + \frac{e^{\eta}}{2}$$
(3.4)

 $'u[\sim 1]'=expand(u[\sim 1])$ 

$$u^{1} = \frac{V_{x}}{c} \tag{3.5}$$

 $u[\sim 2] = expand(u[\sim 2])$ 

$$u^2 = \frac{V_y}{c} \tag{3.6}$$

 $'u[\sim 3]'=expand(u[\sim 3])$ 

$$u^{3} = \frac{e^{\eta} V_{y}^{2}}{2 c^{2}} + \frac{e^{\eta} V_{x}^{2}}{2 c^{2}} + \frac{e^{\eta}}{2} - \frac{1}{2 e^{\eta}}$$
(3.7)

Enforce the laboratory time parametrization (w must be fixed as shown below for t to become the laboratory time)

 $w := expand\left(\frac{1}{2} \cdot Trace(U)\right)$ 

$$w := \frac{1}{2 e^{\eta}} + \frac{e^{\eta} V_x^2}{2 c^2} + \frac{e^{\eta} V_y^2}{2 c^2} + \frac{e^{\eta}}{2}$$
 (3.8)

 $expand(u[\sim 0] - w);$ 

### Circular trajectories in plane waves

$$CircTrajParams := \left\{ V_{x} = diff(r \cdot \cos(\omega \cdot t), t), V_{y} = diff(r \cdot \sin(\omega \cdot t), t), \eta = -\ln\left(\operatorname{sqrt}\left(1 + \left(\frac{r \cdot \omega}{c}\right)^{2}\right)\right) \right\};$$

$$CircTrajParams := \left\{ V_x = -r \omega \sin(\omega t), V_y = r \omega \cos(\omega t), \eta = -\frac{\ln\left(1 + \frac{r^2 \omega^2}{c^2}\right)}{2} \right\}$$
 (4.1)

Find trajectroy

 $dsolve(subs(CircTrajParams, w \cdot diff(x(t), t) = u[\sim 1]))$  assuming c > 0;

$$x(t) = \frac{r\cos(\omega t)}{\sqrt{r^2 \omega^2 + c^2}} + CI$$
 (4.2)

 $dsolve(subs(CircTrajParams, w \cdot diff(y(t), t) = u[\sim 2]))$  assuming c > 0 and omega > 0;

$$y(t) = \frac{r\sin(\omega t)}{\sqrt{r^2 \omega^2 + c^2}} + CI$$
 (4.3)

 $dsolve(subs(CircTrajParams, w \cdot diff(z(t), t) = u[\sim 3]));$ 

$$z(t) = \_C1 \tag{4.4}$$

Notice that the equation for critical time does not depend on t; thus, there will be no pari creation.  $simplify(subs(CircTrajParams, u[\sim 0]));$ 

$$\frac{\sqrt{r^2 \omega^2 + c^2}}{|c|} \tag{4.5}$$

## Pair creation in the homogeneous electric field

EParams := 
$$\{V_x = 0, V_y = 0\}$$
;  
EParams :=  $\{V_x = 0, V_y = 0\}$  (5.1)

Calculate the fields

 $F := simplify(subs(EParams, (GetFarayday(subs(EParams, \Lambda), IdentityMatrix(2)))))$  assuming c > 0 and m > 0;

$$F := \begin{bmatrix} \frac{\left(e^{-\eta(t)} + e^{\eta(t)}\right) m c \dot{\eta}(t)}{2 e} & 0\\ 0 & -\frac{\left(e^{-\eta(t)} + e^{\eta(t)}\right) m c \dot{\eta}(t)}{2 e} \end{bmatrix}$$

$$(5.2)$$

GetMagneticField(F);

GetElectricField(F)

$$\frac{\left(e^{-\eta} + e^{\eta}\right) m c \left(\frac{\partial}{\partial t} \eta\right) \hat{k}}{2 e}$$
(5.4)

Find  $\eta$  to make electric homogeneous time independent

$$odetest\left(\eta = \arcsin\left(\frac{E_0 e t}{m \cdot c}\right), coeff\left(GetElectricField(F), \_k\right) = E_0\right);$$

$$0$$
(5.5)

Update params

$$EParams := EParams union \left\{ \eta = \operatorname{arcsinh} \left( \frac{E_0 e t}{m \cdot c} \right) \right\};$$

EParams := 
$$\left\{ V_x = 0, \ V_y = 0, \ \eta = \operatorname{arcsinh}\left(\frac{E_0 \ e \ t}{m \ c}\right) \right\}$$
 (5.6)

Find crytical time

t0 := solve(subs(EParams, w), t);

$$t0 := \frac{\operatorname{Im} c}{E_0 e} \tag{5.7}$$

Calculate the pair creation rate notice  $2 \cdot mc^2$  in front of w

$$\exp\left(-\frac{2\cdot\operatorname{Im}\left(\operatorname{int}\left(\operatorname{subs}\left(\operatorname{EParams},2\cdot\boldsymbol{m}\cdot\boldsymbol{c}^{2}\cdot\boldsymbol{w}\right),t=0..t0\right)\right)}{\hbar}\right);$$

$$e^{-\frac{\pi m^2 c^3 \Re\left(\frac{1}{E_0}\right)}{e \hbar}}$$
(5.8)

### Pair creation in two-color laser field ( $\omega + 3\omega$ )

$$LaserParams := \left\{ V_x = \frac{c}{\operatorname{sqrt}(g)} \cdot \cos(\omega \cdot t), V_y = 0, \eta = 0 \right\}$$

$$LaserParams := \left\{ V_x = \frac{c \cos(\omega t)}{\sqrt{g}}, V_y = 0, \eta = 0 \right\}$$
(6.1)

Calculate the fields

 $F := simplify(subs(LaserParams, (GetFarayday(subs(LaserParams, \Lambda), IdentityMatrix(2)))))$  assuming c > 0 and m > 0;

$$F := \begin{bmatrix} 0 & -\frac{\left(c^2 \cos(\omega t)^2 + 2c^2 g\right) m \omega \sin(\omega t)}{e c g^{3/2}} \\ 0 & 0 \end{bmatrix}$$
 (6.2)

 $B_{-} \coloneqq \textit{simplify}(\textit{GetMagneticField}(F)) \text{ assuming } \omega > 0 \text{ and } g > 1;$ 

$$\overrightarrow{B} := -\frac{\sin(\omega t) \, \omega \, \widehat{j} \, m \, (\cos(2 \, \omega t) + 1 + 4 \, g)}{4 \, e \, g^3 \, |^2}$$

$$\tag{6.3}$$

 $E_{-} \coloneqq \textit{simplify}(\textit{GetElectricField}(F)) \text{ assuming } \omega > 0 \text{ and } g > 1;$ 

$$\vec{E} := -\frac{\sin(\omega t) \,\omega \,\hat{i} \,c \,m \,(\cos(2\,\omega t) + 1 + 4\,g)}{4\,e\,g^{3/2}}$$
(6.4)

Find crytical time

solve(subs(LaserParams, w), t) assuming g > 0;

$$\frac{\frac{\pi}{2} - I \operatorname{arcsinh}(\sqrt{2} \sqrt{g})}{\omega}, \frac{\frac{\pi}{2} + I \operatorname{arcsinh}(\sqrt{2} \sqrt{g})}{\omega}$$
(6.6)

t0 := convert(combine(solve(subs(LaserParams, w), t)[2]), arcsinh) assuming g > 0;

$$t\theta := \frac{\frac{\pi}{2} - I \operatorname{arcsinh}(I\sqrt{-2g})}{\omega}$$
 (6.7)

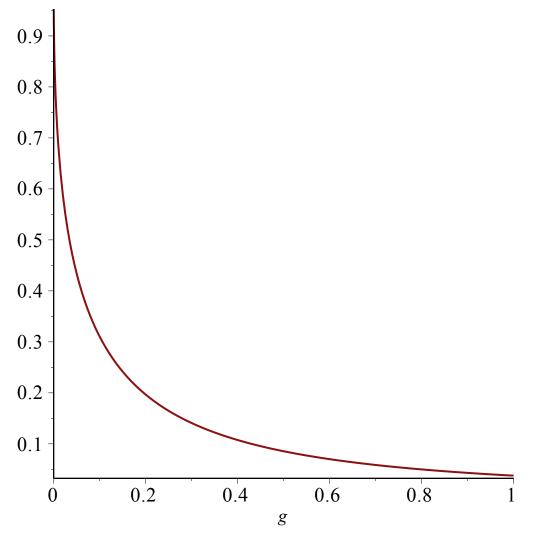
Calculate the pair creation rate notice  $2 \cdot mc^2$  in front of w

$$transition := \exp\left(-\frac{2 \cdot \operatorname{Im}\left(int\left(subs\left(LaserParams, 2 \cdot m \cdot c^2 \cdot w\right), t = 0 ...t0\right)\right)}{\hbar}\right) \text{ assuming } r > 0$$

and c > 0 and  $\omega > 0$  and g > 1;

$$-\frac{mc^2\left(-2\sqrt{2}\sqrt{1+2g}g+8\arcsin\left(\sqrt{2}\sqrt{g}\right)g^3\mid^2+2\arcsin\left(\sqrt{2}\sqrt{g}\right)\sqrt{g}\right)}{2\omega g^3\mid^2\hbar}$$
transition := e (6.8)

 $plot(subs(m = 1, c = 1, omega = 1, \hbar = 1, (transition)), g = 0..1, );$ 



Get the asymptotic expression for large g

$$asympt \left(-\frac{m c^2 \left(-2 \sqrt{2} \sqrt{1+2 g} g+8 \operatorname{arcsinh} \left(\sqrt{2} \sqrt{g}\right) g^{3 / 2}+2 \operatorname{arcsinh} \left(\sqrt{2} \sqrt{g}\right) \sqrt{g}\right)}{2 \omega g^{3 / 2} \hbar}, g, 2\right);$$

$$-\frac{m c^{2} \left(-4+8 \ln \left(2 \sqrt{2}\right)+4 \ln \left(g\right)\right)}{2 \omega \hbar}-\frac{m c^{2} \left(2 \ln \left(2 \sqrt{2}\right)+\ln \left(g\right)\right)}{2 \omega \hbar g}+O\left(\frac{1}{g^{2}}\right)$$
 (6.9)

simplify 
$$\left(\exp\left(-\frac{mc^2(4\ln(g))}{2\omega\hbar}\right)\right)$$
;

simplify 
$$\left(\exp\left(-\frac{mc^2\left(-4+8\ln\left(2\sqrt{2}\right)\right)}{2\omega\hbar}\right)\right)$$
;
$$e^{-\frac{2mc^2\left(-1+3\ln(2)\right)}{\omega\hbar}}$$
(6.11)

$$evalf(-1+3 \ln(2));$$

$$1.079441542$$
(6.12)