

# Characterizing Computational Techniques for Experimentally Feasible Non-Local Evolution of Two-Qubit Systems

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**Abstract:** This paper investigates computational methods and presents a computational framework for simulating non-local evolution of two-qubit systems. This evolution may then easily be combined with finite local operations to give control of near-arbitrary precision. Current research has focused on developing the theoretical background which we hope to computationally implement. We will characterize how useful operators (particularly the *CNOT* gate) are generated in this process, and they will be approximated and performed computationally in optimal time utilizing a adaptive-algorithm approach. We will also analyze the non-local piece of the system Hamiltonian, including determining experimental means of generating it. Particularly, we will be evaluating its action on the Poincaré-Bloch-Sphere.

## Mutation Protocol

Mutations are designated on the basis of chance and magnitude. We set a particular threshold value,  $\alpha$ , which determines whether or not a particular organism's parameters will be mutated. The parameters are the  $\beta'$  rotation coefficients which are specified in the exponential form of the local operations. The other changeable parameter,  $\mu$ , specifies the degree to which the angles will be mutated.

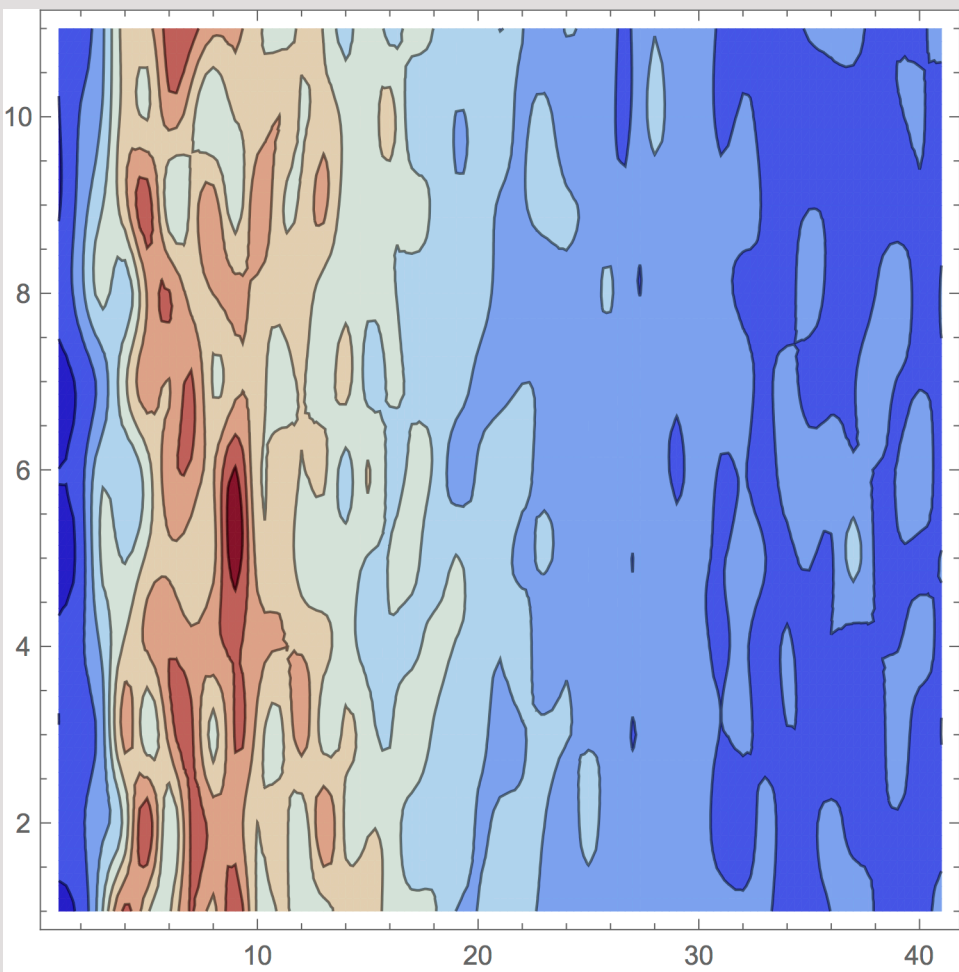


Figure 1: This plot shows the fidelity of the best organism in a ten-organism set after 100 generations with respect to mutation chance and mutation magnitude. The vertical axis, when the value is multiplied by 0.1, gives the chance in any generation for one of the angles defining an organism to be mutated. The horizontal axis, when multiplied by 0.025, gives the mutation magnitude (this ranges between 0 and 1 radian(s)). The plot shows clear correlation between mutation magnitude and organism success.

## The Space as a Whole

A key part to this algorithm's success is knowledge of the search space, and how mobility through it responds to variation in the parameters of its constituent functions. Figure 2 shows the fidelity, on average, of organisms in the *XX* problem.

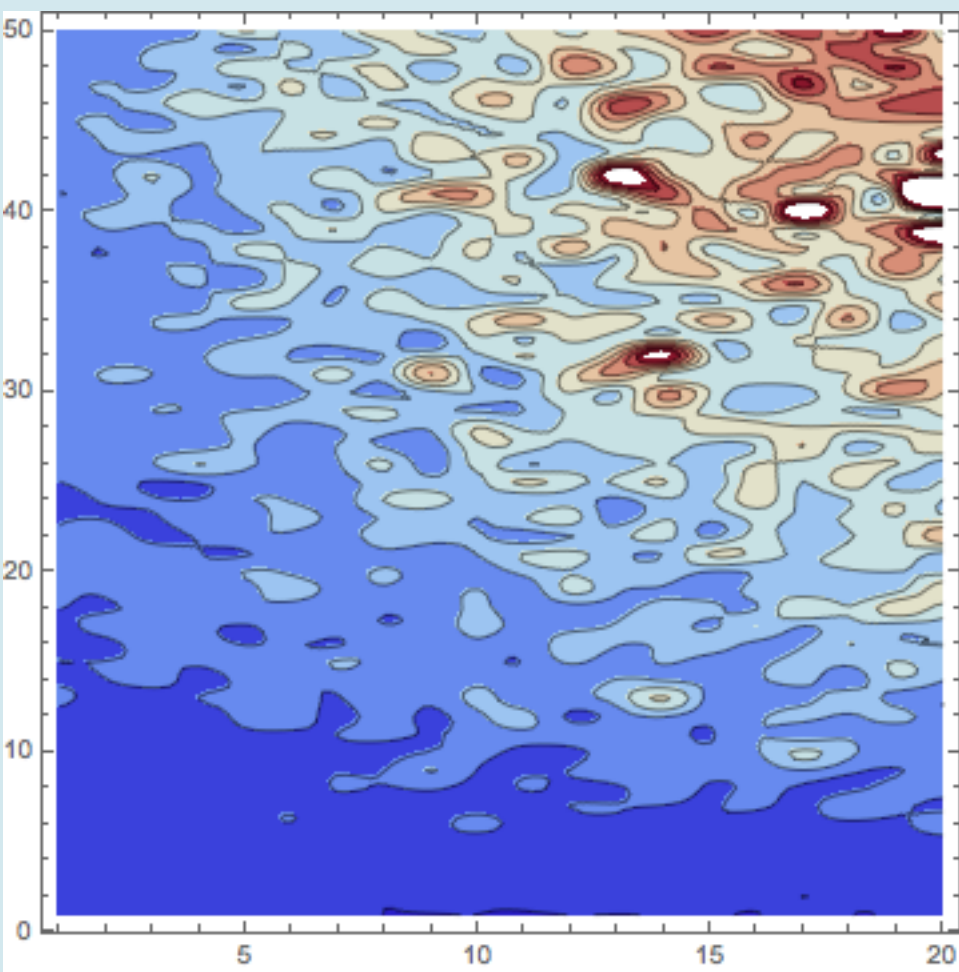
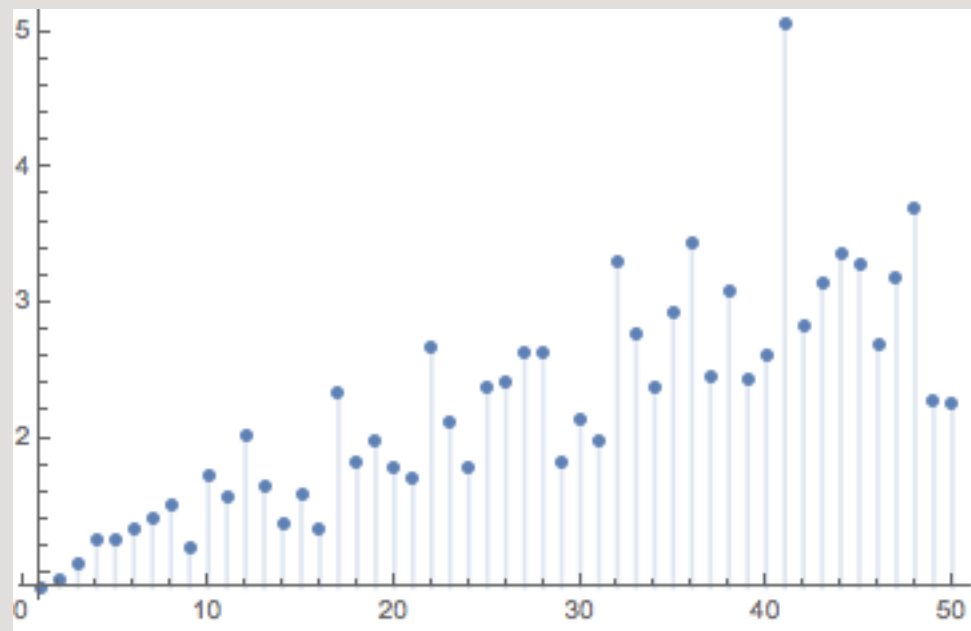


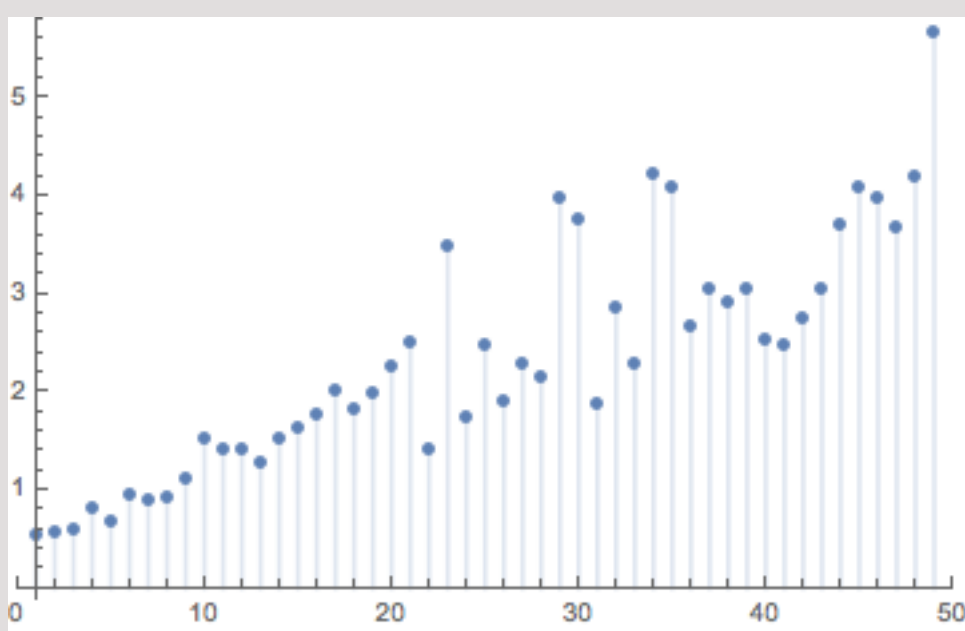
Figure 2: This is a representation of the nature of the search space. Along the horizontal axis is generation number for the set of organisms, and along the vertical is the population of the set. The plot is taken as an average of 20 runs across the entire grid, with the contour plot interpolated between the data at each point; values shown are the fidelity of the best organism on average, and range from 0 to 6.

## Variation in Population and Generation Number

A series of methods is used to determine fidelity in a populations of a set size, generation number, and gene length. We can then produce a series of partially-evaluated functions that can be mapped over a list of integers. The plotted results of this mapping gives knowledge of a slice of the search space, analogous to partial derivatives, which can help determine which variables produce the best solutions with the smallest computational cost.



(a) Lorem ipsum



(b) Lorem ipsum, lorem ipsum, Lorem ipsum, lorem ipsum, Lorem ipsum

## The *XX* Problem

This problem represents a proof-of-concept. They purpose of this scheme is to generate complex unitary evolution operators, and not only is the problem

$$\exp\{\sigma_x \otimes \sigma_x\} \mapsto \exp\{\sigma_x \otimes \sigma_x\} \quad (1)$$

attempting to represent a simple unitary evolution operator, it is attempting to represent the one we began with! Ideally, the results of a high-fidelity simulation in this case would yield identity matrices which produces the same result with two genes. To demonstrate trends in efficiency, we have produced a small table of fidelities, for example.

Pop.	20 Gen.	50 Gen.
10	0.63	0.95
25	0.86	1.71
50	1.17	2.17
100	4.79	2.31

Table 1: A chart of fidelities for the *XX* problem. These values are not averaged, and give some feel as to the variations inherent to each trial.

## The *XY* Problem

In the interest of a less trivial problem, we attempt to move to another simple represented Hamiltonian ( $\sigma_y \otimes \sigma_y$ ). We can represent this concisely by

$$\exp\{\sigma_x \otimes \sigma_x\} \mapsto \exp\{\sigma_y \otimes \sigma_y\} \quad (2)$$

The fidelities can be charted as done in the previous problem, or we may produce a larger chart of the space as in Figure 2.

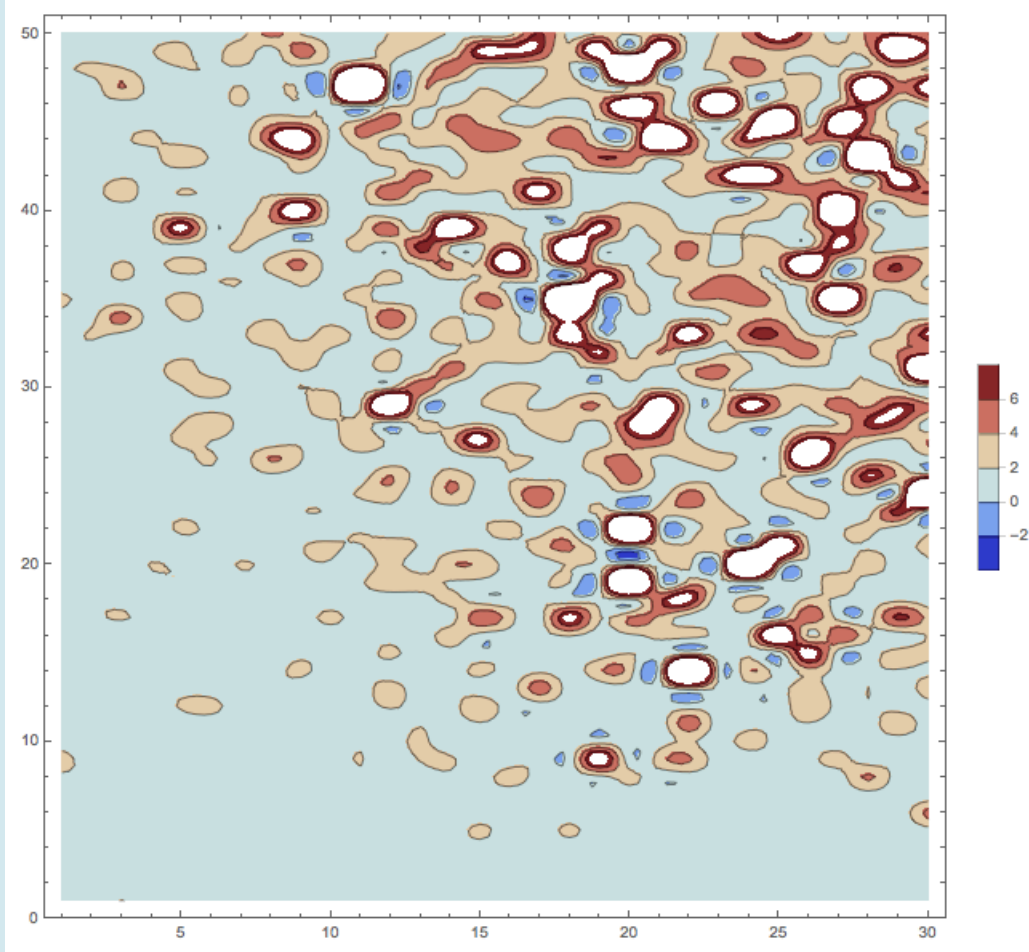


Figure 4: This shows the same plot as Figure 2, but in the context of the *XY* problem. A scale is included to show the fidelities. Note the absence of well defined contours.

As mentioned previously, we can also look at the operators generated from the best organism's set of angles (i.e.  $A_1 \otimes B_1$  and  $A_2 \otimes B_2$ ). Even general trends in these can be helpful in recognizing traits of solutions.

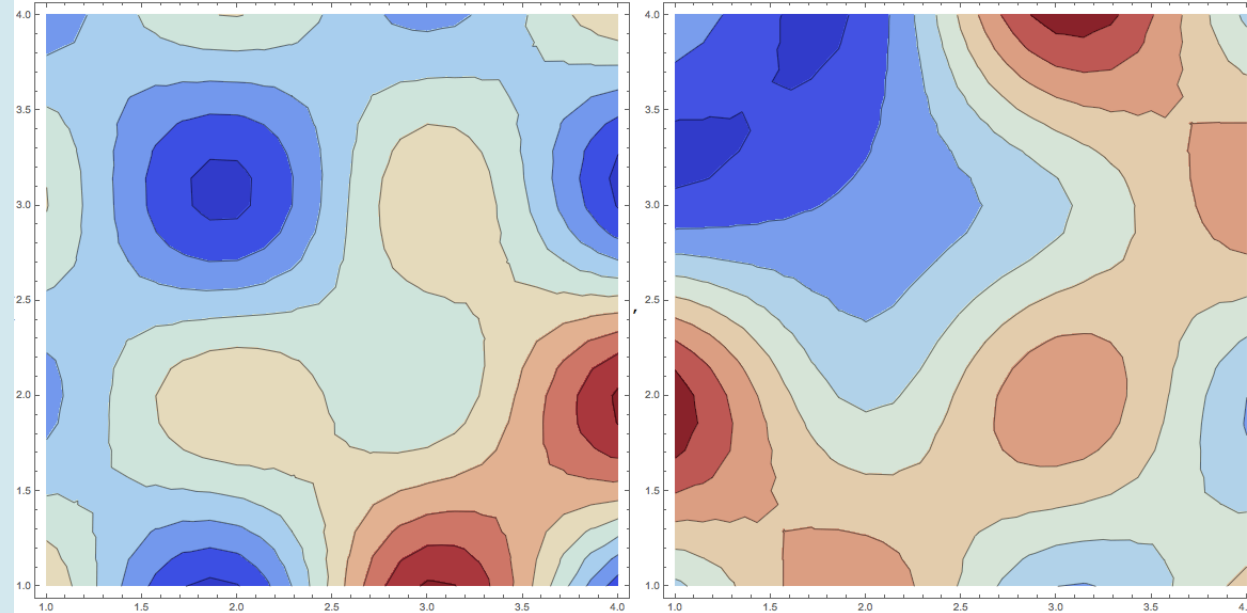


Figure 5: A contour plot of the local operations ( $A_i \otimes B_i$ ) performed by an organism in the *XY* problem. Note that this is an interpolation of a matrix's elements (a four by four grid) and is therefore a useful heuristic, rather than a concrete tool for analysis. The modulus value of the entries were taken, and thus phase factors are lost.

In these types of problems, high fidelities can be reached, but the algorithm is less stable, and generally requires a larger population to produce similar quality.



## Background

The ultimate goal of Control Theory is the ability to arbitrarily manipulate the state of an ensemble of qubits. In the context of quantum computation, this problem is reduced further: basic quantum gates allow for modular systems [3]. However, the generation of these gates is non-trivial, and the arbitrary Hamiltonian is, in general, not readily constructible. Often one will wish to achieve a specific unitary evolution operator or gate  $U$  corresponding to the solitary time evolution of a Hamiltonian  $H$ . When  $H$  may not be engineered easily, the evolution can be achieved in other ways [3], often involving the use of local operations (LO), repeated measurements on the system [2], along with standard Hamiltonian-driven evolution drawing from some collection (possibly a one element collection)  $H_1, H_2, \dots, H_n$ , all of which *can* be lucratively generated.

In the simplest case, a two particle system, the Hamiltonians involved may contain non-local terms, usually corresponding to particle-particle interactions. This paper will propose a framework for taking the general non-local Hamiltonian [2]

$$H = A \otimes \mathbb{1} + \mathbb{1} \otimes B + \sum_{ij} M_{ij} \sigma_i \otimes \sigma_j \quad (3)$$

and interspersing its non-local part's evolution with LO to achieve a desired  $U$ . It is only the non-local part of the Hamiltonian which must be studied, and through decomposition of this piece into its canonical form, and some insight into Lie Theory, we are able to formulate a computational approach [4]. Devising a scheme to generate a  $U$  from a collection of available operators and Hamiltonians has direct applicability to a variety of problems, but cannot be done purely analytically on a massive scale, thus making a computational method necessary.

## Adaptive Algorithm Theory

Adaptive algorithms attempt to mimic the process of evolution by the mechanisms of mutation and natural selection. A set of mathematical objects dubbed *organisms*, each containing some set of data, are subjected to tests against some metric depending on the organism's features [6]. This is considered one *generation*, and the best organisms of each generation are bred together (by some deterministic algorithm) and mutated to compete in the next generation.

$\Theta$  : Arguments

$H$  : Constant

$N$  : Constant

$\{A_i\} \quad \{B_i\}$

$\{t_i\}$

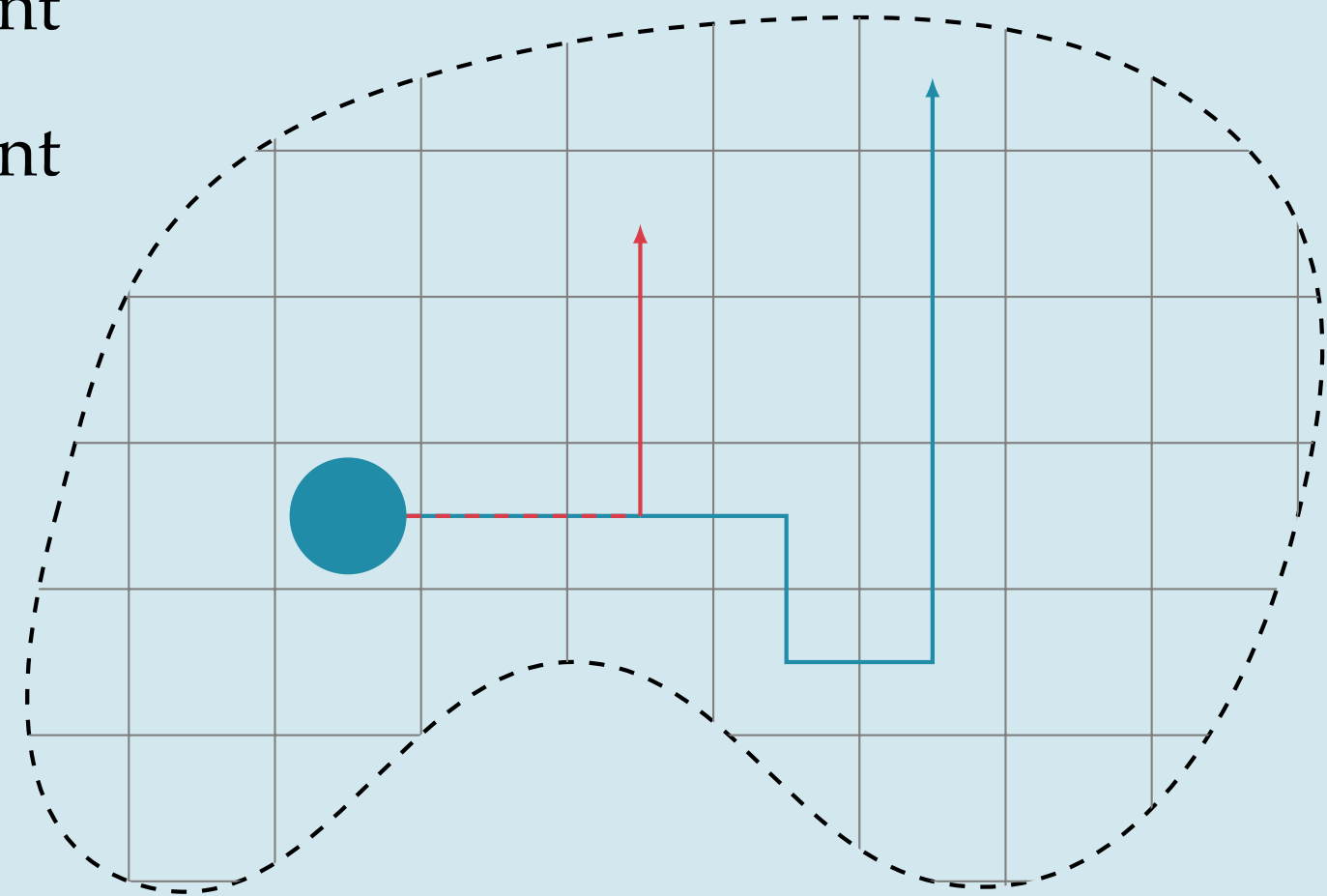


Figure 6: This is a simplification; the organism being tested is represented by the circle in blue, while the search-space (which has a defined metric represented by the grid) surround it. The paths of different color shown represent the non-deterministic paths possible as the organism mutates and is bred with other successful organisms. The actual space being searched is not two- but many-dimensional.

## Group Theory Exploitation

The entirety of this simulation scheme is done with the pauli basis, and this allows their wonderful properties to be exploited. Given further research, the algorithm would implement group theory concepts in its heuristic, but currently, this idea is only used in data compression, and the breeding/mutation of matrices.

The local operations described in the preceding section can be realized as unitary matrices, and described by three angles. Because of this, we can imagine all possible unitary operators as lying on the surface of the 3-Sphere[1].

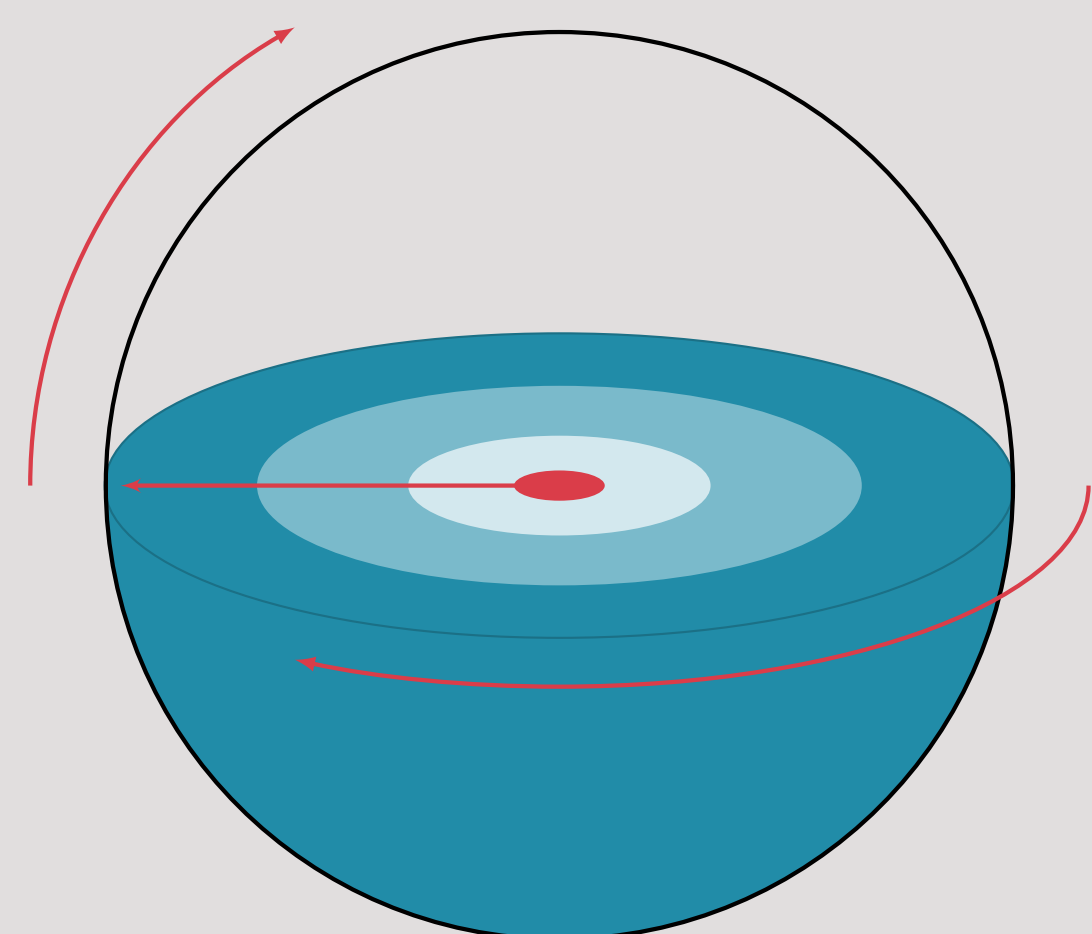


Figure 7: A pictorial representation of the 3-sphere: the upper half has been removed for clarity. The three arrows represent the degrees of freedom along the 3-sphere's surface, the points of which have been projected to the interior of a 2-sphere. (The arrow along the traditional 'radius', if extended, gives degeneracy). Despite implication there is no preferred direction along which to move across the 3-sphere.

## The CNOT Problem

$$\exp\{\sigma_x \otimes \sigma_x\} \mapsto CNOT \quad (4)$$

The purpose of this simulation scheme is to generate methods to create useful quantum operators. These manifest in the canonical quantum gates. Among the most notorious of these (and from which all others can be constructed [? ]) is the *CNOT* gate, represented in matrix form as such

$$CNOT = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} \quad (5)$$

This operator allows one qubit's state to affect the operation applied to another. This is incredibly pertinent in control theory. By the nature of the correlation it implies between particles, this operator is not the tensor product of two local operators. Due to this, the non-local piece of the initial Hamiltonian must take the brunt of the work in producing *CNOT*.

Initial runs yield exotic matrix plots

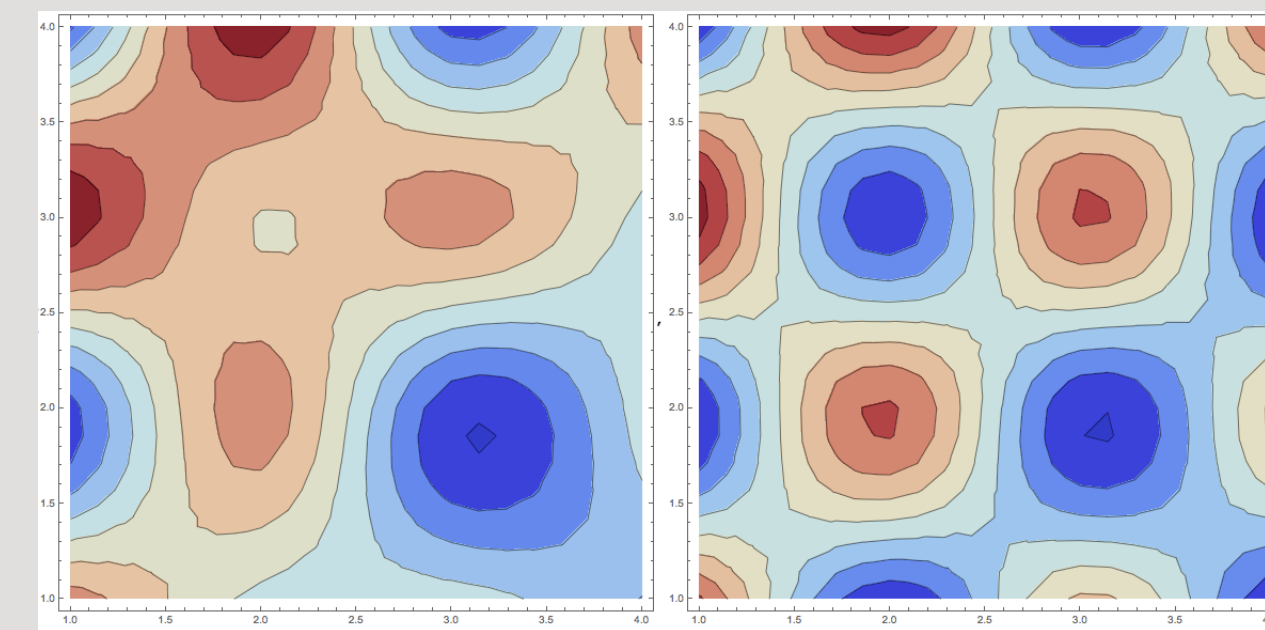


Figure 8: Plots of the LO for the *CNOT* problem. Note that the operators are markedly different, but maintain a concentration of value in the upper off-diagonals.

However, the difficulty of simulating this operator manifests itself in relatively low fidelities given the same generation and population numbers used in the previous problems. There are a few possible ways to increase fidelity (or rather, the potential fidelity) of *CNOT* simulations.

## The 0 Problem

Although there are many simulations possible that are useful to control theory, the final example will be in the simulation of the zero-Hamiltonian, or the stasis of a system. Engineering dynamic Hamiltonians has distinct uses, but preserving the state of the system is unique in two respects: it represents the negation of the original Hamiltonian, and it allows the state of a system to be preserved indefinitely. The problem is represented as such (the zero-Hamiltonian produces the identity as its unitary evolution operator).

$$\exp\{\sigma_x \otimes \sigma_x\} \mapsto \mathbb{1} \quad (6)$$

To give a feel for the fidelity, below is a sampled unitary evolution operator approximation from an average run at 100 organisms and 20 generations.

$$\tilde{U} = \begin{pmatrix} \mathbf{0.86} & 0.07 & 0.07 & 0.10 \\ 0.07 & \mathbf{0.86} & 0.09 & 0.07 \\ 0.17 & 0.10 & \mathbf{1.08} & 0.15 \\ 0.10 & 0.17 & 0.14 & \mathbf{1.08} \end{pmatrix} \quad (7)$$

## Conclusion

We have identified various aspects of using adaptive algorithms that are beneficial and detrimental in determining local operation sets to simulate an arbitrary gate. Although Our results have consistently agreed mathematically, with fairly large fidelity, the run-time grows exponentially with the number of organisms, generations, and gene number. This can only be countered well by running this program with multiple processors (i.e. in parallel, which this algorithm is well-suited for) [5]. Moreover, a lot of lag in the code is derived from matrix multiplication, which is a daunting task to optimize. Perhaps, using a more effective algorithm will improve the run-time efficiency of the code as well. We can however partially ease matrix multiplication by taking advantage of particular expansions of matrix exponentials, which can be truncated depending upon our desired accuracy. Using commutation relations and Lie algebras may immensely help with the reduction of run-time as well (this involves both engineering the heuristic and the initial populations). The methods forming the basis of this scheme are naive, and base little of their workings off of the defining tricks of linear algebra and group theory. There are properties to be exploited, but we chose to produce a functional process before an optimal one. We also wish to represent better ways of representing the resulting operators, to give more intuition about their structure.

Although we compromise time for efficiency, our efforts in implementing adaptive algorithms to solve this quantum computation problem has yielded high success in producing a viable solution sets. With complex systems, as are necessary to produce non-trivial quantum computations, analytical methods for generating the correct unitary evolution operator are simply not feasible. With the proper improvements, non-deterministic methods such as the one here are vital to the success of quantum computation.

## References

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