# 2d to 2.5d transform

We assume that in nature there is no lateral variation along the y axis  $(\partial_y = 0)$ , and we model the synthetic dc electric potential in a true 2-dimensional setting. In 3d, we want to model

$$-\nabla \cdot \sigma(x,z)\nabla u(x,y,z) = s(x,y,z). \tag{1}$$

In the Fourier  $k_y$ -domain we have,

$$-\nabla \cdot \sigma \nabla \tilde{u}(x, k_y, z) + k_y^2 \sigma \tilde{u}(x, k_y, z) = \frac{1}{2} s(x, y, z).$$
 (2)

The 3d solution on the xz-plane is thus

$$u(x, y = 0, z) = \frac{2}{\pi} \int_0^\infty \tilde{u} \, dk_y.$$
 (3)

Discretized, we have

$$u = \frac{2}{\pi} \sum_{i} \tilde{u}(k_i) \,\omega_i,\tag{4}$$

but what are  $k_i$  and  $\omega_i$ ? We follow Pidlisecky and Knight [2008] and proceed to find them by noting that the Green's function solution (for homogeneous  $\sigma$ ) of (1) on the half xz-plane is

$$u(x, y = 0, z) = \frac{\mathbf{i}}{2\pi\sigma} \underbrace{\left(\underbrace{\frac{1}{||x - s_{+}||_{2}} - \underbrace{\frac{1}{||x - s_{-}||_{2}}}_{r/R}}\right)}_{1/R},$$
 (5)

bringing back to the Fourier domain,

$$\tilde{u} = \int_0^\infty u \cos(y \, k_y) \, \mathrm{d}y = \frac{\mathbf{i}}{2\pi\sigma} (B_o(k_y r_+) - B_o(k_y r_-)), \tag{6}$$

where  $B_o$  is the zero order modified Bessel function of the second kind. By plugging in equations 5 and 6 into equation 4 we discretize by

$$1 \approx \sum_{j} \underbrace{\frac{2R}{\pi} \{B_o(k_j r_+) - B_o(k_j r_-)\}}_{K_{i,i}} \omega_j \tag{7}$$

$$K = \frac{2}{\pi} R \left\{ B_o(k r_+) - B_o(k r_-) \right\}$$
 (8)

$$v \approx K\omega,$$
 (9)

where K=K(k,s) is a matrix of size  $(n_R\times n_k)$ , v is a vector of length  $n_R$  whose entries should approximate 1, and  $k=(k_{yi})$ ,  $\omega=(\omega_i)$  are vectors of length  $n_k$ . We minimize

$$\Phi(k) = ||1 - K \underbrace{(K^T K)^{-1} K^T}_{\omega}||_2^2 = ||1 - v(k)||_2^2,$$

using a regularized Newton method. Note that both k and  $\omega$  are geometry dependent and not parameter dependent.

## Finding $k_y$ and $\omega$ for a given s

- 1. initial guess for  $k = (\text{some numbers}) \cdot \Delta x$  and build K(k, s).
- 2.  $v \leftarrow K(K^{\top}K)^{-1}K^{\top} \cdot 1$ .
- 3. compute  $J = \nabla_k v$  using  $n_k$  finite differences,
- 4.  $\nabla_k \Phi^{\top} \leftarrow J^{\top} (1 v) + \beta k$ ,
- 5.  $\Delta k = (J^{\top}J + \beta I)^{-1} \cdot \nabla_k \Phi^{\top}$ ,
- 6.  $k \leftarrow k + \alpha \Delta k$ .
- 7. build K(k,s)
- 8. check if v is almost 1,
- 9. repeat 2-8
- 10.  $\omega = (K^{\top}K)^{-1}K^{\top} \cdot 1$ ,
- 11. correct for flatness  $k \leftarrow k \cdot \Delta x$
- 12. return k and  $\frac{2}{\pi}\omega$ .

# Finding the 2.5d electric potential u for given s and $\sigma$

Given a source s and conductivity  $\sigma$ , the forward model is computed as follows,

1. retrieve from memory (or compute) k and  $\omega$ ,

2. choose  $k_i \in k$  and build  $L_i$  with the right boundary conditions for that  $k_i$ ,

$$L^{i} \approx -\nabla \cdot \sigma \nabla,$$
  
$$L_{i} = L^{i} + k_{i}^{2} \sigma,$$

- 3. solve  $L_i \tilde{u}_i = \frac{s}{2}$  for  $\tilde{u}_i$  and store,
- 4. repeat 3-4 until all k has been used,
- 5.  $u = \sum_{i} \omega_i \, \tilde{u}_i \to d = Mu \to e = d d^o$ .

## 2.5d inversion

Each 2d forward model is,

$$L_i = L^i + k_i^2 \sigma,$$
  

$$L_i \tilde{u}_i = \frac{s}{2},$$
  

$$\tilde{d}_i = M \tilde{u}_i.$$

We can write the 2.5d data and its Jacobian as a linear combination of each 2d problem,

$$d = M \underbrace{\sum_{i} \omega_{i} \tilde{u}_{i}}_{u} = \sum_{i} \omega_{i} \underbrace{M \tilde{u}_{i}}_{\tilde{d}_{i}} = \sum_{i} \omega_{i} \tilde{d}_{i},$$

$$\nabla_{\sigma} d = \underbrace{\sum_{i} \omega_{i} J_{i}}_{I},$$

where

$$J_i = ML_i^{-1}S_i^{\mathsf{T}}, \quad \text{and} \quad S_i = -\left((\nabla_{\sigma}L^i)\tilde{u}_i\right)^{\mathsf{T}} - k_i^2 \operatorname{diag}(\tilde{u}_i)^{\mathsf{T}}.$$

We can now write the gradient of the 2.5d data as a linear combination of the 2d gradients,

$$g = \left(\sum_{i} \omega_{i} J_{i}\right)^{\top} e = \sum_{i} \omega_{i} \underbrace{J_{i}^{\top} e}_{g_{i}} = \sum_{i} \omega_{i} g_{i},$$

where

$$g_i = S_i a_i, \qquad \text{and} \qquad L_i^{\top} a_i = M^{\top} e.$$

#### Finding g for given s and $\sigma$

Given a source s, conductivity  $\sigma$  and weights  $\{k,\omega\}$  the gradient is computed as follows,

- 1. compute the 2.5d forward model to get  $\{\tilde{u}_i\}$  and e,
- 2. choose  $k_i \in k$ ,
- 3. build  $L_i$  and  $S_i = -\left((\nabla_{\sigma}L^i)\tilde{u}_i\right)^{\top} k_i^2 \operatorname{diag}(\tilde{u}_i)^{\top}$ ,
- 4. solve  $L_i^{\top} a_i = M^{\top} e$  for  $a_i$ ,
- 5. compute and store  $g_i = S_i a_i$ ,
- 6. repeat 2-5 until all k has been used,
- 7.  $g = \frac{2}{\pi} \sum_{i} \omega_i g_i$ .

## **Routines**

- dc\_kfourier.m for given source s outputs  $\{k, \omega\}$ .
- dc\_fwd2\_5d.m for a given source s and for each  $k_i \in k$  performs dc\_fwd\_k.m and solves u by weight-stacking  $\{\tilde{u}_{k_i}\}$ . Outputs L, u,  $\{\tilde{u}_{k_i}\}$ , d=Mu and  $e=d-d^o$ .
  - dc\_fwd\_k.m for a given source s and a given  $k_i \in k$  solves the 2d fwd problem  $L_{k_i}\tilde{u}_{k_i} = s$ . Outputs L and  $\tilde{u}_{k_i}$ .
- dc\_gradient2\_5d.m for a given source s, its weights  $\{k,\omega\}$ , its 2d potentials  $\{\tilde{u}_{k_i}\}$ , its matrix L and its 2.5d error e, outputs g as a weighted stack of  $\{g_{k_i}\}$ .

# Verification with analytical models

anal-homo and anal-bi.

## References

Adam Pidlisecky and Rosemary Knight. Fw2\_5d: A matlab 2.5-d electrical resistivity modeling code. *Computers & Geosciences*, 34(12):1645–1654, 2008.