- Discussion of Hugh Thomas 'talk
A probabilistic perspective on \$3 theory

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Amplituhedron Day 8 April 2021 · STOCHASTIC QUANTISATION goes back to [PARISI & WU, 1981]:

Consider Euclidean QFT measure as stationary measure of a stochastic process

- $v(d\phi) := exp(-S(\phi)) \frac{d\phi}{7}$ (formel!) ► static picture:
- functional derivative $\partial_{\epsilon} \phi = \left(\frac{\delta S(\phi)}{\delta \phi} \right) + 3$ Langevin dynamics:

 $\phi = \phi(t, x)$, $t \stackrel{\triangle}{=} fictitious time$ x = space-time

space - time white noise $E[3(4,x)3(s,y)] = \delta(4-s)\delta^{(d)}(x-y)$

- $S(\phi) := \int_{\mathbb{R}^4} \frac{1}{2} |\nabla \phi(x)|^2 + \frac{1}{3} \phi(x)^3 d^{(4)} x$

 $\rightarrow \partial_t \phi = \Delta \phi - \phi^2 + 3 \quad (\phi^3 \text{ equation on } R_+ \times \mathbb{T}^d, \ \phi(0,\cdot) = 0)$

• Fixed - point argument for
$$\square$$
: $\phi = P * (-\phi^2 + 3)$

NB:
$$3 \in C^{-\frac{d+2}{2}-\kappa} as. \Rightarrow 9 \in C^{-3+2-\kappa} a.s. (Schander)$$

· Cure: consider abstract FP problem [Hairy]:

$$\Phi = \mathcal{P}\left(-\Phi^2 + 0\right)$$

$$\frac{1+\sqrt{2}}{2} \quad \Phi(z) = 1 - 2 + \sqrt{2} + \sqrt{2} - 2 + \sqrt{2} + \sqrt{2} + \sqrt{2} = 1 + \sqrt{2} + \sqrt{2} = 1 + \sqrt{2}$$

►
$$\pi^{\epsilon}$$
 := $P * \pi^{\epsilon} \tau$

$$\underline{Ex.:}$$
 $\Pi^{\varepsilon} \Upsilon^{\varepsilon} = P * \Pi^{\varepsilon} \mathcal{P} = P * (\Pi^{\varepsilon})^{2} = P * (\mathcal{E})^{\varepsilon} \stackrel{\triangle}{=} regularised at scale $\varepsilon$$

· Problem in regularity structures:

$$\left(\partial_{\xi} - \Delta\right) \hat{\phi}_{\varepsilon} = -\hat{\phi}_{\varepsilon}^{2} + \vec{\beta}_{\varepsilon} + \sum_{\text{deg}(\varepsilon) < 0} C_{\tau}(\varepsilon) \Upsilon_{\tau}^{\text{RHS}}(\hat{\phi}_{\varepsilon})$$

— Counterterms to ensure ex. of $\hat{\phi} := \lim_{\epsilon \to 0} \hat{\phi}_{\epsilon}$.

* [Bruned, Chandra, Chevyrev, Hairer]

with $C_{\tau}(\varepsilon) := \mathbb{E} \left[\pi^{\varepsilon} A_{-\tau}(0) \right]$

• Explicitly for the ϕ^3 eq. [Beglind & Bruned]:

$$\partial_t \hat{\phi}_{\varepsilon} = \Delta \hat{\phi}_{\varepsilon} - \hat{\phi}_{\varepsilon}^2 + \mathcal{Z}_{\varepsilon} + C_o(\varepsilon) + C_1(\varepsilon) \hat{\phi}_{\varepsilon}$$

• $C_i(\varepsilon) = \mathbb{E} \left[\prod^{\varepsilon} A_{-\tau_i}(0) \right]$ for some $\tau_i \in A_{-\tau_i}(0)$

Example:

- ► Each of the shaded trees gives a (sum of) vaccoum Feynman diagram!

 (TT = acts multiplicatively on forests)
- $= \mathbb{E}\left[\mathbf{T}^{\varepsilon}\right] = \frac{1}{8} \left(\mathbf{0}\right) + \frac{1}{4} \left(\mathbf{0}\right) + \frac{1}{4}$

Let's look at this one in some more detail!

$$\mathbb{E}\left[\mathbf{T}^{\varepsilon} \overset{\circ}{\bigvee} (0)\right] = \mathbb{E}\left[\left(\mathbf{T}^{\varepsilon} \overset{\circ}{\bigvee} (0)\right)^{2}\right] = \mathbb{E}\left[\left(\mathbf{T}^{\varepsilon} \overset{\circ}{\bigvee} (0)\right)^{2}\right] + \text{use Wick's thm!}$$

- ► These computations become unwieldy very quickly!
- We need to exploit cancellations for $\hat{\phi} := \lim_{\varepsilon \downarrow 0} \hat{\phi}_{\varepsilon}$ to exist!

In the spirit of the original motivation:

Can we use AMPLITUHEDRON - like structures to organise our Fynnam diagram computations?

(RS allow to treat locally subcritical = superenormalisable stochastic PDEs, i.e. # [T: deg(T) < 0] < 0.

Ever more important when one wants to study critical SPDEs!)

THANK YOU!