#### Jacobs University Bremen

### Sensors & Filtering



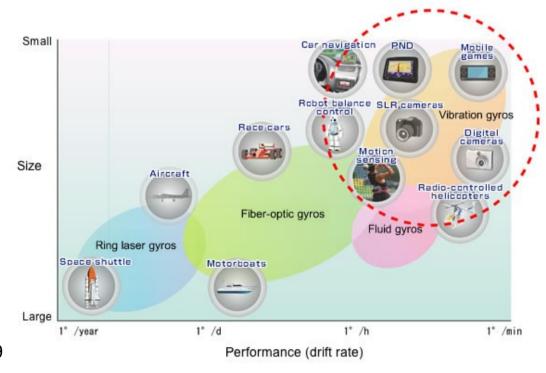
Input to output analysis of a CVG



Automation CO23-320203

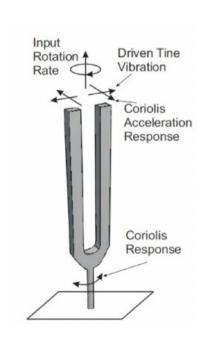
#### Back to gyroscopes

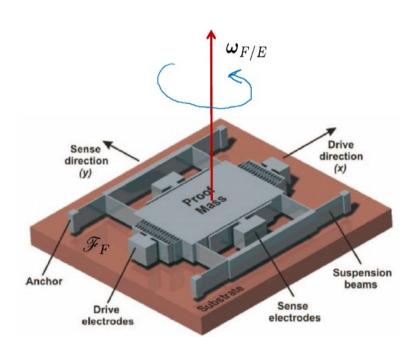
Many different existing gyro technologies strive for their market share



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## Coriolis Vibratory Gyroscope

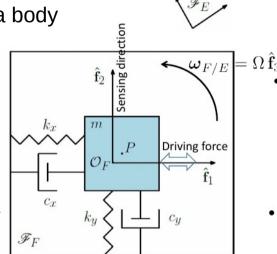




# Coriolis Vibratory Gyroscope idealisation

- Using Coriolis acceleration, which appears on a body in motion in a rotating reference frame, we can estimate the magnitude of rotation!
  - We need to force motion of this mass along one axis (through some drive force F(t))
- The elements of the sensor:
  - Mass m, free to move in the x and y directions
  - (Massless) springs with elastic constants  $k_x$  and  $k_y$
  - Dampers with damping coefficients  $c_x$  and  $c_y$
- Two reference frames will be necessary in the analysis
  - Inertial E
  - Mobile F, attached to the base slab of the sensor, with unit basis vectors

$$\hat{f}_1, \hat{f}_2, \hat{f}_3$$



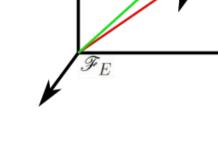
# CVG reference frames setup

- Let us analyse the dynamics of the center of mass P in the intertial frame of reference
- Define the vector  $r_{OEP}$  as a sum of two other characteristic vectors:

$$r_{\mathcal{O}_E P} = r_{\mathcal{O}_E \mathcal{O}_F} + r_{\mathcal{O}_F P}$$

 The quantities in ref. frame F can be expressed in such way:

$$\mathbf{r}_{\mathcal{O}_F P} = x \,\hat{\mathbf{f}}_1 + y \,\hat{\mathbf{f}}_2$$
$$\mathbf{v}_{P/F} = \dot{x} \,\hat{\mathbf{f}}_1 + \dot{y} \,\hat{\mathbf{f}}_2$$
$$\mathbf{a}_{P/F} = \ddot{x} \,\hat{\mathbf{f}}_1 + \ddot{y} \,\hat{\mathbf{f}}_2$$



(notation:  $a_{P/F}$  – acceleration of point P in frame F, etc)

#### CVG dynamics

• The rotation around axis  $\hat{\mathbf{f}}_3$  which we want to measure can be expressed in the following way:

$$\boldsymbol{\omega}_{F/E} = \Omega \,\hat{\mathbf{f}}_3$$

 The only way to use Newton's laws is to apply them in an inertial reference frame. For mass m:

$$m \mathbf{a}_{P/E} = \mathbf{F}_{ext}$$

Let us derive this acceleration. We must start with velocity

$$\mathbf{v}_{P/E} \triangleq \left(\frac{d\mathbf{r}_{\mathcal{O}_E P}}{dt}\right)_E$$

$$\mathbf{a}_{P/E} \triangleq \left(\frac{d\mathbf{v}_{P/E}}{dt}\right)_E$$

$$v_{P/E} = \dot{r}_{\mathcal{O}_E \mathcal{O}_F} + \dot{r}_{\mathcal{O}_F P} = v_{F/E} + \frac{d}{dt} \left( |r_{\mathcal{O}_F P}| \hat{r}_{\mathcal{O}_F P} \right)$$

$$= v_{F/E} + \frac{d}{dt} \left( |r_{\mathcal{O}_F P}| \right) \hat{r}_{\mathcal{O}_F P} + |r_{\mathcal{O}_F P}| \frac{d\hat{r}_{\mathcal{O}_F P}}{dt} = v_{F/E} + \underbrace{\hat{f}_1 \dot{x} + \hat{f}_2 \dot{y}}_{v_{P/F}} + \omega_{F/E} \times r_{\mathcal{O}_F P}$$

$$a_{P/E} = \dot{v}_{P/E} = \dot{v}_{F/E} + \dot{v}_{P/F} + \frac{d}{dt} \left( \omega_{F/E} \times r_{\mathcal{O}_F P} \right)$$

$$= a_{F/E} + \frac{d}{dt} \left( |v_{P/F}| \hat{v}_{P/F} \right) + \frac{d\omega_{F/E}}{dt} \times r_{\mathcal{O}_F P} + \omega_{F/E} \times \frac{dr_{\mathcal{O}_F P}}{dt}$$

$$= a_{F/E} + a_{P/F} + \alpha_{F/E} \times r_{\mathcal{O}_F P} + \omega_{F/E} \times v_{P/F} +$$

$$+ \omega_{F/E} \times \omega_{F/E} \times r_{\mathcal{O}_F P} + \omega_{F/E} \times v_{P/F}$$

with  $\alpha_{F/E}$  – the angular acceleration of frame F in frame E

- Collection of terms of the last line yields a term  $2\omega_{F/E} \times v_{P/F}$  called the Coriolis acceleration
- (see https://www.youtube.com/watch?v=RBVi-4aftNo and related videos to see details of derivation)

 We can assume that the accelerations perceived by our gyroscope in E are not very high:

$$a_{F/E} \approx 0$$

#### **CVG** Coriolis effect

 We have obtained the term which expresses the Coriolis acceleration:

$$2\omega_{F/E} \times \mathbf{v}_{P/F}$$

 Substitute the x/y expressions into the terms of the acceleration equation:

$$2\boldsymbol{\omega}_{F/E} \times \mathbf{v}_{P/F} = 2(\Omega\,\hat{\mathbf{f}}_3) \times (\dot{x}\,\hat{\mathbf{f}}_1 + \dot{y}\,\hat{\mathbf{f}}_2) = 2(\Omega\,\dot{x}\,\hat{\mathbf{f}}_2 - \Omega\,\dot{y}\,\hat{\mathbf{f}}_1)$$

$$\boldsymbol{\omega}_{F/E} \times \boldsymbol{\omega}_{F/E} \times \mathbf{r}_{O_FP} = \Omega^2 \,\hat{\mathbf{f}}_3 \times (\hat{\mathbf{f}}_3 \times (x \,\hat{\mathbf{f}}_1 + y \,\hat{\mathbf{f}}_2)) = -\Omega^2 (x \,\hat{\mathbf{f}}_1 + y \,\hat{\mathbf{f}}_2)$$

Let us look at the external force in more detail now:

$$\mathbf{F}_{ext} = -(k_x x + c_x \dot{x})\hat{\mathbf{f}}_1 - (k_y y + c_y \dot{y})\hat{\mathbf{f}}_2 + F(t)\hat{\mathbf{f}}_1$$

After all the substitutions and plugging into Newton's law we obtain:

$$a_{P/E} = \frac{1}{m} F_{ext}$$

• When we collect the terms containing  $\hat{\mathbf{f}}_1$  and  $\hat{\mathbf{f}}_2$ 

$$(\ddot{x} - 2\Omega\dot{y} - \dot{\Omega}y - \Omega^2x)\hat{\mathbf{f}}_1 + (\ddot{y} + 2\Omega\dot{x} + \dot{\Omega}x - \Omega^2y)\hat{\mathbf{f}}_2$$
$$= -(\frac{k_x}{m}x + \frac{c_x}{m}\dot{x})\hat{\mathbf{f}}_1 - (\frac{k_y}{m}y + \frac{c_y}{m}\dot{y})\hat{\mathbf{f}}_2 + \frac{1}{m}F(t)\hat{\mathbf{f}}_1$$

 $\rightarrow$  forced motion along the  $\hat{\mathbf{f}}_1$  axis

$$\ddot{x} - 2\Omega\dot{y} - \dot{\Omega}y - \Omega^2x + \frac{k_x}{m}x + \frac{c_x}{m}\dot{x} = \frac{1}{m}F(t)$$

→ dynamics of the output axis

$$\ddot{y} + 2\Omega\dot{x} + \dot{\Omega}x - \Omega^2y + \frac{k_y}{m}y + \frac{c_y}{m}\dot{y} = 0$$

some cross-terms are present

For the input axis:

$$\ddot{x} - 2\Omega\dot{y} - \dot{\Omega}y - \Omega^2x + \frac{k_x}{m}x + \frac{c_x}{m}\dot{x} = \frac{1}{m}F(t)$$

• Simplification #2: let us assume that  $\Omega$  is slowly varying, thus  $\dot{\Omega}\approx 0$ . We obtain:

$$\ddot{x} + (\frac{k_x}{m} - \Omega^2)x + \frac{c_x}{m}\dot{x} - 2\Omega\dot{y} = \frac{1}{m}F(t)$$

Simplification #3: equal coefficients c and k on both axes:

$$k_x = k_y = k$$
,  $c_x = c_y = c$ 

• Let us write the characteristic frequency of this system as  $\omega_n$  and the quality factor as Q. Observe that

$$\omega_n = \sqrt{\frac{k}{m}}$$
  $Q = \frac{\sqrt{km}}{c}$   $\frac{\omega_n}{Q} = \sqrt{\frac{k}{m}} \times \frac{c}{\sqrt{km}} = \frac{c}{m}$ 

We can rewrite the input axis equation as:

$$\ddot{x} + (\omega_n^2 - \Omega^2)x + \frac{\omega_n}{Q}\dot{x} - 2\Omega\dot{y} = \frac{1}{m}F(t)$$

• Simplification #4: suppose that the measured angular rate ( $\sim 1$  rad/s) will be greatly inferior to the natural frequency of the system ( $\sim 10^4 - 10^5$  rad/s)

$$\ddot{x} + \omega_n^2 x + \frac{\omega_n}{Q} \dot{x} - 2\Omega \dot{y} = \frac{1}{m} F(t)$$

• If we "pump" a sinusoidal force at the natural frequency in the x-axis of the form  $F(t) = F_D \sin(\omega_n t)$ , we can expect a steady oscillation in this direction:

$$x(t) = x_D \sin(\omega_n t + \phi)$$

#### CVG - input to output

What about the output equation?

$$\ddot{y} + 2\Omega\dot{x} + \dot{\Omega}x - \Omega^2y + \frac{k}{m}y + \frac{c}{m}\dot{y} = 0$$

Using simplifications #2 and #4 we obtain:

$$\ddot{y} + \frac{\omega_n}{Q}\dot{y} + \omega_n^2 y = -2\Omega\dot{x}$$

 By applying Laplace transform to both sides, we can look at the frequency characteristics of the system:

$$\frac{Y(s)}{X(s)} = \frac{-2\Omega s}{s^2 + \frac{\omega_n}{Q}s + \omega_n^2}$$

• From here, we are close to the result of the input vs the output amplitude:  $\frac{Y_D}{X_D} = \frac{2\,\Omega\,Q}{\omega_n}$