

The Local Density of States

There's this thing called the "Single Particle Green's Function" or "Feynman Propagator" or "Correlation Function". It is defined as

$$G(\mathbf{r}', t', \mathbf{r}, t) = -i \langle 0 | T \{ \Psi(\mathbf{r}, t) \Psi^\dagger(\mathbf{r}', t') \} | 0 \rangle \quad (1)$$

Here, \mathbf{r} and \mathbf{r}' are position vectors, $|0\rangle$ is the many body ground state with N electrons, and $\Psi^\dagger(\mathbf{r}', t')$ creates an electron while $\Psi(\mathbf{r}, t)$ destroys an electron. Ignore the T thing; it's called "time ordering" and is not relevant here. I will let $t' = 0$ and $t > 0$ just 'cause I can, and I will set $\hbar = 1$. In the Heisenberg picture, operators evolve as

$$\Psi(\mathbf{r}, t) = e^{iHt} \Psi(\mathbf{r}) e^{-iHt} \quad (2)$$

$$\Psi^\dagger(\mathbf{r}, t) = e^{-iHt} \Psi^\dagger(\mathbf{r}) e^{iHt} \quad (3)$$

so, putting this all together

$$G(\mathbf{r}', \mathbf{r}, t) = -i \langle 0 | \Psi(\mathbf{r}) e^{-iHt} \Psi^\dagger(\mathbf{r}') | 0 \rangle = -i \langle \mathbf{r} | e^{-iHt} | \mathbf{r}' \rangle \quad (4)$$

This has a physical interpretation. It is the probability amplitude that an electron starts at position \mathbf{r}' , then is subject to time evolution for time t , and ends up at position \mathbf{r} . Now, lets suppose we have a complete set of $N + 1$ electron energy eigenstates denoted by $\{|\lambda\rangle\}$. They have energies $\{E_\lambda\}$. The resolution of the identity says $1 = \sum_\lambda |\lambda\rangle \langle \lambda|$.

$$G(\mathbf{r}', \mathbf{r}, t) = -i \sum_{\lambda\lambda'} \langle \mathbf{r} | \lambda \rangle \langle \lambda | e^{-iHt} | \lambda' \rangle \langle \lambda' | \mathbf{r}' \rangle = -i \sum_\lambda e^{-iE_\lambda t} \psi_\lambda^*(\mathbf{r}') \psi_\lambda(\mathbf{r}) \quad (5)$$

Now let $\mathbf{r}' = \mathbf{r}$ because STM injects/removes electrons at the same position that it makes its measurement. Also, an STM measurement is "slow" in that a measurement timescale is much longer than the time it takes for dynamics on the surface to happen. Thus, we should Fourier Transform the Green's Function.

$$G(\mathbf{r}, \omega) = -i \sum_\lambda \int_0^\infty e^{i(\omega - E_\lambda)t} e^{-\delta t} |\psi_\lambda(\mathbf{r})|^2 dt \quad (6)$$

Note that I put in $e^{-\delta t}$ to help the integral converge. Physically, $1/\delta$ is the lifetime of a quasiparticle excitation. We will take the limit wherer δ goes to zero at the end of the calculation.

$$G(\mathbf{r}, \omega) = -i \sum_\lambda \left[\frac{e^{i(\omega - E_\lambda)t} e^{-\delta t}}{i(\omega - E_\lambda) - \delta} \right]_0^\infty |\psi_\lambda(\mathbf{r})|^2 = - \sum_\lambda \left[\frac{1}{\omega - E_\lambda + i\delta} \right] |\psi_\lambda(\mathbf{r})|^2 \quad (7)$$

Define the "Spectral Function"

$$A(\mathbf{r}, \omega) = \frac{1}{\pi} \text{Im}[G(\mathbf{r}, \omega)] = \frac{1}{\pi} \sum_\lambda \frac{\delta}{(\omega - E_\lambda)^2 + \delta^2} |\psi_\lambda(\mathbf{r})|^2 \quad (8)$$

You might notice the similarity to the ARPES spectra function

$$A(k, \omega) = \frac{1}{\pi} \text{Im}[G(k, \omega)] = \frac{1}{\pi} \frac{\text{Im}\Sigma}{(\omega - E_k - \text{Re}\Sigma)^2 + (\text{Im}\Sigma)^2} \quad (9)$$

Anyway, take the $\delta \rightarrow 0$ limit, and $\frac{1}{\pi} \sum_{\lambda} \frac{\delta}{(\omega - E_{\lambda})^2 + \delta^2} \rightarrow \delta(\omega - E_{\lambda})$, where $\delta(\omega - E_{\lambda})$ is the Dirac Delta Function. Finally, we get the Local Density of States

$$A(\mathbf{r}, \omega) = \sum_{\lambda} |\psi_{\lambda}(\mathbf{r})|^2 \delta(\omega - E_{\lambda}) = \text{LDOS}(\mathbf{r}, \omega) \quad (10)$$

I probably lost some proportionality factors, but they don't matter because dI/dV is only proportional to LDOS anyway.