Asymmetric [Ellipsoid] active particle:

An asymmetric active particle undergoes self-propulsion with a velocity $v_p \hat{\mathbf{n}}(t)$ along its longer axis in 2-dimension. The angle between the x-axis of the lab frame and long axis of the ellipsoid at time t is represented by $\theta(t)$. The component of the orientation vector $\hat{\mathbf{n}}(t)$ w.r.t the lab frame can be expressed as $\hat{\mathbf{n}}(t) \equiv \Big(\cos\theta(t), \sin\theta(t)\Big)$. The coupled Langevin equation in the presence of an external force F and torque τ are

$$\partial_t r_i = \Gamma_{ij}(\theta(t)) F_j + v_p \hat{n}_i(t) + \xi_i(t),$$

$$where \Gamma_{ij}(\theta(t)) = \bar{\Gamma} \delta_{ij} + \frac{\Delta \Gamma}{2} M_{ij}(\theta(t))$$

$$\partial_t \theta = \Gamma_{\theta} \tau + \xi_{\theta}(t)$$

$$(1)$$

$$ar{\Gamma} = rac{\Gamma_{\parallel} + \Gamma_{\perp}}{2}, \Delta \Gamma = \Gamma_{\parallel} - \Gamma_{\perp} \ and \ M_{ij}(heta(t)) = egin{bmatrix} \cos 2 heta(t) & sin2 heta(t) \ sin2 heta(t) & -\cos 2 heta(t) \end{bmatrix}$$

 $\xi_i(t)$ and ξ_{θ} are the Gaussian random noise with zero mean and

$$\langle \xi_{ heta}(t) \xi_{ heta}(t\prime)
angle = 2K_B T \Gamma_{ heta} \delta(t-t\prime) = 2D_r \delta(t-t\prime)$$

for a fixed angle $\theta(t)$,

$$\langle \xi_i(t)\xi_i(t\prime)\rangle = 2K_BT \Gamma_{ij}(\theta(t)) \delta(t-t\prime)$$

1. Calculation of $\langle \hat{n}_i(t_1)\hat{n}_j(t_2)\rangle_{\theta_0}$, the average is taken for fixed initial angle θ_0

$$\theta(t_1) \equiv \theta_1, \theta(t_2) \equiv \theta_2$$

$$\langle \hat{n}_i(t_1)\hat{n}_j(t_2)\rangle_{\theta_0} = \begin{bmatrix} \langle \cos\theta_1\cos\theta_2\rangle & \langle \cos\theta_1\sin\theta_2\rangle \\ \langle \sin\theta_1\cos\theta_2\rangle & \langle \sin\theta_1\sin\theta_2\rangle \end{bmatrix}$$
(3)

$$2\langle\cos\theta_1\cos\theta_2\rangle = e^{-D_r(t_1+t_2-2\min(t_1,t_2))} + \cos 2\theta_0 e^{-D_r(t_1+t_2+2\min(t_1,t_2))}$$

$$2\langle \sin \theta_1 \sin \theta_2 \rangle = e^{-D_r(t_1 + t_2 - 2\min(t_1, t_2))} - \cos 2\theta_0 e^{-D_r(t_1 + t_2 + 2\min(t_1, t_2))}$$

$$2\langle\cos\theta_1\sin\theta_2\rangle = \sin 2\theta_0 e^{-D_r(t_1+t_2+2\min(t_1,t_2))}$$

$$2\langle\cos heta_2\sin heta_1
angle=\sin2 heta_0e^{-D_r(t_1+t_2+2\min(t_1,t_2))}$$

Substituting, we have

$$\langle \hat{n}_i(t_1)\hat{n}_j(t_2)\rangle_{\theta_0} = \frac{\delta_{ij}}{2}e^{-D_r(t_1+t_2-2\min(t_1,t_2))} + \frac{M_{ij}(\theta_0)}{2}e^{-D_r(t_1+t_2+2\min(t_1,t_2))}$$
(4)

Without Trap

2. Langevin equation:

$$\frac{\partial r_i}{\partial t} = v_p \hat{n}_i(t) + \xi_i(t)
\Longrightarrow r_i(t) - r_i(0) = \int_0^t v_p \hat{n}_i(t') dt' + \int_0^t \xi_i(t') dt'
\Longrightarrow \Delta r_i(t) = \int_0^t v_p \hat{n}_i(t') dt' + \int_0^t \xi_i(t') dt'$$
(5)

Mean Square Displacement (MSD):-

Similarly,

$$\langle \Delta y^{2}(t) \rangle_{\theta_{0}} = 2\tau_{r}v_{p}^{2}[t - \tau_{1}(t)] - \frac{2\tau_{r}v_{p}^{2}}{3}\cos 2\theta_{0}[\tau_{1}(t) - \tau_{4}(t)] + 2K_{B}T\,\bar{\Gamma}t - K_{B}T\Delta\Gamma\,\cos 2\theta_{0}\tau_{4}(t) \quad (7)$$

Therefore,

$$\begin{split} \langle \Delta r^2(t) \rangle_{\theta_0} &= \langle \Delta x^2(t) \rangle_{\theta_0} + \langle \Delta y^2(t) \rangle_{\theta_0} \\ \langle \Delta r^2(t) \rangle_{\theta_0} &= 4 (K_B T \,\bar{\Gamma} + \tau_r v_p^2) t - 4 \tau_r v_p^2 \tau_1(t) \end{split} \tag{8}$$

3. The correlation $\langle \Delta x(t_1) \Delta x(t_2) \rangle_{\theta_0}$

$$\begin{split} \langle \Delta x(t_{1})\Delta x(t_{2})\rangle_{\theta_{0}} &= v_{p}^{2} \int_{0}^{t_{1}} dt_{1}' \int_{0}^{t_{2}} dt_{2}' \left\langle \hat{n}_{x}(t_{1}') \hat{n}_{x}(t_{2}') \right\rangle_{\theta_{0}} + \int_{0}^{t_{1}} dt_{1}' \int_{0}^{t_{2}} dt_{2}' \left\langle \xi_{x}(t_{1}') \xi_{x}(t_{2}') \right\rangle_{\theta_{0}} \\ \langle \Delta x(t_{1})\Delta x(t_{2})\rangle_{\theta_{0}} &= v_{p}^{2} \int_{0}^{t_{1}} dt_{1}' \int_{0}^{t_{2}} dt_{2}' \left[\frac{1}{2} e^{-D_{r}(t_{1}'+t_{2}'-2\min(t_{1}',t_{2}'))} + \frac{M_{xx}(\theta_{0})}{2} e^{-D_{r}(t_{1}'+t_{2}'+2\min(t_{1}',t_{2}'))} \right] \\ &+ 2K_{B}T \int_{0}^{t_{1}} dt_{1}' \int_{0}^{t_{2}} dt_{2}' \left\langle \Gamma_{xx}(\theta(t_{1}')) \right\rangle_{\theta_{0}} \delta(t_{1}'-t_{2}') \\ \langle \Delta x(t_{1})\Delta x(t_{2})\rangle_{\theta_{0}} &= v_{p}^{2} \int_{0}^{t_{2}} dt_{2}' \int_{t_{2}'}^{t_{2}} dt_{1}' \frac{1}{2} e^{-D_{r}(t_{1}'-t_{2}')} + v_{p}^{2} \int_{0}^{t_{2}} dt_{2}' \int_{t_{2}'}^{t_{1}} dt_{1}' \frac{1}{2} e^{-D_{r}(t_{1}'-t_{2}')} \\ &+ v_{p}^{2} \frac{M_{xx}(\theta_{0})}{2} \int_{0}^{t_{2}} dt_{2}' \int_{t_{2}'}^{t_{2}} dt_{1}' e^{-D_{r}(t_{1}'+3t_{2}')} + v_{p}^{2} \frac{M_{xx}(\theta_{0})}{2} \int_{0}^{t_{2}} dt_{2}' \int_{t_{2}'}^{t_{1}} dt_{1}' e^{-D_{r}(t_{1}'+3t_{2}')} \\ &+ 2K_{B}T \int_{0}^{t_{1}} dt_{1}' \int_{0}^{t_{2}} dt_{2}' \left[\overline{\Gamma} + \frac{\Delta\Gamma}{2} \left\langle \cos 2\theta(t_{1}') \right\rangle \right)_{\theta_{0}} \right] \delta(t_{1}' - t_{2}') \\ \langle \Delta x(t_{1})\Delta x(t_{2})\rangle_{\theta_{0}} &= v_{p}^{2}\tau_{r} \left[2t_{2} - \tau_{1}(t_{2}) - 2\tau_{r} \left(e^{\frac{-(t_{1}-t_{2})}{2\tau_{r}}} - e^{\frac{-t_{1}}{2\tau_{r}}} \right) \right] + \cos 2\theta_{0}\tau_{r}v_{p}^{2} \left[\frac{2\tau_{4}(t_{2})}{3} - \frac{\tau_{1}(t_{2})}{3} - e^{\frac{-t_{1}}{2\tau_{r}}} \tau_{3}(t_{2}) \right] \\ &+ 2K_{B}T \, \overline{\Gamma}t_{2} + K_{B}T\Delta\Gamma \cos 2\theta_{0}\tau_{4}(t_{2}) \end{split}$$

In the presence of harmonic trap

Potential confinement, $U(x,y) = \kappa(x^2 + y^2)/2$

Langevin Equation in the presence of harmonic trap,

$$\frac{\partial x}{\partial t} = -\kappa x \left[\bar{\Gamma} + \frac{\Delta \Gamma}{2} \cos 2\theta(t) \right] - \frac{\kappa y}{2} \Delta \Gamma \sin 2\theta(t) + v_p \hat{n}_x(t) + \xi_x(t)$$

$$\frac{\partial y}{\partial t} = -\frac{\kappa x}{2} \Delta \Gamma \sin 2\theta(t) - \kappa y \left[\bar{\Gamma} - \frac{\Delta \Gamma}{2} \cos 2\theta(t) \right] + v_p \hat{n}_y(t) + \xi_y(t)$$
(10)

Perturbative Expansion:

Define $R \equiv (x, y)^T$, the above equation can be reduce as

$$\dot{R}(t) = -\kappa \Big[\bar{\Gamma} \mathbb{I} + \frac{\Delta \Gamma}{2} \mathbb{M}(t) \Big] R(t) + v_p \hat{\mathbf{n}}(t) + \xi(t)$$
where $\mathbb{M}(t) = \begin{bmatrix} \cos 2\theta(t) & \sin 2\theta(t) \\ \sin 2\theta(t) & -\cos 2\theta(t) \end{bmatrix}$ (11)

Perturbative expansion is

$$R(t) = R_0(t) - \left(\frac{\kappa \Delta \Gamma}{2}\right) R_1(t) + \left(\frac{\kappa \Delta \Gamma}{2}\right)^2 R_2(t) + \mathcal{O}\left(\frac{\kappa \Delta \Gamma}{2}\right)^3$$
(12)

Substituting equation 12 in 11, we obtain

$$\dot{R}_{0}(t) = -\kappa \bar{\Gamma} R_{0}(t) + \xi(t) + v_{p} \hat{\mathbf{n}}(t)
\dot{R}_{1}(t) = -\kappa \bar{\Gamma} R_{1}(t) + \mathbb{M}(t) R_{0}(t)
\dot{R}_{2}(t) = -\kappa \bar{\Gamma} R_{2}(t) + \mathbb{M}(t) R_{1}(t)$$
(13)

Taking R(0)=0, solving the above equations

$$R_{0}(t) = \int_{0}^{t} e^{-\kappa \bar{\Gamma}(t-t')} \xi(t') dt' + v_{p} \int_{0}^{t} e^{-\kappa \bar{\Gamma}(t-t')} \hat{\mathbf{n}}(t') dt'$$

$$R_{1}(t) = \int_{0}^{t} e^{-\kappa \bar{\Gamma}(t-t')} \mathbb{M}(t') R_{0}(t') dt'$$

$$R_{2}(t) = \int_{0}^{t} e^{-\kappa \bar{\Gamma}(t-t')} \mathbb{M}(t') R_{1}(t') dt'$$

$$(14)$$

The equal time correlation matrix

$$\langle R_{i}(t)R_{j}(t)\rangle_{\xi,\theta_{0}} = \langle R_{0,i}(t)R_{0,j}(t)\rangle_{\xi,\theta_{0}} - \left(\frac{\kappa\Delta\Gamma}{2}\right)\langle R_{0,i}(t)R_{1,j}(t)\rangle_{\xi,\theta_{0}} + \left(\frac{\kappa\Delta\Gamma}{2}\right)^{2} \left[\langle R_{1,i}(t)R_{1,j}(t)\rangle_{\xi,\theta_{0}} + 2\langle R_{0,i}(t)R_{2,j}(t)\rangle_{\xi,\theta_{0}}\right] + \mathcal{O}\left(\frac{\kappa\Delta\Gamma}{2}\right)^{3}$$

$$(15)$$

4. Calculation of $\langle R_0(t)R_0(t)\rangle_{\xi,\theta_0}$

$$\langle R_{0}(t)R_{0}(t)\rangle_{\xi,\theta_{0}} = \int_{0}^{t} dt' \int_{0}^{t} dt'' e^{-\kappa\bar{\Gamma}(t-t')} e^{-\kappa\bar{\Gamma}(t-t'')} \langle \xi(t')\xi(t'')\rangle_{\xi,\theta_{0}}$$

$$+ v_{p}^{2} \int_{0}^{t} dt' \int_{0}^{t} dt'' e^{-\kappa\bar{\Gamma}(t-t')} e^{-\kappa\bar{\Gamma}(t-t'')} \langle \hat{\mathbf{n}}(t')\hat{\mathbf{n}}(t'')\rangle_{\xi,\theta_{0}}$$

$$\langle R_{0}(t)R_{0}(t)\rangle_{\xi,\theta_{0}} = 2K_{B}Te^{-2\kappa\bar{\Gamma}t} \int_{0}^{t} dt' \int_{0}^{t} dt'' e^{\kappa\bar{\Gamma}(t'+t'')} \langle \bar{\Gamma}\mathbb{I} + \frac{\Delta\Gamma}{2}\mathbb{M}(\theta(t'))\rangle_{\xi,\theta_{0}} \delta(t'-t'')$$

$$+ v_{p}^{2}e^{-2\kappa\bar{\Gamma}t} \int_{0}^{t} dt' \int_{0}^{t} dt'' e^{\kappa\bar{\Gamma}(t'+t'')} \left[\frac{\mathbb{I}}{2}e^{-D_{r}(t'+t''-2\min(t',t''))} + \frac{\mathbb{M}(\theta_{0})}{2}e^{-D_{r}(t'+t''+2\min(t',t''))} \right]$$

$$\text{Considering the case of } t' > t'',$$

$$\langle R_{0}(t)R_{0}(t)\rangle_{\xi,\theta_{0}} = \frac{K_{B}T}{\kappa} \mathbb{I} \left(1 - e^{-2\kappa\bar{\Gamma}t} \right) + K_{B}T\Delta\Gamma\mathbb{M}(\theta_{0}) \left(\frac{e^{-4D_{r}t} - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma} - 4D_{r}} \right)$$

$$+ \frac{v_{p}^{2}e^{-2\kappa\bar{\Gamma}t}}{2} \mathbb{I} \times 2 \int_{0}^{t} dt'' \int_{t''}^{t} dt' e^{\kappa\bar{\Gamma}(t'+t'')} e^{-D_{r}(t'-t'')}$$

$$+ \frac{v_{p}^{2}e^{-2\kappa\bar{\Gamma}t}}{2} \mathbb{M}(\theta_{0}) \times 2 \int_{0}^{t} dt'' \int_{t''}^{t} dt' e^{\kappa\bar{\Gamma}(t'+t'')} e^{-D_{r}(t'+3t'')}$$

$$\langle R_{0}(t)R_{0}(t)\rangle_{\xi,\theta_{0}} = \frac{K_{B}T}{\kappa} \mathbb{I} \left(1 - e^{-2\kappa\bar{\Gamma}t} \right) + K_{B}T\Delta\Gamma\mathbb{M}(\theta_{0}) \left(\frac{e^{-4D_{r}t} - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma} - 4D_{r}} \right)$$

$$+ v_{p}^{2} \mathbb{I} \left[\frac{1 - e^{-(D_{r}+\kappa\bar{\Gamma})t}}{(\kappa\bar{\Gamma})^{2} - D_{r}^{2}} - \frac{1 - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma}(\kappa\bar{\Gamma} - D_{r})} \right] + \frac{v_{p}^{2}\mathbb{M}(\theta_{0})}{(\kappa\bar{\Gamma} - D_{r})} \left[\frac{e^{-4D_{r}t} - e^{-(D_{r}+\kappa\bar{\Gamma})t}}{(\kappa\bar{\Gamma} - 3D_{r})} - \frac{e^{-4D_{r}t} - e^{-2\kappa\bar{\Gamma}t}}{2(\kappa\bar{\Gamma} - 2D_{r})} \right]$$

$$\langle x_0^2(t) \rangle_{\xi,\theta_0} = \frac{K_B T}{\kappa} \left(1 - e^{-2\kappa\bar{\Gamma}t} \right) + K_B T \Delta\Gamma \cos 2\theta_0 \left(\frac{e^{-4D_r t} - e^{-2\kappa\Gamma t}}{2\kappa\bar{\Gamma} - 4D_r} \right)$$

$$+ v_p^2 \left[\frac{1 - e^{-(D_r + \kappa\bar{\Gamma})t}}{(\kappa\bar{\Gamma})^2 - D_r^2} - \frac{1 - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma}(\kappa\bar{\Gamma} - D_r)} \right] + \frac{v_p^2 \cos 2\theta_0}{(\kappa\bar{\Gamma} - D_r)} \left[\frac{e^{-4D_r t} - e^{-(D_r + \kappa\bar{\Gamma})t}}{(\kappa\bar{\Gamma} - 3D_r)} - \frac{e^{-4D_r t} - e^{-2\kappa\bar{\Gamma}t}}{2(\kappa\bar{\Gamma} - 2D_r)} \right]$$
and
$$\langle y_0^2(t) \rangle_{\xi,\theta_0} = \frac{K_B T}{\kappa} \left(1 - e^{-2\kappa\bar{\Gamma}t} \right) - K_B T \Delta\Gamma \cos 2\theta_0 \left(\frac{e^{-4D_r t} - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma} - 4D_r} \right)$$

$$+ v_p^2 \left[\frac{1 - e^{-(D_r + \kappa\bar{\Gamma})t}}{(\kappa\bar{\Gamma})^2 - D_r^2} - \frac{1 - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma}(\kappa\bar{\Gamma} - D_r)} \right] - \frac{v_p^2 \cos 2\theta_0}{(\kappa\bar{\Gamma} - D_r)} \left[\frac{e^{-4D_r t} - e^{-(D_r + \kappa\bar{\Gamma})t}}{(\kappa\bar{\Gamma} - 3D_r)} - \frac{e^{-4D_r t} - e^{-2\kappa\bar{\Gamma}t}}{2(\kappa\bar{\Gamma} - 2D_r)} \right]$$

The cross correlation function

$$\langle x_0(t)y_0(t)\rangle_{\xi,\theta_0} = \Delta D \sin 2\theta_0 \left(\frac{e^{-4D_r t} - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma} - 4D_r}\right) + \frac{v_p^2 \sin 2\theta_0}{(\kappa\bar{\Gamma} - D_r)} \left[\frac{e^{-4D_r t} - e^{-(D_r + \kappa\bar{\Gamma})t}}{(\kappa\bar{\Gamma} - 3D_r)} - \frac{e^{-4D_r t} - e^{-2\kappa\bar{\Gamma}t}}{2(\kappa\bar{\Gamma} - 2D_r)}\right]$$
(18)

where $\Delta D = K_B T \Delta \Gamma$

5. Calculation of $\langle R_{0,i}(t)R_{1,j}(t)\rangle_{\xi,\theta_0}$

$$\langle R_{0,i}(t)R_{1,j}(t)\rangle_{\xi,\theta_{0}} = \left\langle R_{0,i}(t)\int_{0}^{t} e^{-\kappa\bar{\Gamma}(t-t')} \sum_{k} M_{jk}(t')R_{0}(t')dt' \right\rangle_{\xi,\theta_{0}}$$

$$\langle R_{0,i}(t)R_{1,j}(t)\rangle_{\xi,\theta_{0}} = \left\langle \int_{0}^{t} dt' e^{-\kappa\bar{\Gamma}(t-t')} \sum_{k} M_{jk}(t')\langle R_{0,i}(t)R_{0,j}(t')\rangle_{\xi} \right\rangle_{\theta_{0}}$$

$$(19)$$

$$\langle R_{0,i}(t)R_{0,j}(t')\rangle_{\xi} = \frac{K_B T}{\kappa} \delta_{ij} \left[e^{-\kappa\bar{\Gamma}|t-t'|} - e^{-\kappa\bar{\Gamma}(t+t')} \right]$$

$$+ K_B T \Delta \Gamma e^{-\kappa\bar{\Gamma}(t+t')} \int_0^{\min(t,t')} dt'_1 e^{-\kappa\bar{\Gamma}t'_1} M_{ij}(t'_1) + v_p^2 e^{-\kappa\bar{\Gamma}(t+t')} \int_0^t dt'_1 \int_0^{t'} dt'_2 e^{\kappa\bar{\Gamma}(t'_1+t'_2)} \langle \hat{n}_i(t'_1)\hat{n}_j(t'_2) \rangle_{\xi}$$
(20)

Substitute the value of $\langle R_{0,i}(t)R_{0,j}(t')\rangle_{\xi}$ in equation 19,

$$\langle R_{0,i}(t)R_{1,j}(t)\rangle_{\xi,\theta_{0}} = \frac{K_{B}T}{\kappa}M_{ji}(\theta_{0})e^{-2\kappa\bar{\Gamma}t}\int_{0}^{t}dt' \left[e^{(2\kappa\bar{\Gamma}-4D_{r})t'} - e^{-4D_{r}t'}\right]$$

$$+K_{B}T\Delta\Gamma e^{-2\kappa\bar{\Gamma}t}\int_{0}^{t}dt'\int_{0}^{t'}dt'_{1}\left\langle \sum_{k}M_{jk}(t')M_{ik}(t'_{1})\right\rangle_{\theta_{0}}$$

$$+v_{p}^{2}e^{-2\kappa\bar{\Gamma}t}\int_{0}^{t}dt'\int_{0}^{t}dt'_{1}\int_{0}^{t'}dt'_{2}e^{\kappa\bar{\Gamma}(t'_{1}+t'_{2})}\left\langle \sum_{k}M_{jk}(t')\hat{n}_{i}(t'_{1})\hat{n}_{k}(t'_{2})\right\rangle_{\theta_{0}}$$
(21)

For the MSD along x and y directions, i=j,

$$\left\langle \sum_{k} M_{ik}(t') M_{ik}(t'_{1}) \right\rangle_{\theta_{0}} = e^{-4D_{r}(t'-t'_{1})}$$

$$\left\langle \sum_{k} M_{ik}(t') \hat{n}_{i}(t'_{1}) \hat{n}_{k}(t'_{2}) \right\rangle_{\theta_{0}} = e^{-D_{r}(4t'-3t'_{1}-t'_{2})}$$
(22)

where we use the case of $\ t'>t_1'>t_2'$ to solve the above equation.

The third term of equation 21:

$$\begin{split} v_{p}^{2}e^{-2\kappa\bar{\Gamma}t} \int_{0}^{t} dt' \int_{0}^{t} dt'_{1} \int_{0}^{t'} dt'_{2} e^{\kappa\bar{\Gamma}(t'_{1}+t'_{2})} e^{-D_{r}(4t'-3t'_{1}-t'_{2})} &= v_{p}^{2}e^{-2\kappa\bar{\Gamma}t} \int_{0}^{t} dt' e^{-4D_{r}t'} \int_{0}^{t} dt'_{2} e^{(\kappa\bar{\Gamma}+3D_{r})t'_{1}} e^{(\kappa\bar{\Gamma}+D_{r})t'_{2}} \\ &= v_{p}^{2}e^{-2\kappa\bar{\Gamma}t} \int_{0}^{t} dt' e^{-4D_{r}t'} \left[\int_{0}^{t'} dt'_{2} e^{(\kappa\bar{\Gamma}+D_{r})t'_{2}} \int_{t'_{2}}^{t} dt'_{1} e^{(\kappa\bar{\Gamma}+3D_{r})t'_{1}} + \int_{0}^{t'} dt'_{2} e^{(\kappa\bar{\Gamma}+D_{r})t'_{2}} \int_{t'_{2}}^{t'} dt'_{1} e^{(\kappa\bar{\Gamma}+3D_{r})t'_{1}} \right] \\ &= v_{p}^{2}e^{-2\kappa\bar{\Gamma}t} \int_{0}^{t} dt' e^{-4D_{r}t'} \left[\int_{0}^{t'} dt'_{2} e^{(\kappa\bar{\Gamma}+D_{r})t'_{2}} \int_{t'_{2}}^{t} dt'_{1} e^{(\kappa\bar{\Gamma}+3D_{r})t'_{1}} - \int_{0}^{t} dt'_{2} e^{(\kappa\bar{\Gamma}+3D_{r})t'_{1}} - \frac{(e^{(2\kappa\bar{\Gamma}+3D_{r})t'_{1}} - 1)}{(\kappa\bar{\Gamma}+3D_{r})} (\kappa\bar{\Gamma}+D_{r})} - \frac{e^{(2\kappa\bar{\Gamma}+3D_{r})t'_{1}} - e^{(2\kappa\bar{\Gamma}+3D_{r})t'_{1}} - e^{(2\kappa\bar{\Gamma}+3D_{r})t'_$$

The contribution to the mean square displacement along the x direction is

$$\langle x_{0}(t)x_{1}(t)\rangle_{\xi,\theta_{0}} = \frac{K_{B}T}{\kappa}\cos 2\theta_{0}\left(\frac{e^{-4D_{r}t} - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma} - 4D_{r}} - \frac{e^{-2\kappa\bar{\Gamma}t} - e^{-(2\kappa\bar{\Gamma}+4D_{r})t}}{4D_{r}}\right)$$

$$+ \Delta D\left(\frac{1 - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma}(2\kappa\bar{\Gamma} + 4D_{r})} - \frac{e^{-2\kappa\bar{\Gamma}t} - e^{-(2\kappa\bar{\Gamma}+4D_{r})t}}{4D_{r}(2\kappa\bar{\Gamma} + 4D_{r})}\right) + \frac{v_{p}^{2}}{(\kappa\bar{\Gamma} + 3D_{r})}\left[\frac{1}{(\kappa\bar{\Gamma} - 3D_{r})}\left(\frac{3(\kappa\bar{\Gamma} - D_{r})}{(\kappa\bar{\Gamma} + D_{r})2\kappa\bar{\Gamma}} - \frac{e^{-(\kappa\bar{\Gamma}-3D_{r})t}}{4D_{r}}\right)\right]$$

$$+ \frac{1}{2(\kappa\bar{\Gamma} - D_{r})}\left(\frac{e^{-(\kappa\bar{\Gamma}+D_{r})t}(\kappa\bar{\Gamma} - 5D_{r})}{2D_{r}(\kappa\bar{\Gamma} + D_{r})}\right) - \frac{1}{(\kappa\bar{\Gamma} + 2D_{r})}\left(\frac{1 - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma}} - \frac{e^{-2\kappa\bar{\Gamma}t} - e^{-2(\kappa\bar{\Gamma}+2D_{r})t}}{4D_{r}}\right)\right]$$

and to the y direction is

$$\langle y_{0}(t)y_{1}(t)\rangle_{\xi,\theta_{0}} = -\frac{K_{B}T}{\kappa}\cos2\theta_{0}\left(\frac{e^{-4D_{r}t} - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma} - 4D_{r}} - \frac{e^{-2\kappa\bar{\Gamma}t} - e^{-(2\kappa\bar{\Gamma}+4D_{r})t}}{4D_{r}}\right)$$

$$+ \Delta D\left(\frac{1 - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma}(2\kappa\bar{\Gamma} + 4D_{r})} - \frac{e^{-2\kappa\bar{\Gamma}t} - e^{-(2\kappa\bar{\Gamma}+4D_{r})t}}{4D_{r}(2\kappa\bar{\Gamma} + 4D_{r})}\right) + \frac{v_{p}^{2}}{(\kappa\bar{\Gamma} + 3D_{r})}\left[\frac{1}{(\kappa\bar{\Gamma} - 3D_{r})}\left(\frac{3(\kappa\bar{\Gamma} - D_{r})}{(\kappa\bar{\Gamma} + D_{r})2\kappa\bar{\Gamma}} - \frac{e^{-(\kappa\bar{\Gamma}-3D_{r})t}}{4D_{r}}\right)\right]$$

$$+ \frac{1}{2(\kappa\bar{\Gamma} - D_{r})}\left(\frac{e^{-(\kappa\bar{\Gamma}+D_{r})t}(\kappa\bar{\Gamma} - 5D_{r})}{2D_{r}(\kappa\bar{\Gamma} + D_{r})}\right) - \frac{1}{(\kappa\bar{\Gamma} + 2D_{r})}\left(\frac{1 - e^{-2\kappa\bar{\Gamma}t}}{2\kappa\bar{\Gamma}} - \frac{e^{-2\kappa\bar{\Gamma}t} - e^{-2(\kappa\bar{\Gamma}+2D_{r})t}}{4D_{r}}\right)\right]$$

6. Calculation of $\langle R_{1,i}(t)R_{1,j}(t)\rangle_{\xi,\theta_0}$