Unit-2 Convolution, Correlation of Signals and Fourier Series

Unit-2

Part-1:

(Convolution and Correlation of Signals)

- Convolution and their Properties
- Correlation of Signals
- Cross correlation and Auto Correlation Functions
- Properties of Correlation functions
- Relation between Convolution and Correlation

Part-2:

(Fourier Series)

- Fourier Series Properties
- Dirchlet t Conditions
- Trigonometric Fourier Series
- Exponential Fourier Series
- Complex Fourier Spectrum

Part-1: (Convolution and Correlation of Signals)

Introduction:

Convolution is a mathematical way of combining two signals to form a third signal. It is an important operation because it relates the input signal and impulse response of the system to the output of the system.

Correlation is also a mathematical operation that uses two signals to form a third signal. It compares two signals inorder to determine the degree of similarity between them. It is very widely used in communication ensmeering

correlation may be cross correlation or auto correlation. When one signal is correlated with another signal to form a third signal, it is called cross correlation. When a signal is correlated with itself to form another signal, it is called auto correlation.

Convolution:

Convolution: Convolution is a mathematical operation which is used to express the input-output relationship of an LTI System. It is an important operation in LTI continuous-time systems. It relates the input and impulse response of the system to the output.

Consider an LTI system which is intially relaxed at t=0. If the input to the system is an impulse, then the output of the system is denoted by htt) and is called the impulse response of the ness system.

The impulse response is denoted by

We know that any arbitrary signal alto can be represented by 2(t) = \(\int a(2) \int (t-2) de - \(\omega \) The system output is given by y(t) = T[x(t)] -- (3) : y(t) = T[sa(2) 8(t-2) d2] yet) = [x(2) T [614-2)] de - 4 If the response of the system due to impulse 8(t) is h(t), The response of the system due to delayed input 81t-2) is h(t, 2) = T[5(t-2)] -- 3 From egns (46) y(+) = 52(2) b(+, 2)d2

For a Time invariant system
$$h(t;t) = h(t-t) - \emptyset$$
 from eqn \emptyset & \emptyset & \emptyset

$$y(t) = \int_{-\infty}^{\infty} \chi(t)h(t-t)dt$$
This is called convolution integral or simply convolution. The convolution of two sequent signals $\chi(t)$ and $\chi(t)$ can be represented as
$$\chi(t) = \chi(t) + h(t)$$

This is called convolution integral or simply convolution. The convolution of two sequent signals alt and hit can be represented as

 $y(t) = \int_{-\infty}^{\infty} \chi(z)h(t-z) dz \quad \text{if both } \chi(t) \leq h(t) \text{ one non-causal}$ $= \int_{-\infty}^{\infty} \chi(z)h(t-z)dz \quad \text{if } \chi(t) \text{ is non-causal } \leq h(t) \text{ is causal}$ $= \int_{-\infty}^{\infty} \chi(z)h(t-z)dz \quad \text{if } \chi(t) \text{ is causal } \leq h(t) \text{ is non-causal}$ $= \int_{-\infty}^{\infty} \chi(z)h(t-z)dz \quad \text{if } \chi(t) \leq \text{causal} \leq h(t) \text{ is non-causal}$ $= \int_{-\infty}^{\infty} \chi(z)h(t-z)dz \quad \text{if both } \chi(t) \leq h(t) \text{ are causal}.$

Properties of Convolution:

Let us consider two signals $x_1(t)$ and $x_2(t)$. The convolution of two signals $x_1(t)$ and $x_2(t)$ is given by

 $x_1(t) * x_2(t) = \int x_1(\tau) x_2(t-\tau) d\tau = \int x_2(\tau) x_1(t-\tau) d\tau$

The properties of convolution are as follows:

Commutative property The commutative property of convolution states that

$$x_1(t) * x_2(t) = x_2(t) * x_1(t)$$

Distributive property The distributive property of convolution states that

$$x_1(t) * [x_2(t) + x_3(t)] = [x_1(t) * x_2(t)] + [x_1(t) * x_3(t)]$$

Associative property The associative property of convolution states that

$$x_1(t) * [x_2(t) * x_3(t)] = [x_1(t) * x_2(t)] * x_3(t)$$

Shift property The shift property of convolution states that if

$$x_1(t) * x_2(t) = z(t)$$

Then $x_1(t) * x_2(t-T) = z(t-T)$

Similarly, $x_1(t-T) * x_2(t) = z(t-T)$

and $x_1(t-T_1) * x_2(t-T_2) = z(t-T_1-T_2)$

Convolution with an impulse \tilde{z} Convolution of a signal x(t) with a unit impulse is the signal itself. That is,

$$x(t) * \delta(t) = x(t)$$

Width property Let the duration of $x_1(t)$ and $x_2(t)$ be T_1 and T_2 respectively. Then the duration of the signal obtained by convolving $x_1(t)$ and $x_2(t)$ is $T_1 + T_2$.

Problems:

EXAMPLE 1 Find the convolution of the following signals:

(i)
$$x_1(t) = e^{-2t} u(t)$$
; $x_2(t) = e^{-4t} \dot{u}(t)$ (i) denoted by

(ii)
$$x_1(t) = t u(t);$$
 $x_2(t) = t u(t)$

(iii)
$$x_1(t) = \cos t u(t)$$
; $x_2(t) = u(t)$

(iv)
$$x_1(t) = e^{-3t} u(t)$$
; $x_2(t) = u(t+3)$

(v)
$$x_1(t) = r(t);$$
 $x_2(t) = e^{-2t} u(t)$

Solutions:

Solution:

(i) Given

$$x_1(t) = e^{-2t} u(t); x_2(t) = e^{-4t} u(t)$$

We know that
$$x_1(t) * x_2(t) = \int_{-\infty}^{\infty} x_1(\tau) x_2(t-\tau) d\tau$$

 $\{x(z) h(t-z) dz\}$ if z(t) is causal and h(t) is non-causal

$$\therefore x_1(t) * x_2(t) = \int_{-\infty}^{\infty} e^{-2\tau} u(\tau) e^{-4(t-\tau)} u(t-\tau) d\tau$$
with the property of the contract of the property of the

 $u(\tau) = 1$ for $\tau > 0$ and $u(t - \tau) = 1$ for $(t - \tau) \ge 0$ or for $\tau < t$. Hence $u(\tau)$ $u(t - \tau) = 1$ only for $0 < \tau < t$. For all other values of τ , $u(\tau)$ $u(t - \tau) = 0$.

$$x_1(t) * x_2(t) = \int_0^t e^{-2\tau} e^{-4(t-\tau)} d\tau$$

$$=e^{-4t}\int_{0}^{t}e^{2\tau}d\tau=e^{-4t}\left[\frac{e^{2\tau}}{2}\right]_{0}^{t}=e^{-4t}\left(\frac{e^{2t}-1}{2}\right)=\frac{e^{-2t}-e^{-4t}}{2}\quad (\text{for } t\geq 0)_{+}=\frac{e^{-2t}-e^{-4t}}{2}u(t)$$

(ii) Given

$$x_1(t) = t u(t); x_2(t) = t u(t)$$

We know that

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$$x_1(t) * x_2(t) = \int_0^\infty x_1(\tau) x_2(t-\tau) d\tau$$

$$x_1(t) * x_2(t) = \int_{-\infty}^{\infty} \tau u(\tau) (t - \tau) u(t - \tau) d\tau$$

 $u(\tau) = 1$ for $\tau > 0$ and $u(t - \tau) = 1$ for $(t - \tau) \ge 0$ or for $\tau < t$.

Hence $u(\tau)$ $u(t-\tau)=1$ only for $0<\tau< t$. For all other values of τ , $u(\tau)$ $u(t-\tau)=0$.

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$$\therefore x_1(t) * x_2(t) = \int_0^t \tau(t-\tau) d\tau = \int_0^t t\tau d\tau - \int_0^t \tau^2 d\tau = t \left[\frac{\tau^2}{2}\right]_0^t - \left[\frac{\tau^3}{3}\right]_0^t$$

$$= t \left(\frac{t^2}{2} - 0\right) - \left(\frac{t^3}{3} - 0\right) = \frac{t^3}{2} - \frac{t^3}{3} = \frac{t^3}{6} \quad \text{(for } t \ge 0) \quad \therefore \quad x_1(t) * x_2(t) = \frac{t^3}{6} u(t)$$

(iii) Given
$$x_1(t) = \cos t \ u(t); \ x_2(t) = u(t)$$

We know that $x_1(t) * x_2(t) = \int_{-\infty}^{\infty} x_1(\tau) \ x_2(t-\tau) \ d\tau : x_1(t) * x_2(t) = \int_{-\infty}^{\infty} \cos \tau \ u(\tau) \ u(t-\tau) \ d\tau$

 $u(\tau) = 1$ for $\tau > 0$ and $u(t - \tau) = 1$ for $(t - \tau) \ge 0$ or for $\tau < t$. Hence $u(\tau)$ $u(t - \tau) = 1$ only for $0 < \tau < t$. For all other values of τ , $u(\tau)$ $u(t - \tau) = 0$.

$$x_1(t) * x_2(t) = \int_0^t \cos \tau d\tau = \left[\sin \tau\right]_0^t = \sin t \quad \text{for } t \ge 0$$

$$x_1(t) * x_2(t) = \sin t u(t)$$

(iv) Given $x_1(t) = e^{-3t} u(t); x_2(t) = u(t+3)$ We know that $x_1(t) * x_2(t) = \int_{-\infty}^{\infty} x_1(\tau) x_2(t-\tau) d\tau$

$$\therefore x_1(t) * x_2(t) = \int_{-\infty}^{\infty} e^{-3\tau} u(\tau) u(t+3-\tau)$$

In this case, $u(\tau) = 0$ for $\tau < 0$ and $u(t+3-\tau) = 0$ for $\tau > t+3$. $u(\tau)$ $u(t+3-\tau) = 1$ only for $0 < \tau < t+3$. For all other values of τ , $u(\tau)$ $u(t+3-\tau) = 0$.

$$\therefore x_1(t) * x_2(t) = \int_0^{t+3} e^{-3\tau} d\tau = \left[\frac{e^{-3\tau}}{-3} \right]_0^{t+3} = \frac{e^{-3(t+3)} - 1}{-3} = \frac{1 - e^{-3(t+3)}}{3}$$

$$y(t) = 0 (for t < -3) = \frac{1 - e^{-3(t+3)}}{3} (for t > -3)$$

(v) Given
$$x_1(t) = r(t) = tu(t); x_2(t) = e^{-2t} u(t)$$

We know that $x_1(t) * x_2(t) = \int_{-\infty}^{\infty} x_1(\tau) x_2(t-\tau) d\tau$
 $x_1(t) * x_2(t) = \int_{-\infty}^{\infty} \tau u(\tau) e^{-2(t-\tau)} u(t-\tau) d\tau$

 $u(\tau) = 1$ for $\tau > 0$ and $u(t - \tau) = 1$ for $(t - \tau) \ge 0$ or for $\tau < t$. Hence $u(\tau)$ $u(t - \tau) = 1$ only for $0 < \tau < t$. For all other values of τ , $u(\tau)$ $u(t - \tau) = 0$.

$$z_1(t) * x_2(t) = \int_0^t \tau e^{-2(t-\tau)} d\tau = e^{-2t} \int_0^t \tau e^{2\tau} d\tau = e^{-2t} \left\{ \left[\frac{\tau e^{2\tau}}{2} \right]_0^t - \int_0^t \frac{e^{2\tau}}{2} d\tau \right\}$$

$$= e^{-2t} \left\{ \left[\frac{te^{2t}}{2} - \left[\frac{e^{2\tau}}{4} \right]_0^t \right] \right\} = e^{-2t} \left(\frac{te^{2t}}{2} - \frac{e^{2t}}{4} + \frac{1}{4} \right) = \frac{t}{2} - \frac{1}{4} + \frac{e^{-2t}}{4} \quad \text{(for } t \ge 0\text{)}$$

$$\therefore x_1(t) * x_2(t) = \left(\frac{t}{2} - \frac{1}{4} + \frac{e^{-2t}}{4}\right) u(t)$$

Convolution Theorems:

Convolution of signals may be done either in time domain or in frequency domain. So there are following two theorems of convolution associated with Fourier transforms:

- 1. Time convolution theorem
- 2. Frequency convolution theorem

Fime Convolution Theorem

The time convolution theorem states that convolution in time domain is equivalent to multiplication of their spectra in frequency domain. Mathematically, if

$$x_1(t) \longleftrightarrow X_1(\omega)$$
 and
$$x_2(t) \longleftrightarrow X_2(\omega)$$
 Then
$$x_1(t) * x_2(t) \longleftrightarrow X_1(\omega) X_2(\omega)$$

Proof:
$$F[x_1(t) * x_2(t)] = \int_{-\infty}^{\infty} [x_1(t) * x_2(t)] e^{-jax} dt$$
We have
$$x_1(t) * x_2(t) = \int_{-\infty}^{\infty} x_1(\tau) x_2(t-\tau) d\tau$$

$$F[x_1(t) * x_2(t)] = \int_{-\infty}^{\infty} \left\{ \int_{-\infty}^{\infty} [x_1(\tau) x_2(t-\tau) d\tau] \right\} e^{-j\alpha t} dt$$

Interchanging the order of integration, we have

$$F[x_1(t) * x_2(t)] = \int_{-\infty}^{\infty} x_1(\tau) \left[\int_{-\infty}^{\infty} x_2(t-\tau) e^{-j\omega t} dt \right] d\tau$$

Letting $t - \tau = p$, in the second integration, we have

$$t = p + \tau$$
 and $dt = dp$

$$F[x_1(t) * x_2(t)] = \int_{-\infty}^{\infty} x_1(\tau) \left[\int_{-\infty}^{\infty} x_2(p) e^{-j\omega t} dp \right] d\tau$$

$$= \int_{-\infty}^{\infty} x_1(\tau) \left[\int_{-\infty}^{\infty} x_2(p) e^{-j\omega p} dp \right] e^{-j\omega \tau} d\tau$$

$$= \int_{-\infty}^{\infty} x_1(\tau) X_2(\omega) e^{-j\omega \tau} d\tau = \int_{-\infty}^{\infty} x_1(\tau) e^{-j\omega \tau} d\tau X_2(\omega)$$

$$= X_1(\omega) X_2(\omega)$$

Frequency Convolution Theorem

The frequency convolution theorem states that the multiplication of two functions in time domain is equivalent to convolution of their spectra in frequency domain. Mathematically, if

$$x_1(t) \longleftrightarrow X_1(\omega)$$

and

$$x_2(t) \longleftrightarrow X_2(\omega) \xrightarrow{\pi_1} x_2(t) \longleftrightarrow X_2(\omega)$$

Then

$$x_1(t) x_2(t) \longleftrightarrow \frac{1}{2\pi} [X_1(\omega) * X_2(\omega)]$$

Proof:

$$\begin{split} \mathbf{F}[x_1(t) \ x_2(t)] &= \int\limits_{-\infty}^{\infty} \left[x_1(t) \ x_2(t)\right] e^{-j\epsilon nt} \, dt \, , \\ &= \int\limits_{-\infty}^{\infty} \left[\frac{1}{2\pi} \int\limits_{-\infty}^{\infty} X_1(\lambda) e^{j\lambda t} \, d\lambda\right] x_2(t) \, e^{-j\alpha t} \, dt \, \, . \end{split}$$

Interchanging the order of integration, we get

$$\begin{split} \mathbf{F}[x_1(t) \; x_2(t)] &= \frac{1}{2\pi} \int_{-\infty}^{\infty} X_1(\lambda) \left[\int_{-\infty}^{\infty} x_2(t) \, e^{-j\omega t} \, e^{j\lambda t} \, dt \right] d\lambda \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} X_1(\lambda) \left[\int_{-\infty}^{\infty} x_2(t) \, e^{-j(\omega t - \lambda)t} \, dt \right] d\lambda \end{split}$$

bmu.

Then

Proofs

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} X_1(\lambda) X_2(\omega - \lambda) d\lambda$$

$$= \frac{1}{2\pi} [X_1(\omega) * X_2(\omega)]$$

$$\therefore \qquad x_1(t) x_2(t) \longleftrightarrow \frac{1}{2\pi} X_1(\omega) * X_2(\omega)$$
or
$$2\pi x_1(t) x_2(t) \longleftrightarrow X_1(\omega) * X_2(\omega)$$

This is frequency convolution theorem in radian frequency. In terms of frequency, we get

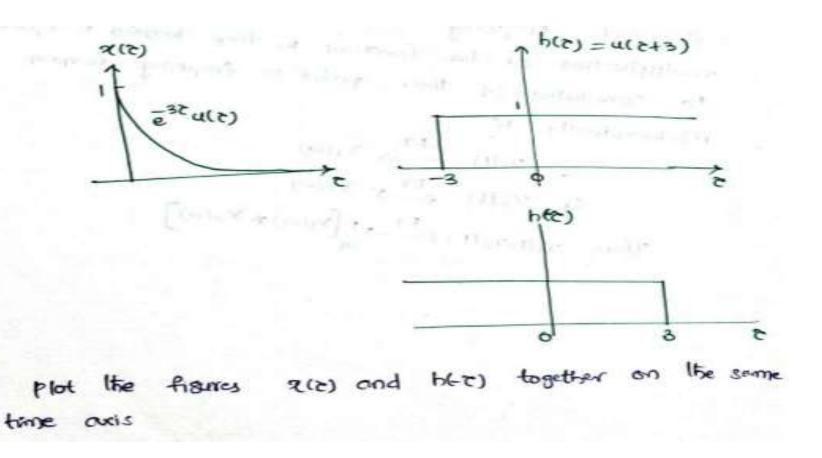
$$F[x_1(t) | x_2(t)] = X_1(f) * X_2(f)$$

Convolution using Graphical Representation:

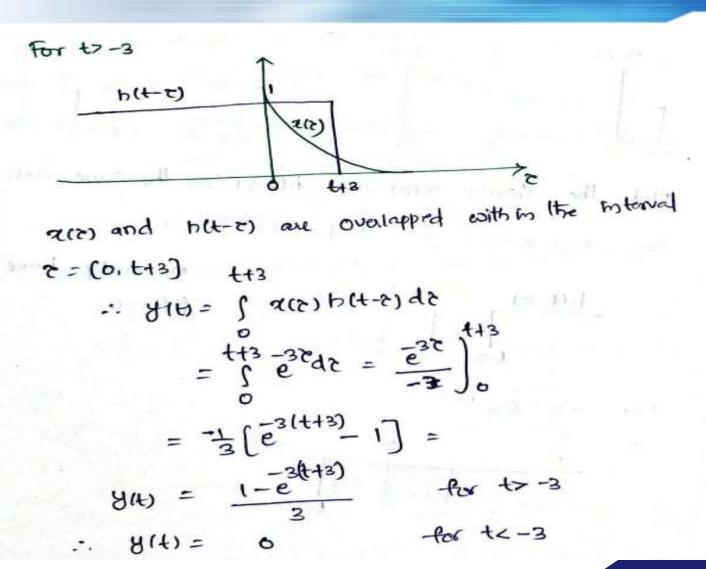
Graphical Representation of Convolution: when two signals with end wells are provided in graphical form, the convolution can be performed by graphical method. It involve the following steps: O for the given signals will and will, draw the signals relied and 21(2) as a function of independent time variable ?. @ Draw the function \$2(-2) which is time reversed function of aday. Then shifting the function by time t to form xelt-es. 3 Draw both signals zell) and 22 (t-z) on the same time z-axis with large time shift. I along the -ve axis (4) Increase the time shift t along the time axis. Multiply the sign of 21(2) 72 (+-1) and integrate over the overlapping period of two signals to obtain convolution at t. 3 Encrease the time shift to step by step and obtain convolution using Step 4 (6) Draw the convolution function all with the values obtained in Step 4 and 5 as a function of t

Problems:

```
find the convolution of the following siznals by graphical
method.
  2(t) = e u(t) and h(t) = u(t+3)
Soln:
given att) = e att)
     & hlt) = (11+3)
    giti = aiti * piti = 2 a(5) pit-6) q5
       q(z) = e^{-3z}u(z) b(z) = u(z+3)
       u(8)=1 for 8>0 u(8+3)=1 for 8>-3
        =0 for 800 =0 for 80-3
35- - 1
                       :. h(c) = u(c+3)=1 for c>-3
```



for
$$t < -3$$
 $t+3$
 $t+3$



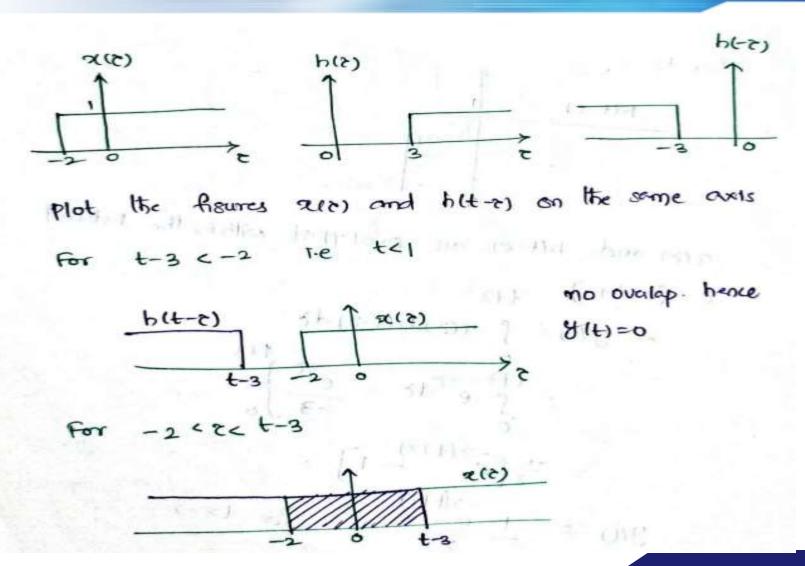
The input and the impulse response to the system are given by
$$x(t) = u(t+2)$$

Est $h(t) = u(t+2)$

Determine the output of the system graphs rally.

Solon: Given $x(t) = u(t+2)$

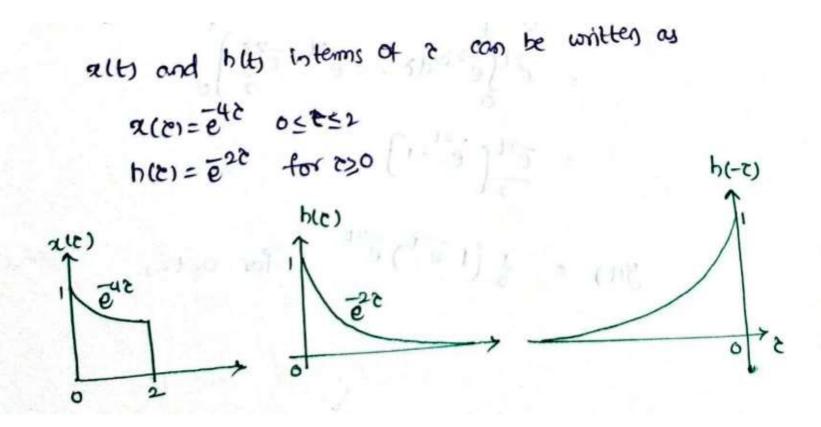
Est $h(t) = u(t+3)$
 $y(t) = x(t) * h(t) = \int_{-\infty}^{\infty} x(t) h(t-t) dt$
 $x(t) = u(t+2) = 1$ for $t > 2$
 $x(t) = u(t+2) = 1$ for $t > 3$
 $x(t) = u(t+2) = 1$ for $t > 3$
 $x(t) = u(t+2) = 1$ for $t > 3$
 $x(t) = u(t+2) = 1$ for $t > 3$
 $x(t) = u(t+2) = 1$ for $t > 3$
 $x(t) = u(t+2) = 1$ for $t > 3$
 $x(t) = u(t+2) = 1$ for $t > 3$



$$y(t) = \int_{-2}^{2} x(t) h(t-t) dt$$

$$= \int_{-2}^{2} (1) (1) dt$$

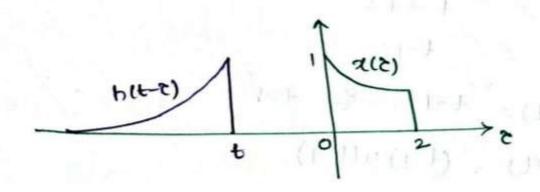
3) The impulse response of the circuit is given as hitt=e^{2t}uit). This circuit is excited by an input of alt) = e (ult)-ult-2)) Determine the output of the circuit. Soln: h(t)= e2tult) = et for t>0 4 21t1 = et (ult) - ult-2)] = et octes gitt= aiti* hitt= \$ a(z) hit-z)de alt) and hith interms of a can be written as 2(0)=et 05+52 h(e) = e20 for 0>0



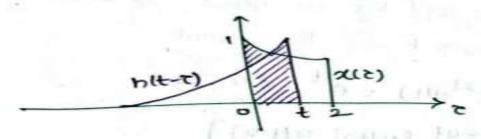
$$h(t-t) = e^{-2(t-t)}u(t-t) = e^{-2(t-t)}$$
 for $t-t>0$

The plots of act) and hit-to drawn on the same time axis as shown below.

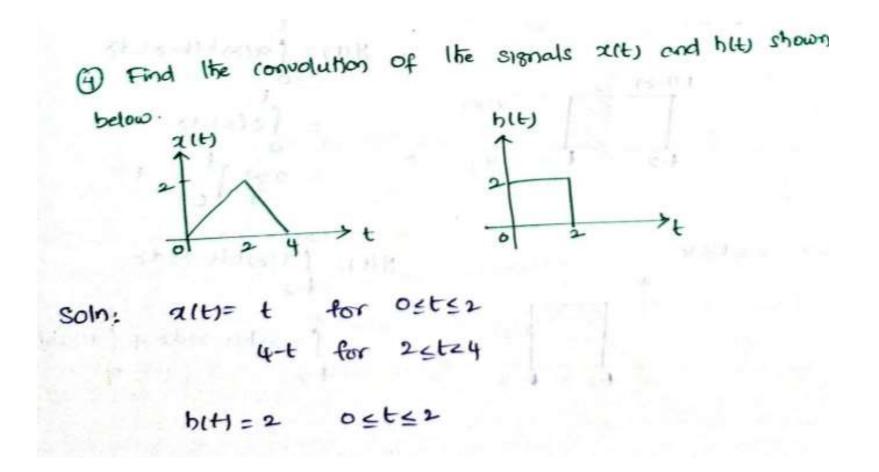
For tzo, The plots do not overlap.



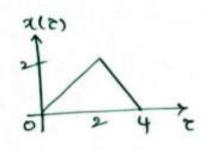
for octs2

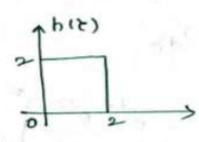


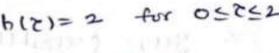
$$\begin{array}{ll}
\vdots & \text{Sith} = \int_{0}^{\infty} a(z) h(t-z) dz \\
&= \int_{0}^{\infty} \left(\overline{e^{2}(t-z)} \right) dz \\
&= \overline{e^{2}t} \left(\overline{e^{2}} \right) \left(\overline{e^{2}(t-z)} \right) dz \\
&= \overline{e^{2}t} \left(\overline{e^{2}} \right) \left(\overline{e^{2}(t-z)} \right) dz \\
&= \overline{e^{2}t} \left(\overline{e^{2}} \right) \left(\overline{e^{2}(t-z)} \right) dz
\end{array}$$

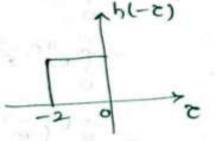


$$a(t) = \begin{cases} t & 0 \le t \le 2 \\ u - t & 2 \le t \le 4 \end{cases}$$
 $b(t) = 2$ for $0 \le t \le 2$

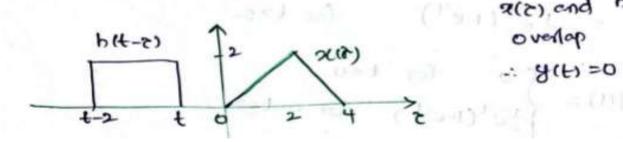


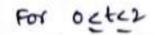


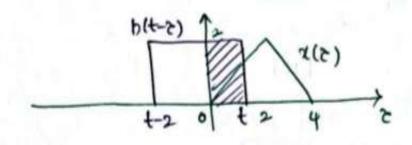




For teo







$$3(t) = \int a(x)h(t-x)dx$$
= $\int a(x)h(t-x)dx$
= $\int a(x)h(t-x)dx$
= $\int a(x)h(t-x)dx$

$$y(t) = \int_{-2}^{2} \pi(z)h(t-z)dz$$

$$= \int_{-2}^{2} \pi(z)h(t-z)dz + \int_{-2}^{2} \pi(z)h(t-z)dz$$

$$= \int_{-2}^{2} \pi(z)h(t-z)dz + \int_{-2}^{2} \pi(z)h(t-z)dz$$

$$\begin{aligned}
y(t) &= \int_{t-2}^{2} 2(x) dx + \int_{t}^{t} (4-x) 2 dx \\
&= 2x^{2} \int_{t-2}^{2} + \int_{t}^{t} (8-2x) dx \\
&= 4 - (t-2)^{2} + \left(8x - 2x^{2}\right)^{\frac{1}{2}} \\
&= 4 - (t^{2} + 4 - 4t) + \left(8t - t^{2}\right) - \left(16 - 4\right) \\
&= 4 - t^{2} + 4x^{2}t + 8t - t^{2} - 12 \\
&= -2t^{2} + 12t - 12
\end{aligned}$$

For t26

$$2 + 2(e)$$

$$3 + 2(e)$$

$$4 + 2($$

Concept of Correlation:

Concept of correlation

The signals may be compared on the basis of similarity of waveforms. Quantitatively, a comparison may be based upon the amount of the component of one waveform contained in the other waveform. If $x_1(t)$ and $x_2(t)$ are two waveforms, then the waveform $x_1(t)$ contains an amount $C_{12}x_2(t)$ of that particular waveform $x_2(t)$ in the interval (t_1, t_2) , where

$$C_{12} = \frac{\int_{t_1}^{t_2} x_1(t) x_2(t) dt}{\int_{t_1}^{t_2} x_2^2(t) dt}$$

 $C_{12} = \frac{\int_{t_1}^{t_2} x_1(t) x_2(t) dt}{\int_{t_2}^{t_2} x_2^2(t) dt}$ The magnitude of the integral in the numerator might be taken as an indication of similarity.

If this integral vanishes, i.e The magnitude of the integral in the numerator If this integral vanishes, i.e

$$\int_{t_1}^{t_2} x_1(t) \ x_2(t) \ dt = 0$$

then the two signals have no similarity over the interval (t_1, t_2) . Such signals are said to be orthogonal over the specified interval.

The integral $\int_{t_1}^{t_2} x_1(t) x_2(t) dt$ forms the basis of comparison of the two signals $x_1(t)$ and $x_2(t)$ over the interval (t_1, t_2) .

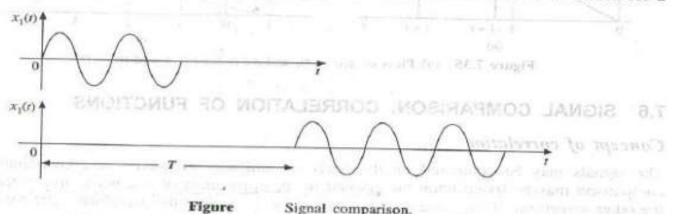
In general we are interested in comparing the two signals over the interval $(-\infty, \infty)$. So the test integral becomes

with the elegan
$$(8-42-5)(301)=\int_0^1 x_1(t)x_2(t)dt=0$$
 (1) a knowledge outs fitted using A

However, there is a difficulty with this test integral which can be illustrated with the example of radar pulse. Figure shows a transmitted pulse and a received pulse which is delayed w.r.t. transmitted pulse by T s. Obviously, the two waveforms are identical except that one

is delayed w.r.t. the other. Yet the test integral $\int_{-\infty}^{\infty} x_1(t) x_2(t) dt$ yields zero because the

product $x_1(t)$ $x_2(t)$ is zero everywhere. This indicates that the two waveforms have no measure of similarity which is obviously a wrong conclusion. Hence in order to search for a



similarity between the two waveforms, we must shift one waveform w.r.t. the other by various amounts and see whether a similarity exists for some amount of shift of one function w.r.t. the other.

w.r.t. the other. Therefore, the test integral is modified as $\int_{-\infty}^{\infty} x_1(t) x_2(t-\tau) dt$

where τ is the searching or scanning parameter. This integral is a function of τ . This integral is known as the cross correlation function between $x_1(t)$ and $x_2(t)$ and is denoted by $R_{12}(\tau)$.

It is immaterial whether we shift the function $x_1(t)$ by an amount of τ in the negative direction or shift the function $x_2(t)$ by the same amount in the positive direction. Thus

Thus the correlation of two functions or signals or waveforms is a measure of similarity between those signals. The correlation is of two types: cross correlation and autocorrelation. The autocorrelation and cross correlation are defined separately for energy (or aperiodic) signals and power (or periodic) signals.

Cross Correlation:

The cross correlation between two different waveforms or signals is a measure of similarity or match or relatedness or coherence between one signal and the time delayed version of another signal. That means the cross correlation between two signals indicates how much one signal is related to the time delayed version of another signal.

Cross correlation of energy signals

Consider two general complex signals $x_1(t)$ and $x_2(t)$ of finite energy. The cross correlation of these two energy signals denoted by $R_{12}(\tau)$ is given by

$$R_{12}(\tau) = \int_{-\infty}^{\infty} x_1(t) x_2 * (t - \tau) dt = \int_{-\infty}^{\infty} x_1(t + \tau) x_2 * (t) dt$$

 $R_{12}(\tau) = \int_{-\infty}^{\infty} x_1(t) x_2 * (t-\tau) dt = \int_{-\infty}^{\infty} x_1(t+\tau) x_2 * (t) dt$ If the two signals $x_1(t)$ and $x_2(t)$ are real, then $R_{12}(\tau) = \int_{-\infty}^{\infty} x_1(t) x_2(t-\tau) dt = \int_{-\infty}^{\infty} x_1(t+\tau) x_2(t) dt$

If $x_1(t)$ and $x_2(t)$ have some similarity, then the cross correlation $R_{12}(\tau)$ will have some finite value over the range of τ . Also if

$$\int_{-\infty}^{\infty} x_1(t) x_2^*(t) dt = 0 \quad \text{i.e. if} \quad R_{12}(0) = 0$$

then the two signals $x_1(t)$ and $x_2(t)$ are called orthogonal signals. That is the cross correlation for orthogonal signals is zero, a smort of regions of the signal signals is zero, a smort of the signal signa Another form of cross correlation between $x_2(t)$ and $x_1(t)$ is defined as: $R_{21}(\tau) = \int x_2(t) x_1^{\tau}(t-\tau) dt$

In the above equations, the cross correlation function $R_{12}(\tau)$ is a function of the variable τ . The variable τ is called the *delay parameter* or the *searching parameter*. It is time delay or time shift of one of the two signals. The delay parameter τ determines the correlation between two signals. Two signals with no cross correlation at $\tau = 0$ can have significant cross correlation by adjusting the parameter τ . Two signals for which the cross correlation is zero for all values of τ are called *uncorrelated* or *incoherent signals*.

Properties of Cross Correlation:

Properties of cross correlation function for energy signals to noiselesson and audit

Following are the properties of cross correlation for energy signals:

1. The cross correlation functions exhibit conjugate symmetry, i.e.) 15 wood ban alangua

In general we are interested in comparing the *to tends over the interval (-50, 50 K integral becomes
$$R_{12}(\tau) = R_{21}(-\tau)$$

That is unlike convolution, cross correlation is not in general commutative, i.e.

The cross correlation between two
$$d(\tau)_{12} = (\tau)_{12} = (\tau)_{12} = 0$$
 is a measure of similarity or match or relatedness or coherence between one signal and the time delayed version of another

i.e. if
$$\int_{0}^{\infty} x_1(t) x_2(t) dt = 0$$
 i.e. if
$$\int_{0}^{\infty} x_1(t) x_2(t) dt = 0$$
 i.e. if

then the two signals are said to be orthogonal over the entire time interval.

3. The cross correlation of two energy signals corresponds to the multiplication of the Fourier transform of one signal by the complex conjugate of Fourier transform of second signal.

i.e.

$$R_{12}(\tau) \longleftrightarrow X_1(\omega) X_2^*(\omega)$$
 but $(\tau)_{12}$ slengts own and if

This is known as correlation theorem.

Cross correlation of power (periodic) signals

The cross correlation function $R_{12}(\tau)$ for two periodic signals $x_1(t)$ and $x_2(t)$ may be defined with the help of average form of correlation. If the two periodic signals $x_1(t)$ and $x_2(t)$ have the same time period T, then cross correlation is defined as:

$$R_{12}(\tau) = \frac{1}{T} \int_{-T/2}^{T/2} x_1(t) x_2^*(t - \tau) dt$$

then the two signals $x_1(t)$ and $x_2(t)$ are called *orthogonal signals*. That is the cross correlation for orthogonal signals is zero. The control of the cross correlation are the cross correlation for orthogonal signals is zero.

Another form of cross correlation between $x_2(t)$ and $x_1(t)$ is defined as:

$$R_{21}(\tau) = \int_{-\infty}^{\infty} x_2(t) x_1^{*}(t-\tau) dt$$

Properties of cross correlation function for power (periodic) signals

Following are the properties of cross correlation for power signals:

 The Fourier transform of the cross correlation of two signals is equal to the multiplication of Fourier transform of one signal and complex conjugate of Fourier transform of other signal.

The subscorrelation function
$$X_1(\omega) = X_1(\omega) =$$

slangis rewood ban slangis
$$T_{T\to\infty}$$
 $T_{T/2}$ $T_{T/2}$

then the signals are said to be orthogonal over the entire time interval.

3. The cross correlation exhibits conjugate symmetry, i.e.

$$R_{12}(\tau) = R_{21}^*(-\tau)$$

4. Unlike convolution, the cross correlation is not commutative, i.e.

where
$$\tau$$
 is called the delay parameter $(\tau)_{12}^{\text{reg}} R \neq (\tau)_{12}^{\text{reg}} R$ afth is shifted by τ in positive direction.

EXAMPLE Prove that $R_{12}(\tau) = R_{21}^*(-\tau)$ i.e. the cross correlation exhibits conjugate symmetry.

Solution: The cross correlation of two signals $x_1(t)$ and $x_2(t)$ is given as:

$$R_{12}(\tau) = \int_{-\infty}^{\infty} x_1(t) x_2^*(t-\tau) dt$$

Let $t - \tau = n$ in the above equation for $R_{12}(\tau)$,

Annique of the section
$$R_{12}(\tau) = \int_0^{\infty} x_1(n+\tau) x_2^{**}(n) dn$$
 with the section of the section $R_{12}(\tau) = \int_0^{\infty} x_1(n+\tau) x_2^{**}(n) dn$

Also we know that add naves at (stretungle grows on to noticisossome and seasons

$$R_{21}(\tau) = \int_{-\infty}^{\infty} x_2(t) \, x_1^*(t - \tau) \, dt$$

Let t = n in the above equation for $R_{21}(\tau)$.

$$R_{21}(\tau) = \int_{-\infty}^{\infty} x_2(n) x_1^*(n-\tau) dn$$

$$\therefore \text{ we then the form of the$$

$$R_{21}^{*}(-\tau) = \int_{-\infty}^{\infty} x_{2}^{*}(n) x_{1}(n+\tau) dn$$

Comparing the above two equations for $R_{12}(\tau)$ and $R_{21}^*(-\tau)$, we can write

sitt to things at stanger over to an
$$R_{12}(\tau)=R_{21}^*(-\tau)$$
 at the intersection property of

Auto Correlation:

Autocorrelation

The autocorrelation function gives the measure of match or similarity or relatedness or coherence between a signal and its time delayed version. This means that the autocorrelation function is a special form of cross correlation function. It is defined as the correlation of a signal with itself.

The autocorrelation is defined separately for energy signals and power signals.

Autocorrelation for energy signals exhibit conjugate symmetry, i.e.

The autocorrelation of an energy signal x(t) is given by

$$R_{11}(\tau) = R(\tau) = \int_0^\infty x(t) x^*(t-\tau) dt$$

where τ is called the delay parameter and the signal x(t) is shifted by τ in positive direction. If x(t) is shifted by τ in negative direction, then

EXAMPLE 7.12 Prove that
$$R_{12}(\tau) = \int_{0}^{\infty} \int_{0}^{\infty$$

Calations The cares correlation of two signals with and

Properties of Auto Correlation function:

Properties of autocorrelation function of energy signals

Following are the properties of autocorrelation for energy signals:

1. The autocorrelation function exhibits conjugate symmetry, i.e.

$$R(\tau) = R(-\tau)$$
 notherpo evods out in $n = 7 - 1$

Thus, it states that the real part of $R(\tau)$ is an even function of τ and the imaginary part of $R(\tau)$ is an odd function of τ .

Proof: The autocorrelation of an energy signal x(t) is given by word an order

$$R(\tau) = \int_{0}^{\infty} x(t)x^{*}(t-\tau)dt$$

Taking the complex conjugate, we have $(7)_{12}$ % and notice avoids and not n=1 to L=1

$$R^*(\tau) = \int_{-\infty}^{\infty} x^*(t) \ x(t-\tau) \ dt$$

$$R^*(-\tau) = \int_0^\infty x^*(t) x(t+\tau) dt = R(\tau)$$

$$R(\tau) = R^*(-\tau)$$

2. The value of autocorrelation function of an energy signal at origin (i.e. at $\tau = 0$) is equal to the total energy of that signal, i.e.

$$R(0) = E = \int_{0}^{\infty} |x(t)|^{2} dt$$

Proof: We have

Proof: We have

$$R(\tau) = \int_{-\infty}^{\infty} x(t)x^{*}(t-\tau) dt$$

Putting $\tau = 0$ gives

$$R(0) = \int_{-\infty}^{\infty} x(t)x^{*}(t) dt = \int_{-\infty}^{\infty} |x(t)|^{2} dt = E$$

3. If τ is increased in either direction, the autocorrelation $R(\tau)$ reduces. As τ reduces autocorrelation, $R(\tau)$ increases and it is maximum at $\tau = 0$, i.e. at the origin. Therefore,

$$|R(\tau)| \le R(0)$$
 (for all τ)

Proof: Consider the functions x(t) and $x(t + \tau)$. $[x(t) \pm x(t + \tau)]^2$ is always greater than or equal to zero since it is squared, i.e.

$$x^{2}(t) + x^{2}(t+\tau) \pm 2x(t) x(t+\tau) \ge 0$$

OF

$$x^{2}(t) + x^{2}(t+\tau) \ge \pm 2x(t) x(t+\tau)$$

Integrating both the sides, we get

$$\int_{-\infty}^{\infty} |x(t)|^2 dt + \int_{-\infty}^{\infty} |x(t+\tau)|^2 dt \ge 2 \int_{-\infty}^{\infty} x(t) x(t+\tau) dt$$

$$E + E \ge 2R(\tau) \quad \text{[If } x(t) \text{ is real valued function]}$$

$$E \geq R(\tau)$$

or
$$R(0) \ge |R(\tau)|$$
 (Since $R(0) = E$)

4. The autocorrelation function $R(\tau)$ and energy spectral density function $\psi(\omega)$ of energy signal form a Fourier transform pair.

$$R(\tau) \longleftrightarrow \psi(\omega)$$

Autocorrelation theorem

The autocorrelation theorem states that the Fourier transform of autocorrelation function $R(\tau)$ yields the energy density function of signal x(t), i.e.

$$F[R(\tau)] = |X(\omega)|^2 = \psi(\omega)$$

Proof: The Fourier transform of autocorrelation function $R(\tau)$ is:

$$F[R(\tau)] = \int_{-\infty}^{\infty} R(\tau) e^{-j\omega\tau} d\tau = \int_{-\infty-\infty}^{\infty} \int_{-\infty-\infty}^{\infty} x(t) x(t-\tau) e^{-j\omega\tau} dt d\tau$$

$$= \int_{-\infty}^{\infty} x(t) e^{-j\omega t} dt \int_{-\infty}^{\infty} x(t-\tau) e^{j\omega(t-\tau)} d\tau$$

$$= X(\omega) \int_{-\infty}^{\infty} x(t-\tau) e^{j\omega(t-\tau)} d\tau$$

Letting $t - \tau = n$ in the second integral, we have

$$F[R(\tau)] = X(\omega) \int_{-\infty}^{\infty} x(n)e^{j\omega n} dn$$

$$= X(\omega) X(-\omega) = |X(\omega)|^2$$

$$= \psi(\omega)$$

Autocorrelation function for power (periodic) signals

The autocorrelation function of a periodic signal with any period T is given by

$$R(\tau) = \operatorname{Lt}_{T \to \infty} \frac{1}{T} \int_{-T/2}^{T/2} x(t) \, x^*(t - \tau) \, dt$$

Properties of autocorrelation function for power signals

Following are the properties of autocorrelation function for power signals:

1. The autocorrelation function exhibits conjugate symmetry, i.e.

$$R(\tau) = R^*(-\tau)$$

Proof: We have

$$(E = (R(\tau) = Lt \int_{T \to \infty}^{T/2} \frac{1}{T} \int_{-T/2}^{T/2} x(t) x^*(t - \tau) dt$$

$$(E = (R(\tau) = Lt \int_{T \to \infty}^{T/2} \frac{1}{T} \int_{-T/2}^{T/2} x(t) x^*(t - \tau) dt$$

$$(E = (R(\tau) = Lt \int_{T \to \infty}^{T/2} \frac{1}{T} \int_{-T/2}^{T/2} x(t) x^*(t - \tau) dt$$

$$(E = (R(\tau) = Lt \int_{T \to \infty}^{T/2} \frac{1}{T} \int_{-T/2}^{T/2} x(t) x^*(t - \tau) dt$$

$$(E = (R(\tau) = Lt \int_{T \to \infty}^{T/2} \frac{1}{T} \int_{-T/2}^{T/2} x(t) x^*(t - \tau) dt$$

$$(E = (R(\tau) = Lt \int_{T \to \infty}^{T/2} \frac{1}{T} \int_{-T/2}^{T/2} x(t) x^*(t - \tau) dt$$

$$(E = (R(\tau) = Lt \int_{T \to \infty}^{T/2} \frac{1}{T} \int_{-T/2}^{T/2} x(t) x^*(t - \tau) dt$$

$$(E = (R(\tau) = Lt \int_{T \to \infty}^{T/2} \frac{1}{T} \int_{-T/2}^{T/2} x(t) x^*(t - \tau) dt$$

$$(E = (R(\tau) = Lt \int_{T \to \infty}^{T/2} \frac{1}{T} \int_{-T/2}^{T/2} x(t) x^*(t - \tau) dt$$

$$R^*(\tau) = Lt_{T \to \infty} \frac{1}{T} \int_{-T/2}^{T/2} x^*(t) \ x(t - \tau) \ dt$$

$$R^*(-\tau) = \operatorname{Lt}_{T \to \infty} \frac{1}{T} \int_{-T/2}^{T/2} x^*(t) \ x(t+\tau) \ dt = R(\tau) \text{ maintaining and a simple problem.}$$

$$R(\tau) = R^*(-\tau)$$

2. The autocorrelation function at origin is equal to the average power of that signal, i.e.

$$R(0) = P = U = \frac{1}{12} |\mathcal{A}(t)|^2 dt$$

2. The autocorrelation function at origin is equal to the average power of that signal, i.e.

$$R(0) = P = U = \frac{1}{15} |\mathcal{M}(t)|^2 dt$$

3. The autocorrelation function $R(\tau)$ has maximum value at the origin, i.e.

$$|R(\tau)| \le R(0)$$
 This is a second of the substitution of the substitution $|R(\tau)| \le R(0)$

The value of autocorrelation reduces as τ increases from origin.

- The autocorrelation function R(τ) and power spectral density S(ω) form a Fourier transform pair, i.e.
 R(τ) ← S(ω)
- The autocorrelation function is periodic with the same period as the periodic signal itself, i.e.

$$R(\tau) = R(\tau \pm nT), \quad n = 1, 2, 3, ...$$

Relation between Convolution and Correlation:

RELATION BETWEEN CONVOLUTION AND CORRELATION

There is a striking resemblance between the operation of convolution and correlation. Indeed the two integrals are closely related. To obtain the cross correlation of $x_1(t)$ and $x_2(t)$

according to the equation
$$R_{12}(\tau) = \int_{-\infty}^{\infty} x_1(t) x_2(t-\tau) dt$$
, we multiply $x_1(t)$ with function $x_2(t)$

displaced by τ sec. The area under the product curve is the cross correlation between $x_1(t)$ and $x_2(t)$ at τ . On the other hand, the convolution of $x_1(t)$ and $x_2(t)$ at $t = \tau$ is obtained by folding $x_2(t)$ backward about the vertical axis at the origin and taking the area under the product curve of $x_1(t)$ and the folded function $x_2(-t)$ displaced by τ . It, therefore, follows that the cross correlation of $x_1(t)$ and $x_2(t)$ is the same as the convolution of $x_1(t)$ and $x_2(-t)$.

The same conclusion can be arrived at analytically as follows:

The convolution of $x_1(t)$ and $x_2(-t)$ is given by

$$x_1(t) * x_2(-t) = \int_{-\infty}^{\infty} x_1(\tau) x_2(\tau - t) d\tau$$

Replacing the dummy variable τ in the above integral by another variable n, we have

$$x_1(t) * x_2(-t) = \int_{-\infty}^{\infty} x_1(n) x_2(n-t) dn$$

Changing the variable from t to τ , we get

$$x_1(\tau) * x_2(-\tau) = \int_{-\infty}^{\infty} x_1(n) x_2(n-\tau) dn = R_{12}(\tau)$$

Hence

$$R_{12}(\tau) = x_1(t) * x_2(-t) \Big|_{t=\tau}$$

Similarly,

$$R_{21}(\tau) = x_2(t) * x_1(-t)|_{t=\tau}$$

All of the techniques used to evaluate the convolution of two functions can be directly applied in order to find the correlation of two functions. Similarly, all of the results derived for convolution also apply to correlation.

If one of the function is an even function of t, let us say $x_2(t)$ is an even function of t, i.e.

$$x_2(t) = x_2(-t)$$

then the cross correlation and convolution are equivalent.

Part-2 Fourier Series

Fourier Series:

```
Fourier Series Representation of Periodic Signals:
       A periodic signal is the one which repeat itself
                -00<t<00.
periodically over
For example, 211) = A smoot is a periodic signal with
        Let us consider a signal zet) which is a sum
 of sine and cosine functions whose frequencies are
 integer multiple of wo as shown below.
  x(t) = ao + a, coswot + a2 cos 2wot + a3 cos 3 wot + ....
               .. + ax coskwet + by sinwet + by smawet +
      basmawat + .... + bk sin kwat -
```

Any periodic stand can be represented as an Infinite sum of sine and cosine functions. This series of sine and cosine terms are known as Prigonometric Fourier sentes and can be represented as alt = ao + & [ancosnwot + bn Smnwot]. where as is de component ao, an, by are called Fourier Sonicy coefficients.

Evaluation of Fourier Series Coefficients

to, -

To evaluate an and by we have to use the following formulae to
$$to$$

To cosmot as must to

To smoot Sin must to

To smoot to

To for to

To to

Evaluation of an: To find an multiply eqn (1) by cosmoot and integrate over one period ie., f orth cosmwotat = f ao cosmwotat + S E [an cosmoot + bn smnwot] cosmoot dt = f an commotat + Ean f commotat + & bn & Sinnwot cosm worldt

In the above fin 1st & grd terms are zero using ean (1) & (1) 50 alt) cosmwotat set) cosmost dt

Evaluation of bn: To find by multiply the equation (1) by sm mont and integrate over one period S x(t) Smmwotht = f as smmwotht t S E [an cos nost + bn sin nost] sin most dt toti
= fao sinm worldt + Ean f cosmoot smmwot dt
n=1 to + & both Sinnwot sin mwotht

From @ Ell 5 x (t) Sin mooth = 0 + 0+ bm = bm = = f xtt) Shimwoot dt to .. bn = 2 f alt) smnwo tot n=1,2,... .. In Trigonometric Fourier sorice att) can be represented alt) = ao + & [ancomwot + bn smnwet] where ao = + 5 alt)dt an = 2 totT xit) cos nwotat bn = 2 toff alt) Sinnwotht

Polar Fourier Series

Polar Founter series representation (or) Cosine Representation (or) Compact Trigonometric Fourser series Representation: The polar Fourier series is the modified form of the trigonometric fourier series. The polar Fourier series is derived from Trigonometric Fourier series as follows. Trigonometric Fourter series is given by xll = aot & [an cosnwot + bn sinnwot] - 0 all) = ao + & Janthi [an cosnwot + Nanthin sin moot Find to to

let
$$\cos \phi = \frac{an}{\sqrt{a_n^2 + b_n^2}}$$
 $g \sin \phi = \frac{bn}{\sqrt{a_n^2 + b_n^2}}$
 $from eqn (3) g (4)$
 $from eqns (3), (3) g (4)$
 $from eqns (4), (4), (4)$
 $from eqns (4), (4), (4)$
 $from eqns (4), (4)$
 $from eqns (4), (4)$
 $from eqns (4), (4)$
 $from eqns$

let
$$D_0 = a_0$$
 $G D_1 = \sqrt{a_1^n + b_1^n}$
 $Alt = D_0 + \frac{c_0}{E} D_1 \cos(n\omega_0 t - \phi)$

where $\phi = tan(\frac{b_1}{a_1})$

Exponential Fourier Series

Exponential Fourier sense to simpler and more compact hence it is mostly used. We will derive an exponential Fourier sense in the following way.

Exponential Fourier sense:

Prigonometric Fourier sense to given by

It = ao +
$$\sum_{n=1}^{\infty} \left[a_n \cos n w_0 t + b_n \sin w_0 t \right] - 0$$

Using Eulers identity

 $i\theta = \cos \theta + i\sin \theta$
 $i\theta = \cos \theta - i\sin$

With
$$= a_0 + \mathcal{E} \left[a_n \cos n\omega_0 + t + b_n \sin \omega_0 t \right] - 0$$

Using Euler's identity
$$e^{j\theta} = \cos \theta + j \sin \theta$$

$$e^{j\theta} = \cos \theta - j \sin \theta$$

$$\cos \theta = \frac{j\theta}{e} - \frac{j\theta}{e} \Rightarrow \cos \theta \cos \theta$$

$$\sin \theta = \frac{j\theta}{e} - \frac{j\theta}{e} \Rightarrow \sin \theta \cos \theta = \frac{j\theta}{e} - \frac{j\theta}{e} \cos \theta$$
From equs 0 , 9 9 9

With $= a_0 + \mathcal{E} \left[a_n \left(\frac{jn\omega_0 t}{e} - \frac{jn\omega_0 t}{e} \right) + \frac{jn\omega_0 t}{e} \right]$

$$b_n \left(\frac{jn\omega_0 t}{e} - \frac{jn\omega_0 t}{e} \right)$$

$$\chi(t) = a_0 + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} + \underbrace{b_n}_{2j} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{b_n}_{2j} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} + \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_{2j^2}}_{2j^2} \right] + \underbrace{\mathcal{E}}_{n=1}^{\infty} e^{jn\omega_0 t} \left[\underbrace{a_n}_{2} - \underbrace{j \underbrace{b_n}_$$

let
$$(0 = 00)$$

$$C_n = \frac{1}{2} [an-jbn]$$

$$C_{m} = \frac{1}{2} [an+jbn]$$

$$2(t) = Cot \underbrace{E}_{n} c_n e + \underbrace{E}_{n-j} C_{-n} e$$

$$Change the index -n to n in the and summation$$

$$2(t) = Cot \underbrace{E}_{n-j} c_n e + \underbrace{E}_{n-j} c_n e$$

$$2(t) = Cot \underbrace{E}_{n-j} c_n e + \underbrace{E}_{n-j} c_n e$$

$$2(t) = Cot \underbrace{E}_{n-j} e$$

$$2(t) = Cot$$

Evaluation of
$$C_n$$
:

$$C_n = \frac{1}{2} \left[a_n - j_{bn} \right]$$
but $a_n = \frac{2}{7} \int_{0}^{\infty} x(t) \cos n\omega_0 t \, dt - 3$

$$S_1 \quad b_n = \frac{1}{2} \int_{0}^{\infty} x(t) \sin n\omega_0 t \, dt - 3$$
From eqns $(0, 0) \in 3$

$$C_n = \frac{1}{2} \left[\frac{1}{7} \int_{0}^{\infty} x(t) \sin n\omega_0 t \, dt - j \right] \int_{0}^{\infty} x(t) \sin n\omega_0 t \, dt$$

$$C_n = \frac{1}{7} \left[\frac{1}{7} \int_{0}^{\infty} x(t) \cos n\omega_0 t \, dt - j \sin n\omega_0 t \, dt \right]$$

$$C_n = \frac{1}{7} \left[\frac{1}{7} \int_{0}^{\infty} x(t) \left[\cos n\omega_0 t - j \sin n\omega_0 t \, dt \right] \right]$$

$$C_n = \frac{1}{7} \left[\frac{1}{7} \int_{0}^{\infty} x(t) \left[\cos n\omega_0 t - j \sin n\omega_0 t \, dt \right] \right]$$

Relation Between Trigonometric and Exponential FS

Derive the relation between Trigonometric Fourier series and exponential Fourier series:

In Trigonometric Fourier series all is given by

$$alt = ao + \frac{ao}{n} [an carnus + bn sin must] - 0$$

The Exponential Fourier series alt is given by

 $alt = \frac{ao}{n} + \frac{tinuot}{n} = \frac{ao}{n} = \frac{tinuot}{n} = \frac{ao}{n}$

$$\chi(t) = co + c_{1}e^{2} + c_{2}e^{2} + c_{3}e^{2} + c_{4}e^{2} + c_{5}e^{2} + c_{5$$

alt =
$$Cot \stackrel{\infty}{\underset{n=1}{E}} \left[(cn+c_n) cosnuot + j(cn-c_n) sinnuot \right]$$

Compane egns (D & B)

 $Co = ao \qquad -B$
 $an = cn+c-n \qquad B$
 $bn = j(cn-c_n) - B$
 $egn (S) \times j$
 $jbn = (-1)(cn-c_n) - B$

add egn (S) & B (S)

 $an + jbn = a - Cn$
 $C-n = \frac{1}{2} \left[an + jbn \right]$

an =
$$c_n + c_n$$

 $jbn = -c_n + c_n$
 $jbn = -c_n + c_n$
 $a_n - jbn = a_n$
 $c_n = \frac{1}{2} [a_n - jb_n]$
 $c_n = \frac{1}{2} [a_n - jb_n]$
 $c_n = \frac{1}{2} [a_n + jb_n]$

Derive the polar Founter series from Exponential Fourner Series & hence prove that $D_n = 2 | C_n |$ an Exponential Fourter Senes all is given by alt) = E Cne Cne $\alpha(t) = co + \underbrace{e}_{n=1}^{\infty} c_n e + \underbrace{e}_{n=$ Co = 90 But we have cn = \frac{1}{2} [an-jbn] - @ cn = 1 [antibon] - 3

From eqns (5), (1) & (1) & (2)

$$a(t) = a_0 + \frac{a_0}{2} + \frac{1}{2} (a_1 - b_1) = a_0 + \frac{a_0}{2} + \frac{1}{2} (a_1 - b_1) = a_0 + \frac{a_0}{2} + \frac{1}{2} (a_1 - b_1) = a_0 + \frac{a_0}{2} + \frac{1}{2} (a_1 - b_1) = a_0 + \frac{a_0}{2} (a_1 - a_1) = a_0 + \frac{a_0}{2} (a_1 - a_2) = a_0 + \frac{a_0}{2}$$

From eqn (1) (8) (1)

$$2(t) = a_0 + \sum_{n=1}^{\infty} \sqrt{a_n^2 + b_n^2} \left(\cos \phi \cosh \phi + \sin \phi + \sin \phi + \sin \phi \right)$$

$$2(t) = a_0 + \sum_{n=1}^{\infty} \sqrt{a_n^2 + b_n^2} \left(\cos (n\omega o t - \phi) - (b) \right)$$

$$bet a_0 = 0_0 \quad \text{s.} \quad \sqrt{a_n^2 + b_n^2} = 0_n \quad -(b)$$

$$2(t) = 0_0 + \sum_{n=1}^{\infty} 0_n \cos(n\omega o t - \phi) - (b)$$

$$2(t) = 0_0 + \sum_{n=1}^{\infty} 0_n \cos(n\omega o t - \phi) - (b)$$

$$2(t) = \frac{1}{2} \left[(a_n - \frac{1}{2}) b_n \right]$$

$$1(c_n) = \frac{1}{2} \left[(a_n - \frac{1}{2}) b_n \right]$$

$$1(c_n) = \frac{1}{2} \left[(a_n - \frac{1}{2}) b_n \right]$$

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$$1(c_n) = \frac{1}{2} \left[(a_n - \frac{1}{2}) b_n \right]$$

Dirchlet Conditions

Convergence of Fourier Series: (Dirchlet Conditions): The Fourier series is convergent if the signal acity eatisfies the tollowing conditions. Those conditions are called Directlet conditions. is single valued property: 2(4) must have only one value at any time instant within the given interval. all) should have finite number of maxima and ti) finite peaks: minima coiting the given merval 2(4) should have at the most finite number of IID Finite Discontinuities: discontinuities within the given interval, W) Absolute integrability: 21t) should be absolute integrable 5 lawlet co i.e.

Symmetry Conditions: Any stand can be represented as a sum of even and odd functions. i.e., alt) = aelt) + aelt) where aett) = = = [2(+)+ 2(-+)] and 2011 = + [214) - 21-4)] symmetry These equations are used to derive the conditions.

is Odd functions have only sine terms:

In Prigonomethic Fourier seves

$$a_{n} = \frac{2}{7} \int_{0}^{7} \chi_{0}(t) cosnwot dt$$

$$\frac{1}{7} \int_{0}^{7} \chi_{0}(t) \int_{0}^{7} (codd + cosnwot dt)$$

$$\frac{1}{7} \int_{0}^{7} (codd + cosnwot dt) \int_{0}^{7} (codd + codd) \int_{$$

.. In Theoremetric F.S.

$$a(t) = a_0 + \sum_{n=1}^{\infty} [a_n rosnwot + b_n simmwot]$$
 $a(t) = a_0 + \sum_{n=1}^{\infty} [a_n rosnwot + b_n simmwot]$
 $a(t) = a_0 + \sum_{n=1}^{\infty} [a_n rosnwot + b_n simmwot]$

From the above eqn of $a(t)$ is odd, it contain only since tems

ti) Even functions have only cosme terms: for even functions a au(t) = 0 all) = delt) an= = S x(t) cosmoot at = Taelt) cosnwotat 4 (xit) cosowot 2 (alt) Sinnwot dt

bn =
$$\frac{2}{7}$$
 (act) Sinnwort at
= $\frac{2}{7}$ (act) Sinnwort at
= $\frac{2}{7}$ (act) Sinnwort at
even $\frac{2}{7}$ (act) Sinnwort at
 $\frac{2}{7}$

Half wave Symmetry:

A periodic stand to said to have a halfwave symmetry If it satisfies the following condition

The fourier senty expansions of such type of periodic signals contain odd harmonius only;

Summary of symmetry Conditions

Type of	cendition	Example	ao	90	bŋ	Property
Symmetry	ત્રામ=ત્રાન)	7(t) 7	20====================================	1 to garthe	0	cosine terms only
edd	α(H=-αH)	Alt)	E STATE OF FA	6	bn= 4 falt)smultdt	Sine terms only
Halfwave	제나=-제나가)	7(t)		an= 4 Sate) Cosmootot	the # Jak Simmootdt	odd n only

Properties of Fourier Series

```
Properties of Fourier Series:
1) Linearity:
 Let us consider two signals alt) and git) with period T.
 If FS [ 21(6)] = Co
    & FS [xx(t)] = Dn
 Then FS [AZILE) + BXILE)] = ACn + Bbn
 Proof = FS[xus] = Co = + suse dt
   FS[AZILE)+BZILE)] = + (AZILE)+BZILE))e
        = Acn + B.Dn
    .. FS[AQILE)+BQL(t)] = ACn +BDn
```

2) Time shifting:

Stmt: If
$$FS[alt] = Cn$$

Then $FS[alt-to] = e^{-jn\omega to} Cn$

Then $FS[alt-to] = e^{-jn\omega t} Cn$

Proof: $FS[alt] = Cn = \frac{1}{1} \int_{17}^{1} alt-to e^{-jn\omega t} dt$

$$FS[alt-to] = \frac{1}{1} \int_{17}^{1} alt-to e^{-jn\omega t} dt$$

$$FS[alt-to] = \frac{1}{1} \int_{17}^{1} a(p) e^{-jn\omega t} dp$$

4) Time scaling =

The time scaled signal of alt) is denoted by anoth. If act, the resulting time scaled stand is expanded version of alt). If an, then anoth is compressed version of alt). If the fundamental frequent period of 21th is T, then the fundamental period of areat) is I and the fundamental frequency is awo. The fourier series coefficients of acat) is some as alt) but the harmonics are now at the frequencies two traws, + 3000, + 4000, ----

5) Frequency shifting: Stmt: If FS[zeti] = Cn Then FS[eswota(t)] = Cn-no FS[xIII] = Cn = + fxII)e

Proof:
$$FS[2(H)] = Cn = \frac{1}{4} \int_{a}^{b} \pi(H)e^{-\frac{1}{2}nwot} dt$$

$$\pi(t) = \sum_{n=-\infty}^{\infty} Cne^{-\frac{1}{2}nwot} (jnwo)$$

$$\frac{d}{dt}\pi(t) = \sum_{n=-\infty}^{\infty} (cne^{-\frac{1}{2}nwot}) \int_{a}^{b} \pi(t)e^{-\frac{1}{2}nwot}$$

$$\frac{d}{dt}\pi(t) = \int_{a}^{\infty} (cne^{-\frac{1}{2}nwot}) \int_{a}^{\infty} \pi(t)e^{-\frac{1}{2}nwot}$$

let
$$t-p=m$$
 $t=m+p$
 $dt=dm$

$$Trick = trick =$$

(payr) serve

9) Parsevals Theorem: Statement: If all) is the periodic signal with fourier coefficients on, then the average power of the periodic signal is given by E |a|2 Proof: P=It + 5 |xit) at = は ナデなばないは (AIN = HIX (AIN = but alt) = S che Take complex conjugate on both sides atly = E chejnwot · P = It + 5 a(t) & che noot dt

Take complex conjugate on both sides

$$\pi^*(t) = \underbrace{\mathcal{E}}_{n=-\infty}^{\infty} C_n e^{jn\omega_0 t}$$

$$\Rightarrow P = \underbrace{It}_{T\to\infty} + \underbrace{\int_{T\to\infty}^{T}}_{T\to\infty} \pi(t) \underbrace{\mathcal{E}}_{n\to\infty}^{\infty} C_n e^{jn\omega_0 t} dt$$

$$= \underbrace{It}_{T\to\infty} + \underbrace{\int_{T\to\infty}^{T}}_{T\to\infty} \pi(t) e^{jn\omega_0 t} dt$$

$$= \underbrace{It}_{T\to\infty} + \underbrace{\mathcal{E}}_{T\to\infty} + \underbrace{\int_{T\to\infty}^{T}}_{T\to\infty} \pi(t) e^{jn\omega_0 t} dt$$

$$= \underbrace{It}_{T\to\infty} + \underbrace{\mathcal{E}}_{T\to\infty} + \underbrace{\int_{T\to\infty}^{T}}_{T\to\infty} \pi(t) e^{jn\omega_0 t} dt$$

$$= \underbrace{It}_{T\to\infty} + \underbrace{\mathcal{E}}_{T\to\infty} + \underbrace{\int_{T\to\infty}^{T}}_{T\to\infty} \pi(t) e^{jn\omega_0 t} dt$$

$$= \underbrace{\mathcal{E}}_{T\to\infty} + \underbrace{\mathcal{E}}_{$$

Start: 1) If
$$FS[x|tt] = Cn$$

Then $FS[x^*(tt)] = C_n^*$

Proof: $x(t) = \mathcal{E}$ on e
 $x^*(t) = \mathcal{E}$ on e

Problems

1 obtain the trigonometric fourter series for the periodic stand shown below え(し) Soln: and zut) = et Trigonometric fourier series representation of xith is xit) = ao + E [an cosnwot + bn 8m nwot] where an = + fattadt an = = falt cosnwotat 되 bn = 목 f xiti Smowoth

Evaluation of ao: ao = + Salt) dt = 10.5 o e dt $= g \left[-\overline{e}^{t} \right]_{n}$ = -2[e-e] 0.7869

Evaluation of an:

$$a_{1} = \frac{2}{T} \int_{CT} T(t) \cos n\omega_{0} t dt$$

$$a_{2} = \frac{2}{0.5} \int_{C} e^{t} \cos n\omega_{0} t dt$$

$$a_{3} = \frac{2}{0.5} \int_{C} e^{t} \cos n\omega_{0} t dt$$

$$= \frac{2\pi}{0.5} \int_{C} e^{t} \cos n\omega_{0} t dt$$

$$= \frac$$

$$a_{n} = 4 e^{\frac{-0.5}{1 + (u_{n} \kappa)^{2}}} \left[-\cos 4n \kappa (o.s) + 4n \kappa \sin 4n \kappa (o.s) \right]$$

$$= \frac{4 \cdot e^{\circ}}{1 + (u_{n} \kappa)^{2}} \left[-1 + 0 \right]$$

$$= \frac{4 \cdot (o.606)}{1 + (u_{n} \kappa)^{2}} \left[-\cos 2n \kappa + 4n \kappa \sin 2n \kappa \right] + \frac{4}{1 + (u_{n} \kappa)^{2}}$$

$$= \frac{4}{1 + (u_{n} \kappa)^{2}} \left[-o.606 + o.41 \right]$$

$$= \frac{4 \cdot (o.244)}{1 + (u_{n} \kappa)^{2}}$$

$$a_{n} = \frac{1.576}{1 + (u_{n} \kappa)^{2}}$$

Evaluation of bn:

$$bn = \frac{2\pi}{T} \int_{TT}^{\infty} x(t) \operatorname{Simnwoth} dt$$

$$= \frac{2\pi}{0.5} \int_{0.5}^{\infty} e^{t} \operatorname{Sim} n \operatorname{an} t dt$$

$$= \frac{2\pi}{0.5} \int_{0.5}^{\infty} e^{t} \operatorname{Sim} n \operatorname{an} t dt$$

$$= \frac{2\pi}{0.5} \int_{0.5}^{\infty} e^{t} \operatorname{Sim} n \operatorname{an} t dt$$

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$$= \frac{2\pi}{0.5} \int_{0.5}^{\infty} e^{t} \operatorname{Sim} n \operatorname{an} t dt$$

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$$= \frac{2\pi}{0.5} \int_{0.5}^{\infty} e^{t} \operatorname{Sim} n \operatorname{an} t dt$$

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$$= \frac{2\pi}{0.5} \int_{0.5}^{\infty} e^{t} \operatorname{Sim} n \operatorname{an} t dt$$

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$$= \frac{2\pi}{0.5} \int_{0.5}^{\infty} e^{t} \operatorname{Sim} n \operatorname{an} t dt$$

$$= \frac{2\pi}{0.5} \int_{0.5}^{\infty} e^{t} \operatorname{An} t dt$$

$$= \frac{2\pi}{0.5} \int_{0.5}^{\infty} e$$

$$bn = 4 \left(\frac{e^{t}}{H(unn)} \left(-\frac{8m}{4n\pi}t - 4m\pi \cos 4n\pi t \right) \right)$$

$$= 4 \cdot \left[\frac{e^{-5}}{1 + (unn)} \left(-\frac{8m}{4n\pi} (o \cdot 5) - 4m\pi \cos 4n\pi (o \cdot 5) - 4m\pi (o \cdot 5)$$

July 9, 2020 111 GKS@CMRIT

3 Obtain the exponential Fourier sens for the figure shown in problem 10

Exponential Fourier series representation of 2111 is

where
$$Cn = T$$
 alther dt

$$C_{n} = \frac{1}{0.5} \int_{0.5}^{0.5} e^{\frac{1}{2}} e^{-\frac{1}{2}n\omega_{0}t} dt$$

$$= \frac{1}{0.5} \int_{0.5}^{0.5} e^{\frac{1}{2}} e^{-\frac{1}{2}n\omega_{0}t} dt$$

$$C_{n} = \frac{1}{0.5} \int_{0}^{0.5} e^{\frac{1}{2} e^{-\frac{1}{2}n\omega_{0}t}} dt$$

$$= a \int_{0.5}^{0.5} e^{\frac{1}{2} e^{-\frac{1}{2}n\omega_{0}t}} dt$$

$$= \frac{-a}{1+j4n\pi} \left(\begin{array}{c} -0.5 - j4n\pi(0.5) \\ e \end{array} \right) - 1$$

$$= \frac{-a}{1+j4n\pi} \left(\begin{array}{c} -0.5 - j2n\pi \\ e \end{array} \right) - 1$$

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$$= \frac{-a}{1+j4n\pi} \left(\begin{array}{c} -0.5 - j2n\pi \\ e \end{array} \right) - 1$$

$$= \frac{-a}{1+j4n\pi} \left(\begin{array}{c} 0.606 - 1 \\ e \end{array} \right)$$

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$$= \frac{-a}{1+j4n\pi} \left(\begin{array}{c} 0.7869 \\ e \end{array} \right) - 1$$

$$= \frac{a}{1+j4n\pi} \left(\begin{array}{c} 0.7869 \\ e \end{array} \right) - 1$$

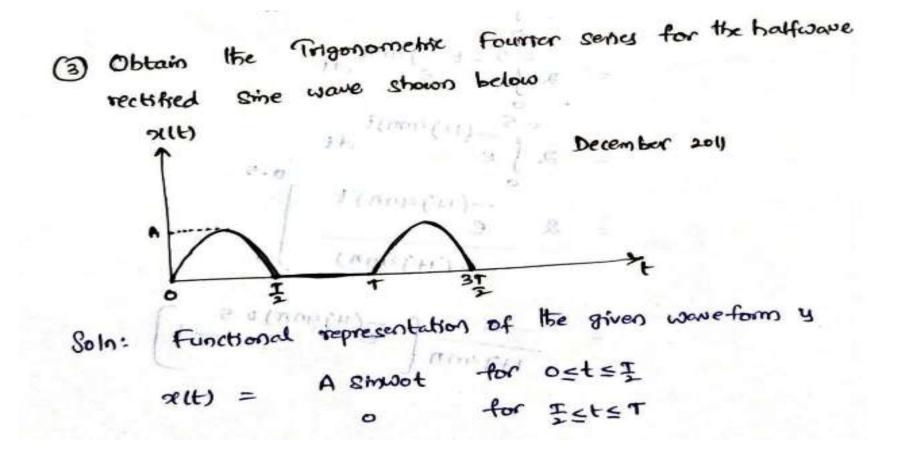
$$= \frac{a}{1+j4n\pi} \left(\begin{array}{c} 0.7869 \\ e \end{array} \right) - 1$$

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$$= \frac{a}{1+j4n\pi} \left(\begin{array}{c} 0.7869 \\ e \end{array} \right) - 1$$



here
$$T = 2\pi$$
 $\therefore \omega_0 = \frac{2\pi}{2\pi} = 1$

Trigonometric fourier series representation of $z(t)$ is

 $z(t) = a_0 + \sum_{n=1}^{\infty} [a_n \cos n\omega_0 t + b_n \sin n\omega_0 t]$

Where $a_0 = \frac{1}{T} \int z(t) dt$
 $z(t) = a_0 + \sum_{n=1}^{\infty} [a_n \cos n\omega_0 t] \int z(t) \sin n\omega_0 t dt$
 $z(t) = a_0 + \sum_{n=1}^{\infty} [a_n \cos n\omega_0 t] \int z(t) \sin n\omega_0 t dt$

Evaluation of as:

$$a_{0} = \frac{1}{T} \int_{CT2}^{T} \pi(t) dt$$

$$= \frac{1}{2\pi} \int_{0}^{\pi} \pi(t) dt + \int_{0}^{\pi} \pi(t) dt$$

$$= \frac{1}{2\pi} \int_{0}^{\pi} A \operatorname{Sin} \omega_{0} t dt + 0$$

$$= \frac{A}{2\pi} \int_{0}^{\pi} \operatorname{Sin} t dt$$

$$= \frac{A}{2\pi} \int_{0}^{\pi} \operatorname{Sin} t dt$$

$$= \frac{A}{2\pi} \int_{0}^{\pi} \operatorname{Cost} \int_{0}^{\pi} = \frac{A}{2\pi} \left(\cos \pi - \cos \phi \right)$$

$$= \frac{A}{2\pi} \left(-1 - 1 \right) = \frac{2A}{2\pi} = \frac{A}{\pi}$$

$$a_{0} = \frac{A}{\pi}$$

Evaluation of an:

$$an = \frac{2}{T} \int_{0}^{T} x(t) \cos n\omega_{0} t dt$$

$$= \frac{2}{2\pi} \int_{0}^{T} A \sin \omega_{0} t \cos n\omega_{0} t dt$$

$$= \frac{A}{T} \int_{0}^{T} 8m t \cos nt dt \qquad : \omega_{0} = 1$$

$$an = \frac{A}{2\pi} \int_{0}^{\pi} 2 \sin t \cos t \, dt$$

$$= \frac{A}{2\pi} \int_{0}^{\pi} (8 \sin (t + nt)) \, dt$$

$$= \frac{A}{2\pi} \int_{0}^{\pi} (8 \sin (t + nt)) \, dt$$

$$= \frac{A}{2\pi} \int_{0}^{\pi} (8 \sin (t + nt)) \, dt$$

$$= -\frac{A}{2\pi} \left[\frac{\cos(Hn)t}{(Hn)} + \frac{\cos(Hn)t}{(Hn)} \right]_{0}^{T}$$

$$= -\frac{A}{2\pi} \frac{1}{(Hn)} \left[\cos \pi(nH) - 1 \right] - \frac{A}{2\pi} \frac{1}{(Hn)} \left[\cos \pi(mh) - 1 \right]$$
For $n = \pm 1, \pm 3, \pm 5, - 1$

$$\cos (Hn)\pi = 1 + \sin \cos (Hn)\pi^{3} = 1$$

$$\cos (Hn)\pi = 1 + \cos (Hn)\pi^{3} = 1$$

$$\sin \pi = -\frac{A}{2\pi} \frac{1}{(Hn)} \left[Hn \right] - \frac{A}{2\pi} \frac{\pi}{(Hn)} \left[Hn \right]$$

$$\cos \pi = -\frac{A}{2\pi} \frac{1}{(Hn)} \left[Hn \right] - \frac{A}{2\pi} \frac{\pi}{(Hn)} \left[Hn \right]$$

$$\sin \pi = -\frac{A}{2\pi} \frac{1}{(Hn)} \left[Hn \right] - \frac{A}{2\pi} \frac{\pi}{(Hn)} \left[Hn \right]$$

$$\sin \pi = -\frac{A}{2\pi} \frac{1}{(Hn)} \left[Hn \right] + \frac{1}{2\pi} \frac{\pi}{(Hn)} \left[Hn \right]$$

For
$$n = \pm 2$$
, ± 4 , ± 6 ,...

 $\cos (Hn)\pi = -1$ $= 4$ $\cos (Hn)\pi = -1$

$$\therefore an = -\frac{A}{2\pi} \cdot \frac{1}{1+n} \left[-1 - 1 \right] - \frac{A}{2\pi} \left(\frac{1}{1-n} \right) \left[-1 - 1 \right]$$

$$= \frac{A}{\pi} \left(\frac{1}{1+n} \right) + \frac{A}{\pi} \left(\frac{1}{1-n} \right)$$

$$= \frac{A}{\pi} \left(\frac{1}{1+n} + \frac{1}{1-n} \right)$$

Evaluation of by:

$$b_{1} = \frac{2}{T} \int_{CT7}^{\infty} x(t) \sin n\omega_{0} t dt$$

$$= \frac{2}{2\pi} \int_{0}^{\infty} A \sin \omega_{0} t \sin n\omega_{0} t dt$$

$$= \frac{A}{\pi} \int_{0}^{\infty} 8 \sin t \sin n t dt$$

$$= \frac{A}{2\pi} \int_{0}^{\infty} (\cos (t - n)t - \cos (t + n)t) dt$$

$$= \frac{A}{2\pi} \int_{0}^{\infty} (\cos (t - n)t) dt$$

$$= \frac{A}{2\pi} \left[\frac{\sin (t - n)t}{t - n} - \frac{\sin (t + n)t}{t - n} \right]_{0}^{\infty}$$

$$= \frac{A}{2\pi} \left[\frac{\cos (t - n)t}{t - n} - \frac{\sin (t + n)t}{t - n} \right]_{0}^{\infty}$$

$$= \frac{A}{2\pi} \left[\frac{\cos (t - n)t}{t - n} - \frac{\sin (t + n)t}{t - n} \right]_{0}^{\infty}$$

$$a_0 = \frac{A}{\pi}$$

$$a_1 = \frac{2A}{\pi(1-n^2)} \quad \text{for } n = 2, 4, 6, 8, \dots$$

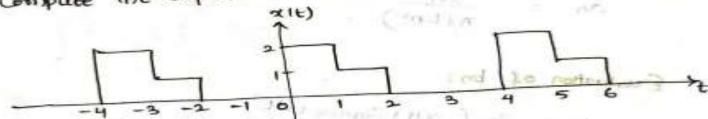
$$b_1 = 0$$

$$x(t) = \frac{A}{\pi} + \frac{2A}{\pi(1-n^2)} cosnoot \quad \text{for } n = 2, 4, 6, 8, \dots$$

$$x(t) = \frac{A}{\pi} + \frac{2A}{\pi(1-n^2)} cosnoot + \frac{2A}{\pi(1-4^2)} cosnoot + \dots$$

$$x(t) = \frac{A}{\pi} + \frac{2A}{\pi(1-n^2)} cosnoot + \frac{2A}{\pi(1-4^2)} cosnoot + \dots$$

$$x(t) = \frac{A}{\pi} + \frac{2A}{3\pi} cosnoot + \frac{2A}{15\pi} cosnoot + \dots$$



$$= \pm \frac{1}{2} \frac{1}{100} + \frac{1}{4} \frac{1}{100} = \frac{1}{100$$

Fourier Spectrum

Fourier Spectrum:

Fourier spectrum of a periodic signal alth com be obtained by plotting the fourier coefficients versus w. The plot of amplitude of fourier coefficients versus w is known as amplitude spectra. The plot of phase of fourier coefficients versus w is known as phase spectra. The two plots together are known as Fourier frequency spectra of 21t). That is in Fourier spectra life amplitude and phase of life fourier coefficients are plotted as a function of frequency. So this type of representation is also known as frequency domain representation of x(t).

Note: The spectrum exists only at discrete frequencies noo where n=0.112,.... Thus like spectrum is mot continuous but exists only at some discrete values of wo and is known as discrete spectrum or line spectrum.

The trigonometric representation of periodic signal alt) contains both sine and cosine terms with the sy-we amplitude coefficients but with no phase angles. In cosine representation all the Fourier coefficients one the with a phase angle on we can plot amplitude spectra (on us w) and phase spectra we can plot amplitude spectra (on us w) and phase spectra (on us w). In this, fourier coefficients present only for the frequencies and may be called as single sided spectra.

En Exponential Fourier series the penalic signal is expressed as sum of exponential function of complex frequencies: 0, ±300, ±200, The amplitudes on are complex and can be represented by magnitude and phase. Therefore we can plot two spectra. The spectra can be plotted for both the and—ve two spectra. Hence the name two sided spectra.

Complex Fourier Spectrum:

The complex Fourier series representation of a function x(t) is equivalent to resolving the function in terms of harmonically related components of the fundamental frequency ω_0 (or f_0). A complex weighting factor F_n (or C_n), called the *spectral amplitude* is assigned to each harmonic component. Graphical representation of a spectral amplitude along with a spectral phase is called the *complex frequency spectrum*.

An amplitude spectrum without phase information does not specify the waveform; because, in general, F_n (or C_n) is a complex quantity. Therefore, such a spectrum is called the *complex frequency spectrum*.

However, if F_n (or C_n) is purely real or purely imaginary, the phase spectrum can be disregarded.

Figure 4.5 below shows a typical amplitude spectrum, where a vertical line has been drawn at each harmonic frequency and the height of the line represents the amplitude at the corresponding harmonic frequency. This spectrum is known as the discrete spectrum or line spectrum and exists only at discrete frequencies that are harmonically related. Figure 4.5(a) represents the spectrum of a trigonometric Fourier series extending from 0 to ∞ , producing a one-sided spectrum as no negative frequencies exist here. Figure 4.5(b) represents the spectrum of a complex exponential Fourier series extending from $-\infty$ to ∞ , producing a two-sided spectrum.

The amplitude spectrum of the exponential series is symmetrical about the vertical axis. This is true for all real periodic functions.

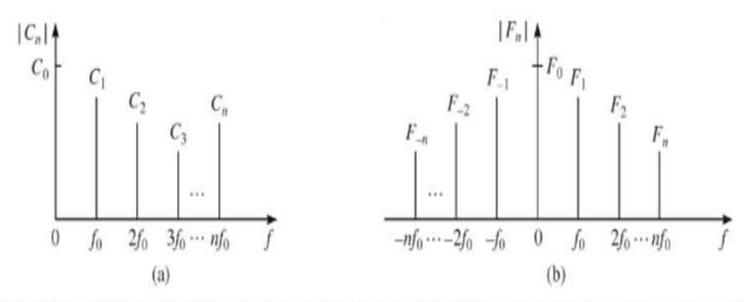


Figure 4.5 Complex frequency spectrum for (a) trigonometric Fourier series and (b) complex exponential Fourier series.

If F_n is a general complex number, then

$$F_n = |F_n| e^{j\theta_n}$$

$$F_{-n} = |F_n| e^{-j\theta_n}$$

$$|F_n| = |F_{-n}|$$

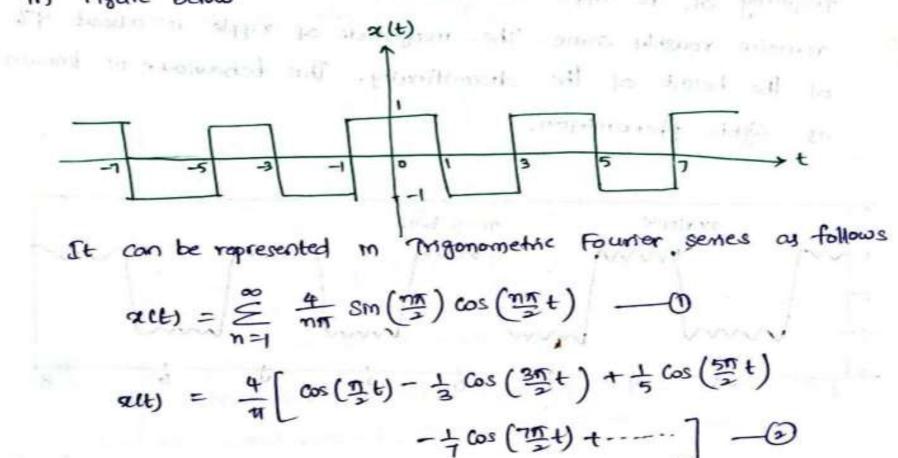
٠.

Hence, the magnitude spectrum is symmetrical about the vertical axis passing through the origin and, thus, it is an even function of ω_n . It is called the *even symmetry* of the magnitude spectrum. Also, θ_n is the phase of F_n and $-\theta_n$ is the phase of F_{-n} . So, the phase spectrum is antisymmetrical about the vertical axis. It is called the *odd symmetry* of phase spectrum.

Accordingly, for a real-valued periodic signal, x(t), the magnitude spectrum is symmetrical and the phase spectrum is anti-symmetrical about the vertical axis passing through the origin. When x(t) is real, then $F_{-n} = F_n^*$, the complex conjugate.

Gibbs Phenominon:

Consider the square wave with time period T as shown in Agure below



The above egn is m on infinite senses. Now let us approximate π it) by a finite value π egn(3) can be truncated as $\pi(11) = \frac{4}{\pi} \left[\cos \left(\frac{\pi}{2}t\right) - \frac{1}{3}\cos \left(\frac{3\pi}{2}t\right) + \frac{1}{5}\cos \left(\frac{5\pi}{2}t\right) - \frac{1}{3}\cos \left(\frac{3\pi}{2}t\right) + \frac{1}{5}\cos \left(\frac{5\pi}{2}t\right) + \frac{1}{5}\cos \left($

The below figure shows the plot of 2014) for n=5 and n=15. We can observe that the truncated Fourier series approaches 214) as n increases. That is an error between 214) and 2014) decreases as n increases.

However at the discontinuity, and) exhibits an ascillatory behaviour and has ripples on both side. As n moreases the frequency of the ripple increases, but the amplitude of the ripples remains roughly some. The magnitude of ripple is about 9% of the height of the discontinuity. This behaviour is known as Gibbs phenominum.

