Extraction of Transverse Single Spin Asymmetry in J/ψ Production in $p\vec{p}$ Interactions at 120 GeV Beam Energy

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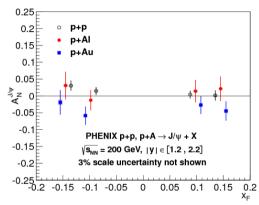
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Transverse Single Spin Asymmetry in J/ψ Production

- In $p\vec{p}$ collisions, the transverse single spin asymmetry (TSSA), A_N , is defined as the amplitude of the azimuthal angular modulation of the outgoing particle's scattering cross section with respect to the transverse spin direction of the polarized proton.
- The asymmetry can be written as function of azimuthal angle ϕ_S^{-1} :

$$A(\phi_S) = \frac{N^{\uparrow}(\phi_S) - N^{\downarrow}(\phi_S)}{N^{\uparrow}(\phi_S) + N^{\downarrow}(\phi_S)} = A_N \sin(\phi_S)$$



■ PHENIX results² shows $A_N^{J/\psi}$ 3 as a function of x_F . In the p+p data a $\sim 2\sigma$ positive A_N in the backward higher x_F bins. The results for other x_F bins are consistent with zero.

 $^{^{1}\}phi_{S}$ is the angle between $\vec{S}_{\mathrm{target}}$ and $\vec{p}_{TJ/\psi}$.

²C. Aidala et al., Phys. Rev. D 98, 012006, arXiv: 1805.01491 (hep-ex) (2018).

³PHENIX convention: x_F is measured w.r.t. p, SpinQuest convention: x_F is measured w.r.t. \vec{p} , q > q

SpinQuest Experiment

- SpinQuest is a fixed-target Dimuon experiment at Fermilab, using an unpolarized 120 GeV proton beam incident on a polarized solid ammonia target.
- SpinQuest measurements will allow us to test models for the internal transverse momentum and angular momentum structure of the nucleon.
- In $p\vec{p}$ collisions, J/ψ particles are primarily produced by strong interaction with $q\bar{q}$ annihilation and gg fusion.
- Our goal is to measure A_N with an absolute error $\mathcal{O}(10^{-2})$ for a few p_T and/or x_F bins.

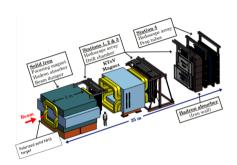


Figure 1: SpinQuest spectrometer.⁴

■ In this presentation, we demonstrate the analysis procedure and extraction of single spin asymmetry (A_N) with kinematics $0.0 \text{ GeV/c} < p_T < 2.0 \text{ GeV/c}$ and $0.4 < x_F < 0.8$.

Data Generation

- Simulated data were generated with kinematics:
 - J/ψ events were considered as signal events.

$$xF = [-0.2, 1.0]$$

where x_E is the Feynman x.

■ Drell-Yan events were considered as background events.

$$xF = [-0.2, 1.0]$$

mass = [1.0, 6.0]

■ Asymmetry was introduced by weighting the data;

$$w_{A_N} = 1 + A_N \sin(\phi_S + \phi_{\text{phase}})$$

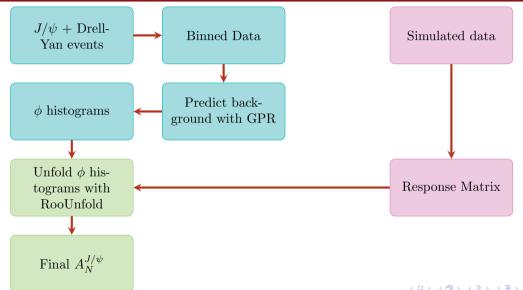
 $w_{\text{Total}} = w_{\text{Gen.}}(mass, x_F) \times w_{A_N}$

where ϕ_S is the angle between \vec{S}_{target} and $\vec{p}_{TJ/\psi}$ and $\phi_{\text{phase}} = 0$. for spin up and $\phi_{\text{phase}} = \pi$ for spin down.

Asymmetry values are set as $A_N^{J/\psi}=0.2$ for J/ψ events and $A_N^{BG}=0.1$ for Drell-Yan events.⁵

⁵Dilution factor of the NH3 was not considered in this study.

Analysis Procedure



Gaussian Process Regression (GPR)

- The Gaussian process model is a probabilistic supervised machine learning technique used in classification and regression tasks. A Gaussian process regression (GPR) model can make predictions incorporating prior knowledge (kernels) and provide uncertainties of the predictions.⁶
- In this analysis, the Radial-Basis Function (RBF) kernel was used as the kernel function in GPR class in sklearn library.

$$k(x_i, x_j) = \exp\left[-\frac{d^2(x_i, x_j)}{2l^2}\right]$$

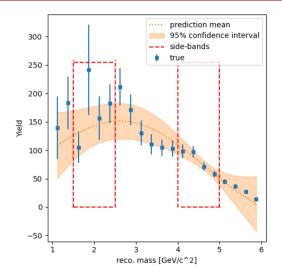
where l is the length scale of the kernel and $d(\cdot, \cdot)$ is the Euclidean distance.⁷

- We fit this kernel in side-band regions on either side of the J/ψ invariant mass peak. Then we used the trained kernel to predict the background in the J/ψ peak region.
- Our first goal is to extract the background under the J/ψ peak using the GPR method with good statistical precision.

⁶C. E. Rasmussen, C. K. I. Williams, Gaussian Processes for Machine Learning, (The MIT Press, Nov. 2005), ISBN: 9780262256834, (https://doi.org/10.7551/mitpress/3206.001.0001).

⁷F. Pedregosa et al., the Journal of machine Learning research 12, 2825-2830 (2011).

Sanity Check



- We used the Drell-Yan mass distribution in different p_T and x_F bins to check the GPR prediction.background
- As shown in figure 9, the prediction from GPR method agrees with the Drell-Yan events with 95% confidence interval in the J/ψ mass region.

Figure 2: GPR prediction for background. The side-bands are given in the red dashed lines.

Predicted Background

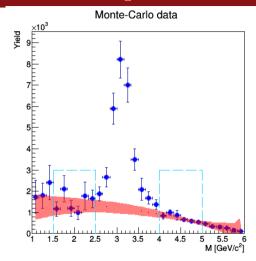


Figure 3: Mass histogram for 1st p_T bin and 1st ϕ bin. Predicted background is given in shaded red region. Side-bands are indicated in dashed blue lines.

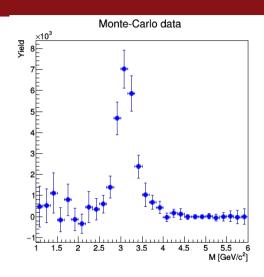


Figure 4: J/ψ signal after subtracting the background.

Unfolding ϕ Distributions

- Unfolding in high energy physics represents the correction of measured spectra in data for the finite detector efficiency, acceptance, and resolution from the detector to particle level.
- The equation of unfolding⁸;

$$\vec{P} = \frac{1}{\epsilon} M^{-1} \eta (\vec{D} - \vec{B})$$

where \vec{D} is the data spectrum, \vec{B} is the background spectrum, η acceptance correction, M^{-1} is the migration matrix, ϵ is the detector efficiency and \vec{P} is the unfolded spectrum.

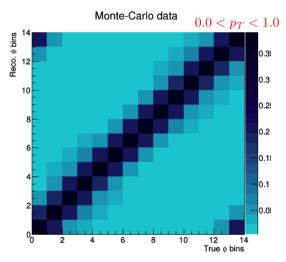
- We calculate the response matrix with Drell-Yan events without any asymmetry included. We used the iterative Bayesian method in ROOUnfold library to unfold the ϕ distributions.⁹
- Our second goal is to correct the bin-by-bin migration using the iterative Bayesian unfolding.



⁸P. Baron, Acta Phys. Polon. B **52**, 863, arXiv: 2104.03036 (hep-ex) (2021).

⁹B. Wynne, arXiv: 1203.4981 (physics.data-an) (Mar. 2012).

Response Matrix for p_T Bins



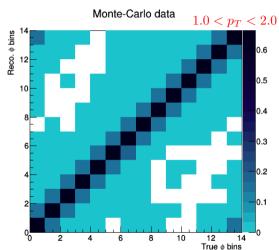


Figure 5: Reco. ϕ vs. true ϕ .

Response Matrix for x_F Bins

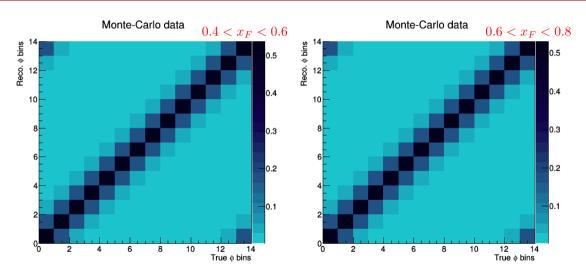
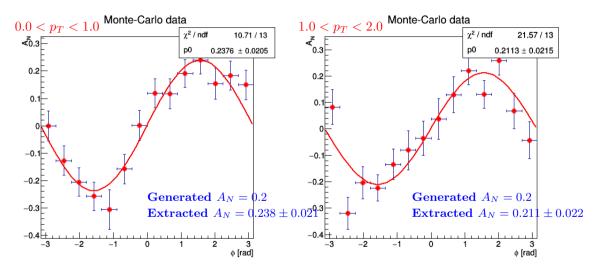


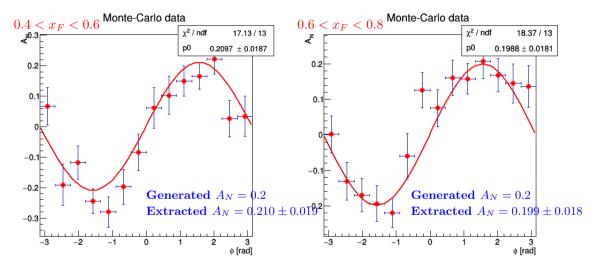
Figure 7: Reco. ϕ vs. true ϕ .

Unfolded $A_N^{J/\psi}$ in p_T Bins



 ${\bf Figure~9:~Unfolded~asymmetry.}$

Unfolded $A_N^{J/\psi}$ in x_F Bins



 ${\bf Figure~11:~Unfolded~asymmetry.}$

Extracted A_N

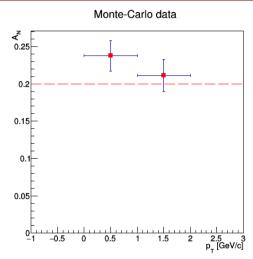


Figure 13: Extracted asymmetry for p_T bins. Generated asymmetry is shown in red dashed line.

Monte-Carlo data

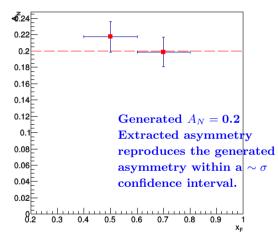


Figure 14: Extracted asymmetry for x_F bins. Generated asymmetry is shown in red dashed line.

Summary

- Gaussian process regression (GPR) is a supervised machine learning method that can be used to predict the background under the J/ψ peak.
- Using GPR method with the RBF kernel, background of the J/ψ mass can be predicted with 95% confidence interval.
- Using iterative Bayesian unfolding we can correct the bin-by-bin migration.
- Using these techniques (GPR+Unfolding), the extracted asymmetry reproduces the generated asymmetry within a $\sim \sigma$ confidence interval.
- SpinQuest does not overlap with PHENIX kinematics.
 - In PHENIX $|x_F| < 0.3$
 - In SpinQuest $|x_F| > 0.4$

SpinQuest will explore a new region of kinematics. Measurement for J/ψ transverse single spin asymmetry can be extracted in a few weeks of data taking with good statistical precision.

- Acknowledgement:
 - This work is supported by the US Department of Energy, Office of Science, Medium Energy Nuclear Physics Program.

Backup Slides

J/ψ Particle

- $\blacksquare J/\psi$ is a vector meson which is a $c\bar{c}$ bound state.
- Discovered by Burton Richter and Samuel Ting in 1974. Awarded Nobel price for the discovery in 1976.
- In $p\vec{p}$ collisions, J/ψ particles are primarily produced by strong interaction with $q\bar{q}$ annihilation and gg fusion.

$J/\psi(1S)$	$I^{G}(J^{PC}) = 0^{-}(1^{-})$
Mass $m = 30$	96.900 ± 0.006 MeV
Full width Γ =	$= 92.9 \pm 2.8 \; {\sf keV} ({\sf S} = 1.1)$
$\Gamma_{ee} = 5.53 \pm$	0.10 keV
Γ_{ee} < 5.4 eV	/, CL = 90%

$J/\psi(1S)$ DECAY MODES		Scale factor/ p Confidence level (MeV/c)	
hadrons	(87.7 ± 0.5)%	_	
virtual $\gamma \to \text{hadrons}$	(13.50 ± 0.30) %	-	
ggg	(64.1 ± 1.0) %	-	
γgg	(8.8 ± 1.1)%	-	
$e^{+}e^{-}$	(5.971± 0.032) %	1548	
$e^+e^-\gamma$	[a] ($8.8 \pm 1.4) \times 10^{-3}$	1548	
$\mu^{+}\mu^{-}$	(5.961± 0.033) %	1545	

Figure 16: J/ψ properties. 11

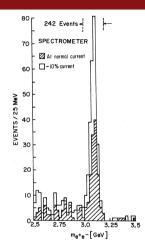


Figure 15: Mass spectrum showing the existence of J/ψ . ¹⁰



¹⁰ J. J. Aubert et al., Adv. Exp. Phys. 5, 128 (1976).

¹¹P. A. Zyla et al., PTEP **2020**, 083C01 (2020).

J/ψ Production

Color evaporation model (CEM), Color Singlet model (CSM) and Color Octet model (COM) are three most prominent models developed to understand the production of J/ψ particle. All three models attempt to factorize the J/ψ production into a relativistic part describing the production of $c\bar{c}$ $d\sigma_{c\bar{c}[n]}$, and a non-relativistic part describing the bound state of two quarks $F_{c\bar{c}[n]}(\Lambda)$;

$$d\sigma(J/\psi + X) = \sum_{n} \int d\Lambda \frac{d\sigma_{c\bar{c}[n]+X}}{d\Lambda} F_{c\bar{c}[n](\Lambda)}$$

where [n] is the quantum state of the $c\bar{c}$ pair and Λ is the energy scale¹².

■ CEM: The non-relativistic part is assumed to be non-zero and constant between $4m_c^2$ and $4m_D^2$ and zero for all other energies, where m_c is the mass of the charm quark and m_D is the mass of D meson.

$$d\sigma(J/\psi + X) = \frac{F_{c\bar{c}[J/\psi]}}{9} \Sigma_n \int_{2mc_c}^{2m_D} dM \frac{d\sigma_{c\bar{c}[n]+X}}{dM}$$



 $^{^{12}\}mathrm{T.}$ Kempel, PhD thesis, Iowa State U., 2011, arXiv: 1107.1293 (nucl-ex).

J/ψ Production

■ CSM: In this model, the $c\bar{c}$ pair emerging from the relativistic scattering diagram is assumed to be in the same quantum state as the produced J/ψ , and the non-relativistic amplitude is the real-space J/ψ wave function evaluated at the origin;

$$d\sigma(J/\psi + X) = \int_0^\infty dM \frac{d\sigma_{c\bar{c}}[^3S_1] + X}{dM} \psi_{J/\psi}(r=0)$$

■ COM: This model attempts to formalize the factorization of relativistic and non-relativistic effects. The model use a generic expansion;

$$d\sigma(J/\psi + X) = \sum_{n} \int_{0}^{\infty} dM \frac{d\sigma_{c\bar{c}}[^{3}S_{1}] + X}{dM} \left\langle \mathcal{O}_{[n]}^{J/\psi} \right\rangle$$

with parameters $\langle \mathcal{O}_{[n]}^{J/\psi} \rangle$, non-relativistic matrix elements associated with the amplitude for producing a J/ψ from a $c\bar{c}$ pair in state [n]. Technique of non-relativistic QCD is apply to calculate the $\langle \mathcal{O}_{[n]}^{J/\psi} \rangle$ parameters in power of v, relative velocity between c and \bar{c} . The model is thus a double expansion, about v^2 and α_S .

Gaussian Process Regression (GPR)

■ Probability density function (PDF) of a multivariate normal distribution (MVN) with dimension D is;

$$\mathcal{N}(x|\mu, \Sigma) = \frac{1}{(2\pi)^{D/2} |\Sigma^{1/2}|} \exp\left[-\frac{1}{2} (x - \mu)^T \sigma^{-1} (x - \mu)\right]$$

where D is the number of dimensions, x is the variable, μ is the mean vector and Σ is the covariance matrix.

■ Gaussian processes are distributions over functions f(x) of which the distribution is defined by a mean function m(x) and positive definite covariance function k(x, x'), with x the function values and x, x' all possible pairs in the input domain;

$$f(x) \sim \mathcal{GP}(m(x), k(x, x'))$$

where for any finite subset $X = x_1,, x_n$ of the domain of x, the marginal distribution is a multivariate Gaussian distribution;

$$f(X) \sim \mathcal{N}(m(X), k(X, X))$$

TSSA

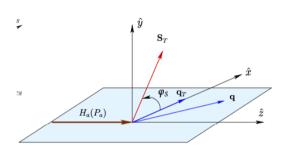


Figure 17: ϕ_S definition in the target rest frame. ¹³

$$\sigma(\phi_S) \propto 1 + PA_N \sin(\phi_S + \phi)$$

$$A(\phi_S) = \frac{\sigma^{\uparrow} - \sigma^{\downarrow}}{\sigma^{\uparrow} + \sigma^{\downarrow}} = A_N \sin \phi_S$$

where P is the target polarization, ϕ_S is the angle between $q_T \& S_T$, ϕ is the spin alignment of the target.

We can extract the A_N using $\sin \phi_S$ modulations.

 $^{^{\}textstyle 13}\mathrm{R.}$ Longo, EPJ Web Conf. $^{\textstyle 137},$ ed. by Y. Foka et al., 05013 (2017).