An Investigation of Charm Quark Jet Spectrum and Shape Modifications in Au+Au Collisions at $\sqrt{s_{\mathrm{NN}}} = 200$ GeV*

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Partons in heavy-ion collisions interact strongly with the Quark-Gluon Plasma (QGP), and hence have their energy and shower structure modified compared to those in vacuum. Theoretical calculations predict that the radiative energy loss, which is the dominant mode of energy loss for gluons and light quarks in the QGP, is suppressed for heavy quarks at low transverse momenta (p_{T}). The excellent secondary vertex resolution provided by the Heavy Flavor Tracker in the STAR experiment at RHIC enables the reconstruction of $D^0(\bar{D}^0)$ mesons at low p_{T} with high signal significance over the background. In this proceeding, we report the first measurements of the $D^0(\bar{D}^0)$ meson tagged jet p_{T} spectra and the $D^0(\bar{D}^0)$ meson radial profile in jets reconstructed from Au+Au collisions at $\sqrt{s_{\mathrm{NN}}} = 200$ GeV, collected by the STAR experiment.

1. Introduction

Relativistic heavy-ion collisions produce Quark-Gluon Plasma (QGP), as predicted by Quantum Chromodynamics (QCD) [1]. Internal probes involving hard scattering processes are used to study the properties of the QGP medium. Jets, one of such probes, manifests as a collimated cluster of final state particles in the detector. The partons which give rise to these jets lose energy to the QGP medium, either through collisions, or through induced gluon *bremsstrahlung* - a phenomenon known as jet quenching [2]. The effects of jet quenching can be seen in measurements of inclusive jets yield suppression [3], and modifications to the jet structure [4]. A study of heavy flavor tagged jets can shed light on the mass and flavor dependence of the parton energy loss and jet structure modifications. The dead-cone effect, as predicted by the QCD, was measured for charm quarks in pp collisions at

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the LHC [5], but remains elusive for heavy-ion collisions. Heavy flavor jets at the LHC have also yet to reveal significant differences with their inclusive counterparts [6, 7], possibly due to having energies much higher than the parton masses. Such studies at the RHIC energies, where lower energy jets are produced, could be the key to better understanding the parton mass dependence. This proceeding will focus on the first measurements of jet transverse momentum p_{T} spectra and the $D^0(\bar{D}^0)$ meson radial profile in tagged jets from Au+Au collisions at $\sqrt{s_{\mathrm{NN}}} = 200$ GeV.

2. Analysis Setup

This work uses Minimum Bias(MB) triggered Au+Au collision events at $\sqrt{s_{\mathrm{NN}}} = 200$ GeV, collected in 2014 by the STAR detector [8] at RHIC. The event selections for this analysis follow the ones used in Ref. [9]. Tracks which pass standard quality cuts at STAR [10], are chosen within the pseudorapidity acceptance of $|\eta| < 1$. The analysis is done in three centrality bins: 0-10 % (central), 10-40 %(mid-central), and 40-80 %(peripheral). $D^0(\bar{D}^0)$ mesons are reconstructed via the decay channel $D^0 \to K^- + \pi^+$ (and its charge conjugate) with a branching ratio of 3.89 % [11]. Several topological selections based on the decay geometry of $D^0(\bar{D}^0)$ are applied to the combinatorial $K\pi$ pairs in an event by using the Heavy Flavor Tracker (HFT), which improves the resolution of tracking from 1 mm at Time Projection Chamber (TPC) to about 30 μm. A more thorough discussion on the selection criteria for the $D^0(\bar{D}^0)$ candidates is available in Ref. [9].

Jets are reconstructed from tracks and towers with $p_{\mathrm{T}} > 0.2$ GeV/c, and transverse energy $E_{\mathrm{T}} > 0.2$ GeV respectively. The jets are defined using the anti-k_{T} clustering algorithm available in the FastJet package [12], with a radius parameter of $R = 0.4$ in the $\eta - \phi$ space. The K and π daughter tracks are replaced with the corresponding $D^0(\bar{D}^0)$ before the jets are reconstructed. A jet area based background subtraction is applied to minimize the effect of the soft background on the jets [13]. Jets with a $D^0(\bar{D}^0)$ constituent with $p_{\mathrm{T,D^0}} \in (5, 10)$ GeV/c are considered for this analysis. In this work, a D^0 tagged jet is a jet with a $D^0(\bar{D}^0)$ candidate as a constituent.

3. $D^0(\bar{D}^0)$ Jet Spectrum and Shape Modifications

To extract the raw yield of $D^0(\bar{D}^0)$ mesons, a method called ${}_s\mathcal{P}lot$ [14] is used. ${}_s\mathcal{P}lot$ calculates per event weights, called sWeights, from an unbinned likelihood fit to the $D^0(\bar{D}^0)$ invariant mass distribution over all kinematics. The weights classify how '*signal-like*' a $D^0(\bar{D}^0)$ candidate is. Figure 1 shows the invariant mass distribution of $K\pi$ candidates in the p_{T} region of 5–10

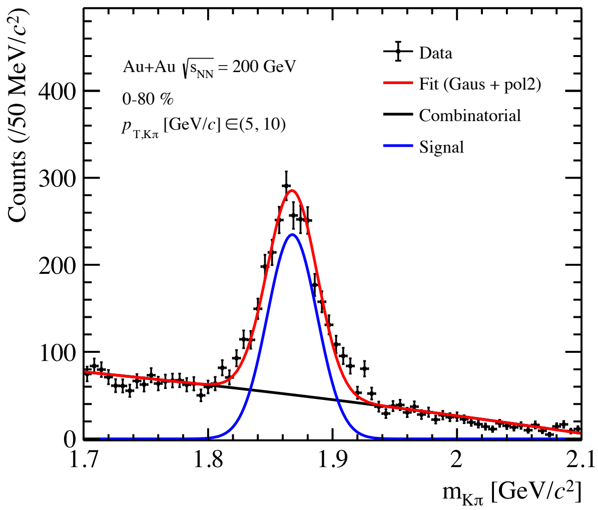

Fig. 1. The invariant mass distribution of $K\pi$ pairs with $p_{\mathrm{T}} \in (5, 10)$ GeV/c. The unlike sign $K\pi$ pairs distribution (*black*) is fit with a Gaussian plus second-order polynomial (*red*) to estimate the $D^0(\bar{D}^0)$ meson yield. The signal after the removal of the background (*blue*) is also shown on the same scale.

GeV/c for 0–80% MB events. The invariant yields of $D^0(\bar{D}^0)$ tagged jets is represented by the formula:

$$\frac{d^2 N_{\mathrm{jet}}}{2\pi\, N_{\mathrm{evt}}\, p_{\mathrm{T,jet}}\, dp_{\mathrm{T,jet}}\, d\eta} = \frac{1}{\mathrm{B.R.}} \times \frac{N_{\mathrm{jet}}^{\mathrm{raw}}}{2\pi\, N_{\mathrm{evt}}\, p_{\mathrm{T,jet}}\, \Delta p_{\mathrm{T,jet}}\, \Delta\eta} \times \frac{1}{\epsilon_{\mathrm{corr}}} \tag{1}$$

where B.R. is the $D^0 \to K^-\pi^+$ decay branching ratio (3.89±0.04%), $N_{\mathrm{jet}}^{\mathrm{raw}}$ is the reconstructed $D^0(\bar{D}^0)$ tagged jets raw counts, N_{evt} is the total numbers of events used in this analysis. The raw yields are corrected for the tracking efficiencies and acceptances of the TPC and HFT, topological cut efficiency, particle identification efficiency, and finite vertex resolution based on the correction factors derived in the STAR study on $D^0(\bar{D}^0)$ production in heavy-ion collisions [9], and the total correction factor is ϵ_{corr}. The nuclear modification factor R_{CP} is defined as the ratio of N_{coll}–normalized yields between central and peripheral collisions where N_{coll} is the number of the binary collisions for a centrality class.

The radial distribution of $D^0(\bar{D}^0)$ mesons in tagged jets is defined by the formula:

$$\frac{1}{N_{\mathrm{jet}}}\frac{dN_{\mathrm{jet}}}{d\mathrm{r}} = \frac{1}{N_{\mathrm{jet}}}\frac{N_{\mathrm{jet}}|_{\Delta\mathrm{r}}}{\Delta\mathrm{r}} \tag{2}$$

where r $= \sqrt{(\eta_{\rm jet} - \eta_{\rm D^0})^2 + (\eta_{\rm jet} - \phi_{\rm D^0})^2}$ is the distance of the $D^0(\eta_{D^0}, \phi_{D^0})$ from the jet axis $(\eta_{\rm jet}, \phi_{\rm jet})$ in the $\eta - \phi$ plane, and $N_{\rm jet}|_{\Delta \rm r}$ is the number of jets with $D^0(\bar{D}^0)$ mesons in the Δr interval.

A Bayesian unfolding procedure [15], with a Monte-Carlo (MC) generated event sample, is used to account for the detector inefficiencies in jet reconstruction. A $D^0(\bar{D}^0)$-enriched sample of pp collision events at $\sqrt{s}$ = 200 GeV is generated using PYTHIA v8.303, with the 'Detroit' tune [16], and propagated through the STAR detector using the GEANT3 package. The FONLL (Fixed Order + Next-to-Leading Logarithms) charm quark spectrum is used as a prior for the unfolding procedure. The charm jet fragmentation function is modeled using PYTHIA, and a systematic study of effects of its variation is in the works. Observables with an asterisk(*), found later in this proceeding, denote this underlying assumption.

The fluctuation due to the heavy-ion background is estimated by embedding one 'single-particle' jet each in MB Au+Au events, and then matching each embedded jet with a reconstructed jet containing the tagged 'single-particle'. The quantity $\Delta p_{\rm T,SPjet} = p^{\rm det}_{\rm T,SPjet} - p^{\rm part}_{\rm T,SPjet}$ models this fluctuation. The superscript 'part' refers to particle-level jets, and 'det' refers to detector-level jets. For the D^0 meson radial profile, a two-dimensional Bayesian unfolding procedure is used to simultaneously correct $N_{\rm jet}$ as a function of $p_{\rm T,jet}$ and Δr.

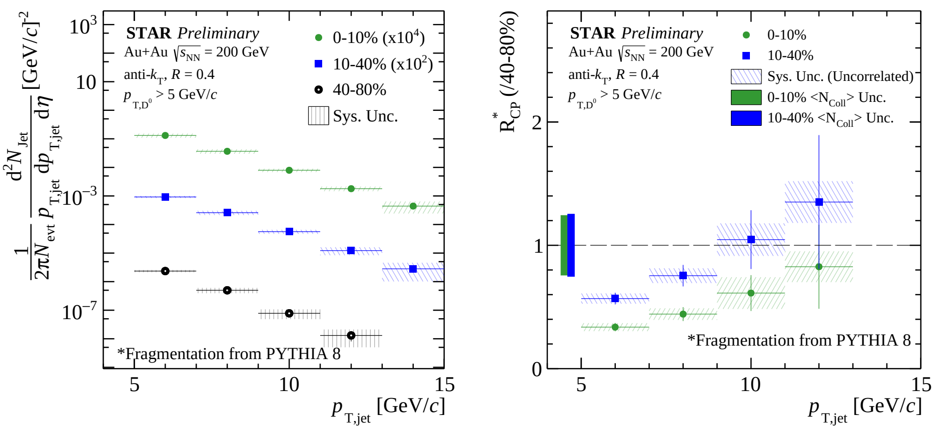

Fig. 2. **Left**: Jet p_T spectra for $D^0(\bar{D}^0)$ jets with $p_{\rm T,D^0} \in (5, 10)$ GeV/c in different centrality classes; **Right**: Nuclear modification factor $R_{\rm CP}*$ for D^0 jets.

The systematic uncertainties across the bins in the reported observables are dominated by the following contributions: a) differences in the invariant yield of D^0 mesons calculated using the $_s\mathcal{P}lot$ method, and a like-sign background subtraction method, and b) systematic uncertainty in $D^0(\bar{D}^0)$

reconstruction taken from Ref. [9]. Systematic variations related to the unfolding procedure are estimated by varying the following: a) the prior from FONLL to the jet distribution for D^0 tagged jets in pp collisions at $\sqrt{s} = 200$ GeV generated by PYTHIA, and b) the regularisation parameter.

The efficiency-corrected invariant yield of $D^0(\bar{D}^0)$ meson tagged jets with $p_{\mathrm{T,D^0}} > 5$ GeV/c is shown in the left panel of Figure 2, as a function of $p_{\mathrm{T,jet}}$ in 0-10%, 10-40%, and 40-80% Au+Au collisions. The spectra in the first two centrality bins are scaled by arbitrary factors for better visibility. The ratios of binary collisions (N_{Coll}) scaled yields in different centrality classes, defined as the nuclear modification factors R^*_{CP}, are shown for the central and the mid-central Au+Au collisions in the right panel of Figure 2, with the peripheral centrality bin as the reference. The bands (blue and green) are uncertainties associated with N_{Coll}. The D^0 jet R_{CP} shows a strong suppression in central collisions than in mid-central collisions at low $p_{\mathrm{T,jet}}$. R^*_{CP} also shows an increasing trend with $p_{\mathrm{T,jet}}$ for both centrality bins which is qualitatively different from the R_{CP} for inclusive jets [3].

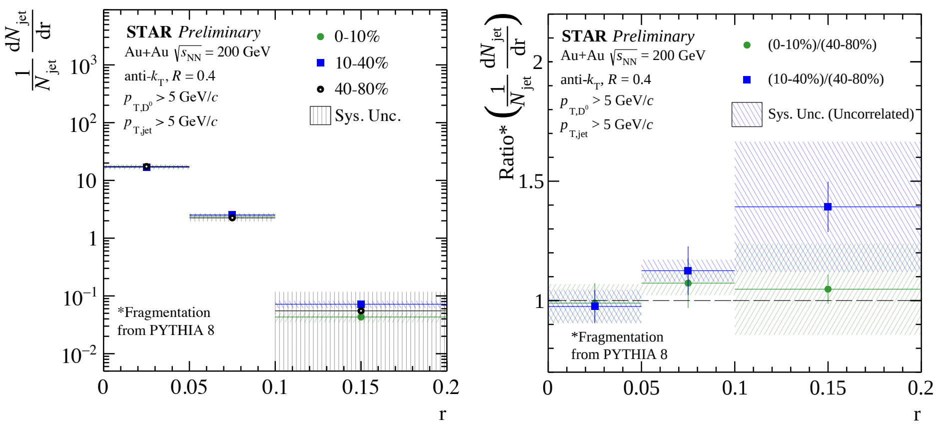

Fig. 3. **Left**: D^0 radial profile for $D^0(\bar{D}^0)$ jets with $p_{\mathrm{T,D^0}} \in (5, 10)$ GeV/c in different centrality classes; **Right**: Ratio of D^0 radial profiles for central and mid-central events with respect to D^0 radial profile for peripheral events.

The radial profile for $D^0(\bar{D}^0)$ mesons with $p_{\mathrm{T,D^0}} \in (5, 10)$ GeV/c in the tagged jets is shown as a function of the distance from the jet axis (r) in 0-10%, 10-40%, and 40-80% Au+Au collisions in the left panel of Figure 3. The ratio of the radial profiles for the central and mid-central events with peripheral events, shown in the right panel of Figure 3, is found to be consistent with unity within the uncertainties. The large uncertainties are dominated by the limited statistics in the peripheral centrality bin.

4. Discussion

The first measurements of D^0 meson tagged jet p_{T} spectra and D^0 meson radial profile is reported for $p_{\mathrm{T,D^0}} \in (5, 10$ GeV/c in Au+Au collisions at $\sqrt{s_{\mathrm{NN}}} = 200$ GeV. The D^0 $p_{\mathrm{T,jet}}$ spectra is found to be suppressed for central and mid-central collisions with the nuclear modification factor showing an increasing trend with $p_{\mathrm{T,jet}}$. This trend is qualitatively different from the inclusive jet measurements at RHIC. The radial profile of D^0 meson in its tagged jets is found to be the same at different centralities. Within the current uncertainties, no hint of differences in charm quark diffusion is observed in the presence of the QGP medium. Further studies are ongoing to extend our measurements to lower $p_{\mathrm{T,D^0}}$, accessible at STAR, to get closer to the charm quark mass. These measurements can constrain theoretical models on parton flavor and mass dependencies of jet energy loss.

REFERENCES

[1] STAR Collaboration. *Nuclear Physics A*, 757(1-2):102–183, Aug 2005.

[2] Megan Connors *et al. Rev. Mod. Phys.*, 90:025005, Jun 2018.

[3] STAR Collaboration. *Phys. Rev. C*, 102:054913, Nov 2020.

[4] CMS Collaboration. *Physics Letters B*, 730:243–263, 2014.

[5] ALICE Collaboration. *Nature*, 605(7910):440–446, May 2022.

[6] CMS Collaboration. *Phys. Rev. Lett.*, 113:132301, Sep 2014.

[7] CMS Collaboration. *Phys. Rev. Lett.*, 125:102001, Sep 2020.

[8] STAR Collaboration. *Nuc. Ins. Methods. A*, 499(2):624–632, 2003.

[9] STAR Collaboration. *Phys. Rev. C*, 99:034908, Mar 2019.

[10] STAR Collaboration. *Phys. Rev. Lett.*, 119:062301, Aug 2017.

[11] Particle Data Group. *Prog. Theor. Exp. Phys*, 2020:083C–84, 2020.

[12] Matteo Cacciari *et al. The Eur. Phys. Jour. C*, 72(3):1896, Mar 2012.

[13] Matteo Cacciari *et al. Physics Letters B*, 659(1):119–126, 2008.

[14] M. Pivk *et al. Nuc. Ins. Methods. A*, 555(1):356–369, 2005.

[15] G. D'Agostini. *Nuc. Ins. Methods. A*, 362(2):487–498, 1995.

[16] Manny Rosales Aguilar and *et al. arXiv*, 2021.