

# Open Heavy Flavor Measurements at RHIC with STAR

Santa Fe Jets and Heavy Flavor Workshop, January 11-13, 2016



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1. University of Illinois at Chicago
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The University of Illinois  
at Chicago

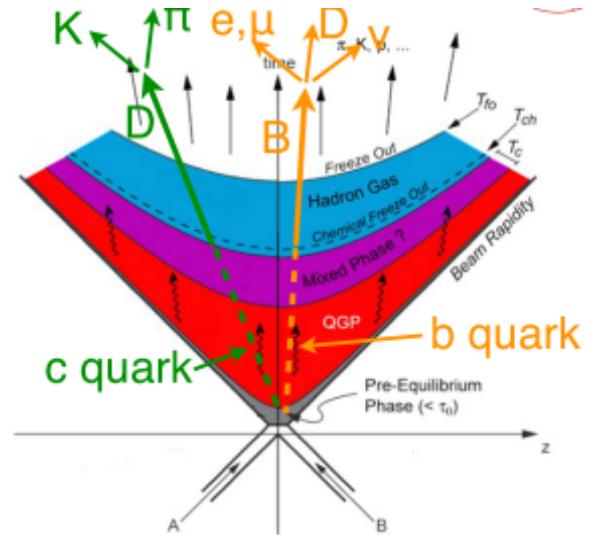


華中師範大學  
Huazhong Normal University

# Heavy Flavor Quarks

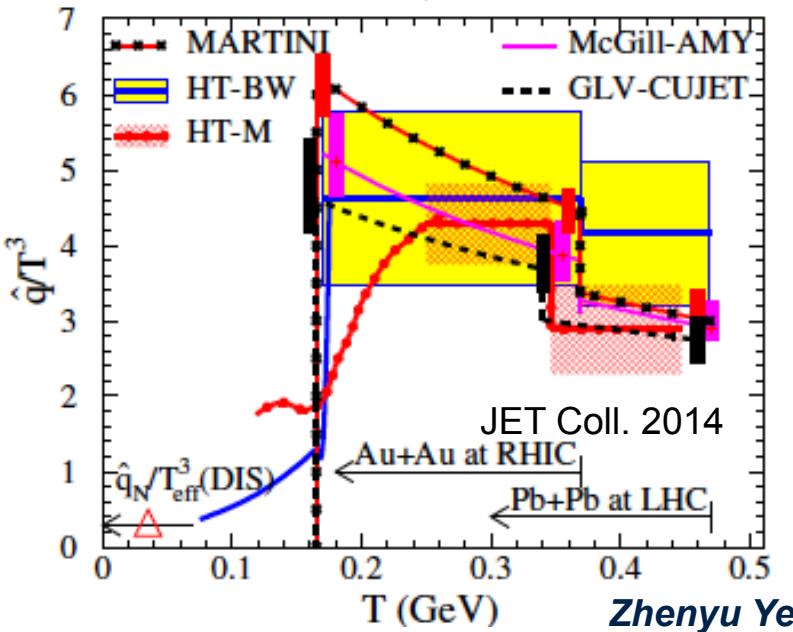
## Heavy quark tomography

- produced mostly from initial hard parton scatterings at RHIC energies; exposed to the whole evolution of the QGP
- total yield or mass not (significantly) altered within the QGP



## Sensitive to parton-medium interactions and medium properties

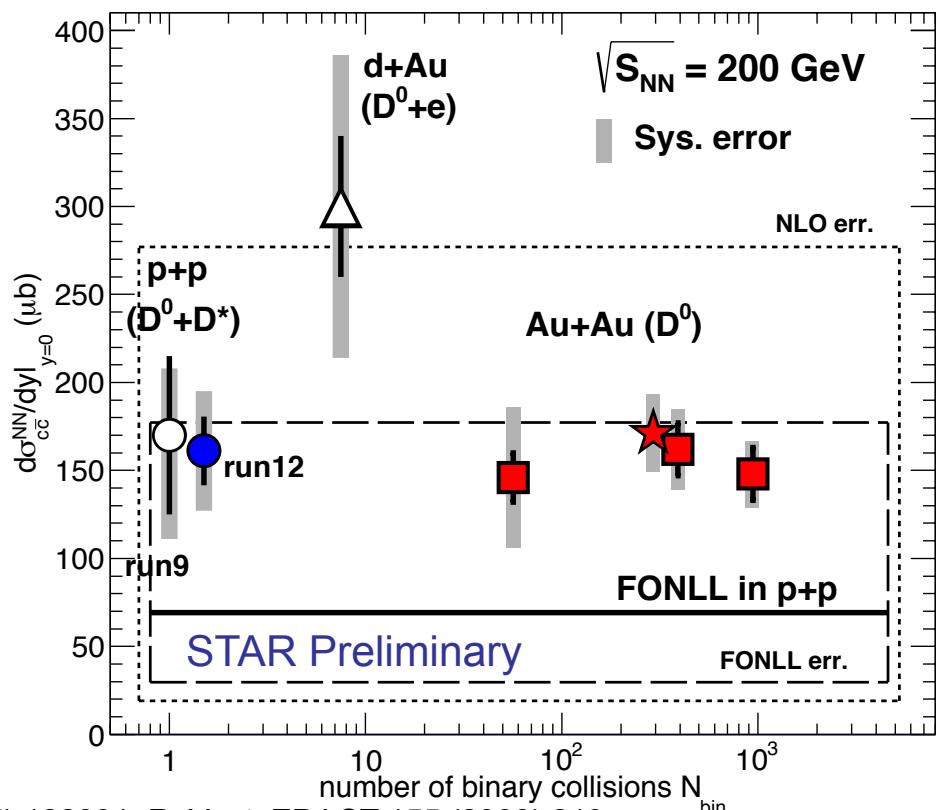
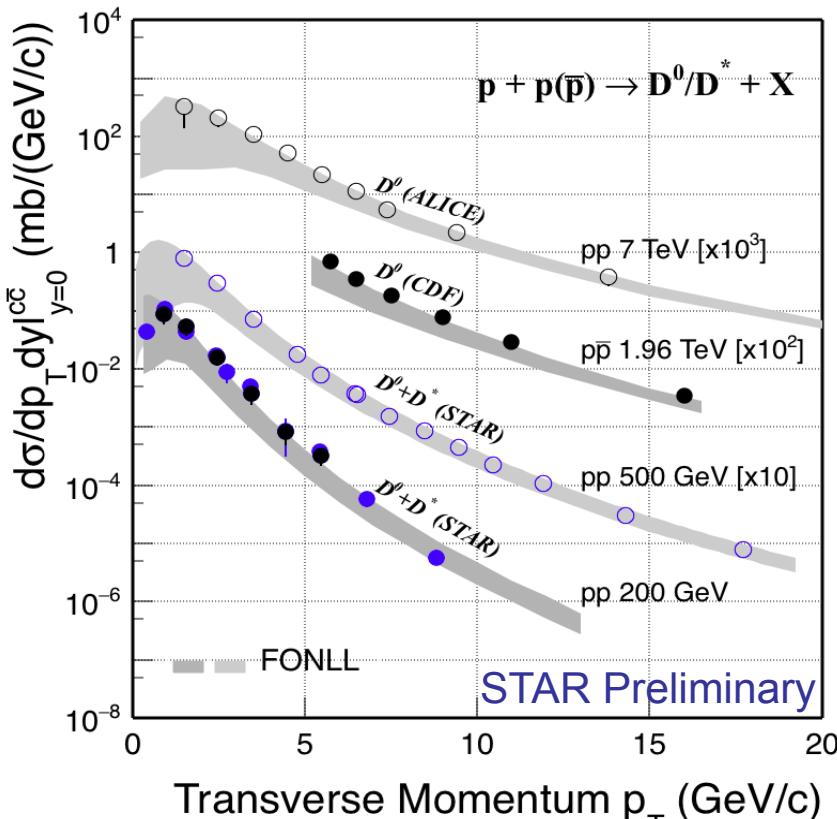
- Comparing light, charm and bottom to disentangle radiative vs collisional energy losses
- Extraction of temperature-dependent parton transport properties needs precise experimental data on heavy flavor production from RHIC



# Open Charm Production at RHIC

Heavy flavor quarks can serve as calibrated probes for the QGP at RHIC:

- production in p+p collisions are described by pQCD calculations
- produced mostly in the initial hard scatterings at RHIC energies



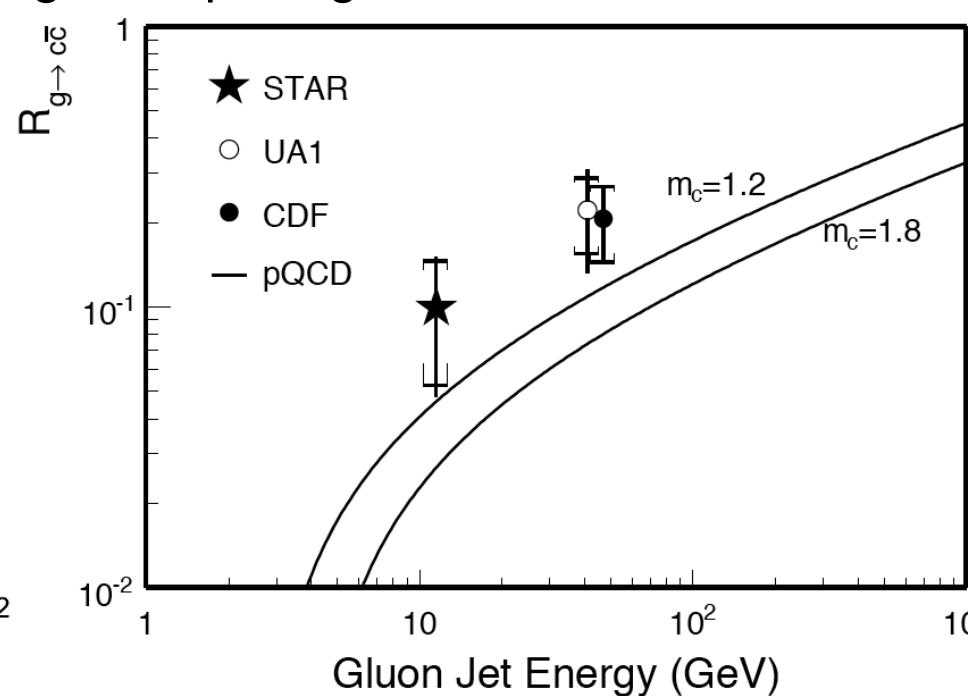
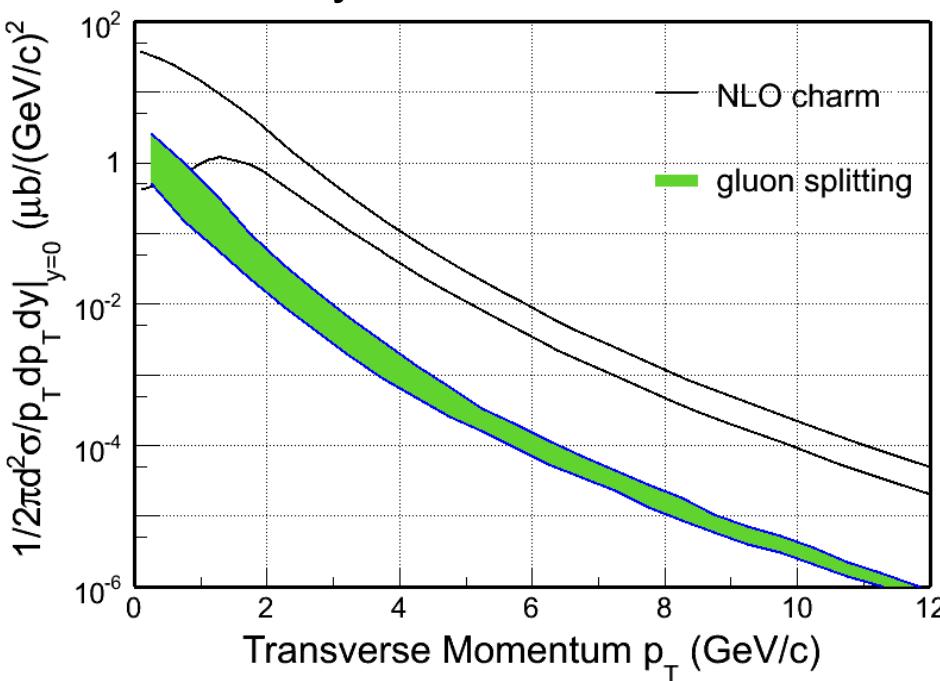
Theory: M. Cacciari et al., PRL 95 (2005) 122001, R. Vogt, EPJ ST 155 (2008) 213

Talk on Bottom production by Zach Miller tomorrow

# Open Charm Production at RHIC

Heavy flavor quarks can serve as calibrated probes for the QGP at RHIC:

- production in p+p collisions are described by pQCD calculations
- produced mostly in the initial hard scatterings at RHIC energies
- has only a small contribution from gluon splitting



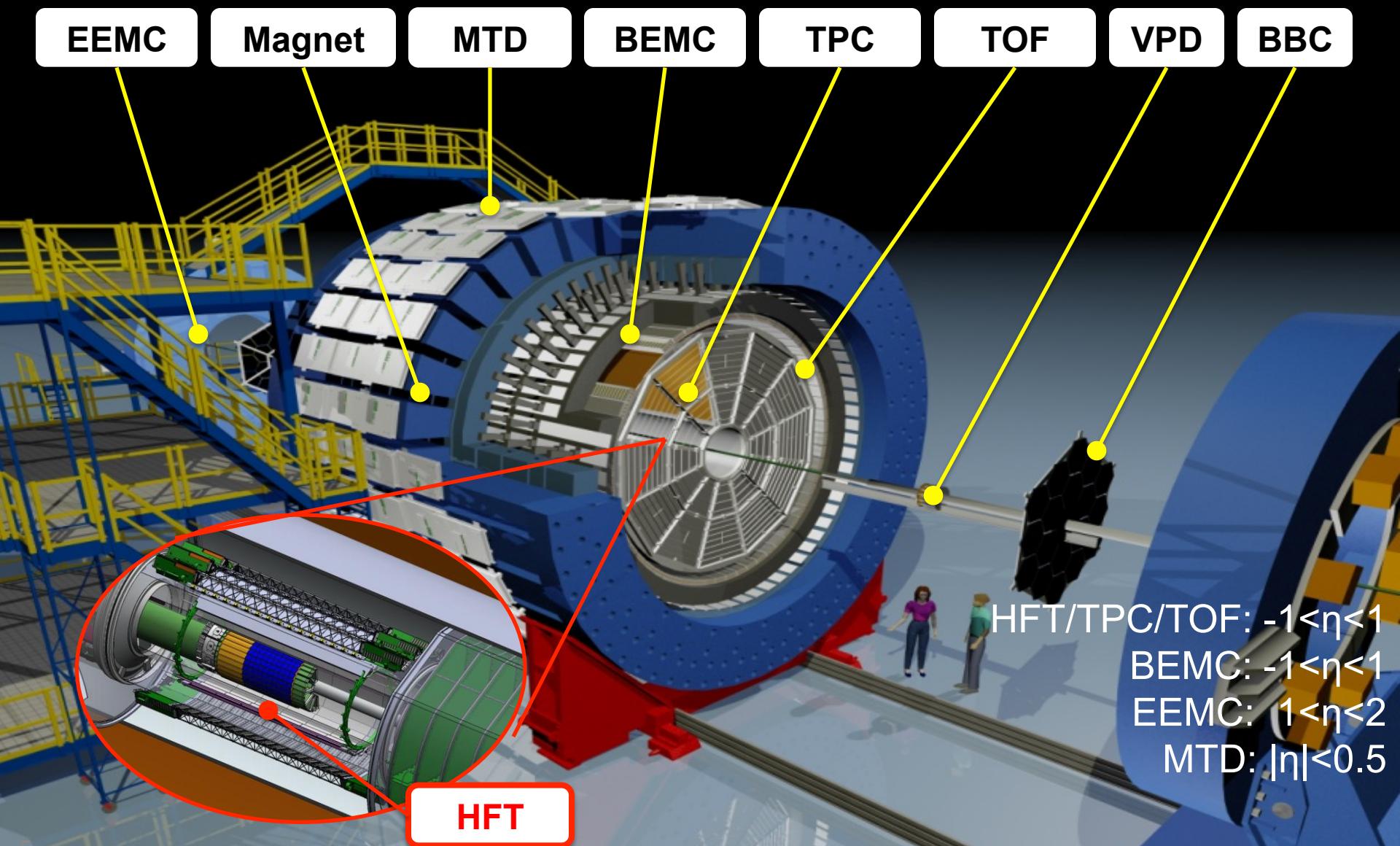
NLO charm: M. Cacciari et al, PRL 95 (2005) 122001

gluon splitting charm: STAR Jet\*pQCD R( $g \rightarrow cc\bar{c}$ )

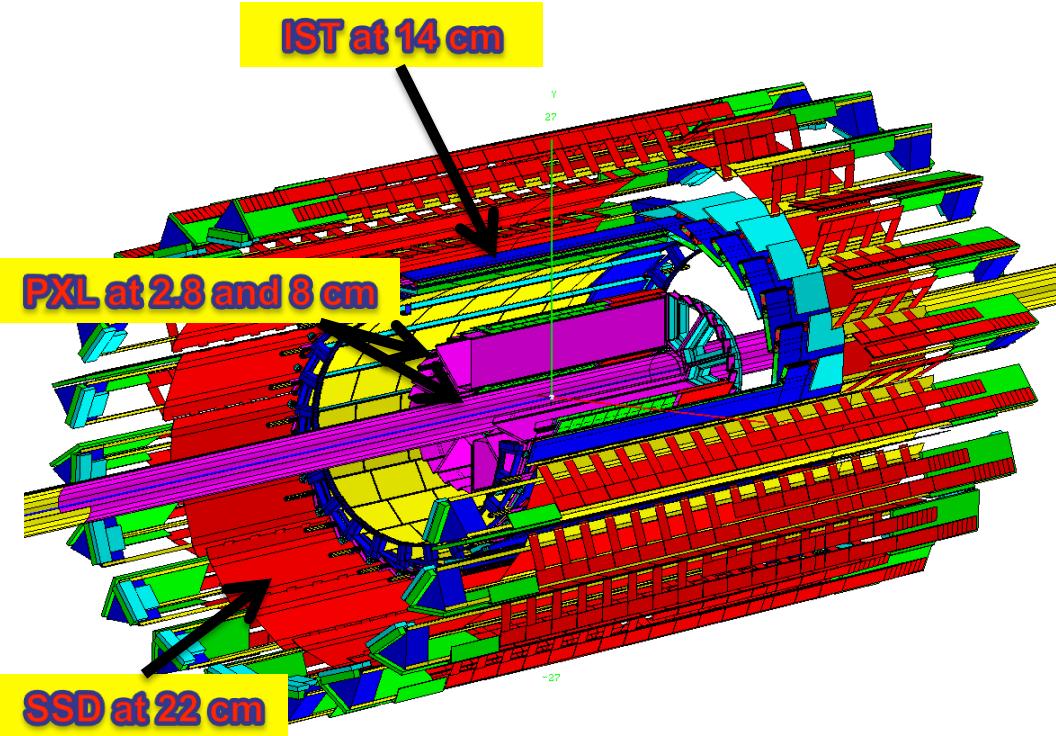
STAR R( $g \rightarrow cc\bar{c}$ ): PRD79 (2009) 112006

pQCD R( $g \rightarrow cc\bar{c}$ ): Mueller & Nason PLB 157 (1985) 226; Mangano & Nason PLB 285 (1992) 160

# STAR Experiment at RHIC



# STAR Heavy Flavor Tracker



## PiXeL detector (PXL)

- two layers of thin Monolithic Active Pixel Sensors with 356M 20.7x20.7  $\mu\text{m}$  pixels
- excellent DCA resolution for HF studies

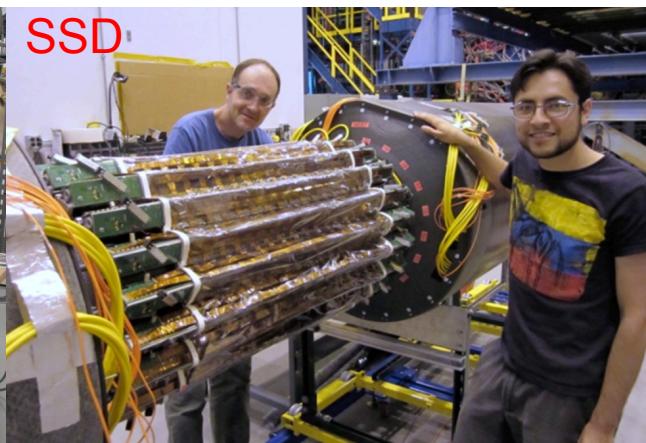
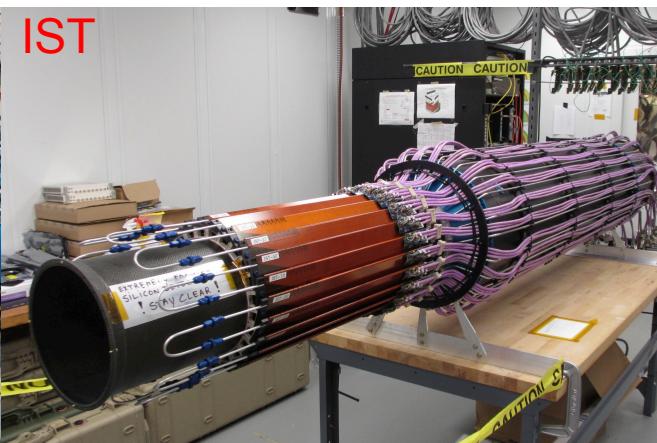
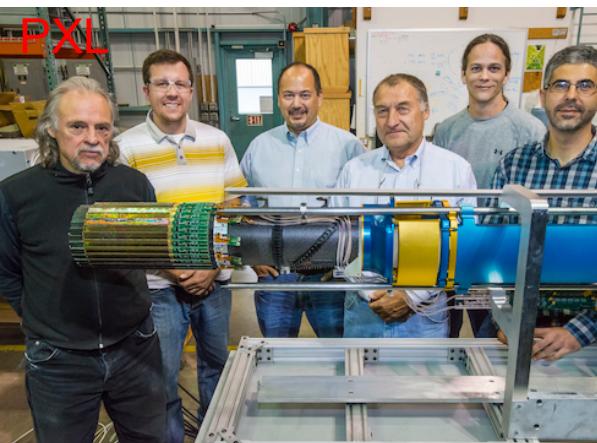
## Intermediate Silicon Tracker (IST)

- one layer of fast readout single-sided double-metal silicon strip detector

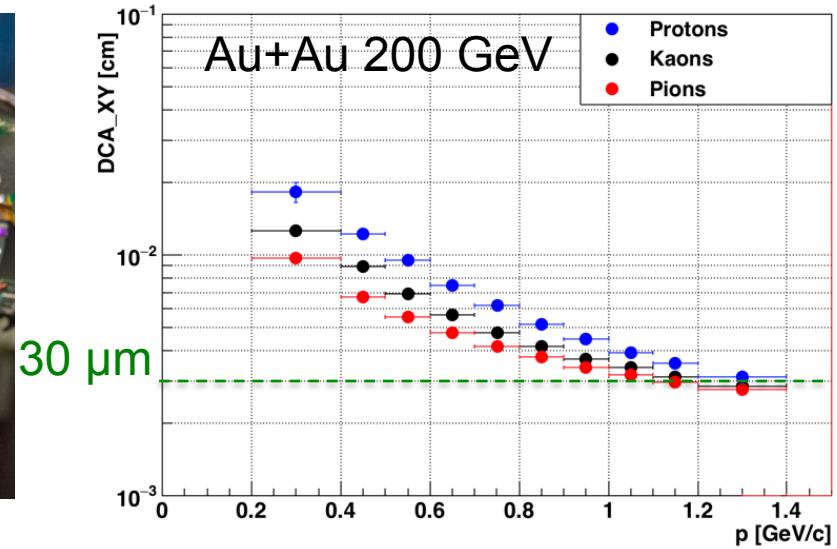
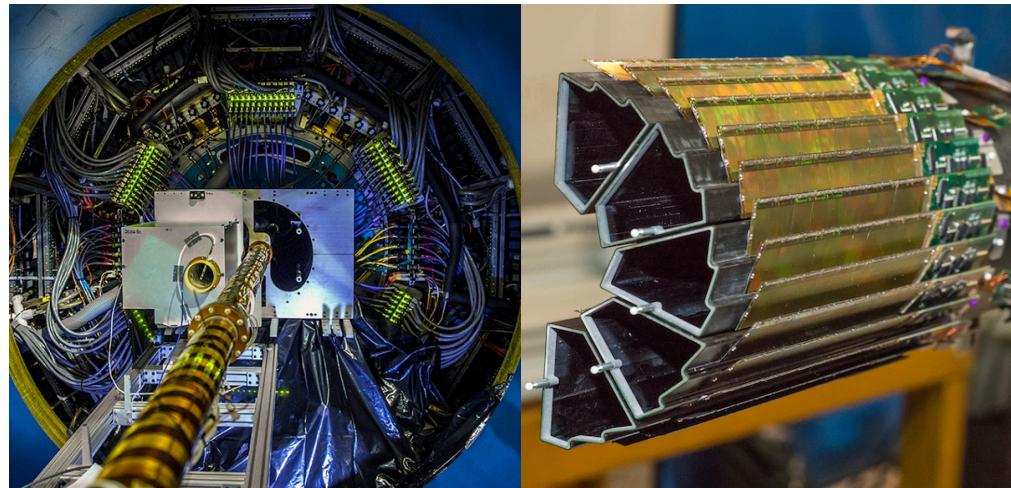
## Silicon Strip Detector (SSD)

- existing one layer of double-sided silicon strip detector with electronic upgrade

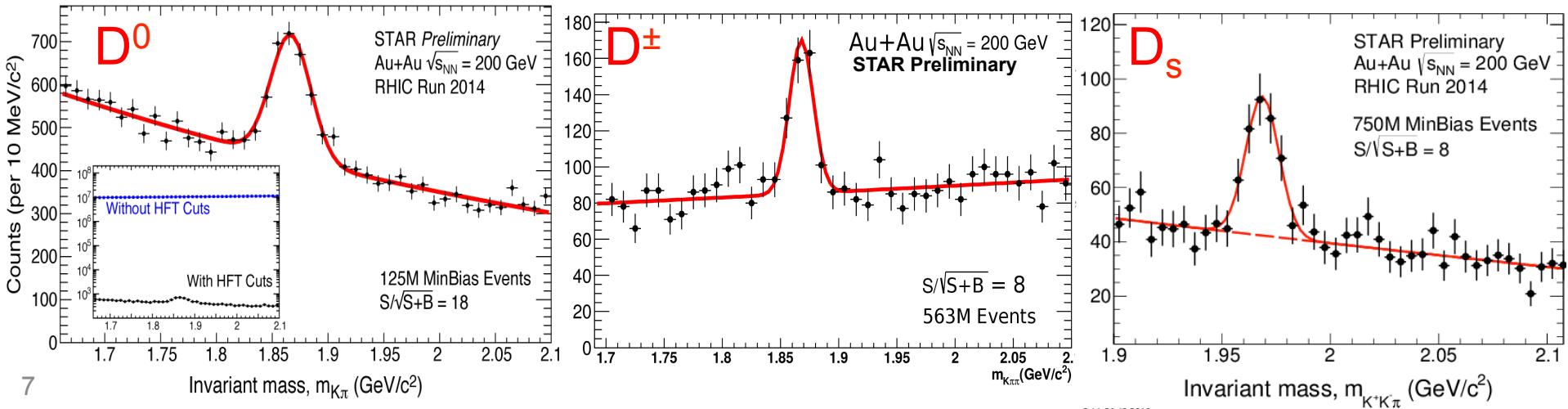
Detector	Radius (cm)	Hit Resolution R/ $\varphi$ - Z ( $\mu\text{m}$ )	Radiation length
SSD	22	20 / 740	1% $X_0$
IST	14	170 / 1800	<1.5 % $X_0$
PXL	2.8/8	6 / 6	$\sim$ 0.4 % $X_0$



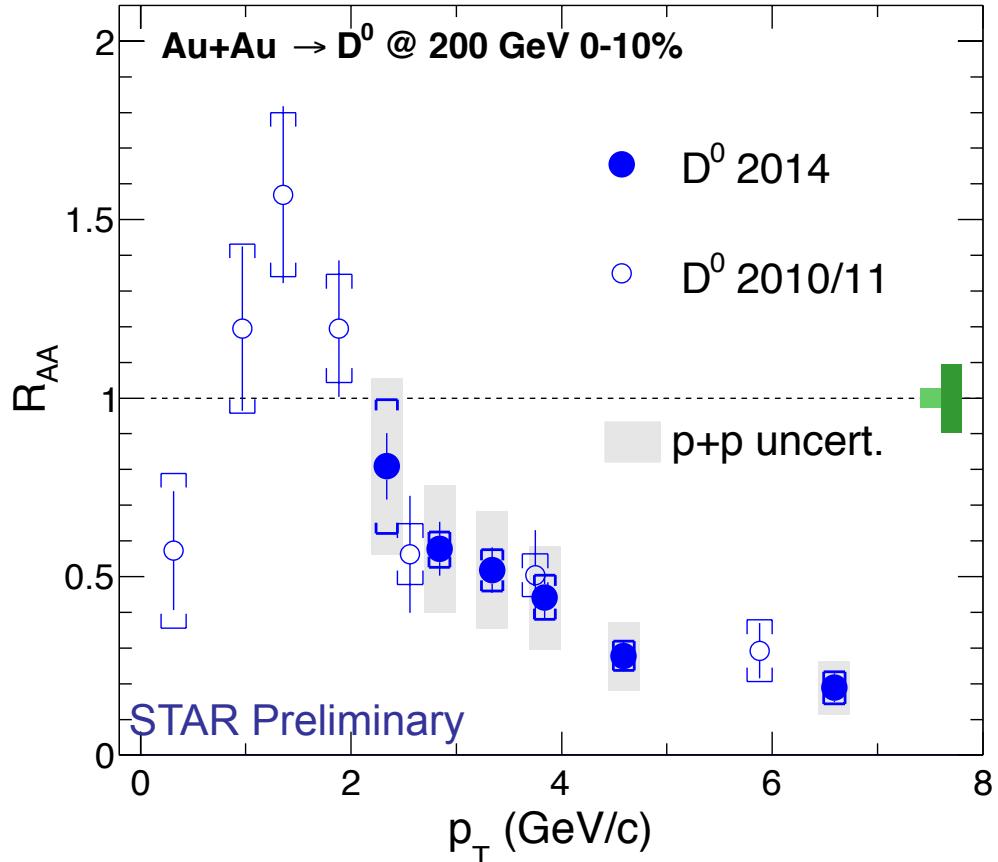
# STAR Heavy Flavor Tracker



A factor of  $\sim 4$  improvement in  $D^0$  significance by the HFT. First results on  $D^\pm$  and  $D_s$ .



# New Results from the HFT – D<sup>0</sup> R<sub>AA</sub>

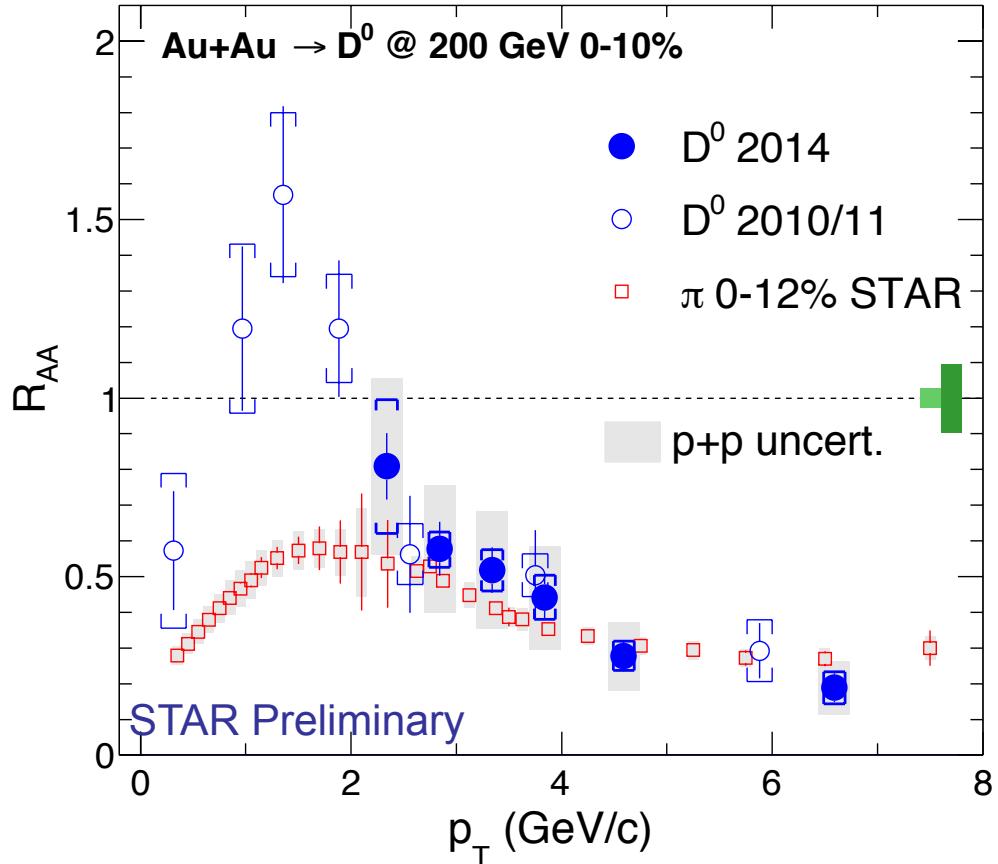


- $R_{AA}(D)>1$  for  $p_T \sim 1.5 \text{ GeV}/c$   
**Charm coalescence with a radially flowing bulk medium**
- High  $p_T$ : significant suppression in central  $Au+Au$  collisions.  
**Strong charm-medium interaction**
- Improved  $Au+Au$  precision at high  $p_T$  thanks to the HFT.  $R_{AA}$  at low  $p_T$  with Run14  $Au+Au$  and Run15  $p+p$  HFT data are underway.

$$R_{AA} = \frac{dN_{AA}/dy}{N_{binary} \cdot dN_{pp}/dy}$$

STAR D<sup>0</sup> 2010/11: PRL 113 (2014) 142301

# New Results from the HFT – D<sup>0</sup> R<sub>AA</sub>

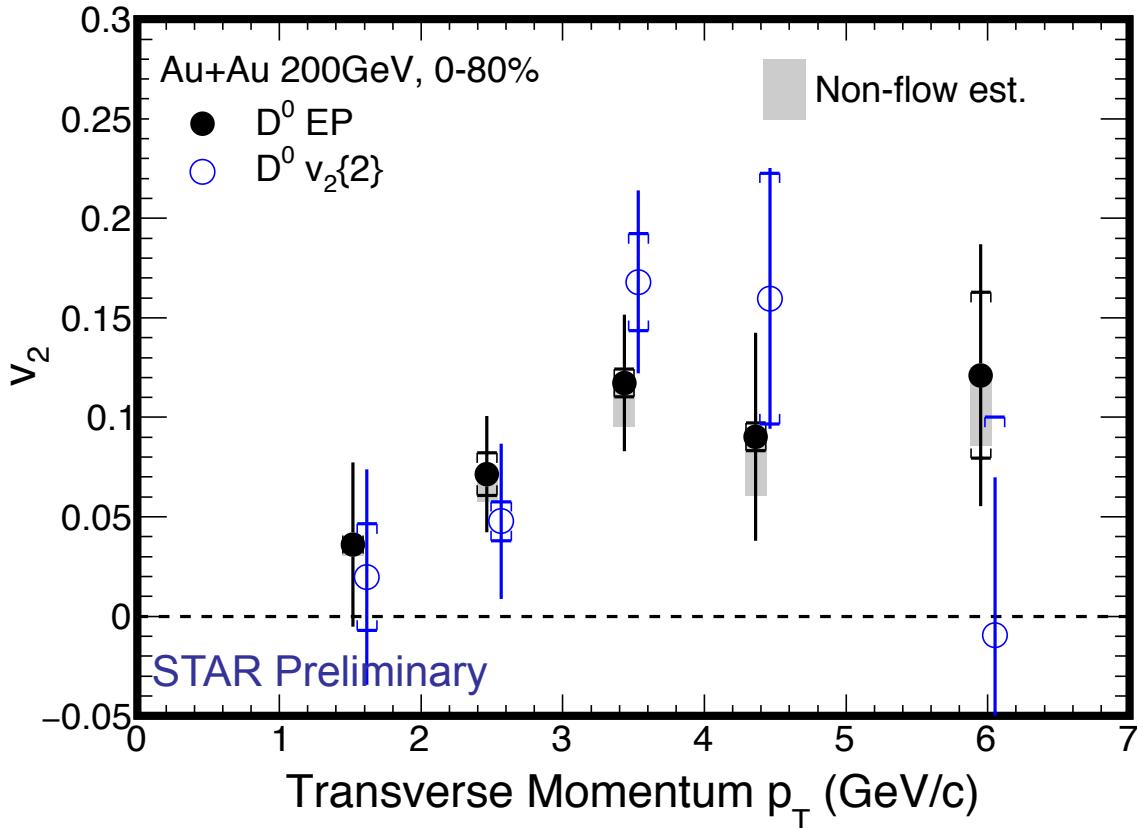


- R<sub>AA</sub>(D)>1 for p<sub>T</sub>~1.5 GeV/c  
**Charm coalescence with a radially flowing bulk medium**
- High p<sub>T</sub>: significant suppression in central Au+Au collisions.  
**Strong charm-medium interaction**
- R<sub>AA</sub>(D) ~ R<sub>AA</sub>(π) at p<sub>T</sub>>4 GeV/c  
**Similar suppression for light partons and charm quarks at high p<sub>T</sub>**

$$R_{AA} = \frac{dN_{AA}/dy}{N_{binary} \cdot dN_{pp}/dy}$$

STAR D<sup>0</sup> 2010/11: PRL 113 (2014) 142301  
STAR π 0-12%: PLB 655 (2007) 104

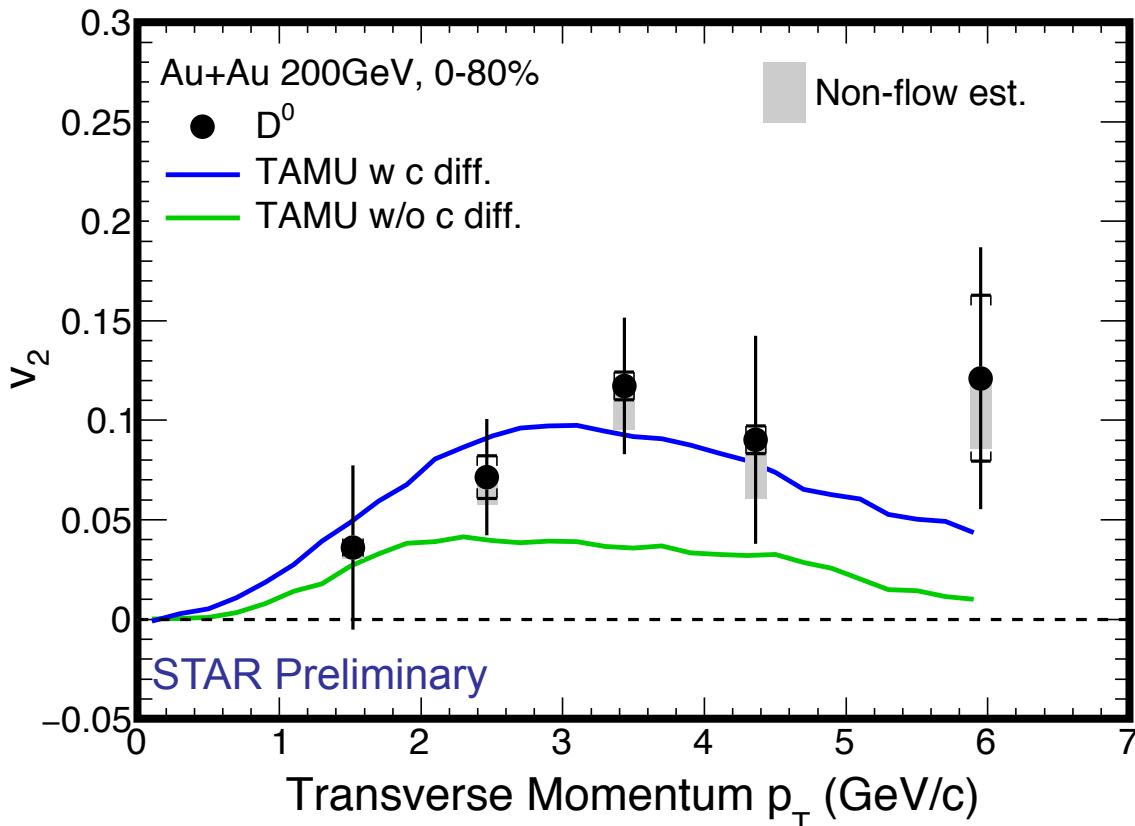
# New Results from the HFT – D<sup>0</sup> v<sub>2</sub>



- Finite  $D^0 v_2$  for  $p_T > 1$  GeV/c

$$\frac{dN}{d\phi} = N_0 \left[ 1 + \sum_n 2v_n \cos n\phi \right]$$

# New Results from the HFT – $D^0 v_2$

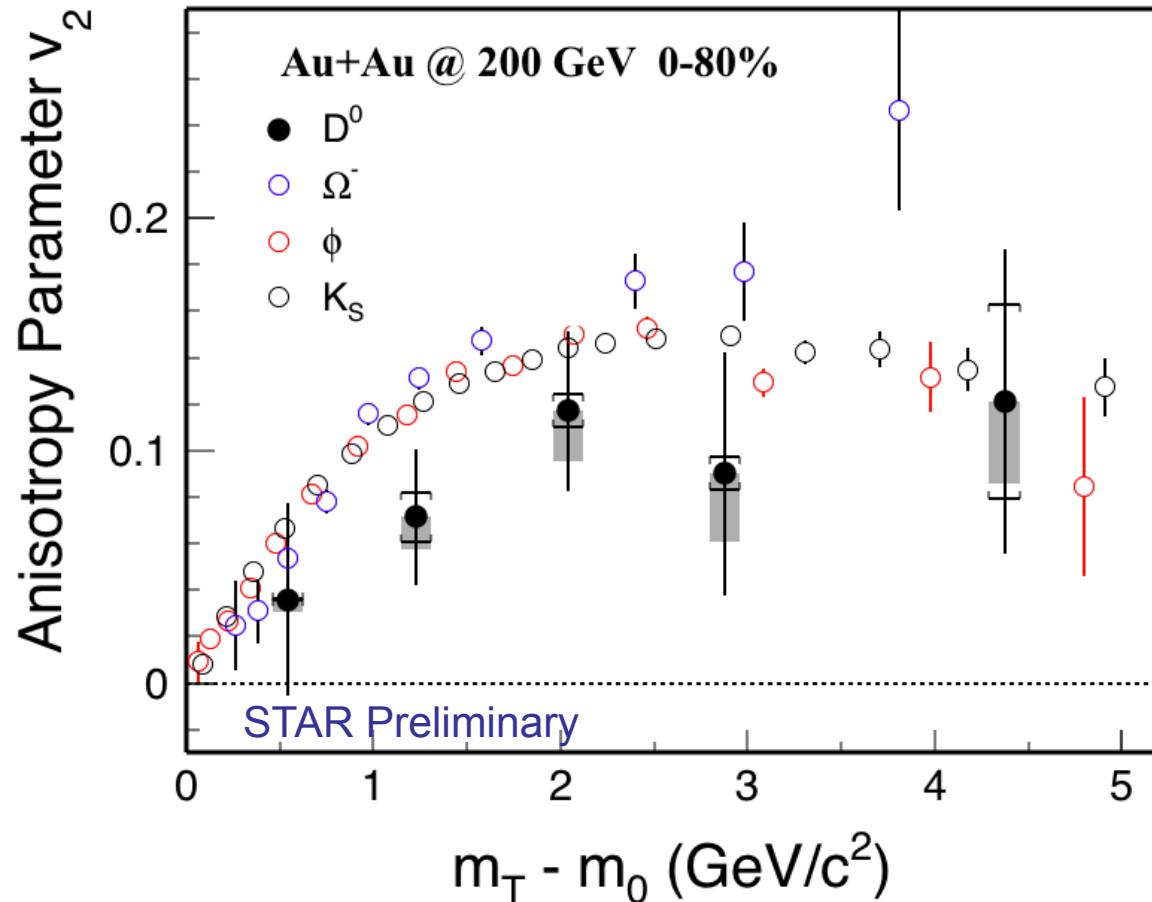


- Finite  $D^0 v_2$  for  $p_T > 1 \text{ GeV}/c$
- Favors charm quark diffusion

$$\frac{dN}{d\phi} = N_0 \left[ 1 + \sum_n 2v_n \cos n\phi \right]$$

Theory curves: latest calculations from private communications  
TAMU: PRC 86 (2012) 014903, PRL 110 (2013) 112301

# New Results from the HFT – D<sup>0</sup> v<sub>2</sub>



- Finite D<sup>0</sup> v<sub>2</sub> for p<sub>T</sub>>1 GeV/c

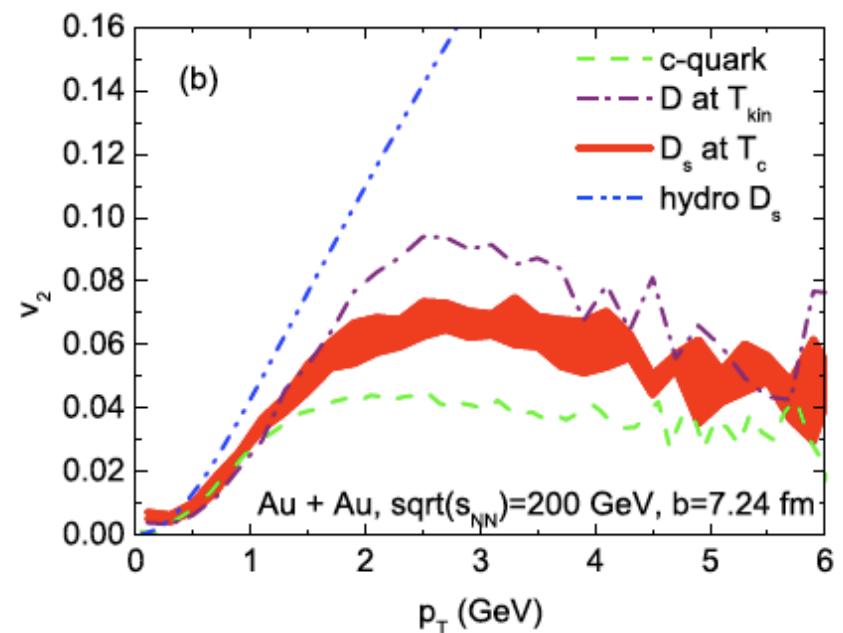
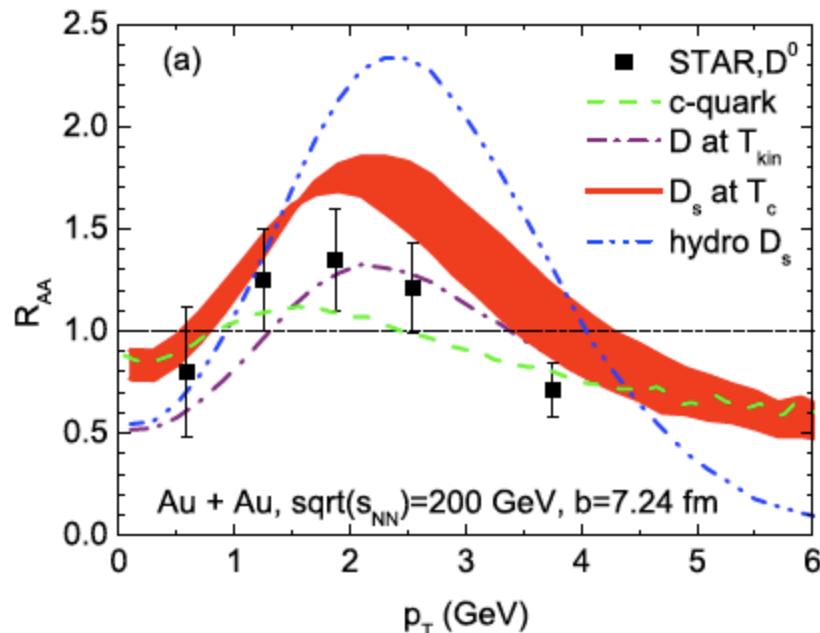
Favors charm quark diffusion

- Lower than light hadron v<sub>2</sub>

Indicates that charm quarks  
are not fully thermalized with  
the medium

$$m_T = \sqrt{p_T^2 + m_0^2}$$

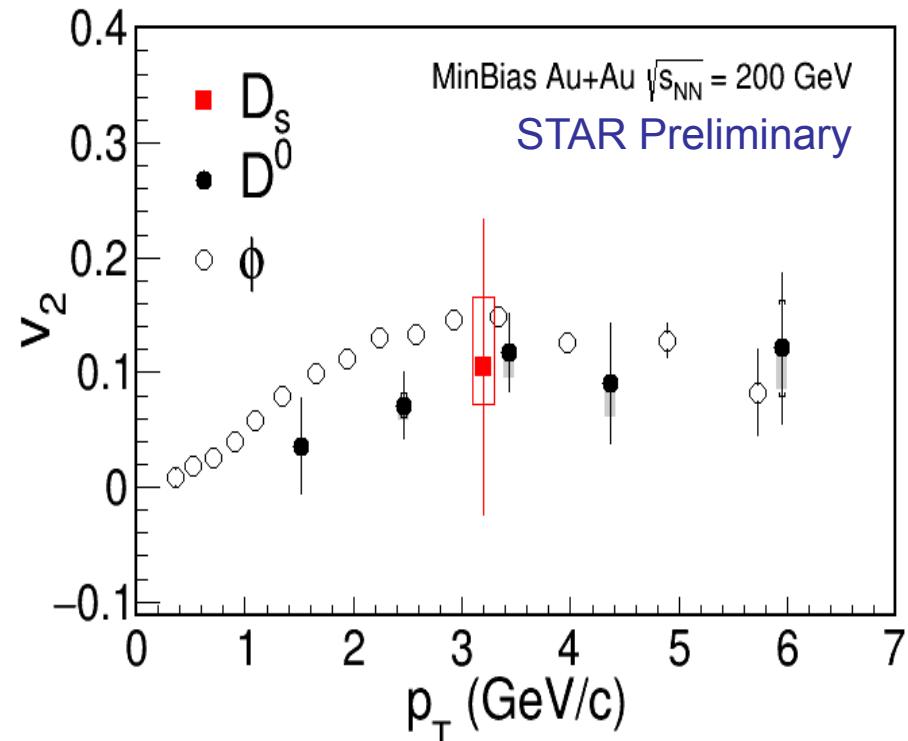
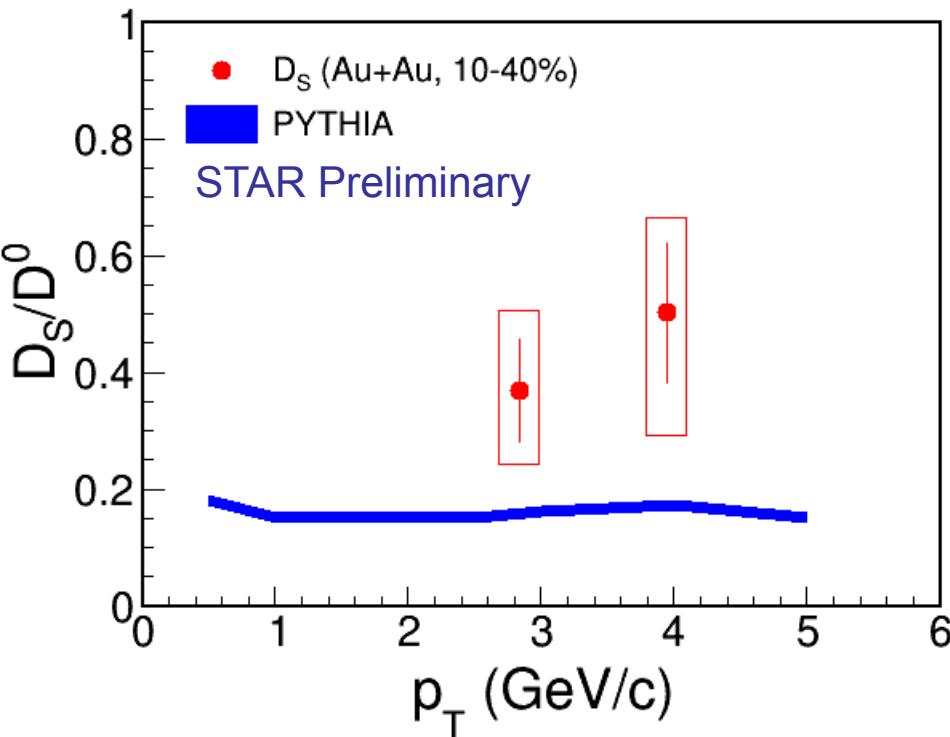
# New Results from the HFT - $D_s$



M. He *et al.*, PRL 110, 112301 (2013)

- Strangeness enhancement in heavy-ion collisions is expected to affect the yield of  $D_s$ : relative increase of  $D_s$  yield than  $D^0$  predicted.
- Elliptic flow of  $D_s < D^0$  is expected due to earlier freeze out of  $D_s$ .

# New Results from the HFT - D<sub>s</sub>



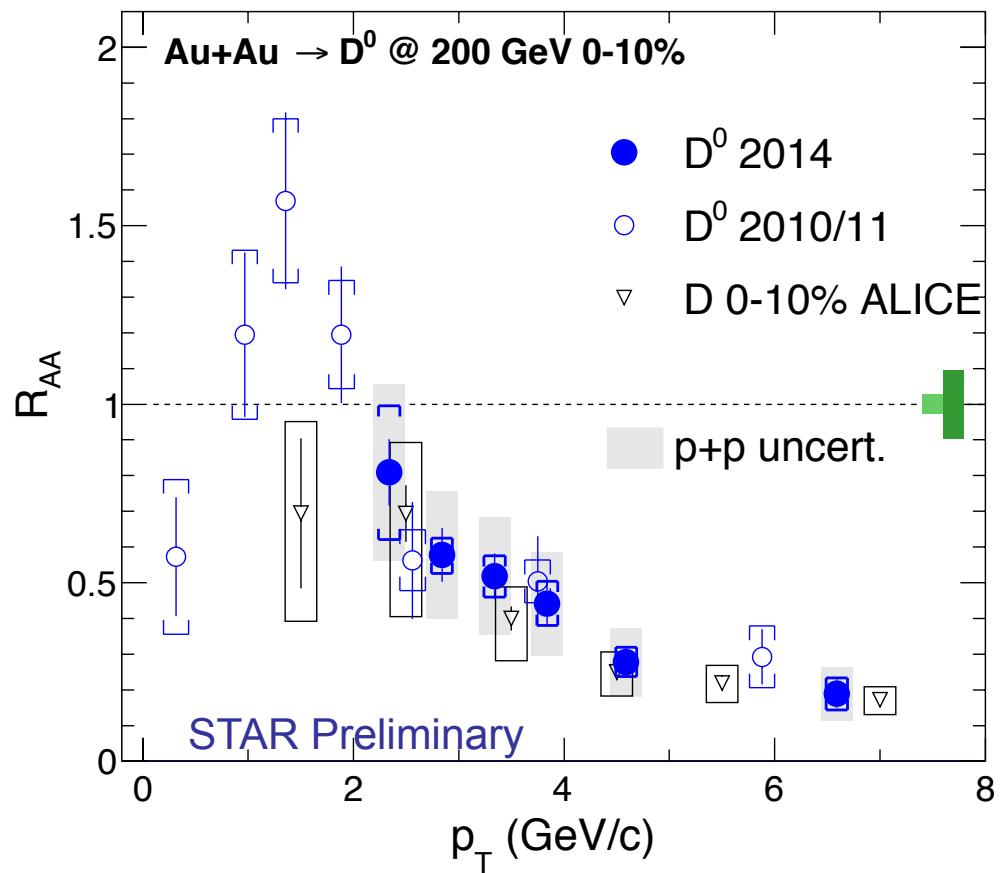
- Strangeness enhancement in heavy-ion collisions is expected to affect the yield of  $D_s$ : relative increase of  $D_s$  yield than  $D^0$  predicted:

The ratio of  $D_s/D^0$  yield measured in Au+Au collisions is found to be higher than that in p+p collisions from PYTHIA

- Elliptic flow of  $D_s < D^0$  is expected due to earlier freeze out of  $D_s$ :

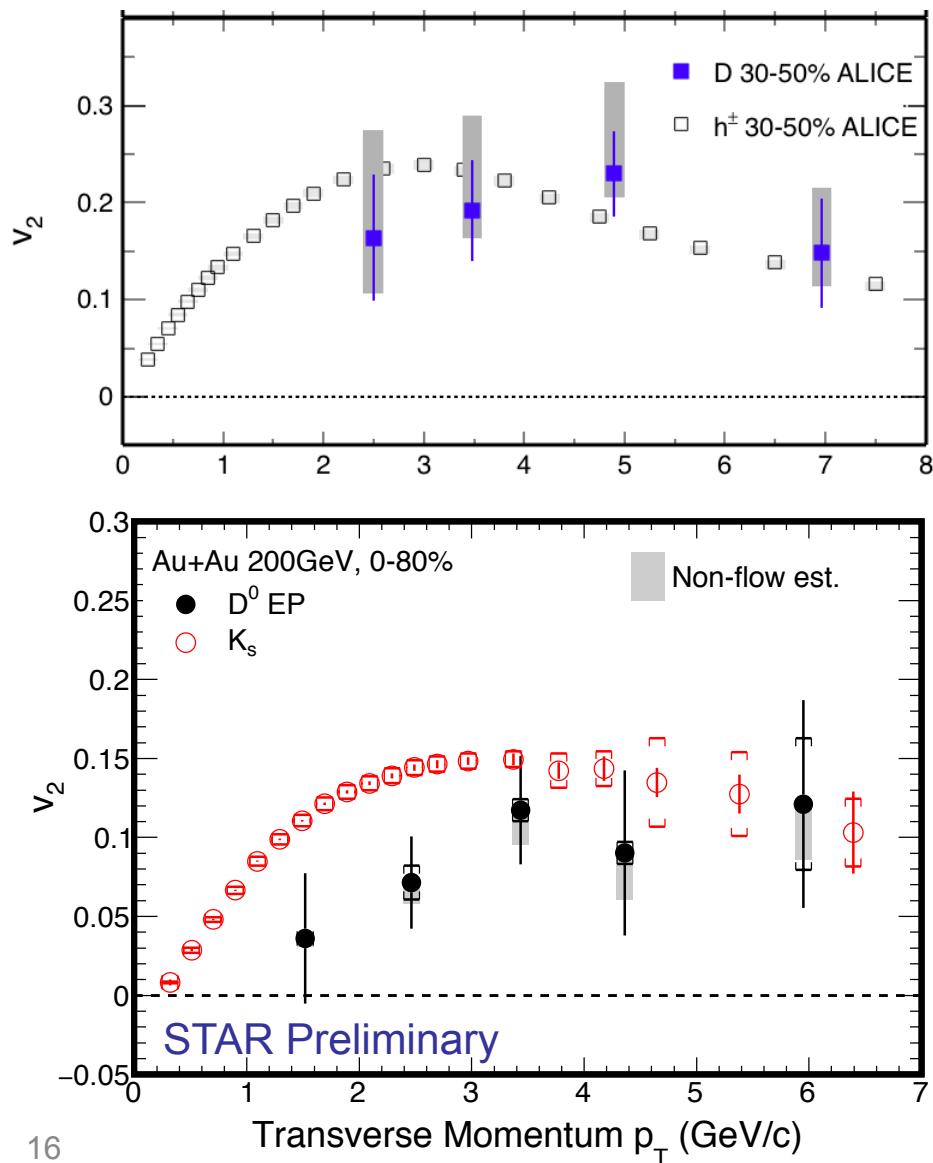
First measurement of  $D_s v_2$  in heavy-ion experiment. More data are needed to draw conclusion.

# Comparison with LHC Results



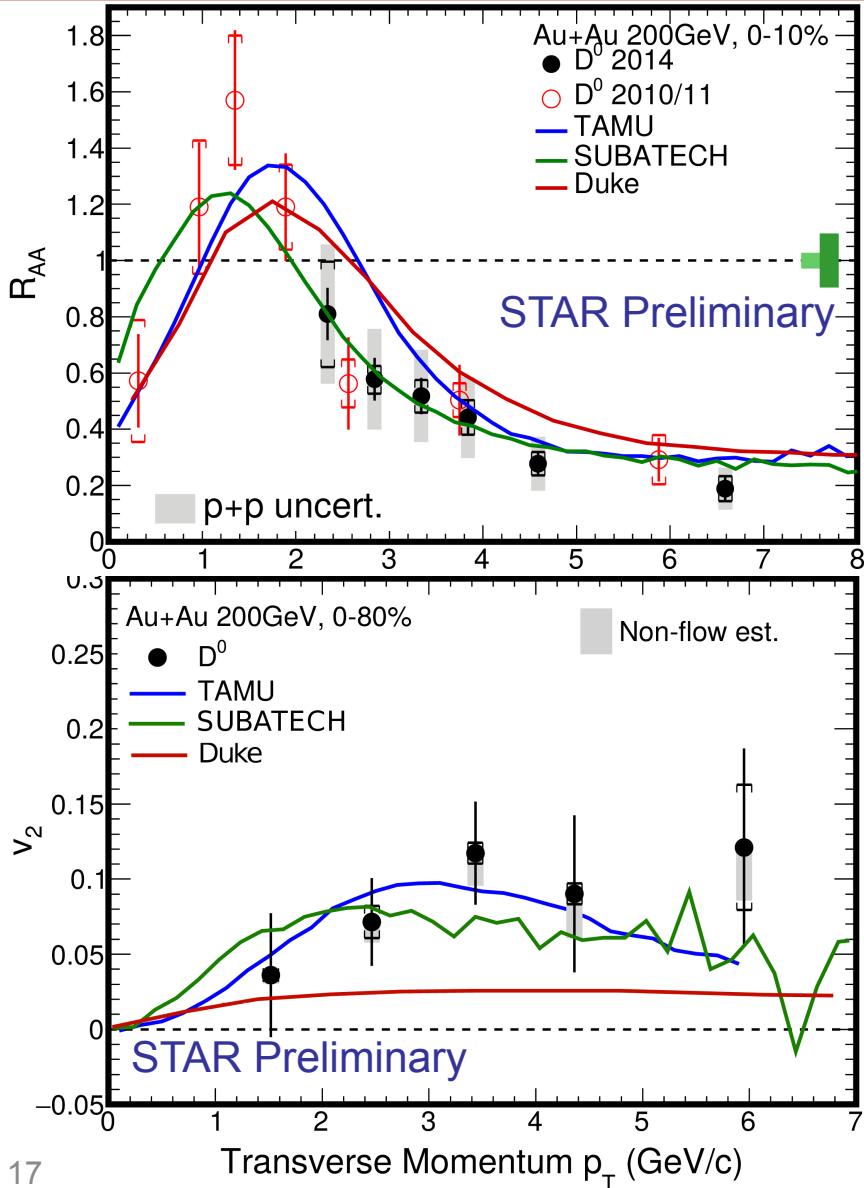
- $D$  meson  $R_{AA}$  @ RHIC  $\sim R_{AA}$  @ LHC  
at  $p_T > 4$  GeV/c
- Strong charm-medium interaction at RHIC and LHC

# Comparison with LHC Results



- $D$  meson  $R_{AA}$  @ RHIC  $\sim R_{AA}$  @ LHC at  $p_T > 4$  GeV/c
- Strong charm-medium interaction at RHIC and LHC
- $D^0 v_2$  LHC results are compatible with light flavor  $v_2$
- $D^0 v_2$  STAR results are lower than light flavor  $v_2$
- Charm thermalized at LHC energy but not fully thermalized at RHIC?
- More precise data and systematic theoretical studies of heavy flavor production at RHIC and LHC will be very helpful.

# Comparison with Theory



TAMU: non-perturb. T-matrix  
 $(2\pi T)D = 2-11$

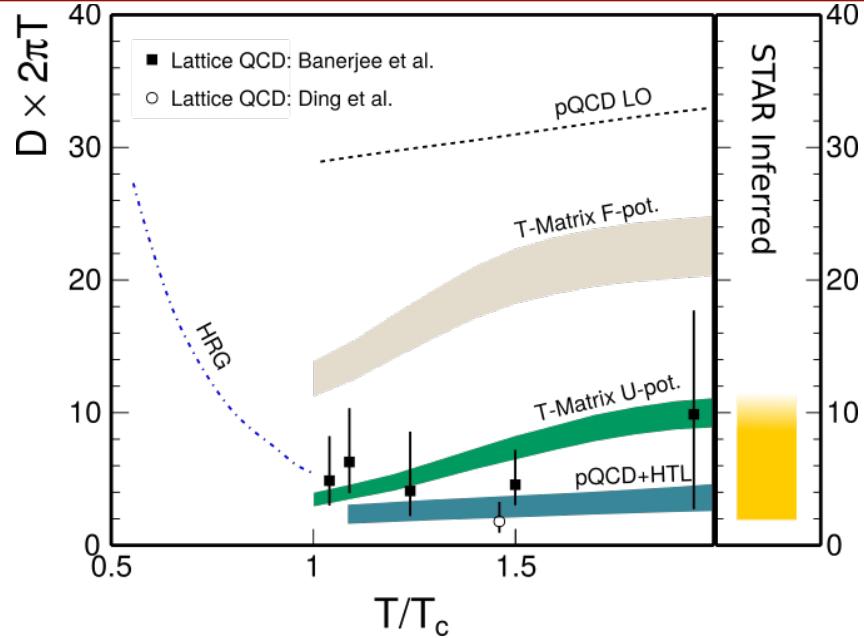
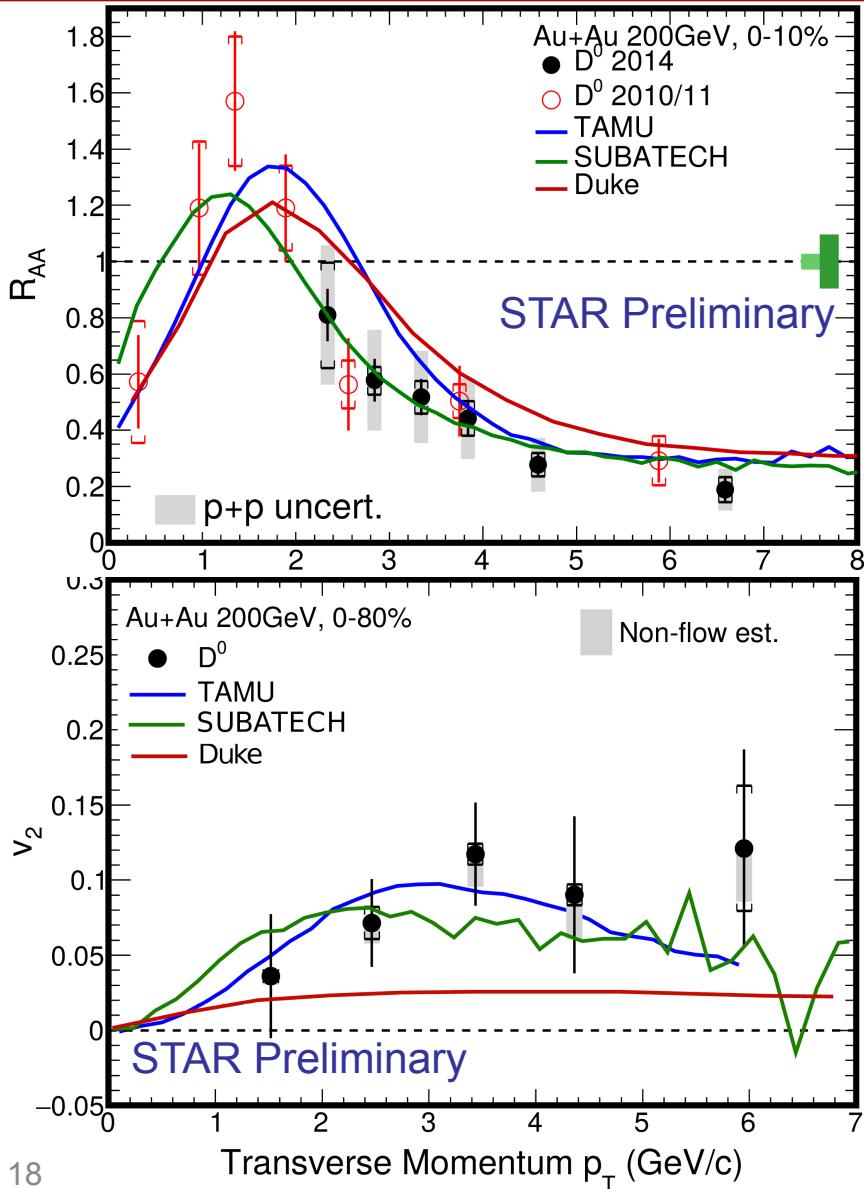
SUBATECH: perturb.+resummation  
 $(2\pi T)D = 2-4$

DUKE: Langevin simulation with input parameter tuned to the LHC data  
 $(2\pi T)D = 7$

	$D \times 2\pi T$	Diff. Calculation
TAMU	2-11	T-Matrix
SUBATECH	2-4	pQCD+HTL
Duke	7	Free parameter

STAR  $D^0$  2010/11: PRL 113 (2014) 142301  
 Theory curves: latest calculations from private communications  
 DUKE: PRC 92 (2015) 024907  
 A. Andronic arXiv:1506.03981(2015)

# Comparison with Theory



Models with charm diffusion coefficient of 2-~10 describe STAR  $D^0 R_{AA}$  and  $v_2$  results.  
 Lattice calculations are consistent with values inferred from data.

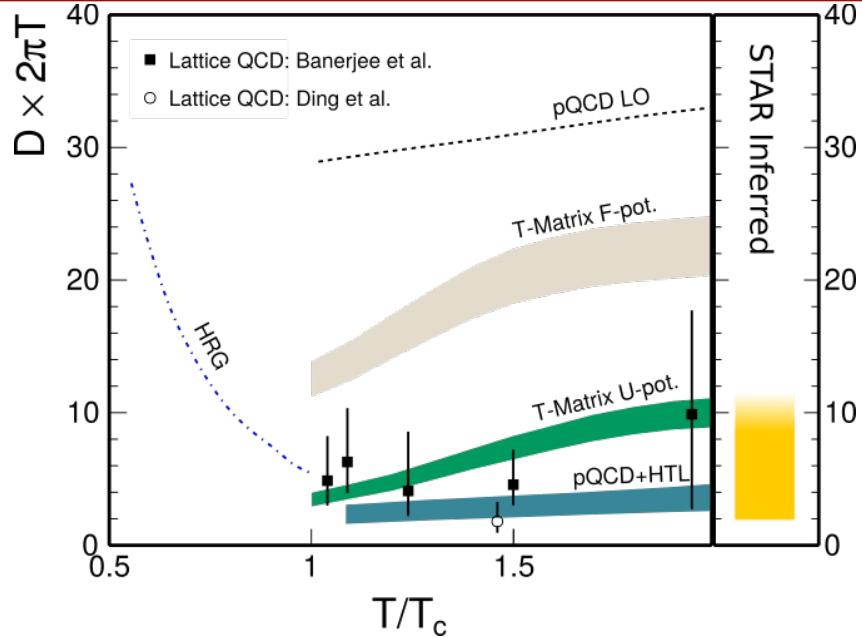
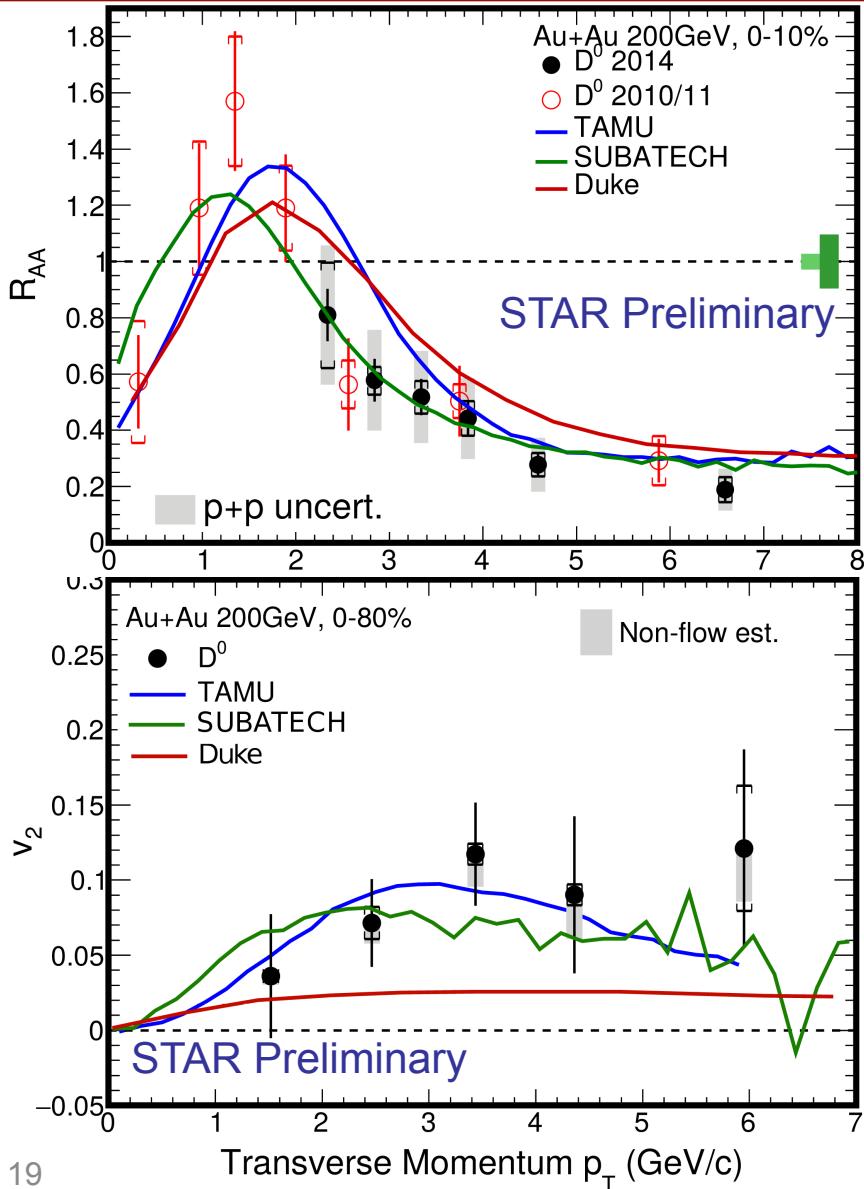
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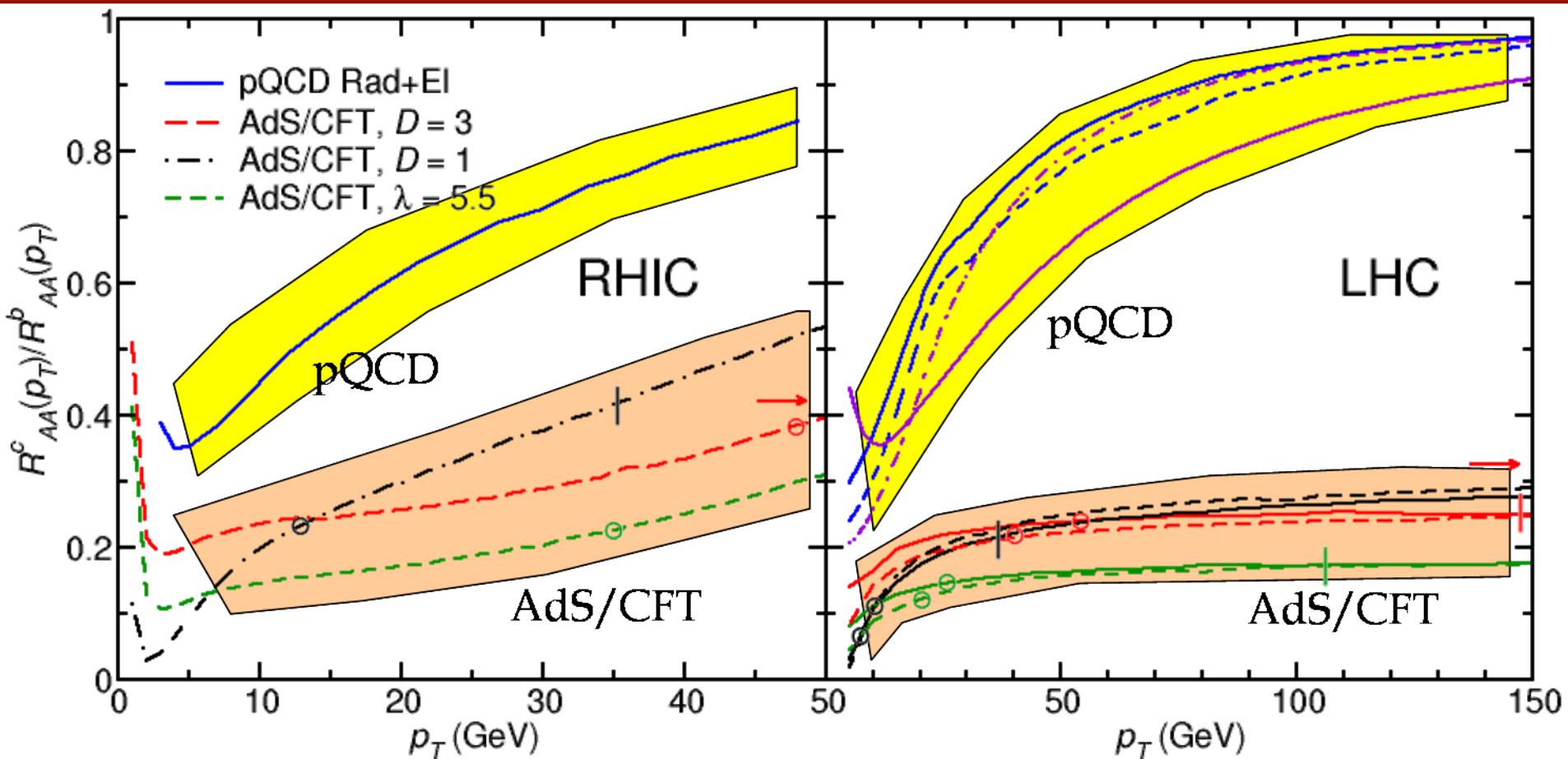
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More precise results expected from STAR Run15 (pp, pAu) and Run16 (AuAu) data: improved p+p baseline, CNM, a factor of ~3 increase in Au+Au data size, improved DCA resolution at low  $p_T$  with AI cables for PXL

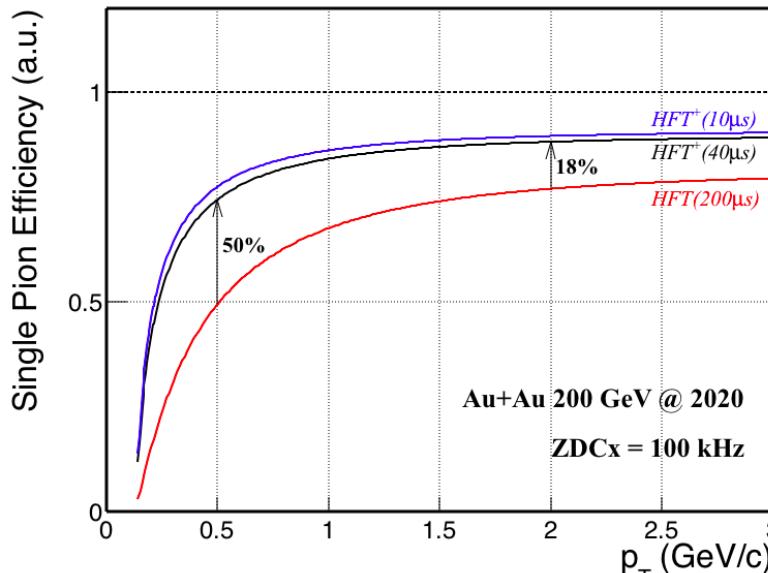
# STAR Heavy Flavor II (2021-2022)



W. Horowitz and M. Gyulassy, arXiv:0710.0703

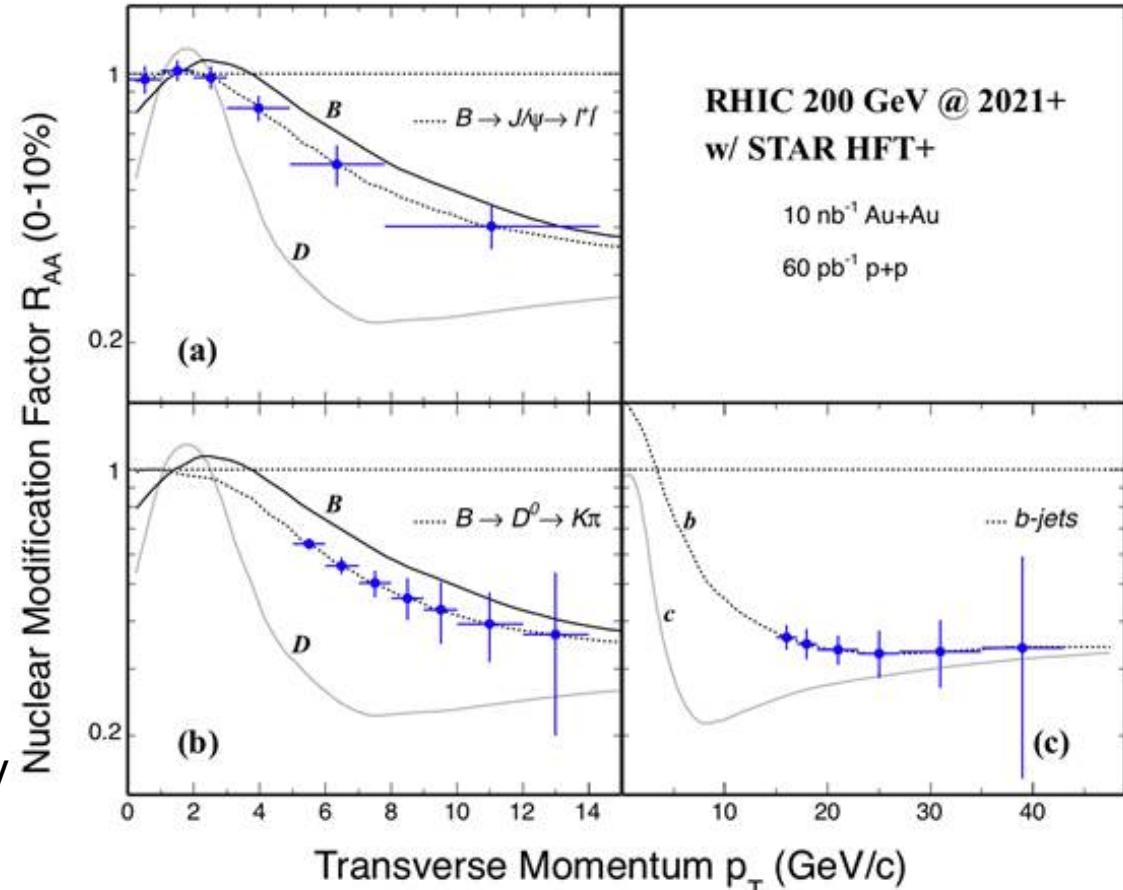
Without Bottom from RHIC, can we claim that we fully understand the energy loss mechanisms, or mass- and temperature-dependent parton transport coefficients of the QGP? Does b quark diffuse in the QGP at RHIC energies and if so how much?

# STAR Heavy Flavor II (2021-2022)



## HFT+ with Faster MAPS sensors

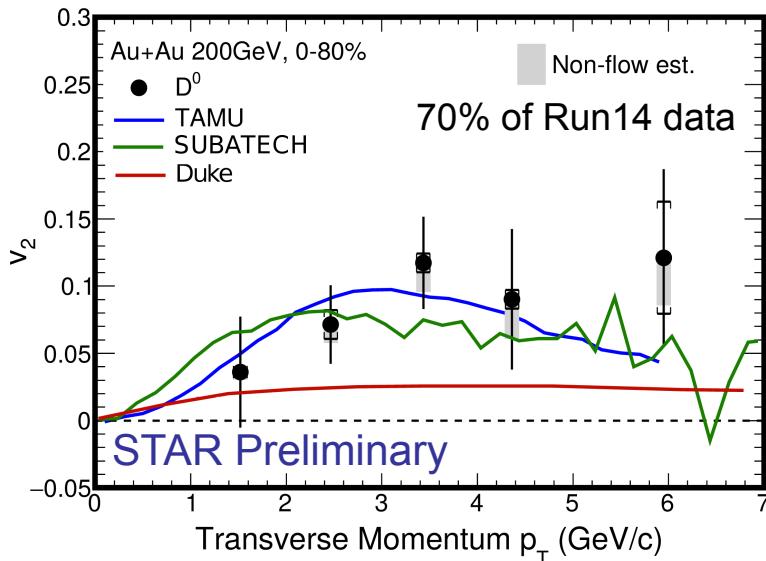
- integration time from  $\sim 185 \mu\text{s}$  to below  $40 \mu\text{s}$  – less pile-up hits and thus better tracking efficiency
- use chips developed for ALICE ITS upgrade and existing HFT infrastructure – cost effective
- experienced team worked on HFT



**Projected  $R_{AA}(0-10\%)$  stat. uncertainty**  
for RHIC pp and AuAu running in 2021-22

Precise bottom measurements with the HFT+ to complete the heavy flavor physics at RHIC. Complementary to ALICE HF and sPHENIX Jet and Upsilon programs.

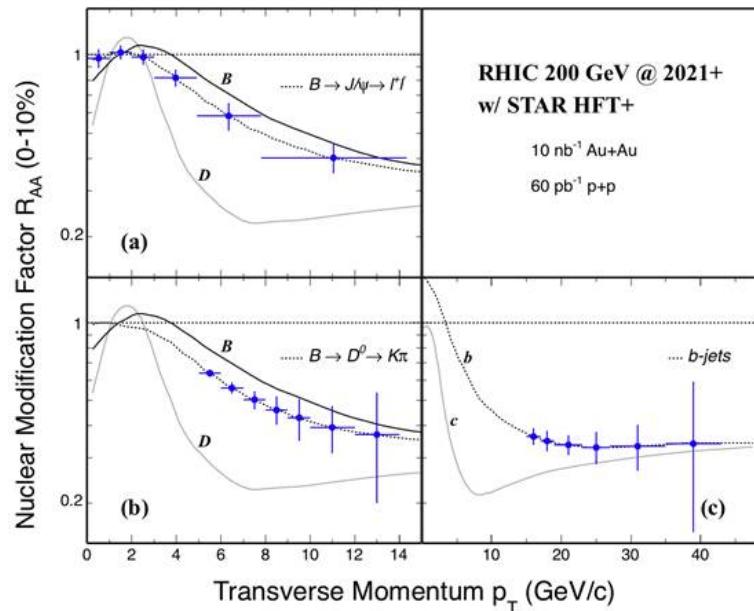
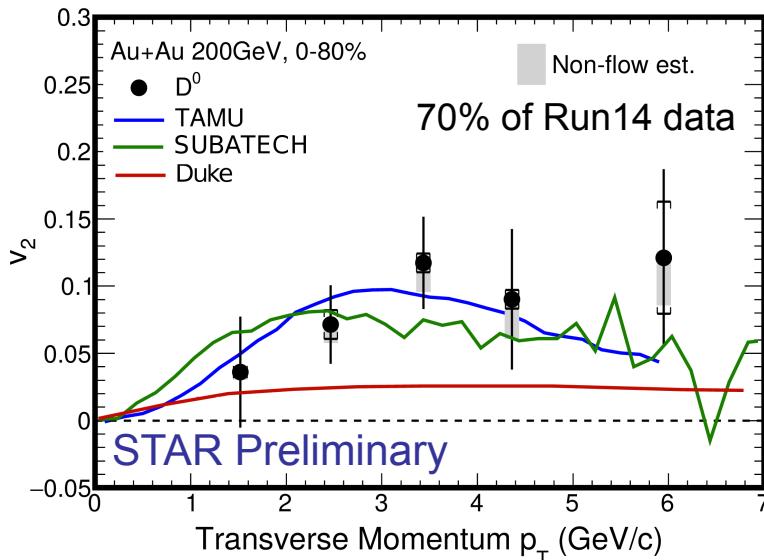
# Summary and Outlook



## STAR HFT in Run14-16

- Run14: Au+Au, results based on ~70% stat.
- Run15: p+p baseline, p+Au for CNM effects
- Run16(+14): x4 Au+Au data size than QM15, inner PXL 0.5->0.4% $X_0$  with AI cables
  - Precise charm results
  - First bottom results

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- Precise charm results
- First bottom results

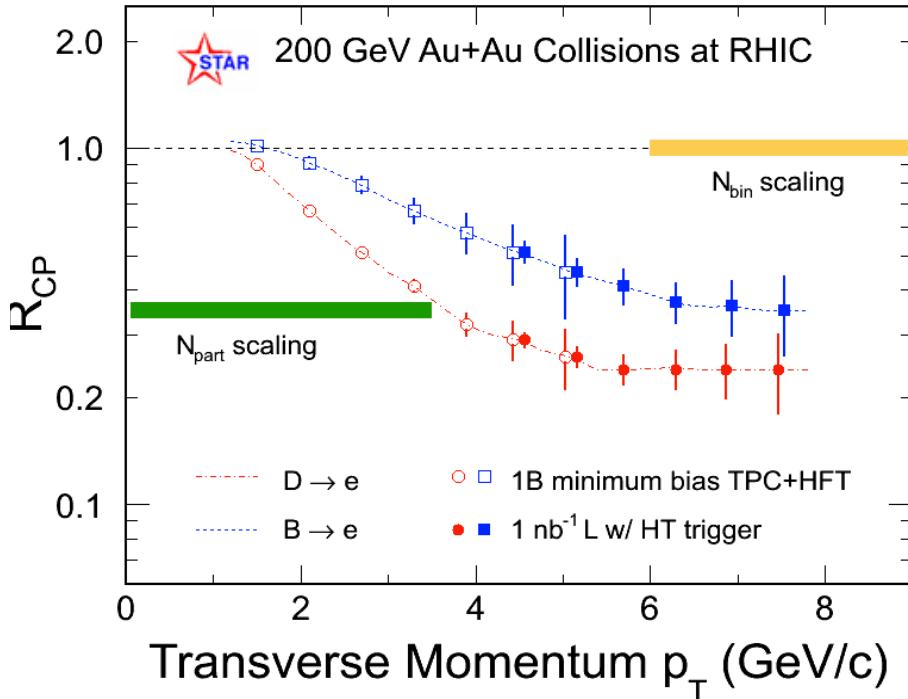
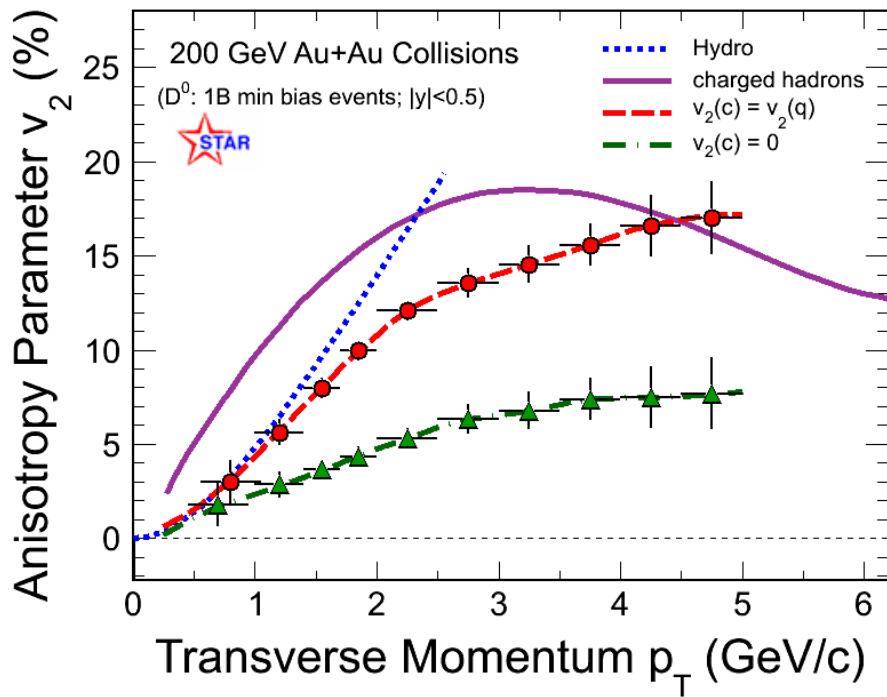
## Upgraded HFT+ in 2020+

HFT+ with faster MAPS sensors will allow precise measurements of bottom quark production at RHIC through B->J/ψ, B->D and b-tagged jets

- Precise bottom results

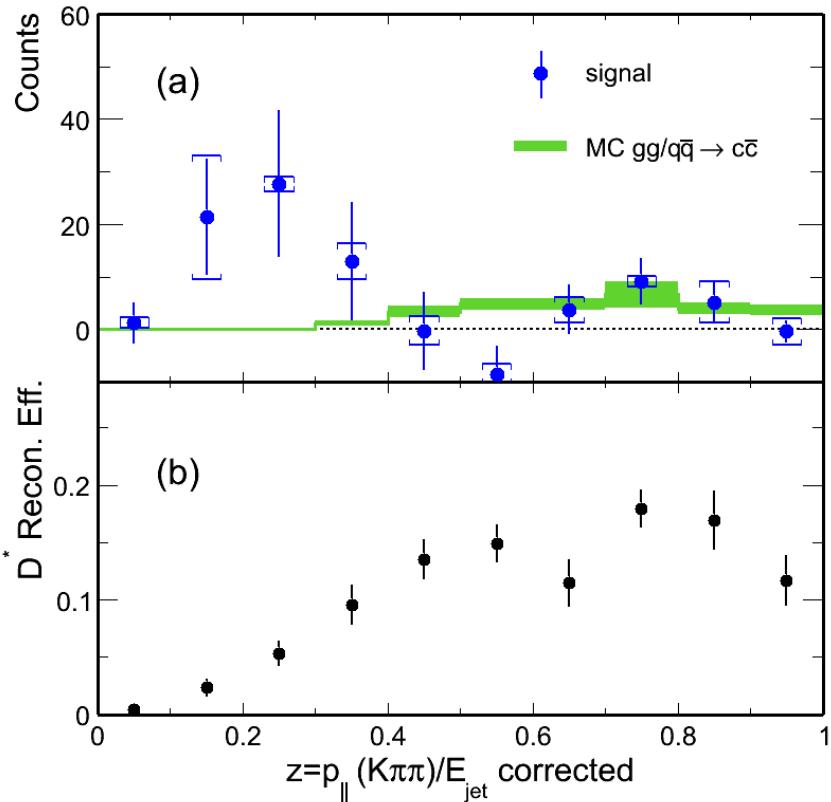
# Backup

# STAR Heavy Flavor Tracker



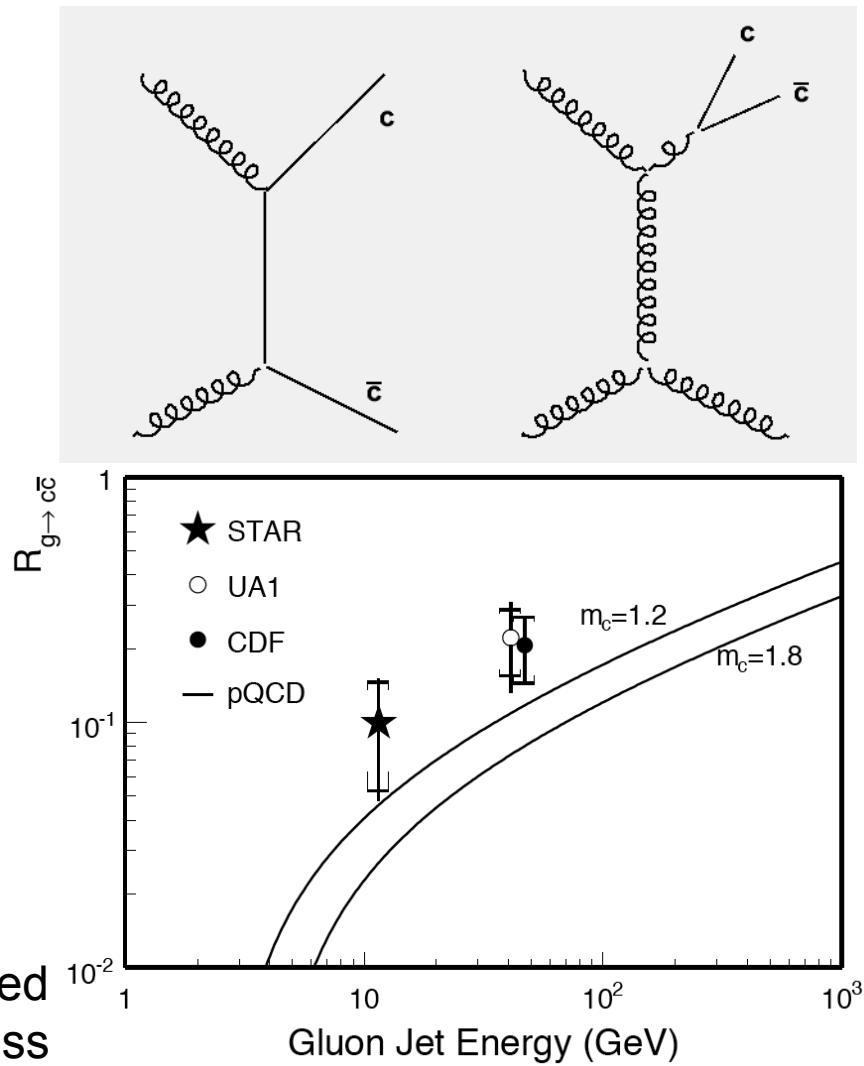
High precision  $R_{AA}$ ,  $R_{pA}$ ,  $v_2$ , correlations results for D mesons and HF leptons;  
Unique at low  $p_T$  → medium thermalization, total charm production

# Charm Production at RHIC



$$N(D^*)/N(\text{jets}) = (1.5 \pm 0.8 \pm 0.7) \times 10^{-2}$$

High  $z$  production is suppressed w.r.t. low  $z$  by trigger bias. The magnitude in data is reproduced by MC with direct flavor creation process. Excess at low  $z$  is from high order processes.



STAR PRD79 (2009) 112006

Zhenyu Ye

# Charm Production at RHIC

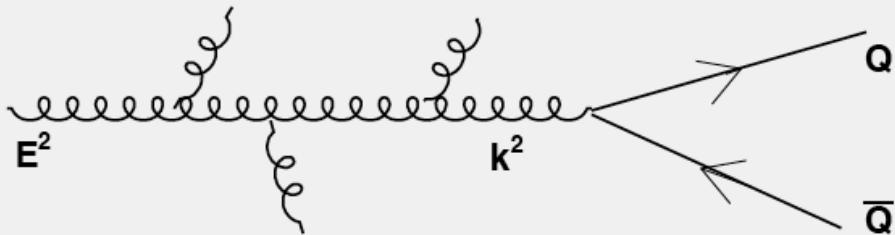


Figure 1: Gluon splitting into  $Q\bar{Q}$

Gluon fragmentation into  $Q\bar{Q}$  pairs is calculable in perturbative QCD. The process is represented in Fig. 1. The gluon multiplicity of those having virtuality  $k^2$  could be calculated as:

$$n_g(E^2, k^2) = \left[ \frac{\ln(E^2/\Lambda^2)}{\ln(k^2/\Lambda^2)} \right]^a \times \exp \left\{ [(2C_A/\pi b) \ln(E^2/\Lambda^2)]^{1/2} \right\} / \exp \left\{ [(2C_A/\pi b) \ln(k^2/\Lambda^2)]^{1/2} \right\}$$

where  $a = -1/4 \times [1 + (2N_f/3\pi b)(1 - C_F/C_A)]$ ,  $b = (11C_A - 2N_f)/12\pi$ .

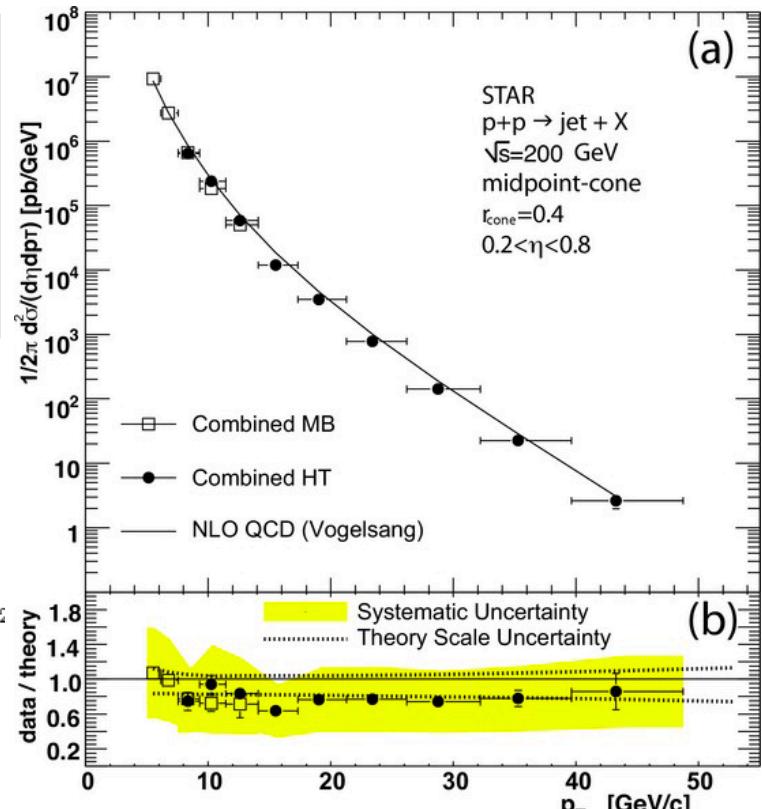
The average number of  $Q\bar{Q}$  pairs in a gluon jet is:

$$R_{Q\bar{Q}}(E) = \int_{4m^2}^{E^2} \frac{dk^2}{k^2} \frac{\alpha_s(k)}{2\pi} \int_{z_-}^{z_+} \frac{1}{2} [z^2 + (1-z)^2 + \frac{2m^2}{k^2}] dz \times n_g(E^2, k^2)$$

where  $z_\pm = (1 \pm \beta)/2$  and  $\beta = \sqrt{1 - 4m^2/k^2}$ .  $m$  is the heavy quark mass.

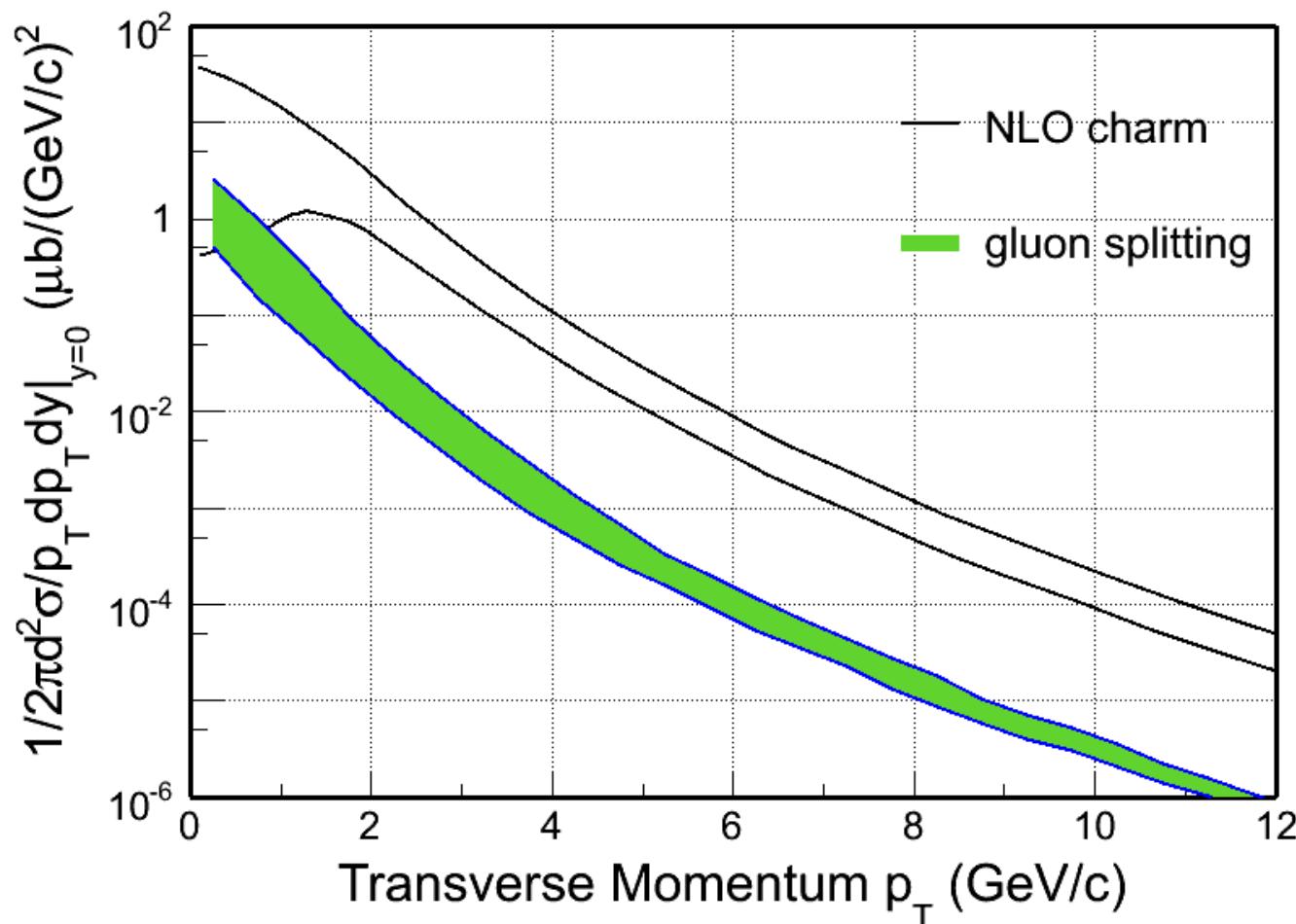
In my calculation for  $z_H$  (charm hadron  $z$ ) coverage, the charm quark fragmentation function is used as the Peterson fragmentation function:

$$D_Q^H(z) \propto \frac{1}{z} \left( 1 - \frac{1}{z} - \frac{\epsilon_Q}{1-z} \right)^{-2}$$



STAR PRL 97 (2006) 252001

# Charm Production at RHIC



Charm production in jets at  $p_T \sim 2-10$  GeV/c has a small contribution from gluon splitting and is dominated by jets initiated by charm quarks

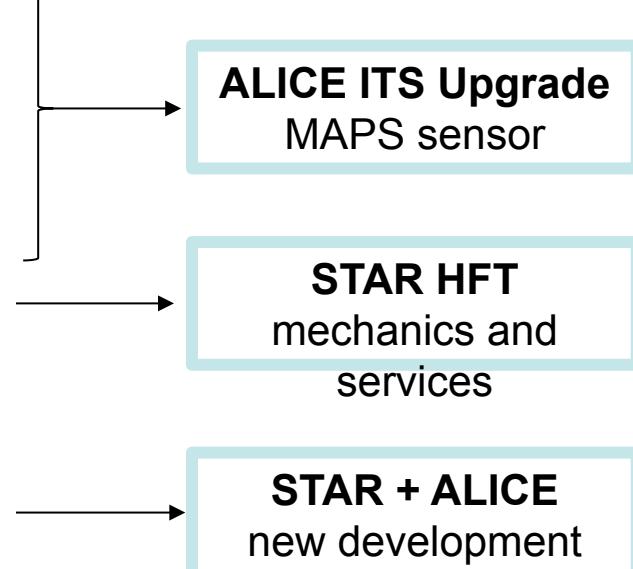
# HFT+ Upgrade plan (2021+)

HFT+ upgrade motivation:

- Measure **bottom quark hadrons** at the RHIC energy
- Take data in **higher luminosity** with high efficiency

HFT+ detector requirements:

- **Faster** frame readout of 40  $\mu\text{s}$  or less
- **Similar or better** pointing resolution  
S/N ratio  
total power consumption  
radiation length
- **Compatible** with the existing insertion mechanism,  
support structure, air cooling system



HFT+ read-out electronics requirements:

- **Compatible** with STAR DAQ system and trigger