Practice: Static and Kinematic Admissibility

Problem 2

Let us consider a sphere with cavity (exterior radius R_e and internal radius R_i). The external surface S_e is free of loading and the interior surface S_i is submitted to a pressure p_0 . The material is assumed to be isotropic linear elastic with young modulus E and Poisson ratio ν .

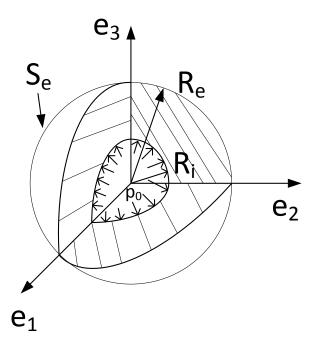


Figure 1: Configuration of problem 2

Question 1: Write the static admissibility equations!

Answer: Find $\underline{\sigma}$ symmetric such that

(a)
$$\underline{\underline{\sigma}} \cdot \underline{n} = \underline{0} \quad \forall M \in S_e$$
 Boundary equation
(b) $\underline{\underline{\sigma}} \cdot \underline{n} = -p_0 \underline{n} \quad \forall M \in S_i$ Boundary equation
(c) $\underline{\underline{\text{div}}\underline{\sigma}} = \underline{0} \quad \forall M \in \Omega$ Internal equilibrium

Question 2: Please expand this equation using spherical coordinates system!

Answer: In spherical coordinates,

$$\underline{\underline{\sigma}} = \begin{bmatrix} \sigma_{rr} & \sigma_{r\theta} & \sigma_{r\varphi} \\ \sigma_{r\theta} & \sigma_{\theta\theta} & \sigma_{\theta\varphi} \\ \sigma_{r\varphi} & \sigma_{\theta\varphi} & \sigma_{\varphi\varphi} \end{bmatrix}$$
(2)

From Eq. (a),

$$\sigma_{rr}(r = R_e, \theta, \varphi) = \sigma_{r\theta}(r = R_e, \theta, \varphi) = \sigma_{r\varphi}(r = R_e, \theta, \varphi) = 0$$
(3)

From Eq. (b),

$$-\sigma_{rr}(r = R_i, \theta, \varphi) = p_0$$

$$-\sigma_{r\theta}(r = R_i, \theta, \varphi) = 0$$

$$-\sigma_{r\varphi}(r = R_i, \theta, \varphi) = 0$$
(4)

From Eq. (c) \rightarrow See Textbook equation 1.11a (divergence with spherical coordinates)

Question 3: Let's assume that the following stress tensor

$$\underline{\underline{\sigma}} = \begin{bmatrix} \sigma_{rr} & 0 & 0 \\ 0 & \sigma_{\theta\theta} & 0 \\ 0 & 0 & \sigma_{\varphi\varphi} \end{bmatrix} \tag{5}$$

with

$$\sigma_{rr} = -A \left[\frac{B}{r^3} - 1 \right] \tag{6}$$

and

$$\sigma_{\theta\theta} = \sigma_{\varphi\varphi} = A \left[\frac{B}{2r^3} + 1 \right] \tag{7}$$

Please calculate the coefficient A and B so that $\underline{\sigma}$ could be a solution of the problem!

Answer: Using the proposed stress tensor, and expanding equation (c), we get

$$\frac{\partial \sigma_{rr}}{\partial r} + \frac{2\sigma_{rr} - \sigma_{\theta\theta} - \sigma_{\varphi\varphi}}{r} = 0$$

$$0 = 0$$

$$0 = 0$$
(8)

We now check if $\underline{\underline{\sigma}}$ ca fulfill the static admissibility equations.

$$\frac{\partial \sigma_{rr}}{\partial r} + \frac{2\sigma_{rr} - \sigma_{\theta\theta} - \sigma_{\varphi\varphi}}{r} = \frac{3AB}{r^4} - \frac{3AB}{r^4} = 0 \rightarrow \text{(c) is OK. Interior equilibrium is satisfied.}$$
(9)

(a)
$$\to -A \left[\frac{B}{R_e^3} - 1 \right] = 0 \to B = R_e^3$$
 (10)

(b)
$$\to -A \left[\frac{R_i^3}{R_e^3} - 1 \right] = p_0 \to A = p_0 \frac{R_e^3}{R_i^3 - R_e^3}$$
 (11)

All the other conditions concerning $\sigma_{\theta r}$ and $\sigma_{r\varphi}$ are automatically satisfied due to the chosen stress tensor. Thus, a statically admissible stress field is:

$$\underline{\underline{\sigma}} = \begin{bmatrix} p_0 \frac{R_e^3}{R_i^3 - R_e^3} \left[\frac{R_i^3}{r} - 1 \right] & 0 & 0 \\ 0 & -p_0 \frac{R_e^3}{R_e^3 - R_i^3} \left[\frac{R_e^3}{2r^3} + 1 \right] & 0 \\ 0 & 0 & -p_0 \frac{R_e^3}{R_e^3 - R_i^3} \left[\frac{R_e^3}{2r^3} + 1 \right] \end{bmatrix}$$
(12)

Question 4: Please calculate the strain tensor $\underline{\underline{\varepsilon}}$? Does it satisfy the compatibility equation? Conclusion?

Answer: Based on the constitutive equation, we can calculate the strains:

$$\underline{\underline{\varepsilon}} = \frac{1+\nu}{E}\underline{\underline{\sigma}} - \frac{\nu}{E} \text{tr}\underline{\underline{\sigma}1} \tag{13}$$

$$\operatorname{tr}\underline{\underline{\sigma}} = -p_0 \frac{R_e^3}{R_e^3 - R_i^3} \tag{14}$$

Thus, we get $\varepsilon_{r\theta} = \varepsilon_{\theta\varphi} = \varepsilon_{\varphi\theta} = 0$. Let us note that

$$A_1 = p_0 \frac{R_e^3}{R_e^3 - R_i^3} \tag{15}$$

We get:

$$\underline{\underline{\varepsilon}} = \begin{bmatrix} \left(\frac{1+\nu}{E}\right) A_1 \left[\frac{R_e^3}{r} - 1\right] + \frac{\nu}{E} A_1 & 0 & 0 \\ 0 & -\left(\frac{1+\nu}{E}\right) A_1 \left[\frac{R_e^3}{2r^3} + 1\right] + \frac{\nu}{E} A_1 & 0 \\ 0 & 0 & -\left(\frac{1+\nu}{E}\right) A_1 \left[\frac{R_e^3}{2r^3} + 1\right] + \frac{\nu}{E} A_1 \end{bmatrix}$$
(16)

In case of complete spherical symmetry, the compatibility equation can be simplified (see Textbook equation 2.14b):

$$\frac{d\varepsilon_{\varphi\varphi}}{dr} + \frac{\varepsilon_{\varphi\varphi} - \varepsilon_{rr}}{r} = 0 \tag{17}$$

Let us verify if this equation is satisfied by $\underline{\varepsilon}$.

$$\frac{d\varepsilon_{\varphi\varphi}}{dr} = 3A_1 \left(\frac{1+\nu}{E}\right) \frac{R_e^3}{2r^4} \tag{18}$$

$$\frac{\varepsilon_{\varphi\varphi} - \varepsilon_{rr}}{r} = -3A_1 \left(\frac{1+\nu}{E}\right) \frac{R_e^3}{2r^4} \tag{19}$$

The two equations above satisfied the compatibility equation. Thus it will be possible to integrate $\underline{\varepsilon}$ to determine a displacement field \underline{u} . Considering there is no additional kinematic condition (no kinematic boundary condition), we got the exact solution of the problem.

Problem 3

Let us consider a domain Ω , simply supported by a rigid basis (Γ) . Ω is submitted to gravity volume forces $f = \rho gy$ and to uniform pressure p > 0 over every vertical side. We assume plane stresses in the plane of x-y and the material is isotropic linear elastic with Young modulus E and Poisson ratio ν .

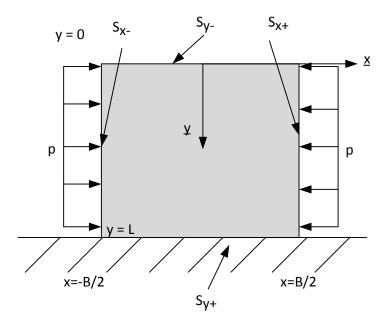


Figure 2: Configuration of problem 3

Question 1: Write down the full problem! (kinematic, static and constitutive equations)

Answer: Kinematic admissibility:

Let us not \underline{u} the displacement vector such that $\underline{u} = \begin{pmatrix} u \\ v \end{pmatrix}$, $\underline{u} = u\underline{x} + v\underline{y}$ $\begin{cases} v = 0, & \forall M \in S_y^+ \\ \underline{\underline{\epsilon}} = \frac{1}{2} \left(\underline{\underline{\nabla}} \underline{u} + \underline{\underline{\nabla}} \underline{u}^t \right), & \text{compatibility} \end{cases}$

$$\begin{cases} v = 0, & \forall M \in S_y^+ \\ \underline{\underline{\epsilon}} = \frac{1}{2} \left(\underline{\underline{\nabla}} \underline{u} + \underline{\underline{\nabla}} \underline{u}^t \right), & \text{compatibility} \end{cases}$$
 (20)

Static admissibility: Find $\underline{\sigma}$ symmetric such that

$$\begin{array}{l}
(1) \ \underline{\underline{\sigma}} \cdot \underline{n} = -p\underline{n} \quad \forall \ M \in S_{x}^{+} \\
(2) \ \underline{\underline{\sigma}} \cdot \underline{n} = -p\underline{n} \quad \forall \ M \in S_{x}^{-} \\
(3) \ \underline{\underline{\sigma}} \cdot \underline{n} = \underline{0} \quad \forall \ M \in S_{y}^{-} \\
(4) \ (\underline{\underline{\sigma}} \cdot \underline{n}) \cdot \underline{t} = \underline{0} \quad \forall \ M \in S_{y}^{+}
\end{array}
\right\}$$
Boundary equation (21)

$$\underline{\operatorname{div}} \sigma + \rho g \underline{y} = \underline{0} \quad \forall \ M \in \Omega \quad \text{Internal equilibrium}$$
 (22)

As plane stress are assumed, the equation can be reduced to:

$$\sigma_{xx,x} + \sigma_{xy,y} = 0
\sigma_{xy,x} + \sigma_{yy,y} + \rho g = 0
\sigma_{xz} = \sigma_{yz} = \sigma_{zz} = 0$$

$$\forall M \in \Omega$$
(23)

Constitutive equation: The 3D constitutive equation is:

$$\underline{\underline{\epsilon}} = \frac{1+\nu}{E}\underline{\underline{\sigma}} - \frac{\nu}{E} \text{tr}\underline{\underline{\sigma}}\underline{\underline{1}}$$
 (24)

Here, we can use the plane stress assumption $(\sigma_{xz} = \sigma_{yz} = \sigma_{zz} = 0)$

$$\begin{bmatrix} \epsilon_{xx} & \epsilon_{xy} & \epsilon_{xz} \\ \epsilon_{xy} & \epsilon_{yy} & \epsilon_{yz} \\ \epsilon_{xz} & \epsilon_{yz} & \epsilon_{zz} \end{bmatrix} = \begin{pmatrix} 1+\nu \\ E \end{pmatrix} \begin{bmatrix} \sigma_{xx} & \sigma_{xy} & 0 \\ \sigma_{xy} & \sigma_{yy} & 0 \\ 0 & 0 & 0 \end{bmatrix} - \frac{\nu}{E} \left(\sigma_{xx} + \sigma_{yy}\right) \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$
(25)

As a consequence, we can also write this equation in 2D:

$$\begin{cases}
\underline{\epsilon}^{2D} = \frac{1+\nu}{E} \underline{\sigma}^{2D} - \frac{\nu}{E} \operatorname{tr} \underline{\sigma}^{2D} \underline{1}^{2D} \\
\epsilon_{xz} = \epsilon_{yz} = 0 \\
\epsilon_{zz} = -\frac{\nu}{E} (\sigma_{xx} + \sigma_{yy}) \quad \text{the out of plane strain is not equal to zero}
\end{cases} \tag{26}$$

Question 2: We assume an Airy function $\phi(x,y) = ax^2 + cy^2 + hy^3$. Find a, c, h so we define a statistically admissible stress field.

Answer: We want to use an approach based on Airy function but we have to be careful in case of volume force.

Let us assume the volume force density \underline{f} derive from a potential:

$$\underline{\mathbf{f}} = \mathrm{grad}V \tag{27}$$

For example here: $V = \rho gy \Rightarrow \underline{\mathbf{f}} = \underline{\mathbf{grad}}V = \begin{pmatrix} 0 \\ \rho g \end{pmatrix}$

In this case:

$$\begin{cases}
\frac{\partial(\sigma_{xx} + V)}{\partial x} + \frac{\partial\sigma_{xy}}{\partial y} = 0 \\
\frac{\partial\sigma_{xy}}{\partial x} + \frac{\partial(\sigma_{yy} + V)}{\partial y} = 0
\end{cases}$$
(28)

and $\sigma_{xx}, \sigma_{xy}, \sigma_{yy}$ can be defined based on the Airy function by the following equation:

$$\sigma_{xx} + V = \frac{\partial^2 \phi}{\partial y^2}; \quad \sigma_{yy} + V = \frac{\partial^2 \phi}{\partial x^2}; \quad \sigma_{xy} = -\frac{\partial^2 \phi}{\partial x \partial y}$$
 (29)

If $\phi(x,y) = ax^2 + cy^2 + hy^3$

$$\sigma_{xx} = \frac{\partial^2 \phi}{\partial y^2} - V = 2C + 6hy - \rho gy \to \sigma_{xx,x} = 0$$

$$\sigma_{yy} = \frac{\partial^2 \phi}{\partial x^2} - V = 2a - \rho gy \to \sigma_{yy,y} = -\rho g$$

$$\sigma_{xy} = -\frac{\partial^2 \phi}{\partial x \partial y} = 0$$
(30)

Interior equation:
$$\begin{cases} \sigma_{xx,x} + \sigma_{xy,y} = 0 \to \text{OK} \\ \sigma_{xy,x} + \sigma_{yy,y} + \rho g = 0 \to \text{OK} \end{cases}$$
(31)

Boundary static equation:

$$(1) \Rightarrow \begin{cases} \sigma_{xx}(x = B/2, y) = -p, & \forall y \to \begin{cases} 2c = -p \\ 6h - \rho g = 0 \end{cases} \to \begin{cases} c = -p/2 \\ h = \frac{\rho g}{6} \end{cases} \\ \sigma_{xy}(x = B/2, y) = 0, & \forall y \text{ OK} \end{cases}$$

$$(2) \Rightarrow \begin{cases} \sigma_{xx}(x = -B/2, y) = -p, & \forall y \text{ OK} \\ \sigma_{xy}(x = -B/2, y) = 0, & \forall y \text{ OK} \end{cases}$$

$$(32) \Rightarrow \begin{cases} \sigma_{yy}(y = 0, x) = 0, & \forall x \to 2a = 0 \to a = 0 \\ \sigma_{xy}(y = 0, x) = 0, & \forall x \text{ OK} \end{cases}$$

$$(4) \Rightarrow \{\sigma_{xy}(y = L, x) = 0, & \forall x \text{ OK} \end{cases}$$

Finally, we have $a=0; c=-p/2; h=\rho g/6$, and we get the following statistically admissible stress field:

$$\begin{cases}
\sigma_{xx} = -p \\
\sigma_{yy} = -\rho gy \\
\sigma_{xy} = 0 \\
\sigma_{xz} = \sigma_{yz} = \sigma_{zz} = 0
\end{cases}$$
(33)

Question 3: Find the related strain field using the constitutive equation. Does that satisfy the compatibility equation? Conclusion?

Answer: As a consequence, using the constitutive equation:

$$\begin{cases}
\epsilon_{xx} = \frac{1}{E}\sigma_{xx} - \frac{\nu}{E}\sigma_{yy} = -\frac{p}{E} + \frac{\nu\rho g}{E}y \\
\epsilon_{yy} = -\frac{\rho g}{E}y + \frac{\nu p}{E} \\
\epsilon_{xy} = \epsilon_{xz} = \epsilon_{yz} = 0 \\
\epsilon_{zz} = \frac{\nu}{E}(p + \rho gy)
\end{cases}$$
(34)

Check the compatibility:

The compatibility equation are obviously satisfied as we just have linear relations (and the compatibility equations involves the second derivation).

Question 4: Integrate the strain field to get the displacement at every point. Can we find the exact solution of the problem using this Airy function?

Answer:

$$u_{,x} = -\frac{p}{E} + \frac{\nu \rho g}{E} y \to u = \left(-\frac{p}{E} + \frac{\nu \rho g}{E} y \right) x + f(y)$$

$$v_{,y} = -\frac{\rho g}{E} y + \frac{\nu p}{E} \to v = -\frac{\rho g}{2E} y^2 + \frac{\nu p}{E} y + g(x)$$

$$\epsilon_{xy} = \frac{1}{2} (u_{,y} + v_{,x}) = 0 \to \frac{\nu \rho g}{E} x + \frac{\partial f(y)}{y} + \frac{\partial g(x)}{x} = 0$$
(35)

$$\frac{\partial f(y)}{y} = \lambda \to f(y) = \lambda y + \beta$$

$$\frac{\nu \rho g}{E} x + \frac{\partial g(x)}{x} = -\lambda \to g(x) = -\frac{\nu \rho g}{2E} x^2 - \lambda x + \gamma$$
(36)

Thus:

$$\begin{cases}
 u = \left(-\frac{p}{E} + \frac{\nu \rho g}{E}y\right)x + \lambda y + \beta \\
 v = -\frac{\rho g}{2E}y^2 + \frac{\nu p}{E}y - \frac{\nu \rho + g}{2E}x^2 - \lambda x + \gamma
\end{cases}$$
(37)

Introducing the kinematic boundary equations:

$$v(x, y = L) \Rightarrow \text{this solution can not be satisfied}$$
 (38)

We do not have the exact solution.

Problem 4

Let us consider a classical test in material engineering that make possible to identify the oedometric modelus \hat{E} . A cylinder material domain Ω is compressed due to the prescribed motion of the top surface $\underline{u} = -U\underline{e}_z$. All surfaces are assumed to be perfect contact without friction. Note that $\underline{P} = -P\underline{e}_z$ the resulting force that is necessary to press the materials (Young modulus E and Poisson ratio ν .

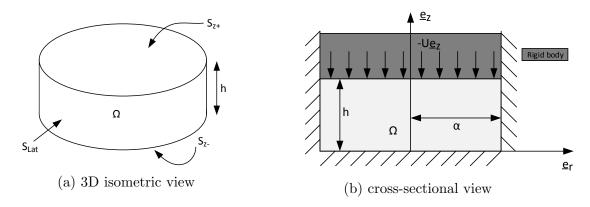


Figure 3: Oedometric compression

Question 1: Write the complete problem: kinematic + static + constitutive equation!

Answer: Kinematic admissibility:

$$\underline{u} \cdot \underline{n} = 0 \quad \forall \underline{M} \in S_{Lat} \to u_r = 0 \quad \forall \underline{M} \in S_{Lat}$$
 (39)

$$\underline{n} \cdot \underline{u} = -U \quad \forall \underline{M} \in S_z^+ \to u_z = -U \quad \forall \underline{M} \in S_z^+ \tag{40}$$

$$\underline{\underline{\varepsilon}}(\underline{u}) = \frac{1}{2} (\underline{\underline{\nabla}}\underline{u} + \underline{\underline{\nabla}}^T\underline{u}) \quad \forall \underline{M} \in \Omega$$
 (41)

$$\underline{u} \cdot \underline{n} = 0 \quad \forall \underline{M} \in S_z^- \to u_z = 0 \quad \forall \underline{M} \in S_z^-$$
 (42)

Static admissibility:

$$\underline{\operatorname{div}}\underline{\sigma} = \underline{0} \quad \forall \underline{M} \in \Omega \tag{43}$$

or in expanded form as below

$$\sigma_{rr,r} + \frac{1}{r} (\sigma_{r\theta,\theta} + \sigma_{rr} - \sigma_{\theta\theta}) + \sigma_{rz,z} = 0$$

$$\sigma_{\theta r,r} + \frac{1}{r} (\sigma_{\theta\theta,\theta} + \sigma_{r\theta} + \sigma_{\theta r}) + \sigma_{\theta z,z} = 0 \qquad \rightarrow \qquad \forall \underline{M} \in \Omega$$

$$\sigma_{zr,r} + \frac{1}{r} (\sigma_{z\theta,\theta} + \sigma_{zr}) + \sigma_{zz,z} = 0$$

$$(44)$$

Contact without friction:

$$(\underline{\sigma} \cdot \underline{t}) \cdot \underline{t} = 0 \quad \forall \underline{M} \in S_z^+ \to \sigma_{rz} = \sigma_{\theta z} = 0 \quad \underline{M} \in S_z^+$$

$$(45)$$

Contact without friction:

$$(\underline{\sigma} \cdot \underline{t}) \cdot \underline{t} = 0 \quad \forall \underline{M} \in S_z^- \to \sigma_{rz} = \sigma_{\theta z} = 0 \quad \underline{M} \in S_z^-$$

$$(46)$$

Contact without friction:

$$(\underline{\sigma} \cdot \underline{t}) \cdot \underline{t} = 0 \quad \forall \underline{M} \in S_{Lat} \to \sigma_{\theta r} = \sigma_{\theta z} = 0 \quad \underline{M} \in S_{Lat}$$
 (47)

Constitutive equations:

$$\underline{\underline{\sigma}} = \lambda \operatorname{tr} \underline{\underline{\varepsilon}} \underline{1} + 2\mu \underline{\underline{\varepsilon}} \quad \forall \underline{M} \in \Omega$$
 (48)

Question 2: Assume that $\underline{u} = -\frac{U}{h}z\underline{e}_z$. Is this kinematically admissible?

Answer: We assume that

$$\underline{u} = \begin{pmatrix} 0 \\ 0 \\ -\frac{U}{h}z \end{pmatrix} = \left(-\frac{U}{h}z\right) \quad \text{Of course } \underline{u} \text{ is kinematically admissible.} \tag{49}$$

By definition (using the gradient of a vector with cylindrical coordinates):

$$\underline{\underline{\nabla}} \cdot \underline{u} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -\frac{U}{h} \end{bmatrix} \to \underline{\varepsilon} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \varepsilon_{zz} \end{bmatrix}$$
 (50)

with $\varepsilon_{zz} = -U/h \rightarrow$ the compatibility equation is obviously satisfied.

Question 3: Find the corresponding stress $\underline{\underline{\sigma}}$. Is it statistically admissible?

Answer:

$$\underline{\underline{\sigma}} = \lambda \operatorname{tr} \underline{\underline{\varepsilon}} \underline{1} + 2\mu \underline{\underline{\varepsilon}} \to \underline{\underline{\sigma}} = \begin{bmatrix} -\lambda \frac{U}{h} & 0 & 0\\ 0 & -\lambda \frac{U}{h} & 0\\ 0 & 0 & -(\lambda + 2\mu) \frac{U}{h} \end{bmatrix}$$
(51)

Eqs. 47, 46, and 45 are obviously satisfied. $\underline{\text{div}\underline{\sigma}} \to \text{obviously}$ equal to 0 as $\underline{\underline{\sigma}}$ is constant.

Question 4: Find oedometric modulus \hat{E} as a function of E and ν .

Answer: The vertical force is equal to

$$\underline{P} = \int_{S_z^+} \underline{\underline{\sigma}} \cdot \underline{e}_z dS_3^+ \tag{52}$$

$$\underline{P} = \int_{S_z^+} -(\lambda + 2\mu) \frac{U}{h} dS_3^+ \underline{e}_z = -\underline{P}\underline{e}_z \quad \text{with} \quad \underline{P} = \frac{(\lambda + 2\mu)US}{h}$$
 (53)

So we can identify the oedometric modulus $\hat{E} = \lambda + 2\mu$.

$$\underline{\underline{\underline{\sigma}}} = \lambda \operatorname{tr} \underline{\underline{\varepsilon}} + 2\mu \underline{\underline{\varepsilon}} \qquad \operatorname{tr} \underline{\underline{\underline{\sigma}}} = (3\lambda + 2\mu) \operatorname{tr} \underline{\underline{\varepsilon}} \\
\underline{\underline{\varepsilon}} = \frac{1 + \nu}{E} \underline{\underline{\underline{\sigma}}} - \frac{\nu}{E} \operatorname{tr} \underline{\underline{\sigma}} \underline{\underline{\underline{1}}} \qquad \operatorname{tr} \underline{\underline{\underline{\varepsilon}}} = (\frac{1 - 2\nu}{E}) \operatorname{tr} \underline{\underline{\underline{\sigma}}} \qquad \to 3\lambda + 2\mu = \frac{E}{1 - 2\nu} \tag{54}$$

Or we can also write as follows:

$$\operatorname{tr}\underline{\underline{\sigma}} = (3\lambda + 2\mu)\operatorname{tr}\underline{\underline{\varepsilon}} \quad \to \quad \underline{\underline{\sigma}} = \frac{\lambda}{3\lambda + 2\mu}\operatorname{tr}\underline{\underline{\sigma}}\underline{\underline{1}} + 2\mu\underline{\underline{\varepsilon}}$$
 (55)

$$\underline{\underline{\varepsilon}} = \frac{1}{2\mu}\underline{\underline{\sigma}} - \frac{\lambda}{2\mu(3\lambda + 2\mu)} \text{tr}\underline{\underline{\sigma}1}$$
(56)

By identification,

$$\frac{1}{2\mu} = \frac{1+\nu}{E} \quad \text{and} \quad \frac{\lambda}{2\mu(3\lambda+2\mu)} = \frac{\nu}{E}$$
 (57)

$$\rightarrow \lambda = \frac{E\nu}{(1+\nu)(1-2\nu)} \quad \text{and} \quad \mu = \frac{E}{2(1+\nu)}$$
 (58)

Thus,

$$\hat{E} = \lambda + 2\mu = \frac{2E\nu + 2E(1 - 2\nu)}{2(1 + \nu)(1 - 2\nu)} = \frac{E(1 - \nu)}{(1 + \nu)(1 - 2\nu)}$$
(59)