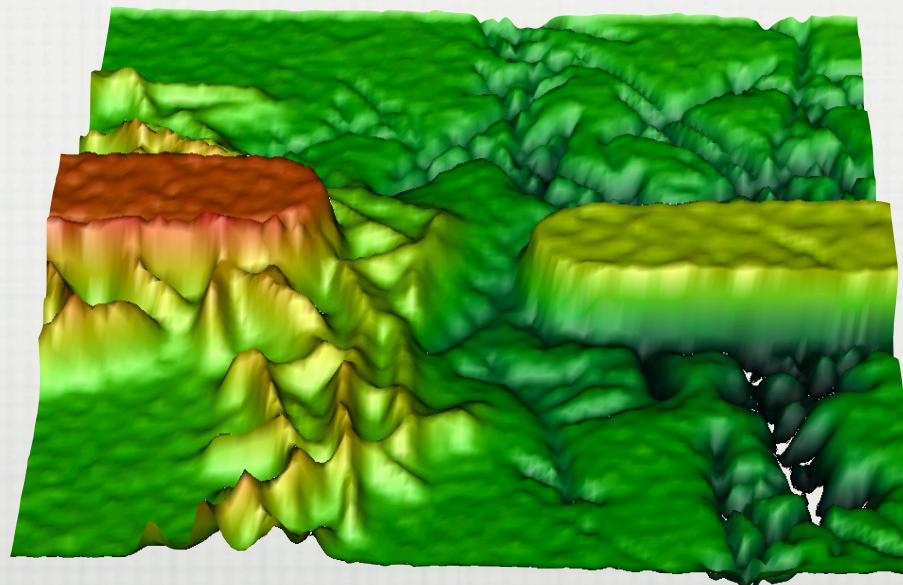


PHY-4810

COMPUTATIONAL PHYSICS

LECTURE 13B: MATRIX ALGEBRA, ADVANCED TOPICS



SUMMARY FROM LAST FRIDAY

$$\mathcal{A} \ x = b$$

- Gauss-Jordan: compute A^{-1} explicitly (N^3 process)
- LU decomposition: $N^3/3$ (does not yield A^{-1})
- Taking advantage of form of the matrix (e.g. Cholesky, etc)

$$\boxed{\begin{bmatrix} \alpha_{00} & 0 & 0 & 0 \\ \alpha_{10} & \alpha_{11} & 0 & 0 \\ \alpha_{20} & \alpha_{21} & \alpha_{22} & 0 \\ \alpha_{30} & \alpha_{31} & \alpha_{32} & \alpha_{33} \end{bmatrix} \cdot \begin{bmatrix} \beta_{00} & \beta_{01} & \beta_{02} & \beta_{03} \\ 0 & \beta_{11} & \beta_{12} & \beta_{13} \\ 0 & 0 & \beta_{22} & \beta_{23} \\ 0 & 0 & 0 & \beta_{33} \end{bmatrix} = \begin{bmatrix} a_{00} & a_{01} & a_{02} & a_{03} \\ a_{10} & a_{11} & a_{12} & a_{13} \\ a_{20} & a_{21} & a_{22} & a_{23} \\ a_{30} & a_{31} & a_{32} & a_{33} \end{bmatrix}}$$

ADVANCED TOPICS

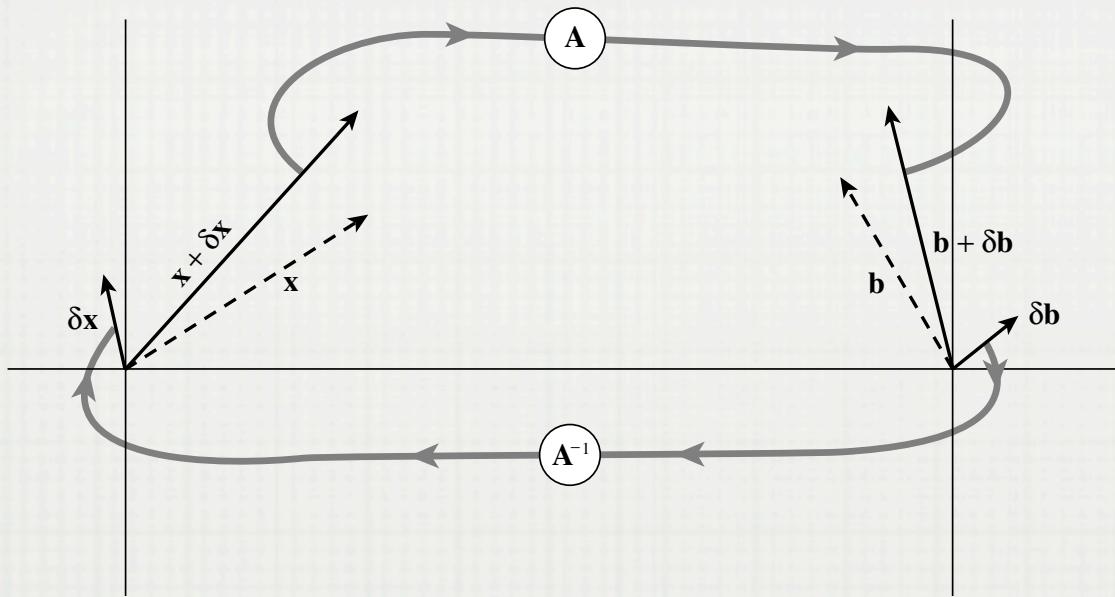
STI: ITERATIVE IMPROVEMENTS

- Suppose we know the solution up to some error: $x + \delta x$
- We can inject the solution in the initial problem to compute how far we are from b: $A . (x + \delta x) = b + \delta b$
- In fact we have: $A . \delta x = \delta b$
- We can now solve this equation to estimate δx :

$$A . \delta x = A . (x + \delta x) - b$$

- Once we know δx , we can then improve solution x
- (*See method mprove from ludcmp.h header file*)

IN PRACTICE (SEE LUDCMP.H FOR IMPLEMENTATION)

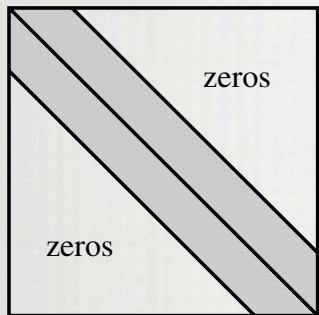


- Each call to `improve` realize one improvement, which is typically sufficient you can run it a few more times to confirm convergence it is a good idea to use this iterative improvement (**only a N^2 cost**)

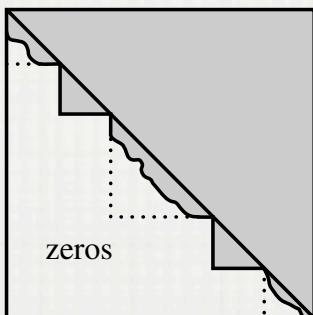
ST2: SINGULAR VALUE DECOMPOSITION

- When problems are too close to singularity
- SVD will **DIAGNOSE** the problem and even provide a **SOLUTION** to it
- SVD also solve linear **LEAST-SQUARE** problems (overdetermined problems)
- This is an advanced topic that could be covered in depth in a project

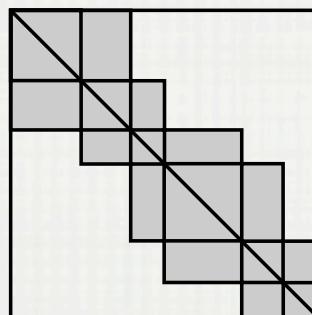
ST2: SPARSE MATRICES



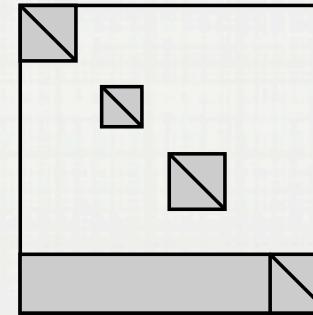
Band-diagonal



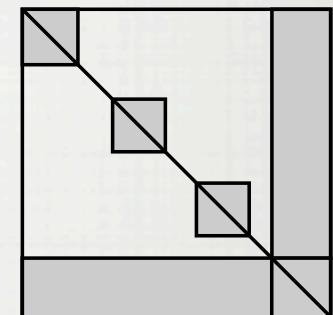
Block triangular



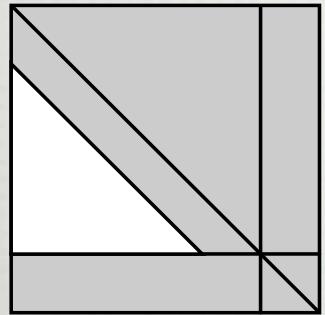
Block tridiagonal



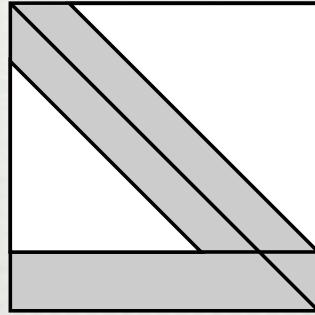
**singly bordered
block diagonal**



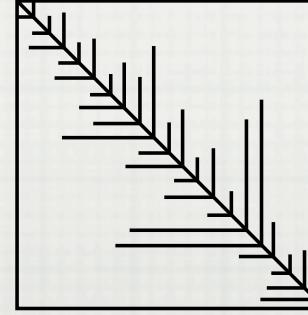
**doubly bordered
block diagonal**



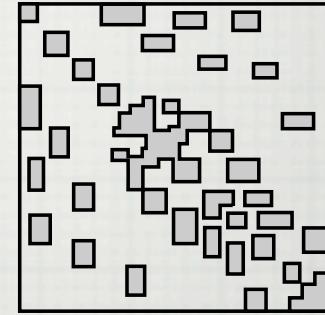
**bordered
band-triangular**



singly and doubly bordered band-diagonal



OTHERS



INDEXED STORAGE: COMPRESSED COLUMN STORAGE

3.0	0.0	1.0	2.0	0.0
0.0	4.0	0.0	0.0	0.0
0.0	7.0	5.0	9.0	0.0
0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	6.0	5.0

index k	0	1	2	3	4	5	6	7	8
val[k]	3.0	4.0	7.0	1.0	5.0	2.0	9.0	6.0	5.0
row_ind[k]	0	1	2	0	2	0	2	4	4

index i	0	1	2	3	4	5
col_ptr[i]	0	1	3	5	8	9

HARWELL-BOEING FORMAT

- 3 vectors are used
 - val : non zero-values
 - row_ind: row index
 - col_ptr: column index ***
 - col_ptr[0]=0;
 - col_ptr[ncol+1]=nsparse
 - col_ptr[i] =>
 - the first nonzero in column i is in col_ptr[i]; the last is at col_ptr[i+1]-1
 - [[answers the question: *How many non-zero elements have we encountered in previous column?*]]

$$\begin{bmatrix} 3.0 & 0.0 & 1.0 & 2.0 & 0.0 \\ 0.0 & 4.0 & 0.0 & 0.0 & 0.0 \\ 0.0 & 7.0 & 5.0 & 9.0 & 0.0 \\ 0.0 & 0.0 & 0.0 & 0.0 & 0.0 \\ 0.0 & 0.0 & 0.0 & 6.0 & 5.0 \end{bmatrix}$$

index k	0	1	2	3	4	5	6	7	8
val[k]	3.0	4.0	7.0	1.0	5.0	2.0	9.0	6.0	5.0
row_ind[k]	0	1	2	0	2	0	2	4	4

index i	0	1	2	3	4	5
col_ptr[i]	0	1	3	5	8	9

3.0	0.0	1.0	2.0	0.0
0.0	4.0	0.0	0.0	0.0
0.0	7.0	5.0	9.0	0.0
0.0	0.0	0.0	0.0	0.0
0.0	0.0	0.0	6.0	5.0

index k	0	1	2	3	4	5	6	7	8
val[k]	3.0	4.0	7.0	1.0	5.0	2.0	9.0	6.0	5.0
row_ind[k]	0	1	2	0	2	0	2	4	4

index i	0	1	2	3	4	5
col_ptr[i]	0	1	3	5	8	9

```

/* in this first example I explicitely define the variables
to be used later to build a sparse matrix */
NRsparseMat A(5,5,9);

//the following should not be needed since they have been
given above
A.nrows=5;
A.ncols=5;
A.nvals=9;

// column index (see lecture notes)
A.col_ptr[0]=0;
A.col_ptr[1]=1;
A.col_ptr[2]=3;
A.col_ptr[3]=5;
A.col_ptr[4]=8;
A.col_ptr[5]=9;

// rows where sparse elements will be located
A.row_ind[0]=0;
A.row_ind[1]=1;
A.row_ind[2]=2;
A.row_ind[3]=0;
A.row_ind[4]=2;
A.row_ind[5]=0;
A.row_ind[6]=2;
A.row_ind[7]=4;
A.row_ind[8]=3; //this is modified compared to NR book
par.2.7 since the matrix there is singular

//now the values themselves
A.val[0] = 3.;
A.val[1] = 4.;
A.val[2] = 7.;
A.val[3] = 1.;
A.val[4] = 5.;
A.val[5] = 2.;
A.val[6] = 9.;
A.val[7] = 6.;
A.val[8] = 5.;
```

```

#include <iostream>
#include "nr3.h"
#include "sort.h"
#include "sparse.h"

int main (int argc, char * const argv[]) {

    int nsparse=0;
    int i, j;
    double eps=0.00001;

    //example from lecture notes
    int n=5;
    MatDoub Full(n,n,0.0);

    Full[0][0]=3.0;
    Full[0][2]=1.0;
    Full[0][3]=2.0;
    Full[1][1]=4.0;
    Full[2][1]=7.0;
    Full[2][2]=5.0;
    Full[2][3]=9.0;
    Full[4][3]=6.0;
    Full[4][4]=5.0;

    //first we need to find out how many non-zero elements we have
    //we store that information in "nsparse"
    for (i= 0; i<n; i++){
        for (j = 0; j<n; j++){
            if(fabs((Full[i][j]))>eps) {
                nsparse++;
            }
        }
    }
    //now we can create a sparse object
    NRsparseMat Sparse(n,n,nsparse);

    nsparse=0;
    for (j = 0; j<n; j++){
        if(j>0) {
            Sparse.col_ptr[j+1]=Sparse.col_ptr[j];
        }
        else
        {
            Sparse.col_ptr[0]=0;
        }
        for (i = 0; i<n; i++){
            if(fabs(Full[i][j])>eps) {
                Sparse.row_ind[nsparse]=i;
                Sparse.val[nsparse]=Full[i][j];
                Sparse.col_ptr[j+1]++;
                nsparse++;
            }
        }
    }
}

```

```

COUT << "NSPARSE=" << NSPARSE << endl;
COUT << "I= \t";
FOR(I=0;I<NSPARSE;I++){
    COUT << I << " \t" ;
}
COUT << endl;
COUT << "VAL= \t";
FOR(I=0;I<NSPARSE;I++){
    COUT << SPARSE.VAL[I] << " \t";
}
COUT << endl;
COUT << "ROW= \t";
FOR(I=0;I<NSPARSE;I++){
    COUT << SPARSE.ROW_IND[I] << " \t";
}
COUT << endl;
COUT << endl;
COUT << "I= \t";
FOR(I=0;I<N+1;I++){
    COUT << I << " \t";
}
COUT << endl;
COUT << "COL= \t";
FOR(I=0;I<N+1;I++){
    COUT << SPARSE.COL_PTR[I] << " \t";
}
COUT << endl;

RETURN 0;
}

```

```

Vinces-MacBook-Pro:Debug vml2$ ./Resistors
nsparse=9
i=      0      1      2      3      4      5      6      7      8
val=    3      4      7      1      5      2      9      6      5
row=   0      1      2      0      2      0      2      4      4
i=      0      1      2      3      4      5
col=   0      1      3      5      8      9
Vinces-MacBook-Pro:Debug vml2$ 

```

IN PRACTICE: SPARSE.H

```
struct NRsparseMat
```

Sparse matrix data structure for compressed column storage.

```
{
```

```
    Int nrows;
```

Number of rows.

```
    Int ncols;
```

Number of columns.

```
    Int nnvals;
```

Maximum number of nonzeros.

```
    VecInt col_ptr;
```

Pointers to start of columns. Length is ncols+1.

```
    VecInt row_ind;
```

Row indices of nonzeros.

```
    VecDoub val;
```

Array of nonzero values.

```
    NRsparseMat();
```

Default constructor.

```
    NRsparseMat(Int m, Int n, Int nnvals);
```

Constructor. Initializes vector to zero.

```
    VecDoub ax(const VecDoub &x) const;
```

Multiply \mathbf{A} by a vector $x[0..ncols-1]$.

```
    VecDoub atx(const VecDoub &x) const;
```

Multiply \mathbf{A}^T by a vector $x[0..nrows-1]$.

```
    NRsparseMat transpose() const;
```

Form \mathbf{A}^T .

```
};
```

SPARSE: EXAMPLE (SEE SPARSE.CPP)

```
#include "sparse.h"

#include "linbcg.h"
#include "asolve.h"

NRsparseMat A(5,5,9);

//here we define A

NRsparseLinbcg C(A);

VecDoub x(5,1.0);
VecDoub b(5,0.0);

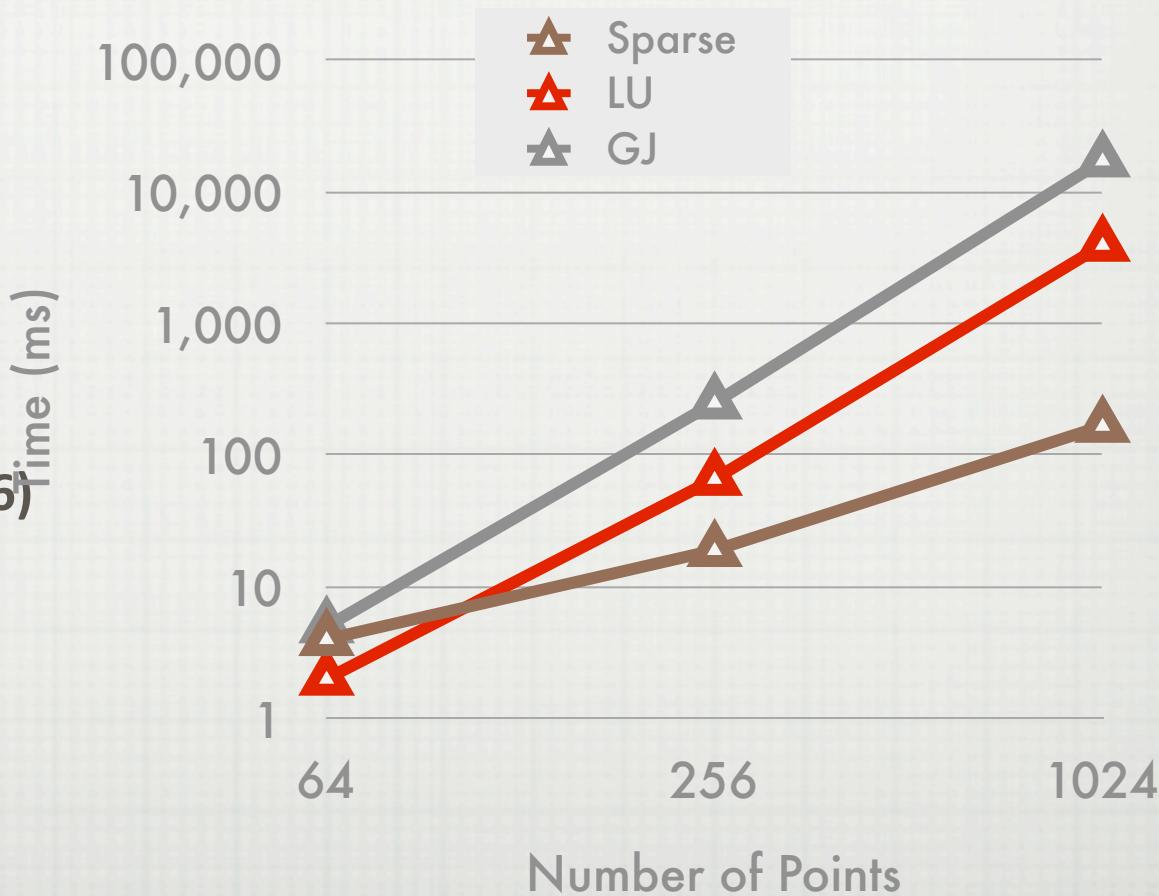
//here we define x and b

Int itol=1;
Doub tol=1.0E-72;;
Int itmax=10000;
Int iter;
Doub err;

C.solve(b,x,itol,tol,itmax,iter,err);
```

SCALING

- Gaussj -> N^3
- LU -> $N^3/3$
- (Cholesky -> $N^3/6$)
- Sparse -> N^2



ST4: IS MATRIX INVERSION AN N^3 PROCESS?

- How many multiplications do you need for the following problem?

$$\begin{pmatrix} a_{00} & a_{01} \\ a_{10} & a_{11} \end{pmatrix} \cdot \begin{pmatrix} b_{00} & b_{01} \\ b_{10} & b_{11} \end{pmatrix} = \begin{pmatrix} c_{00} & c_{01} \\ c_{10} & c_{11} \end{pmatrix}$$

- $c_{00} = a_{00} \times b_{00} + a_{01} \times b_{10}$
- $c_{01} = a_{00} \times b_{01} + a_{01} \times b_{11}$
- $c_{10} = a_{10} \times b_{00} + a_{11} \times b_{10}$
- $c_{11} = a_{10} \times b_{01} + a_{11} \times b_{11}$

- EIGHT, right?

COULD YOU DO IT IN 7 MULTIPLICATIONS?

$$Q_0 \equiv (a_{00} + a_{11}) \times (b_{00} + b_{11})$$

$$Q_1 \equiv (a_{10} + a_{11}) \times b_{00}$$

$$Q_2 \equiv a_{00} \times (b_{01} - b_{11})$$

$$Q_3 \equiv a_{11} \times (-b_{00} + b_{10})$$

$$c_{00} = Q_0 + Q_3 - Q_4 + Q_6$$

$$Q_4 \equiv (a_{00} + a_{01}) \times b_{11}$$

$$c_{10} = Q_1 + Q_3$$

$$Q_5 \equiv (-a_{00} + a_{10}) \times (b_{00} + b_{01})$$

$$c_{01} = Q_2 + Q_4$$

$$Q_6 \equiv (a_{01} - a_{11}) \times (b_{10} + b_{11})$$

$$c_{11} = Q_0 + Q_2 - Q_1 + Q_5$$

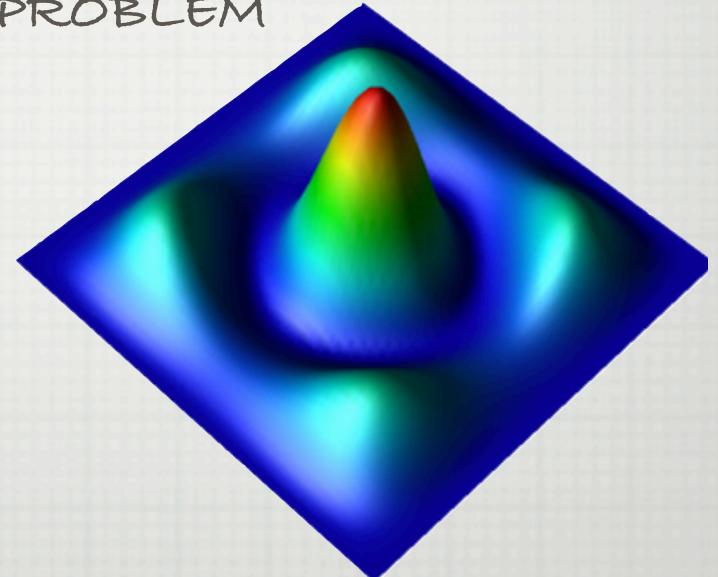
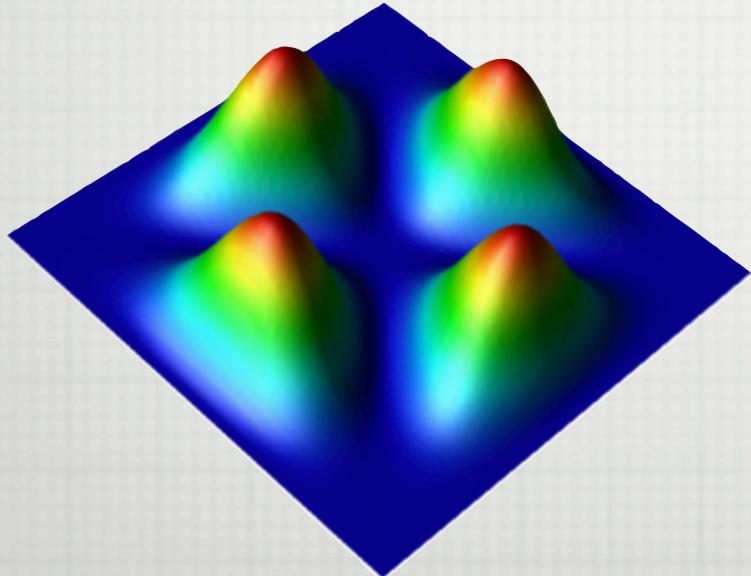
- What's the use of this? There is one fewer multiplications than in full equation, but many more additions and subtractions. It is not clear that anything has been gained.
- But notice that in the a's and b's are never commuted. Therefore these relationships are valid when the a's and b's are themselves matrices.

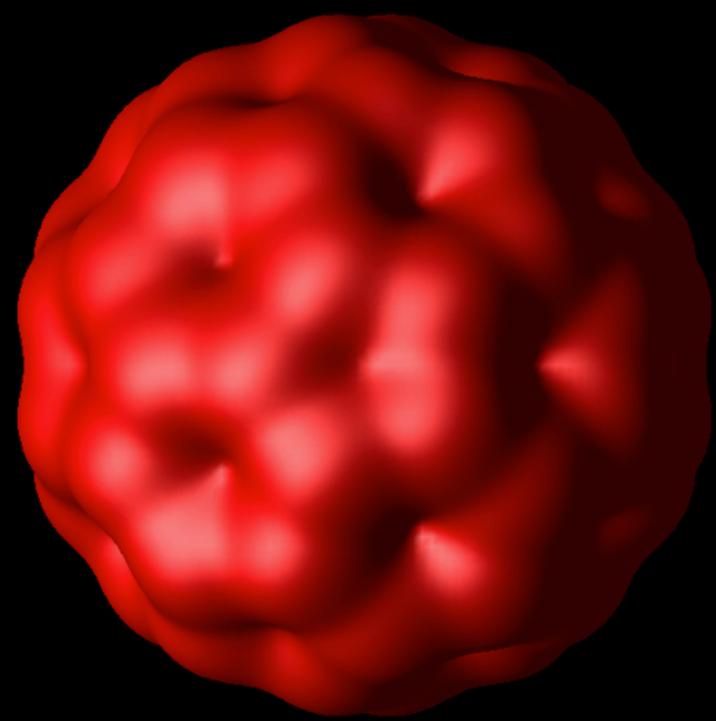
- The problem of multiplying two very large matrices (of order $N = 2^m$ for some integer m) can now be broken down recursively by partitioning the matrices into quarters, sixteenths, etc.
- And note the key point: The savings is not just a factor “7/8”; it is that factor at each hierarchical level of the recursion. In total it reduces the process of matrix multiplication to order $N^{\log_2(7)}$ instead of N^3 . (that's $N^{2.8}$)
- for $n=100$: 411,822 instead of 1,000,000 (41%)
- for $n=10000$: 169,044,093,164 instead of 1,000,000,000,000 (17%)
- What about the additions? For large N , it turns out that there are six times as many additions as multiplications. But, if N is very large, this constant factor is no match for the change in the exponent from N^3 to $N^{\log_2 7}$.
- This type of algorithm starts to show up in high level *blas* routines!

PHY-4810

COMPUTATIONAL PHYSICS

LECTURE 14: EIGENVALUE PROBLEM





EIGENVALUE, EIGENVECTORS?

- An eigenvalue and eigenvector of a square matrix A are a scalar λ and a nonzero vector x so that

$$Ax = \lambda x$$

- The eigenvalue-eigenvector equation for a square matrix can be written

$$(A - \lambda I)x = 0, \quad x \neq 0.$$

- The characteristic equation or characteristic polynomial of A.

$$\det(A - \lambda I) = 0.$$

IN PRACTICE

- Numerical recipes has a collection of methods to compute the eigenvalues/eigenvectors (chapter 11)
- Here we will employ the methods developed for symmetric matrices

```
struct Symmeig {
```

Computes all eigenvalues and eigenvectors of a real symmetric matrix by reduction to tridiagonal form followed by QL iteration.

eigen_sym.h

EXAMPLE: SCHROEDINGER EQUATION

- Time-independent S.E.:

$$-\frac{\hbar^2}{2m} \nabla^2 \psi(x, y, z) + V(x, y, z) \psi(x, y, z) = E \psi(x, y, z)$$

- Or, equivalently:

$$\left(-\frac{\hbar^2}{2m} \nabla^2 + V \right) \psi = E \psi$$

- Finally:

$$\left(\frac{\hbar^2}{2m} = 1 \right)$$

$$\boxed{(-\nabla^2 + V)\psi = E\psi}$$

SOLUTION: DISCRETIZATION OF ∇^2

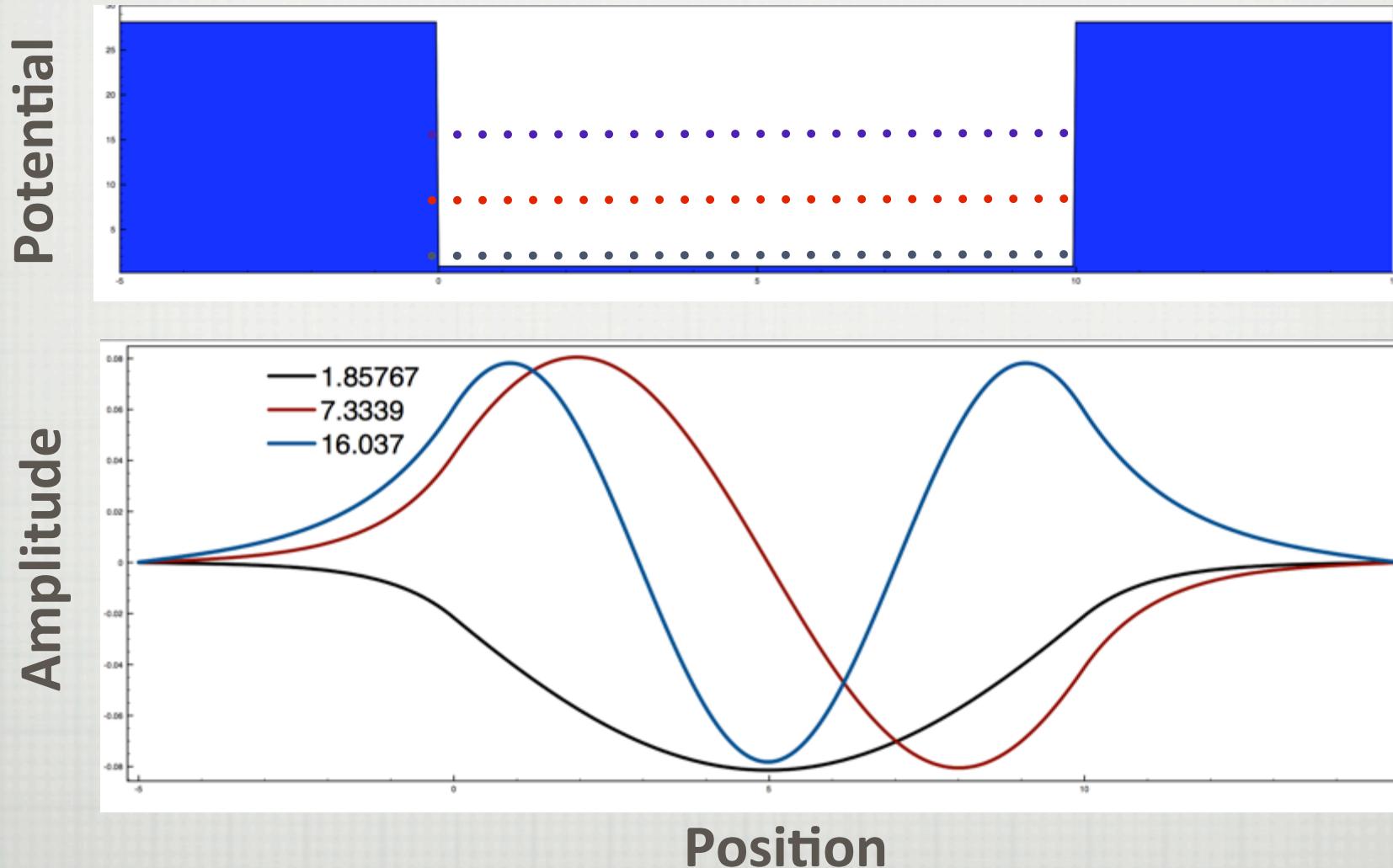
1D:
$$\nabla^2 f = \frac{f_{i-1} - 2f_i + f_{i+1}}{h^2}$$

$$(-\nabla^2 + V)f = \frac{-f_{i-1} + (2 + h^2 V_i) f_i - f_{i+1}}{h^2}$$

2D:
$$\nabla^2 f = \frac{f_{i-1,j} + f_{i+1,j} - 4f_{ij} + f_{i,j-1} + f_{i,j+1}}{h^2}$$

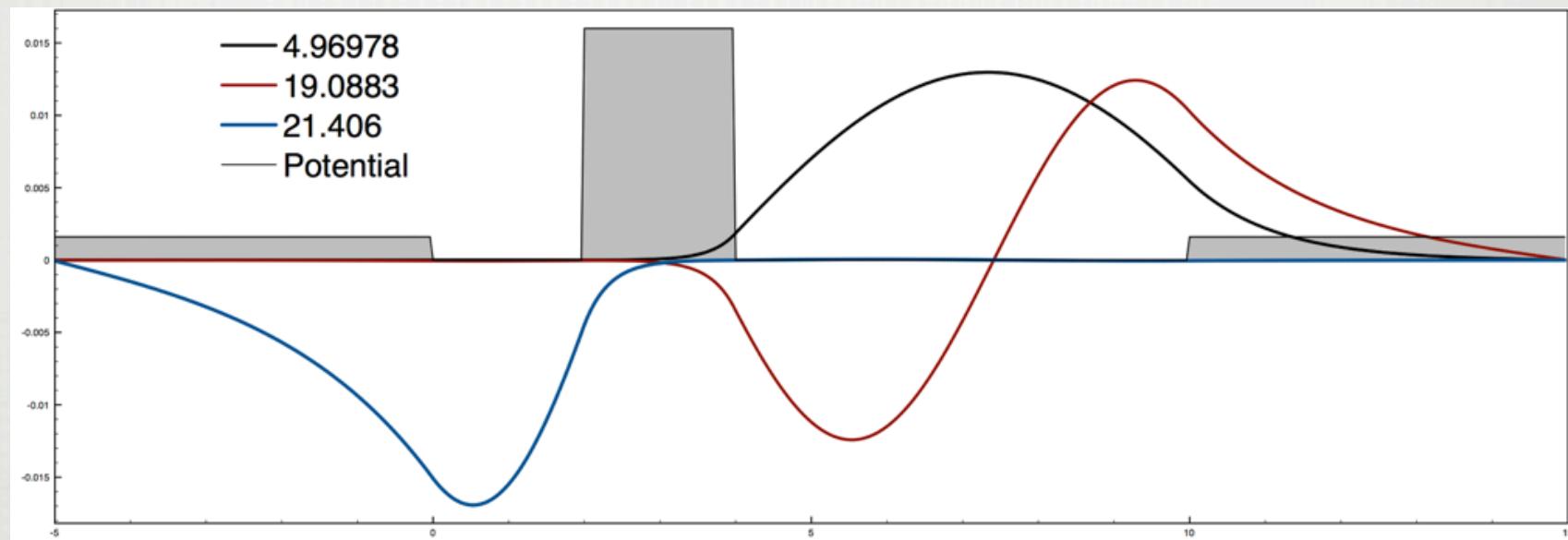
$$((-\nabla^2 + V)f)_{ij} = \frac{-f_{i-1,j} - f_{i+1,j} + (4 + h^2 V_{ij}) f_{ij} - f_{i,j-1} - f_{i,j+1}}{h^2}$$

PARTICLE IN A 1D WELL (BOUND STATES)



PARTICLE IN A MULTIPLE-WELL

Amplitude



Position

HOW?

```
//define Hamiltonian
for(int i=0;i<n;i++){
    Ham[i][i]=2.0+Pot[i];
    if(i<n-1) Ham[i][i+1]=-1.;
    if(i>0) Ham[i][i-1]=-1.;
}

Symmeig Schrodinger(Ham);

for(int i=n-10;i<n;i++){
    cout << "i= " << i << " Eigenvalue (eV)=" << hartree*Schrodinger.d[i]/d2 << endl;
}
cout << endl;

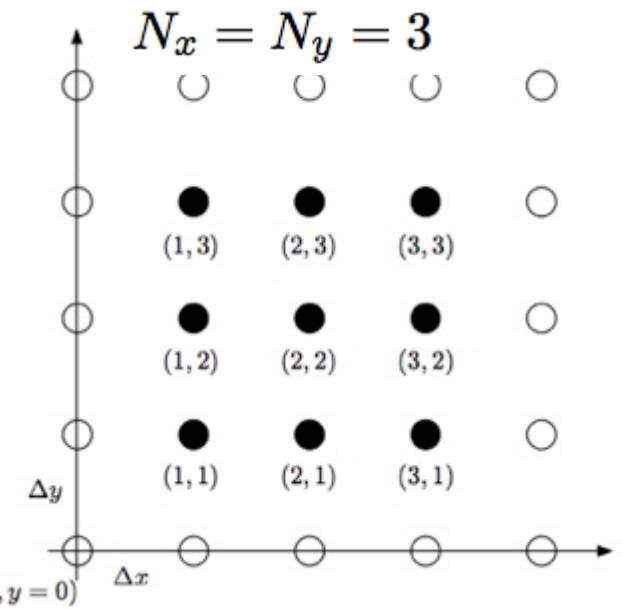
ofstream outfile;
outfile.open("wavefunction.dat");
for(int i=0;i<n;i++){
    outfile << (i-n1)*delta << " ";
    for(int j=0;j<10;j++){
        outfile<< Schrodinger.z[i][n-j-1] << " ";
    }
    outfile << endl;
}
outfile.close();
```

DISCRETIZED S.E. IN 2D

$$[(-\nabla^2 + V)f]_{ij} = \frac{-f_{i-1,j} - f_{i+1,j} + (4 + h^2 V_{ij})f_{ij} - f_{i,j-1} - f_{i,j+1}}{h^2}$$

REMEMBER LAST FRIDAY!

ONE WAY TO BUILD THE MATRIX



$$\partial^2 x \equiv \begin{pmatrix} -2 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & -2 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & -2 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -2 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & -2 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & -2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -2 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & -2 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & -2 & 0 \end{pmatrix}$$

$$\partial^2 y \equiv \begin{pmatrix} -2 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & -2 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & -2 & 0 & 0 & 1 & 0 & 0 & 0 \\ 1 & 0 & 0 & -2 & 0 & 0 & 1 & 0 & 0 \\ 0 & 1 & 0 & 0 & -2 & 0 & 0 & 1 & 0 \\ 0 & 0 & 1 & 0 & 0 & -2 & 0 & 0 & 1 \\ 0 & 0 & 0 & 1 & 0 & 0 & -2 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & -2 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & -2 \end{pmatrix}$$

$$\frac{\partial^2 f(x_n, y_m)}{\partial x^2} \sim \frac{f(x_n - \Delta x, y_m) - 2f(x_n, y_m) + f(x_n + \Delta x, y_m)}{\Delta x^2}$$

PUTTING IT ALL TOGETHER (∇^2)

$$\nabla^2 \equiv \begin{pmatrix} -4 & 1 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 1 & -4 & 1 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & -4 & 0 & 0 & 1 & 0 & 0 & 0 \\ 1 & 0 & 0 & -4 & 1 & 0 & 1 & 0 & 0 \\ 0 & 1 & 0 & 1 & -4 & 1 & 0 & 1 & 0 \\ 0 & 0 & 1 & 0 & 1 & -4 & 0 & 0 & 1 \\ 0 & 0 & 0 & 1 & 0 & 0 & -4 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 1 & -4 & 1 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 1 & -4 \end{pmatrix}$$

$$\nabla^2 \equiv \begin{pmatrix} -4 & 1 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 1 & -4 & 1 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & -4 & 0 & 0 & 1 & 0 & 0 & 0 \\ 1 & 0 & 0 & -4 & 1 & 0 & 1 & 0 & 0 \\ 0 & 1 & 0 & 1 & -4 & 1 & 0 & 1 & 0 \\ 0 & 0 & 1 & 0 & 1 & -4 & 0 & 0 & 1 \\ 0 & 0 & 0 & 1 & 0 & 0 & -4 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 1 & -4 & 1 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 1 & -4 \end{pmatrix}$$

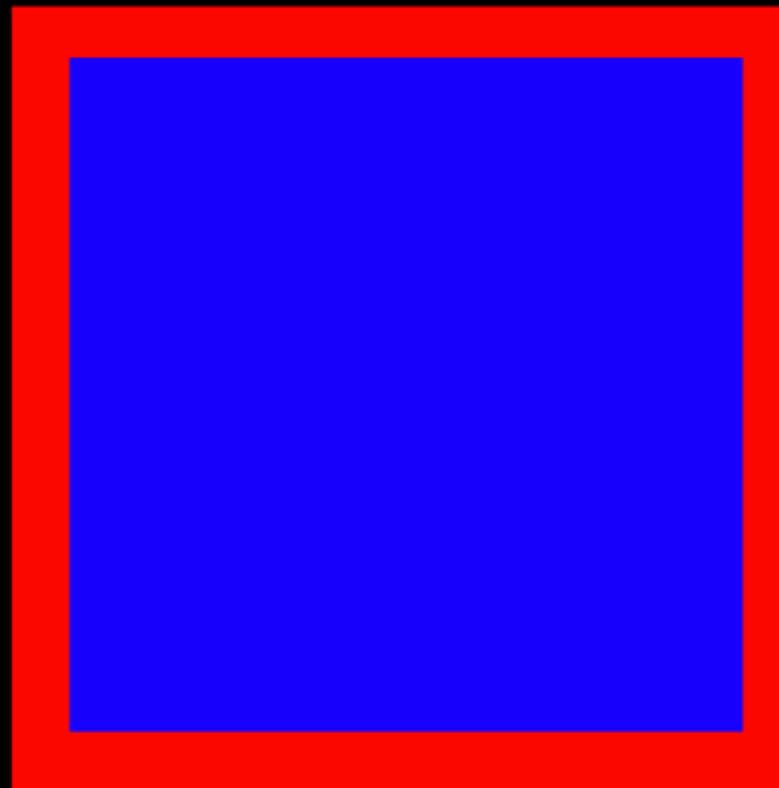
$$\begin{matrix} -4 & 1 & 0 \\ 1 & -4 & 1 \\ 0 & 1 & -4 \end{matrix}$$

$n_y \times (n_x \times n_x)$ blocks along the diagonal

blocks are connected by a diagonal of “1”

500 x 500

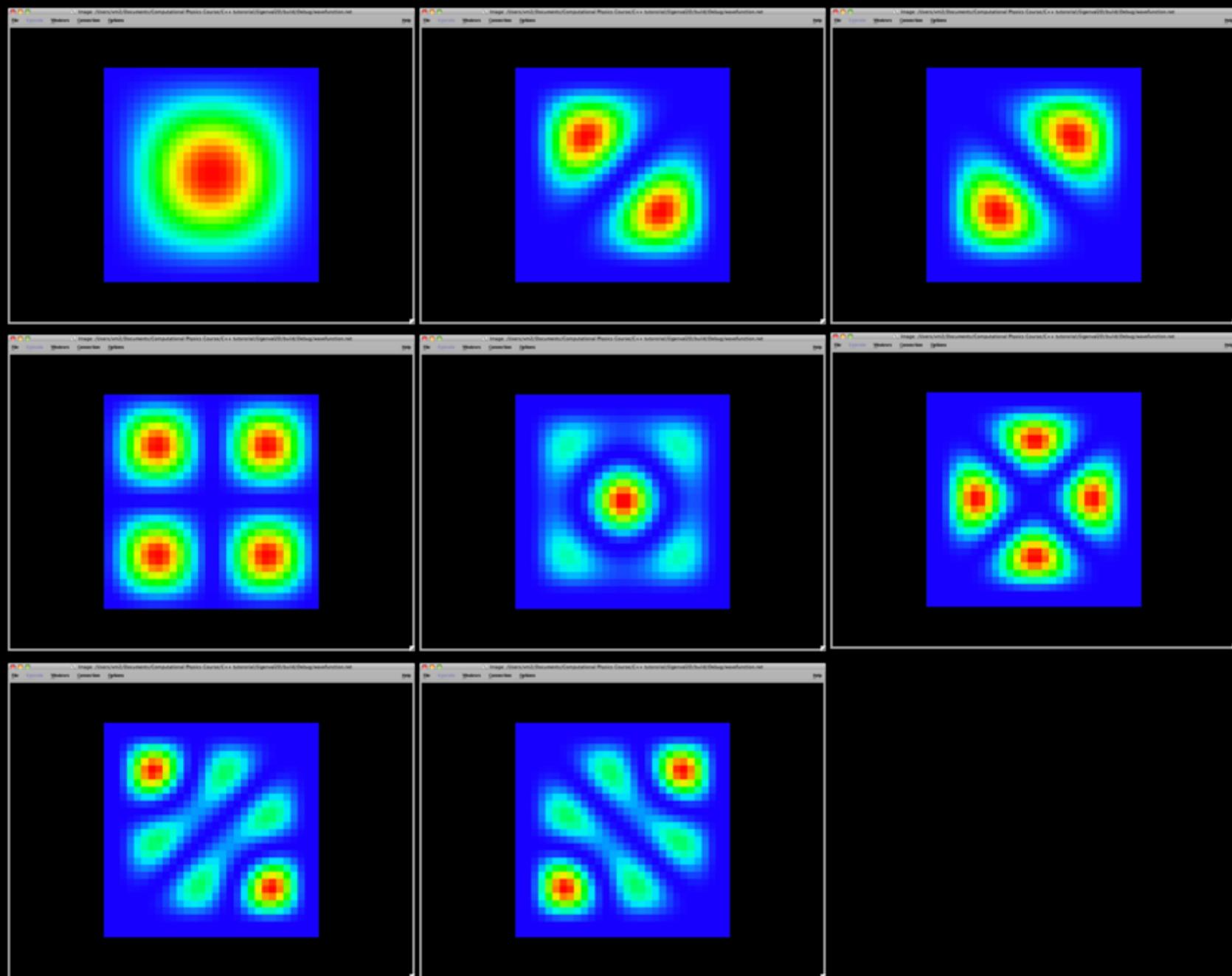
Square well



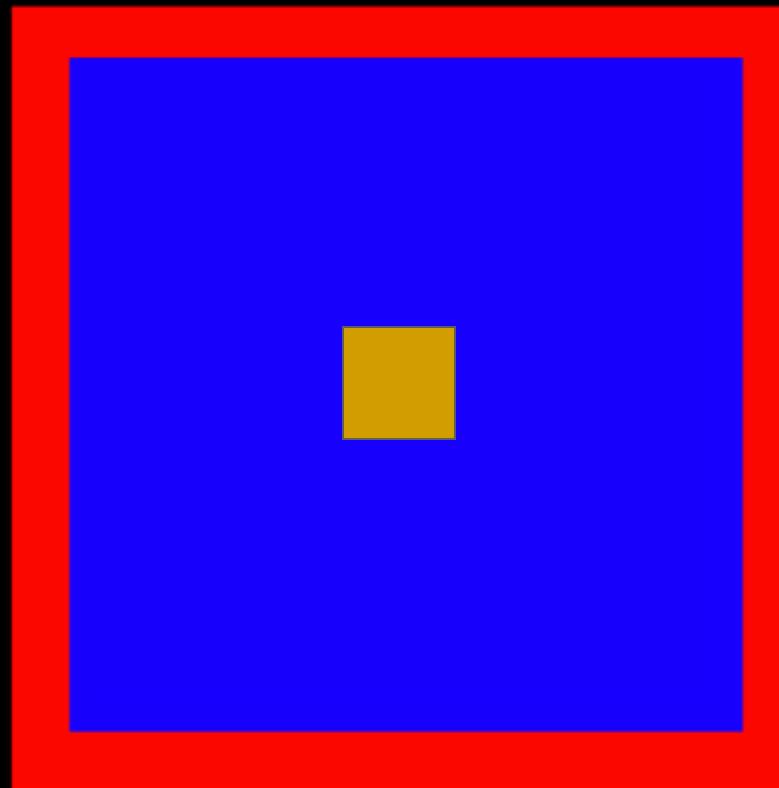
$V=0$

$V=100$

Square well: wavefunctions



Square well with a central pole

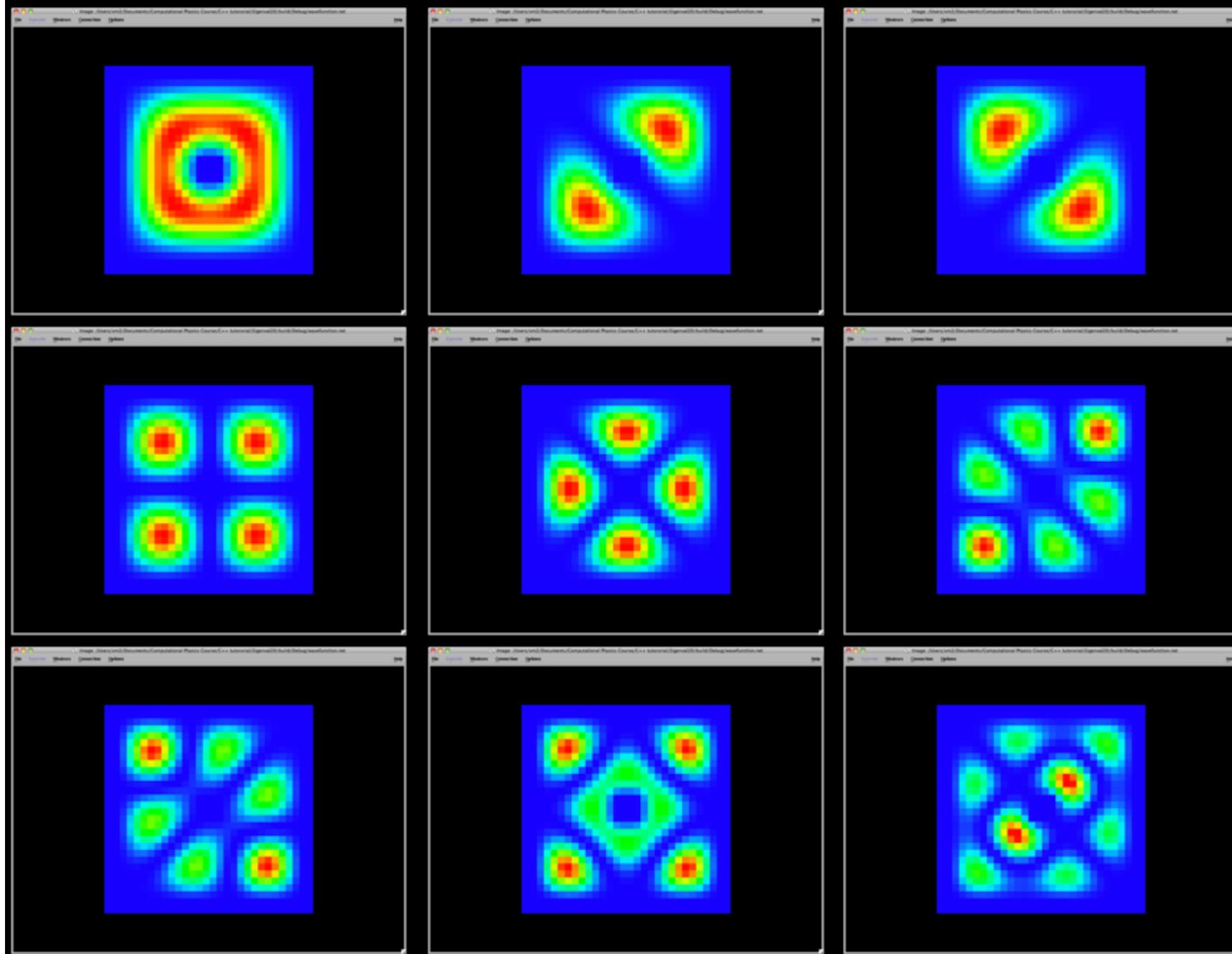


$V=0$

$V=100$

$V=10000$

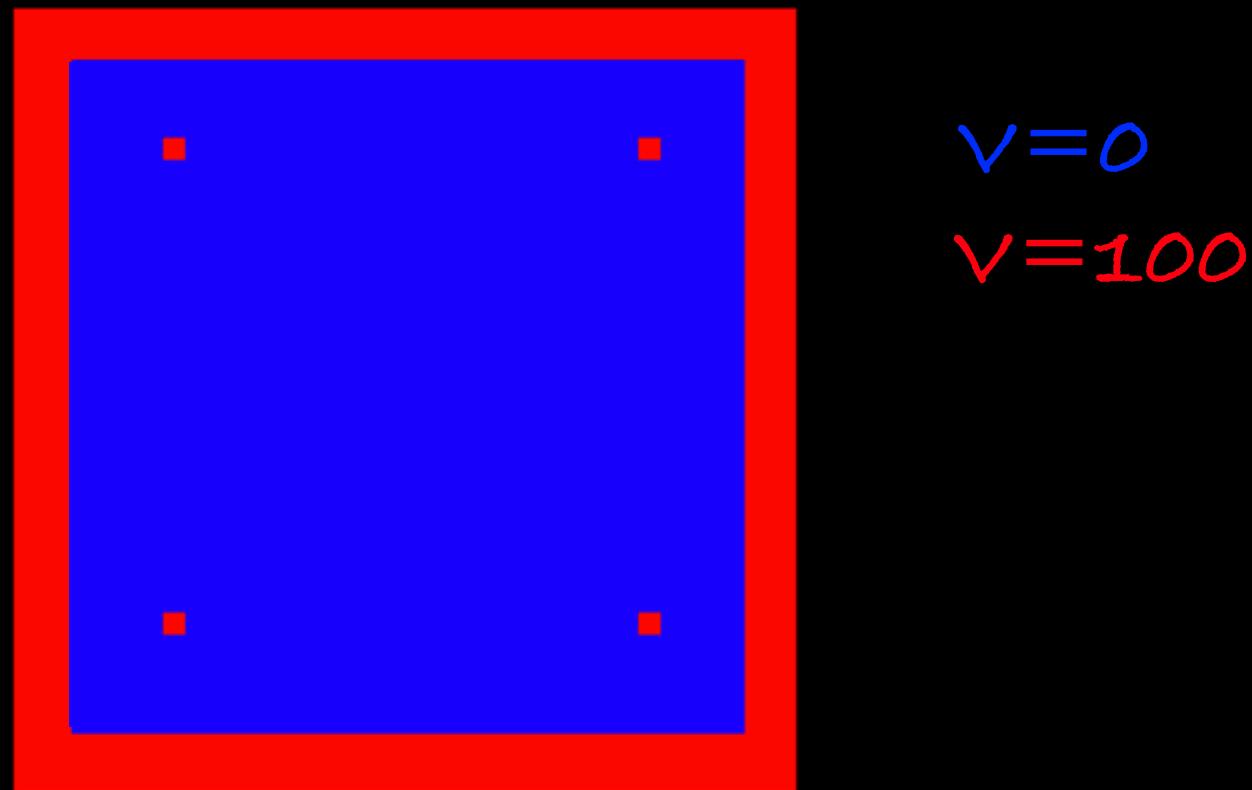
Square well with a pole: wavefunctions

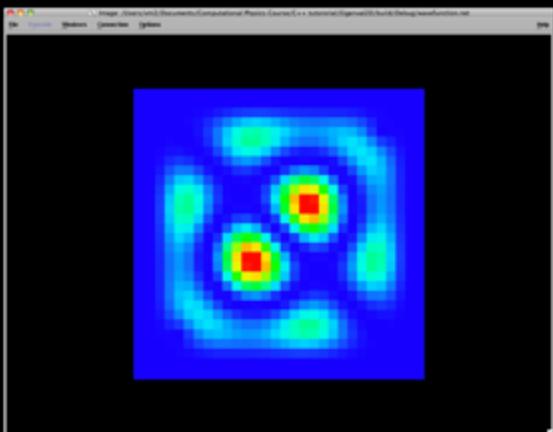
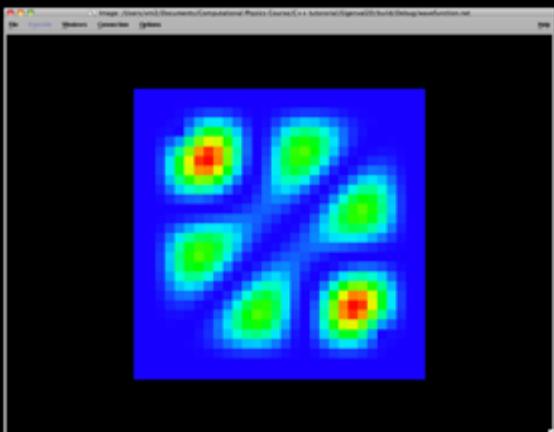
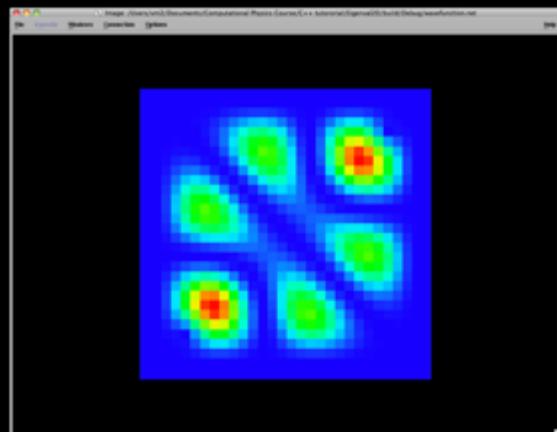
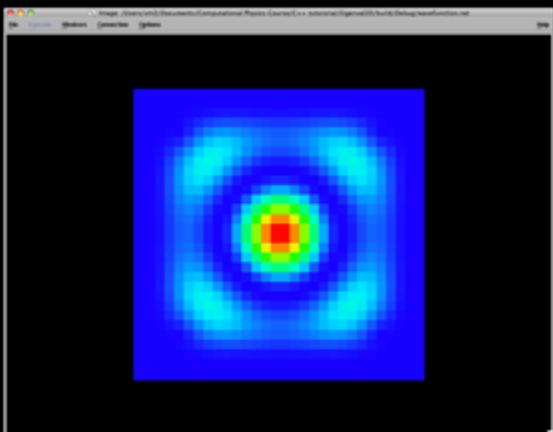
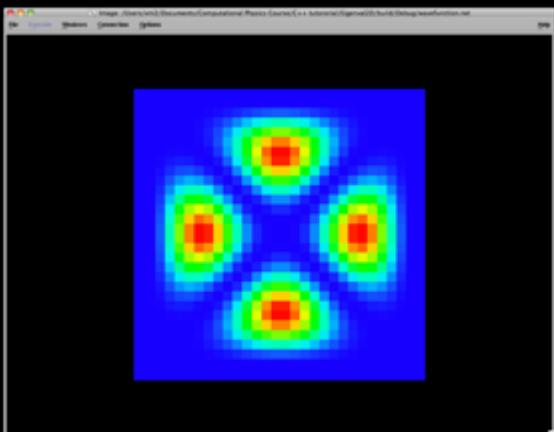
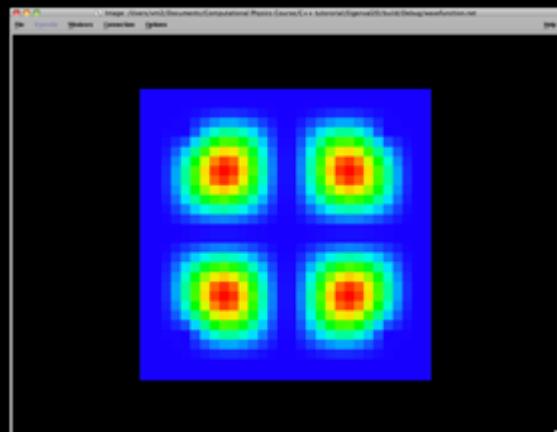
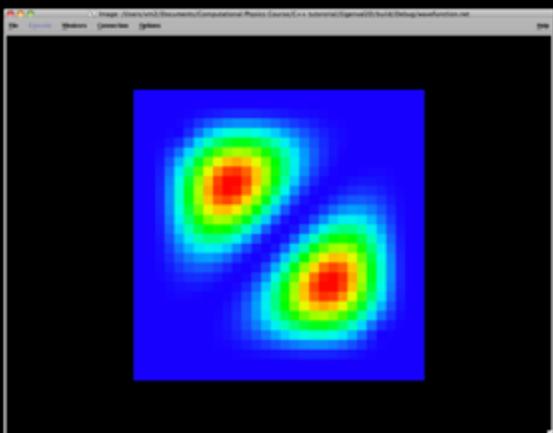
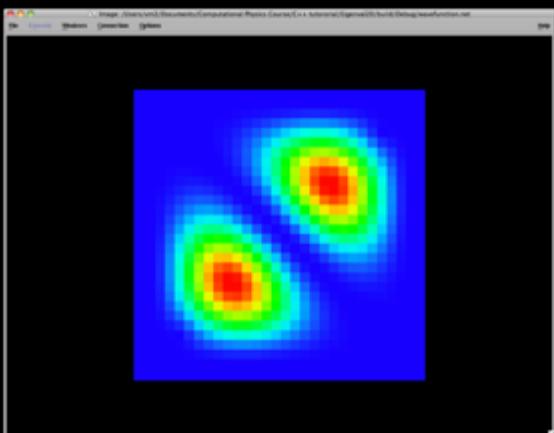
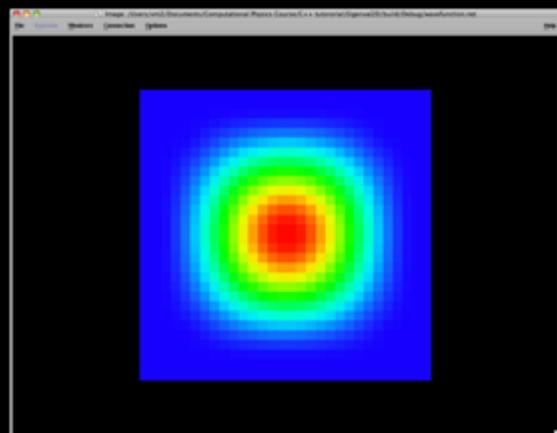


E_l

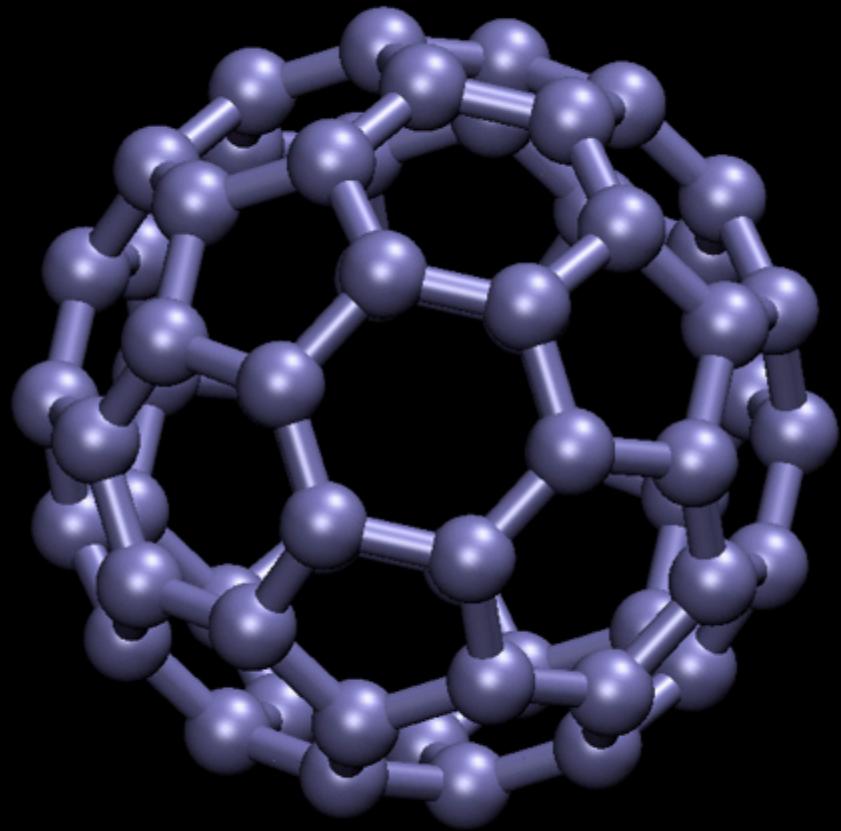
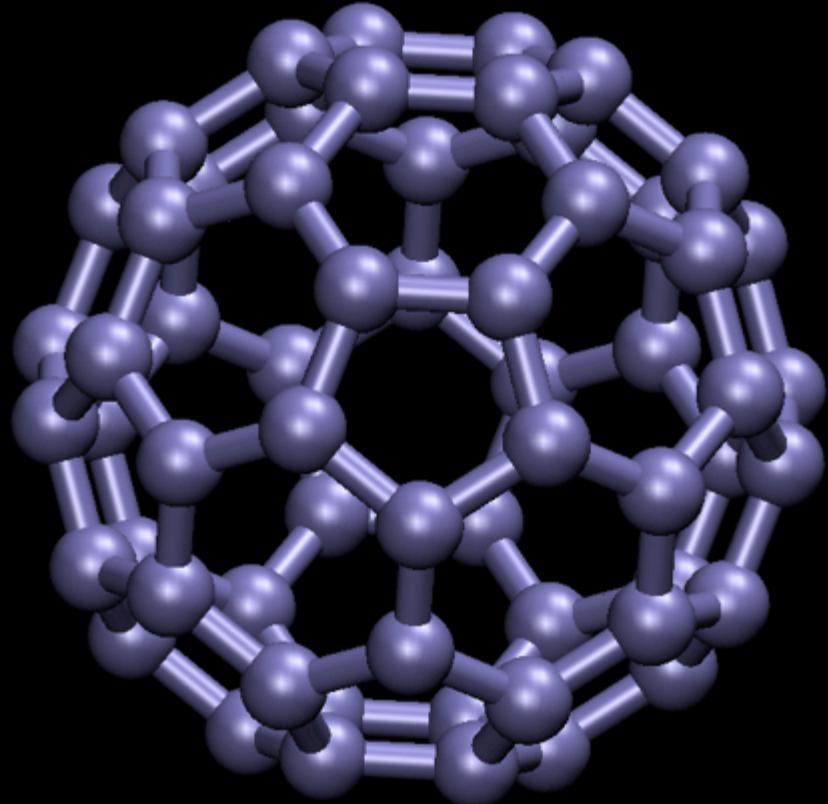
0.0548
0.0733
0.0733
0.108
0.134
0.178
0.178
0.196
0.246

Square well with perturbations





How Schroedinger Equation is solved in
quantum chemistry and computational
nanoscience!



INTRODUCTION TO LCAO METHOD

$$H|\alpha\rangle = E_\alpha|\alpha\rangle$$

$$|\alpha\rangle = \sum_a c_a |\phi_a\rangle$$

We use atomic orbitals as basis: we suppose that the molecule can be approximated by addition of properties of the materials

$$\sum_a |\phi_a\rangle\langle\phi_a| = 1$$

Properties of completeness and orthogonality

$$\langle\phi_a|\phi'_a\rangle = \delta_{aa'}$$

$$|\alpha\rangle = \sum_b |\phi_b\rangle\langle\phi_b|\alpha\rangle \quad c_a = \langle\phi_a|\alpha\rangle$$

$$H|\alpha\rangle = E_\alpha |\alpha\rangle \quad |\alpha\rangle = \sum_a c_a |\phi_a\rangle$$

$$\sum_a c_a H |\phi_a\rangle = E_\alpha \sum_a c_a |\phi_a\rangle$$

$$\langle \phi_{a'} | \sum_a c_a H |\phi_a\rangle = E_\alpha \sum_a c_a \langle \phi_{a'} | \phi_a \rangle$$

$$\sum_a c_a \langle \phi_{a'} | H |\phi_a\rangle = E_\alpha \sum_a c_a \langle \phi_{a'} | \phi_a \rangle$$

$$\sum_a c_a H_{aa'} = E_\alpha \sum_a c_a \delta_{aa'}$$

$$\sum_a H_{a'a} c_a = E_\alpha c'_a \quad \mathcal{H}C = EC$$

HUCKEL THEORY (TIGHT-BINDING)

- All interactions are limited to neighboring atoms, and are identical

$$H_{aa'} = \langle \phi_a | H | \phi'_{a'} \rangle = \begin{cases} 0 \\ \gamma \end{cases}$$

- The overlap between atoms is zero

$$\langle \phi_a | \phi'_{a'} \rangle = \delta_{aa'}$$



IN PRACTICE

```
for(int i=0;i<60;i++){
    for(int j=0;j<60;j++){
        distance=sqrt(pow(r[i][0]-r[j][0],2)+pow(r[i][1]-r[j]
[1],2)+pow(r[i][2]-r[j][2],2));
        if(distance<rcut && distance > 0.1) {
            H[i][j]=gamma0;
        }
    }
}

Symmeig Schrodinger(H);
for(int i=1;i<nat;i++){
    cout << nat-i << " " << Schrodinger.d[i] << endl;
}
```

PERIODIC TABLE OF THE ELEMENTS

<http://www.ktf-split.hr/periodni/en/>

GROUP	1 IA	RELATIVE ATOMIC MASS (1)												18 VIIIA				
PERIOD	1 H HYDROGEN	2 IIA BERYLLIUM	GROUP IUPAC 13 IIIA 5 10.811												2 He HELIUM			
	3 LITHIUM	4 BORON	GROUP CAS												10 4.0026			
1	1 1.0079	2 9.0122	3 6.941	4 9.0122	5 10.811	6 12.011	7 14.007	8 15.999	9 18.998	10 20.180	11 22.990	12 24.305	13 BORON	14 CARBON	15 NITROGEN	16 OXYGEN	17 FLUORINE	18 NEON
2	Li LITHIUM	Be BERYLLIUM	Symbol	BORON	Element Name	Aluminium	Silicon	Phosphorus	Sulphur	Chlorine	Na SODIUM	Mg MAGNESIUM	13 10.811	14 12.011	15 14.007	16 15.999	17 18.998	18 20.180
3	19 39.098	20 40.078	21 44.956	22 47.867	23 50.942	24 51.996	25 54.938	26 55.845	27 58.933	28 58.693	29 63.546	30 65.39	31 69.723	32 72.64	33 74.922	34 78.96	35 79.904	36 83.80
4	K POTASSIUM	Ca CALCIUM	Sc SCANDIUM	Ti TITANIUM	V VANADIUM	Cr CHROMIUM	Mn MANGANESE	Fe IRON	Co COBALT	Ni NICKEL	Cu COPPER	Zn ZINC	Ga GALLIUM	Ge GERMANIUM	As ARSENIC	Se SELENIUM	Br BROMINE	Kr KRYPTON
5	37 85.468	38 87.62	39 88.906	40 91.224	41 92.906	42 95.94	43 (98)	44 101.07	45 102.91	46 106.42	47 107.87	48 112.41	49 114.82	50 118.71	51 121.76	52 127.60	53 126.90	54 131.29
6	Rb RUBIDIUM	Sr STRONTIUM	Y YTTRIUM	Zr ZIRCONIUM	Nb NIOBIUM	Mo MOLOBDENUM	Tc TECHNETIUM	Ru RUTHENIUM	Rh RHODIUM	Pd PALLADIUM	Ag SILVER	Cd CADMIUM	In INDIUM	Sn TIN	Sb ANTIMONY	Te TELLURIUM	I IODINE	Xe XENON
7	55 132.91	56 137.33	57-71 La-Lu Lanthanide	72 178.49	73 180.95	74 183.84	75 186.21	76 190.23	77 192.22	78 195.08	79 196.97	80 200.59	81 204.38	82 207.2	83 208.98	84 (209)	85 (210)	86 (222)
7	Cs CAESIUM	Ba BARIUM	Hf HAFNIUM	Ta TANTALUM	W TUNGSTEN	Re RHENIUM	Os OSMIUM	Ir IRIDIUM	Pt PLATINUM	Au GOLD	Hg MERCURY	Tl THALLIUM	Pb LEAD	Bi BISMUTH	Po POLONIUM	At ASTATINE	Rn RADON	
	87 (223)	88 (226)	89-103 Ac-Lr Actinide	104 (261)	105 (262)	106 (266)	107 (264)	108 (277)	109 (268)	110 (281)	111 (272)	112 (285)	114 (289)	115 (283)	116 (287)	117 (286)	118 (284)	119 (285)

(1) Pure Appl. Chem., 73, No. 4, 667-683 (2001)
Relative atomic mass is shown with five significant figures. For elements have no stable nuclides, the value enclosed in brackets indicates the mass number of the longest-lived isotope of the element.

However three such elements (Th, Pa, and U) do have a characteristic terrestrial isotopic composition, and for these an atomic weight is tabulated.

Editor: Aditya Verdhan (adivar@netlinx.com)

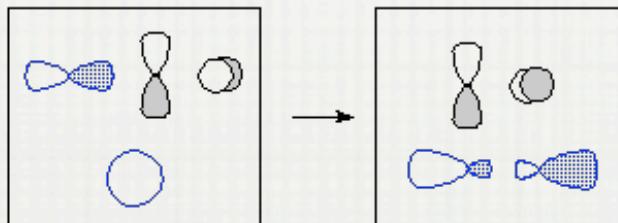
LANTHANIDE

57 138.91	58 140.12	59 140.91	60 144.24	61 (145)	62 150.36	63 151.96	64 157.25	65 158.93	66 162.50	67 164.93	68 167.26	69 168.93	70 173.04	71 174.97
La LANTHANUM	Ce CERIUM	Pr PRASEODYMIUM	Nd NEODYMIUM	Pm PROMETHIUM	Sm SAMARIUM	Eu EUROPIUM	Gd GADOLINIUM	Tb TERBIUM	Dy DYSPROSIUM	Ho HOLMIUM	Er ERBIUM	Tm THULIUM	Yb YTTERBIUM	Lu LUTETIUM

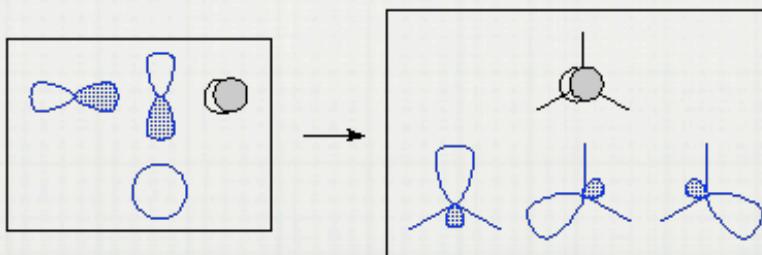
ACTINIDE

89 (227)	90 232.04	91 231.04	92 238.03	93 (237)	94 (244)	95 (243)	96 (247)	97 (247)	98 (251)	99 (252)	100 (257)	101 (258)	102 (259)	103 (262)
Ac ACTINIUM	Th THORIUM	Pa PROTACTINIUM	U URANIUM	Np NEPTUNIUM	Pu PLUTONIUM	Am AMERICIUM	Cm CURIUM	Bk BERKELIUM	Cf CALIFORNIUM	Einsteinium	Fermium	Mendelevium	No Nobelium	Lr Lawrencium

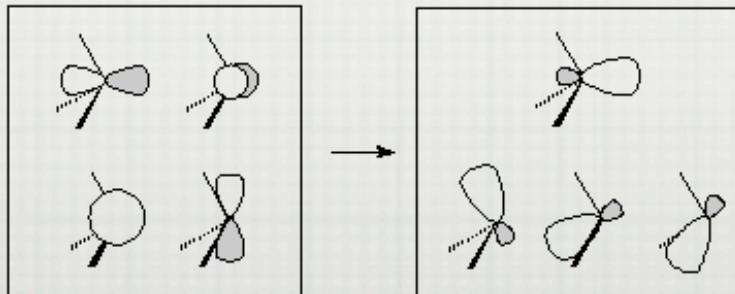
CARBON SP₂



sp hybrid
orbitals

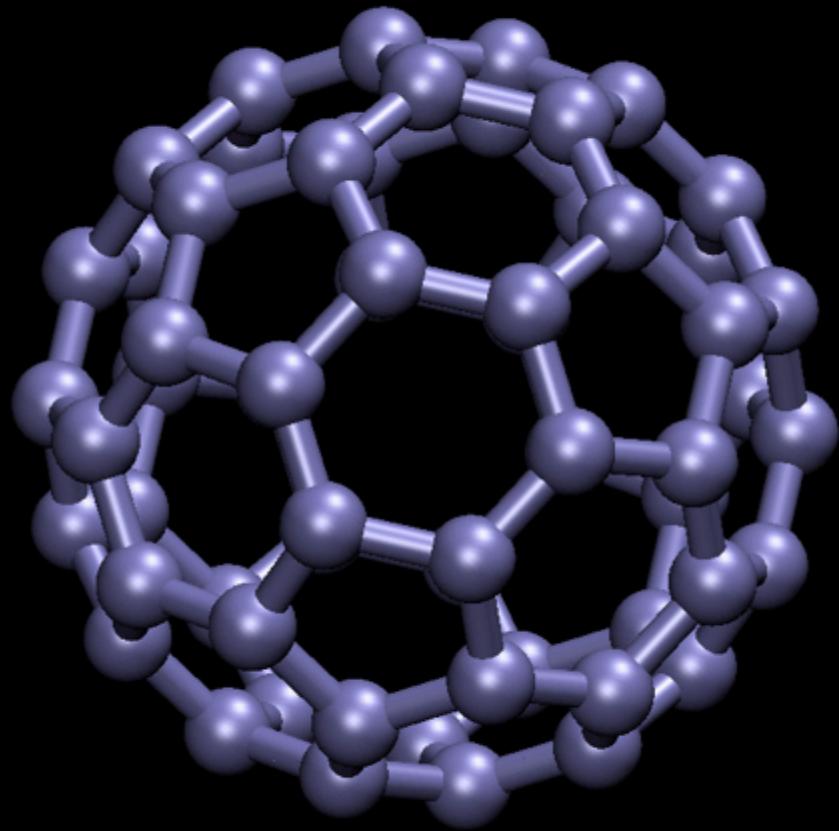
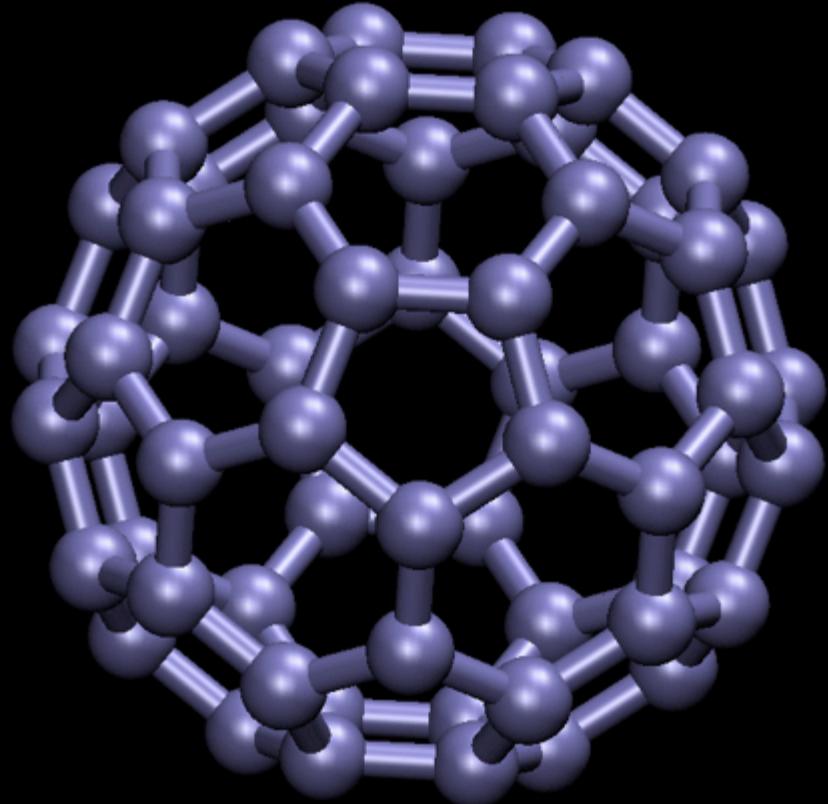


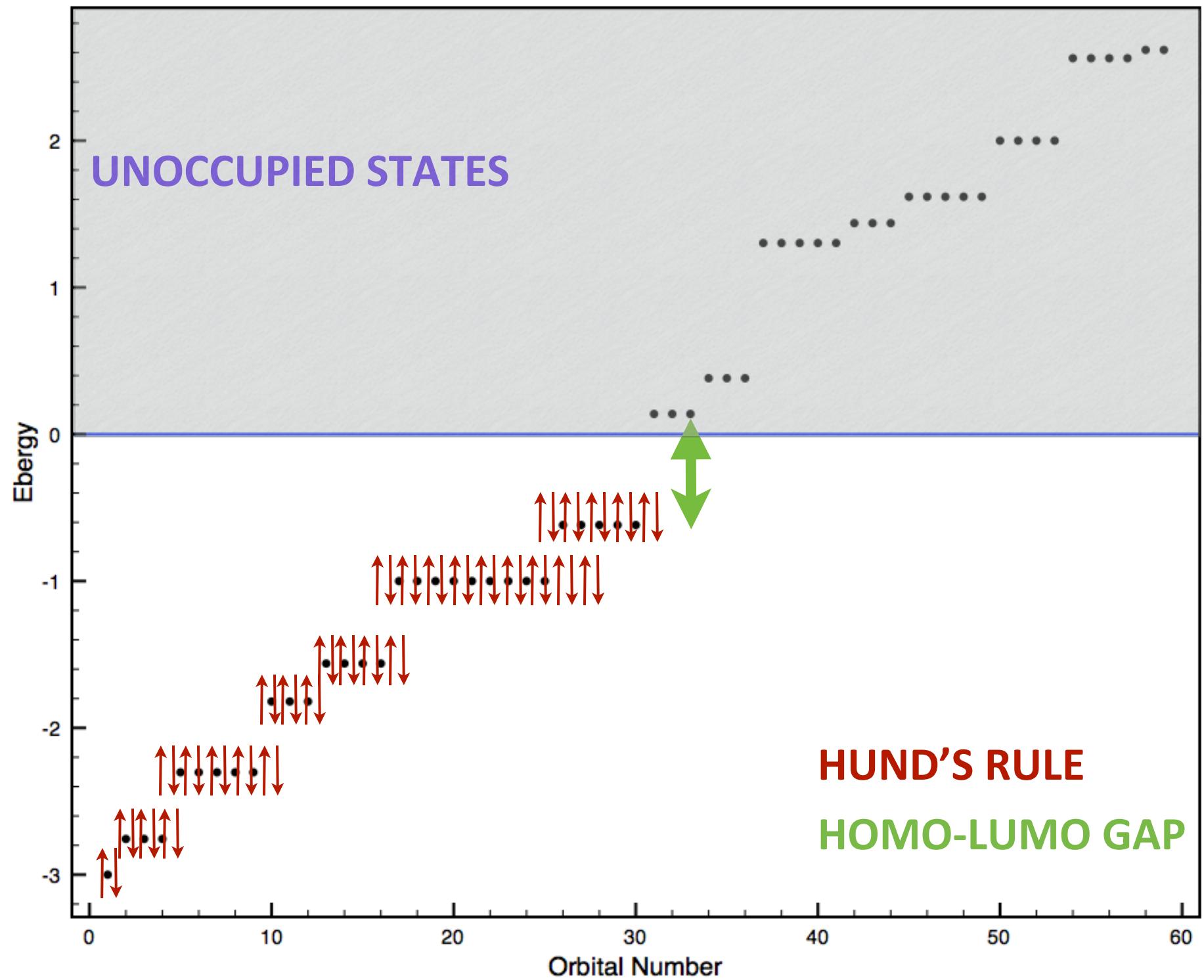
sp^2 hybrid
orbitals



sp^3 hybrid
orbitals



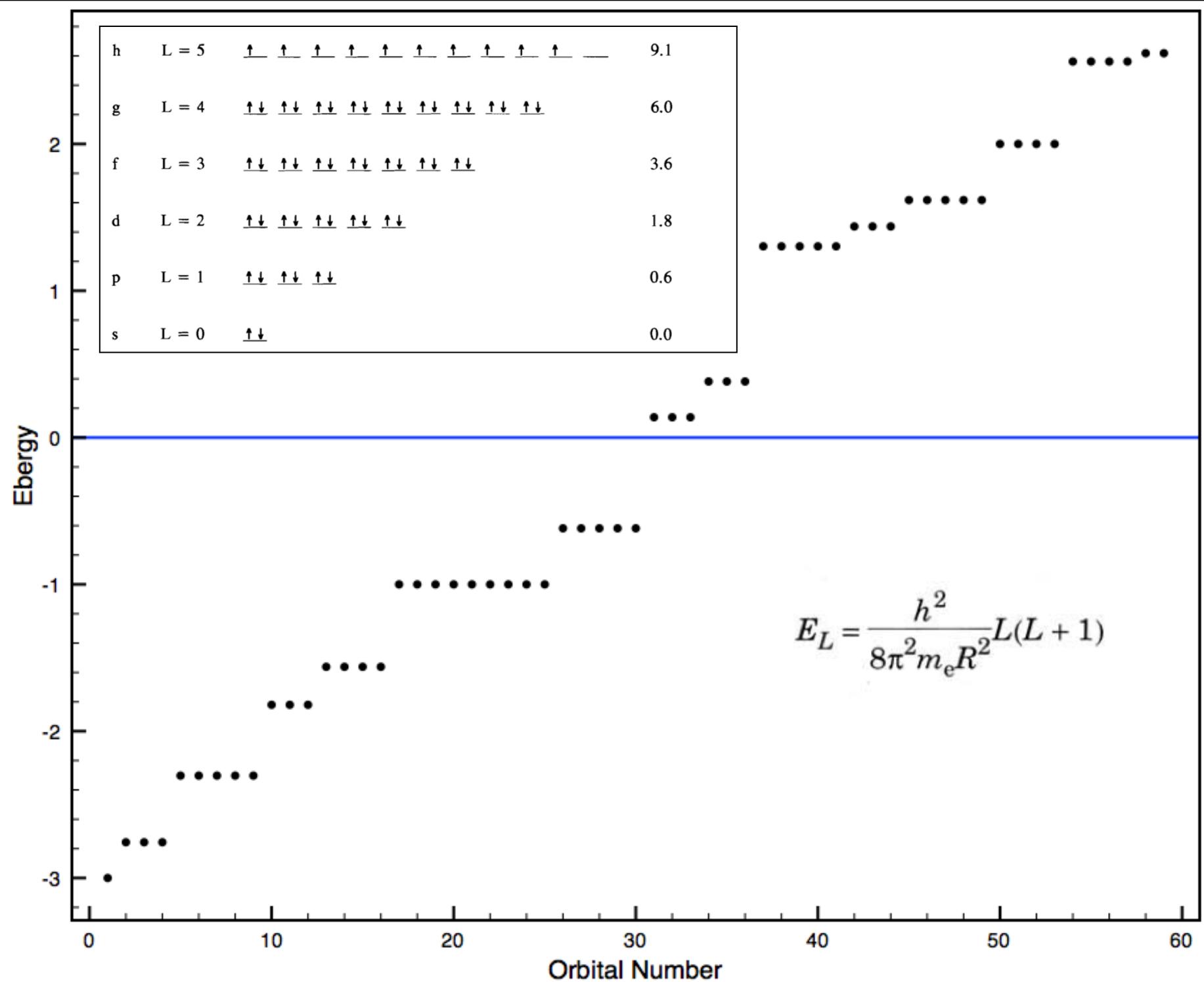




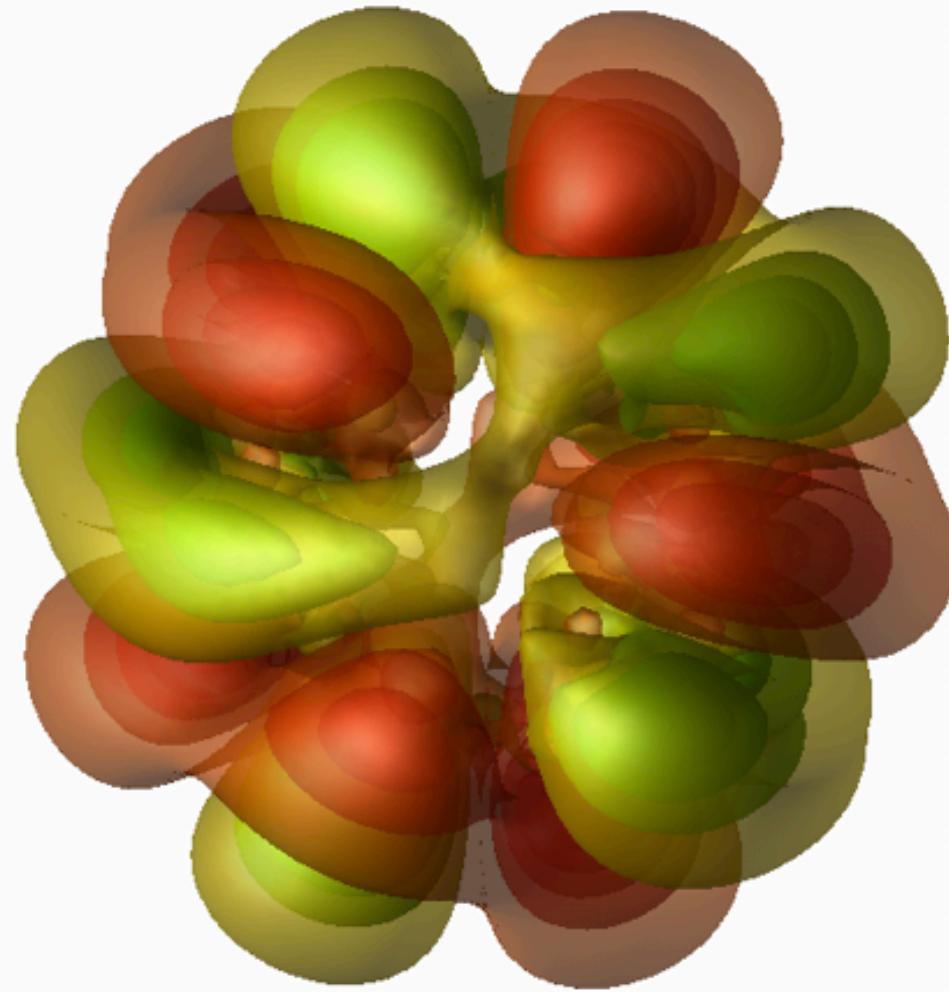
SPHERICAL SHELL MODEL

h	L = 5	$\frac{\uparrow}{\downarrow} \frac{\uparrow}{\downarrow} \frac{\uparrow}{\downarrow} \frac{\uparrow}{\downarrow} \frac{\uparrow}{\downarrow} \frac{\uparrow}{\downarrow} \frac{\uparrow}{\downarrow} \frac{\uparrow}{\downarrow} \frac{\uparrow}{\downarrow} \frac{\uparrow}{\downarrow}$	9.1
g	L = 4	$\frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow}$	6.0
f	L = 3	$\frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow}$	3.6
d	L = 2	$\frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow}$	1.8
p	L = 1	$\frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow} \frac{\uparrow\downarrow}{\uparrow\downarrow}$	0.6
s	L = 0	$\frac{\uparrow\downarrow}{\uparrow\downarrow}$	0.0

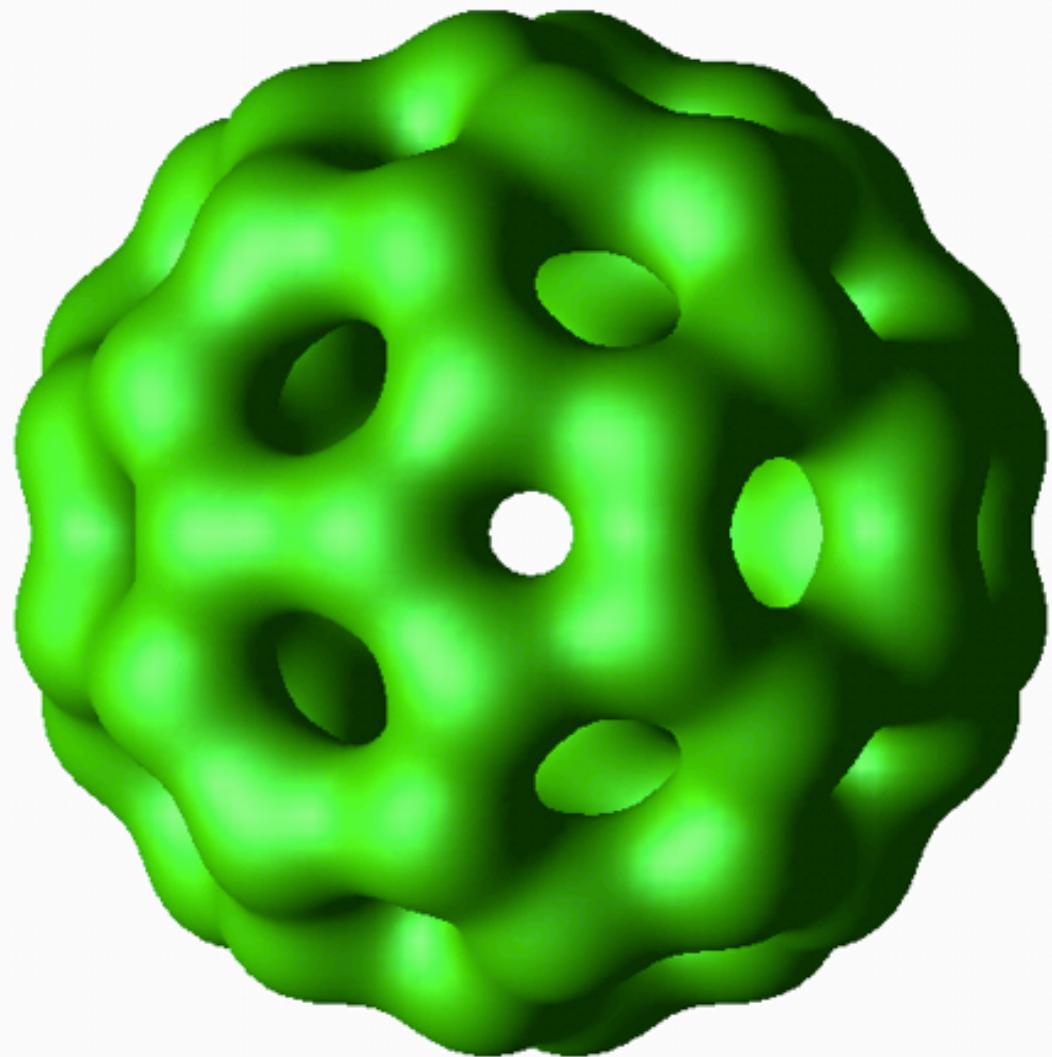
$$E_L = \frac{\hbar^2}{8\pi^2 m_e R^2} L(L+1)$$

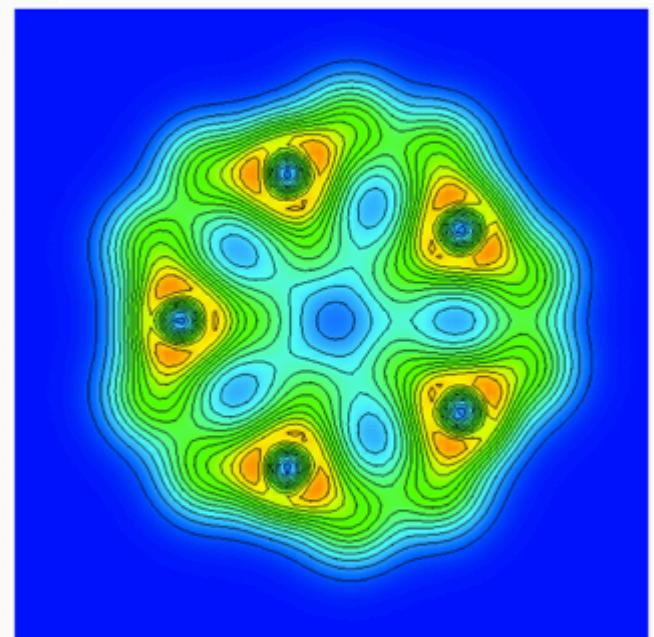
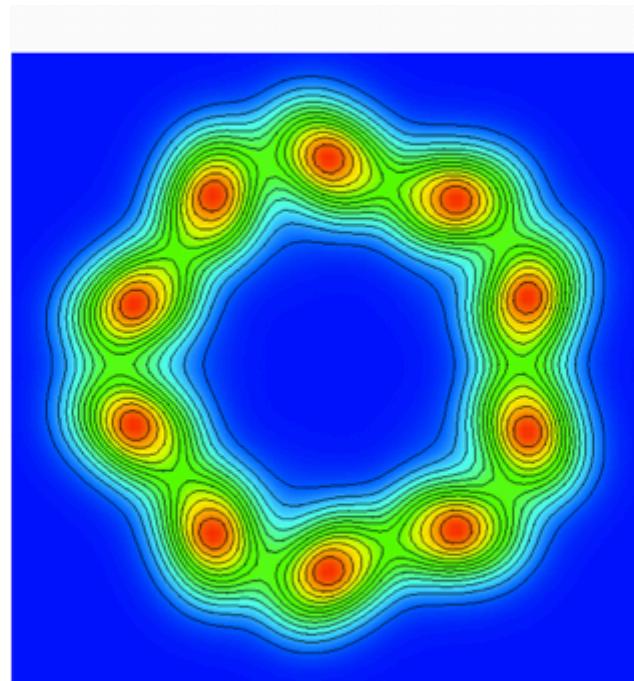
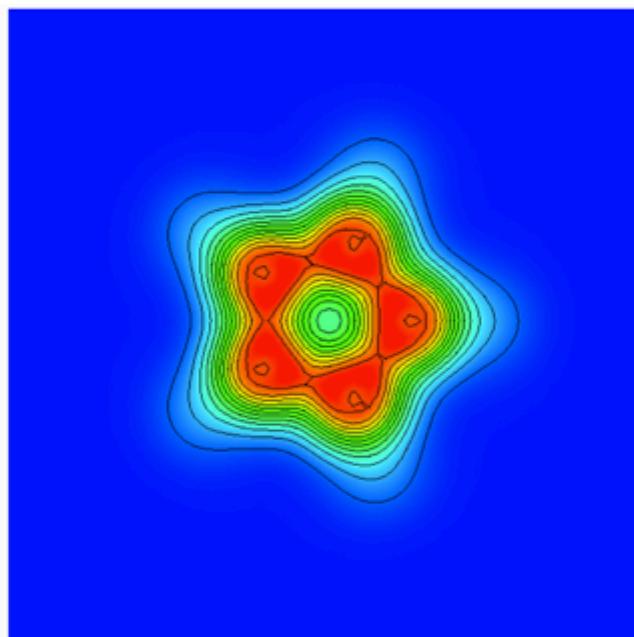


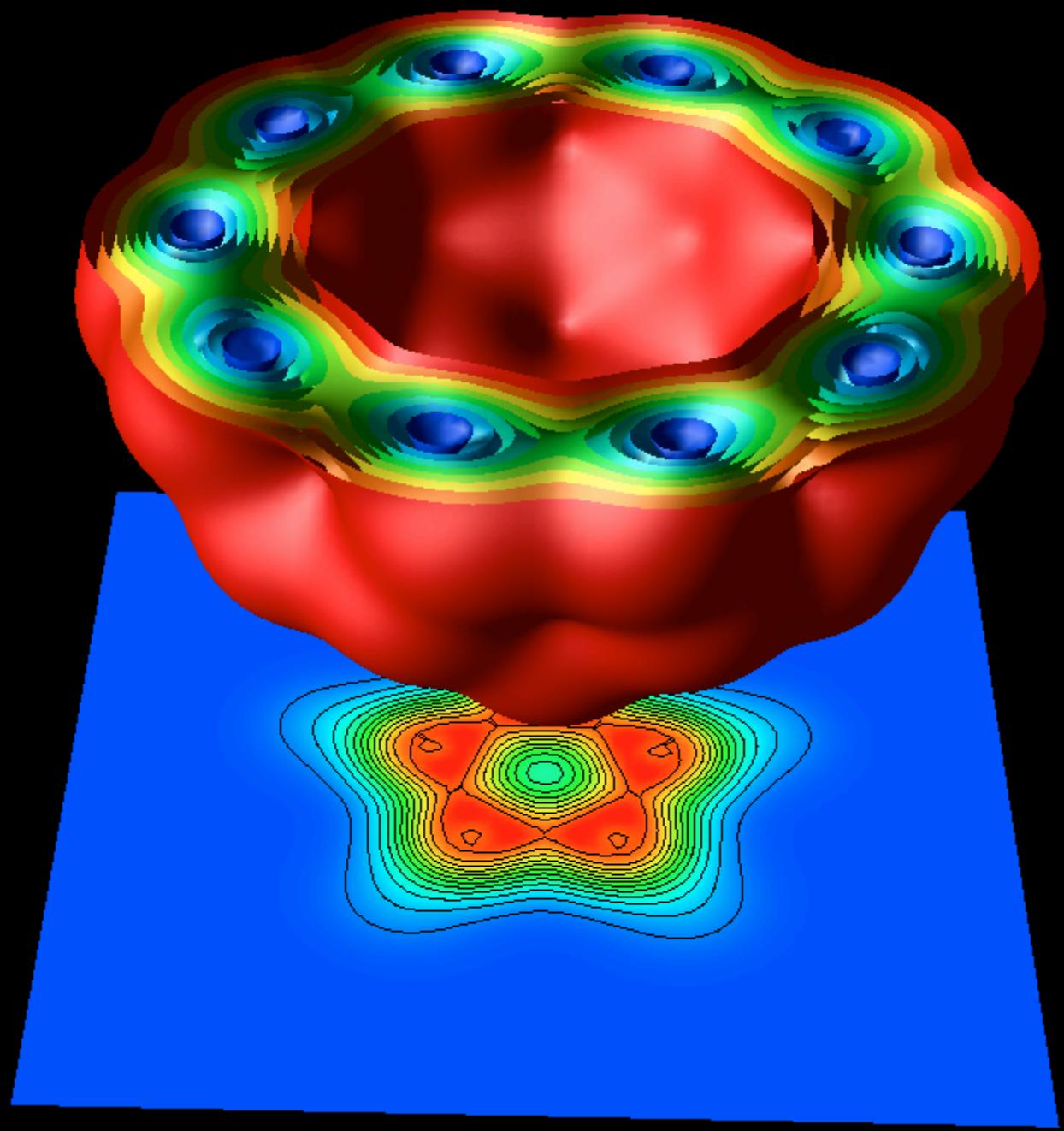
MOLECULAR ORBITAL

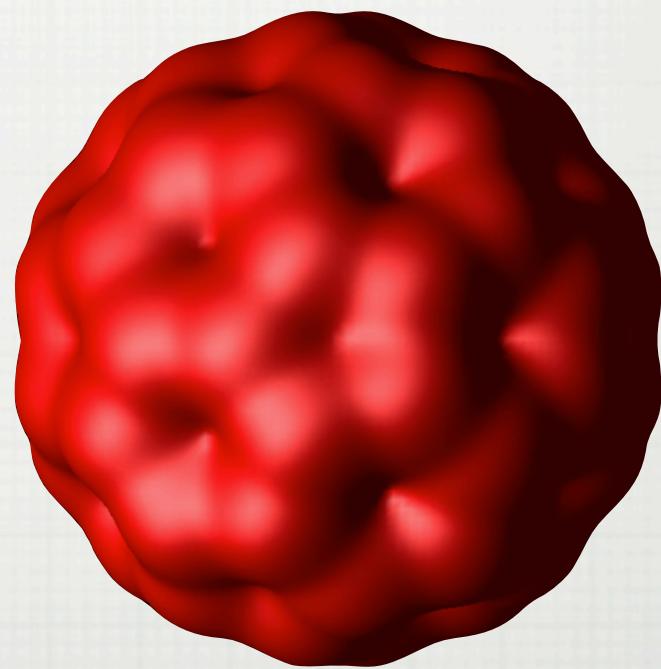


DENSITY OF STATES











HAVE A GOOD DAY