

# ANNOUNCEMENT

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- Creation of New Class**
  - PHYS 6960 ADV COMPUTATIONAL PHYSICS**
  - Will meet on Wednesday 5PM to 6:50 PM**
- Prerequisite: 4810, QM**
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# POSSIBLE TOPICS

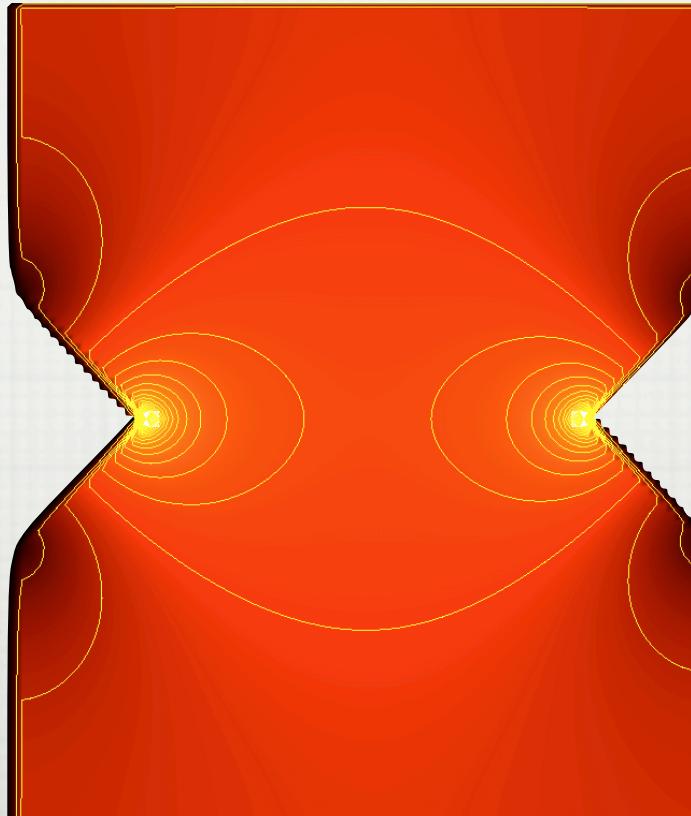
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- (1) finite difference methods;
- (2) Molecular dynamics,
- (3) wavelet transforms;
- (4) Perturbation theory at low orders
- (5) Random numbers and multidimensional integration (importance sampling)
- (6) Dynamical Mean Field Theory for model Hamiltonians
- (7) Local Density Approximation + Dynamical Mean Field Theory (LDA+DMFT) in action
- (8) Parallel programming with MPI
- (9) Recursion techniques
- (10) Sparse ALgebra
- (11) Solving Systems of Equations with Matrices
- (12) Fractals & Statistical Growth
- (13) More GPU
- (14) PDEWaves: String, Quantum Packet, and E&M

# PHY-4810

# COMPUTATIONAL PHYSICS

LECTURE 10: PARTIAL DIFFERENTIAL EQUATIONS (PDE)



# PDE: DEFINITION

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*Partial differential equations (PDEs) are a type of differential equations involving an unknown function (or functions) of **several** independent variables and their **partial derivatives** with respect to those variables.*

$$F(x_1, \dots, x_n, u, \frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_n} u, \frac{\partial^2}{\partial x_1 \partial x_1} u, \dots, \frac{\partial^2}{\partial x_1 \partial x_2} u, \dots) = 0$$

# EXAMPLES

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- 2D Laplace equation
- Heat equation
- Wave equation
- Spherical waves
- Ginzburg Landau (superconductivity)
- Vibrating Membrane

# CLASSIFICATION

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$$A \frac{\partial^2 U}{\partial x^2} + 2B \frac{\partial^2 U}{\partial x \partial y} + C \frac{\partial^2 U}{\partial y^2} + D \frac{\partial U}{\partial x} + E \frac{\partial U}{\partial y} = F$$

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## General Form

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## Types

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## Types

*Elliptic*

*Parabolic*

*Hyperbolic*

$$d = AC - B^2 > 0$$

$$d = AC - B^2 = 0$$

$$d = AC - B^2 < 0$$

$$\nabla^2 U(x) = -4\pi\rho(x)$$

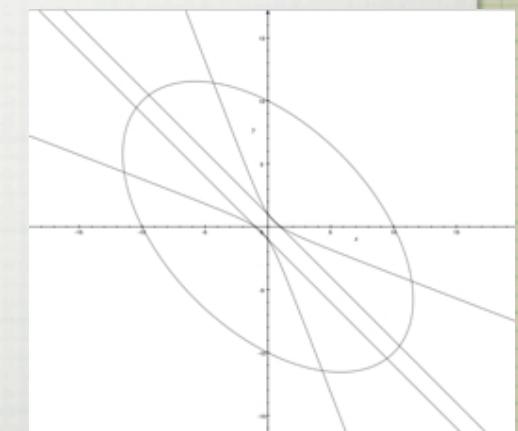
$$\nabla^2 U(\mathbf{x}, t) = a \partial U / \partial t$$

$$\nabla^2 U(\mathbf{x}, t) = c^{-2} \partial^2 U / \partial t^2$$

Poisson's

Heat

Wave



# TYPES OF BOUNDARY CONDITIONS

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- If the boundary condition is the value of the **normal derivative** on the surrounding surface, we have a **Neumann boundary condition**.
- If both the **value** of solution and **its derivative** are specified on a closed boundary, we have a **Cauchy boundary condition** (*these are more initial conditions than boundary conditions!*)



# BOUNDARY CONDITIONS

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Boundary condition	Elliptic (Poisson equation)	Hyperbolic (Wave equation)	Parabolic (Heat equation)
Dirichlet open surface	Underspecified	Underspecified	<i>Unique and stable (1D)</i>
Dirichlet closed surface	<i>Unique and stable</i>	Overspecified	Overspecified
Neumann open surface	Underspecified	Underspecified	<i>Unique and stable (1D)</i>
Neumann closed surface	<i>Unique and stable</i>	Overspecified	Overspecified
Cauchy open surface	Unphysical	<i>Unique and stable</i>	Overspecified
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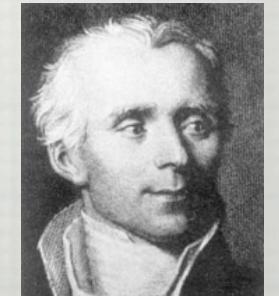
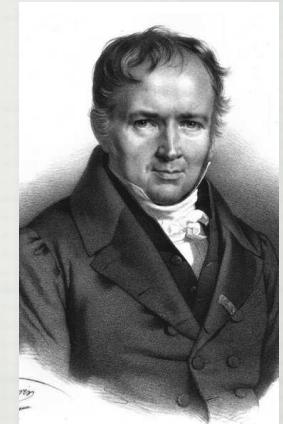
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# LAPLACE AND POISSON EQUATIONS

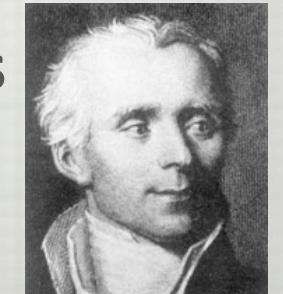
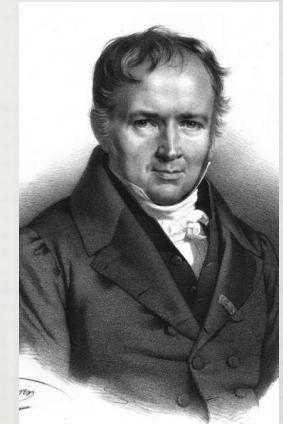
$$\frac{\partial^2 U(x, y)}{\partial x^2} + \frac{\partial^2 U(x, y)}{\partial y^2} = \begin{cases} 0 & \text{Laplace's equation} \\ -4\pi\rho(\mathbf{x}) & \text{Poisson's equation} \end{cases}$$



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- Dirichlet: impose potential  $U$  at the boundaries
- Neuman: impose electric field at the boundaries



# ANALYTICAL SOLUTIONS

---

$$U(x, y) = X(x)Y(y) \quad \Rightarrow \quad \frac{d^2X(x)/dx^2}{X(x)} + \frac{d^2Y(y)/dy^2}{Y(y)} = 0$$

# ANALYTICAL SOLUTIONS

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- Electrodynamics: use of polynomial expansions (for instance, Fourier) using separation of variables.

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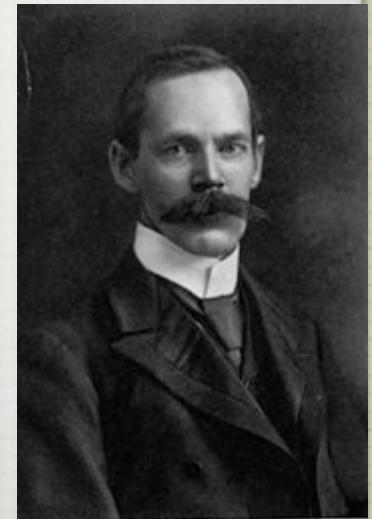
- **Cons:**
  - **Solutions in infinite series end up being numerical!**
  - **Often hard to change specific boundary conditions in complicated geometries**

# NUMERICAL SOLUTION

- Remember Taylor or McLaurin (and homework #1):

$$\begin{array}{l} \cdots \\ \cdot U(x + \Delta x, y) = U(x, y) + \frac{\partial U}{\partial x} \Delta x + \frac{1}{2} \frac{\partial^2 U}{\partial x^2} (\Delta x)^2 + \cdots \cdot \\ \cdot U(x - \Delta x, y) = U(x, y) - \frac{\partial U}{\partial x} \Delta x + \frac{1}{2} \frac{\partial^2 U}{\partial x^2} (\Delta x)^2 - \cdots \cdot \\ \cdots \end{array}$$

- Adding equation yields the central difference approximation:



$$\frac{\partial^2 U(x, y)}{\partial x^2} \simeq \frac{U(x + \Delta x, y) + U(x - \Delta x, y) - 2U(x, y)}{(\Delta x)^2} + \mathcal{O}(\Delta x^4)$$

$$\frac{\partial^2 U(x, y)}{\partial y^2} \simeq \frac{U(x, y + \Delta y) + U(x, y - \Delta y) - 2U(x, y)}{(\Delta y)^2} + \mathcal{O}(\Delta y^4)$$

# LAPLACIAN OF $U(X,Y)$

---

$$\frac{U(x + \Delta x, y) + U(x - \Delta x, y) - 2U(x, y)}{(\Delta x)^2} + \frac{U(x, y + \Delta y) + U(x, y - \Delta y) - 2U(x, y)}{(\Delta y)^2} = -4\pi\rho$$

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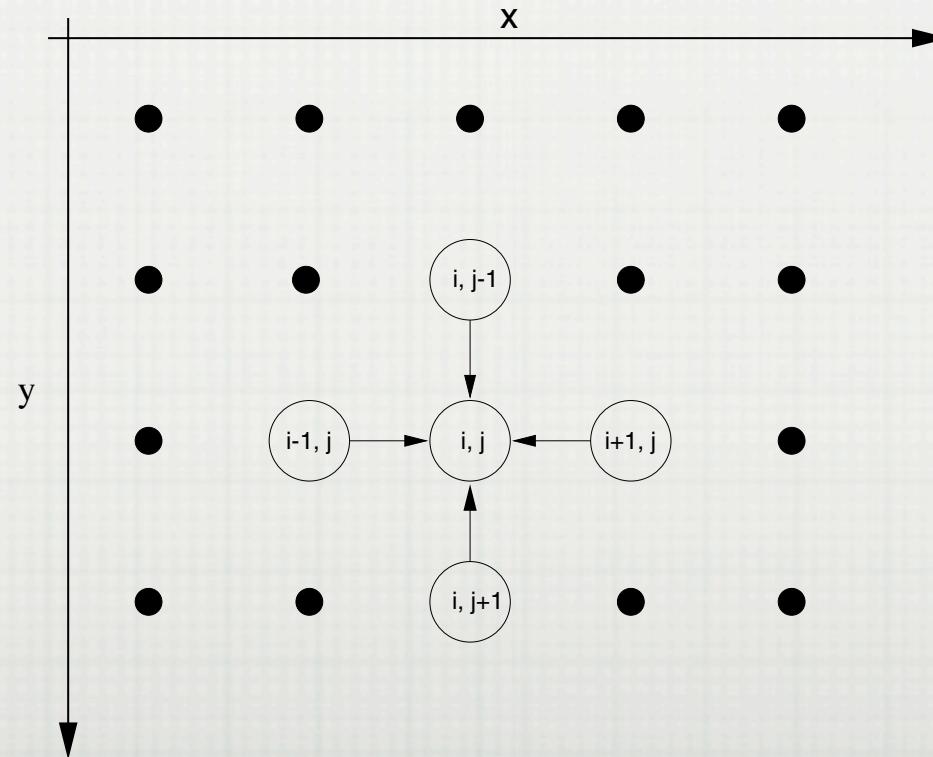
For equal spacing in x and y :

$$U(x + \Delta x, y) + U(x - \Delta x, y) + U(x, y + \Delta y) + U(x, y - \Delta y) - 4U(x, y) = -4\pi\rho$$

# LET'S MEDITATE

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$$U(x + \Delta x, y) + U(x - \Delta x, y) + U(x, y + \Delta y) + U(x, y - \Delta y) - 4U(x, y) = -4\pi\rho$$



# DISCRETIZED VERSION

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$$U(x + \Delta x, y) + U(x - \Delta x, y) + U(x, y + \Delta y) + U(x, y - \Delta y) - 4U(x, y) = -4\pi\rho$$



$$U_{i,j} = \frac{1}{4} [U_{i+1,j} + U_{i-1,j} + U_{i,j+1} + U_{i,j-1}] + \pi\rho(i\Delta, j\Delta)\Delta^2$$

$$x = x_0 + i\Delta, \quad y = y_0 + j\Delta \quad i, j = 0, \dots, N_{\max-1}$$

$$\Delta x = \Delta y = \Delta = L/(N_{\max} - 1)$$

# RELAXATION: JACOBI METHOD (I)

$$U_{i,j} = \frac{1}{4} [U_{i+1,j} + U_{i-1,j} + U_{i,j+1} + U_{i,j-1}] + \pi\rho(i\Delta, j\Delta)\Delta^2$$

**NEW**                                    **OLD**

- IMPOSE BOUNDARY CONDITIONS
- START WITH “REASONABLE GUESS” TO ACCELERATE CONVERGENCE



# RELAXATION: GAUSS-SEIDEL METHOD

$$U_{i,j} = \frac{1}{4} [U_{i+1,j} + U_{i-1,j} + U_{i,j+1} + U_{i,j-1}] + \pi \rho(i\Delta, j\Delta) \Delta^2$$

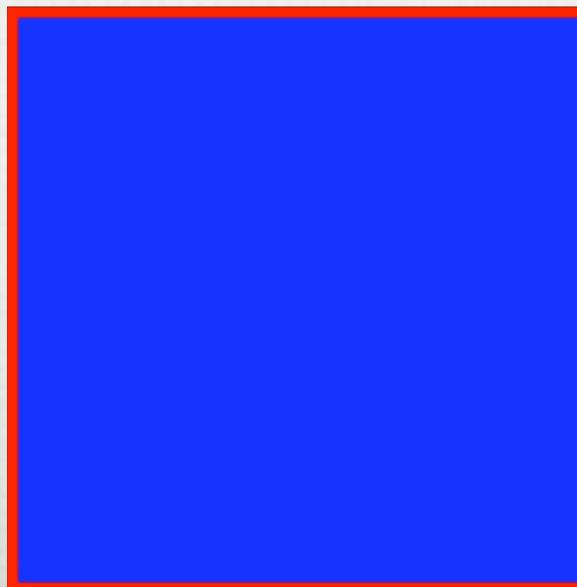
The equation is presented within a dashed orange rectangular frame. Inside the frame, the word "OLD" is written in green above three curly braces that group the terms  $U_{i+1,j}$ ,  $U_{i-1,j}$ , and  $U_{i,j+1}$ . Below these three terms, the word "NEW" is written in red. To the right of the third term, another curly brace groups the terms  $U_{i,j-1}$  and  $\pi \rho(i\Delta, j\Delta) \Delta^2$ . Below this second group, the word "NEW" is also written in red.



# FIRST TEST: BOUNDARY CONDITIONS

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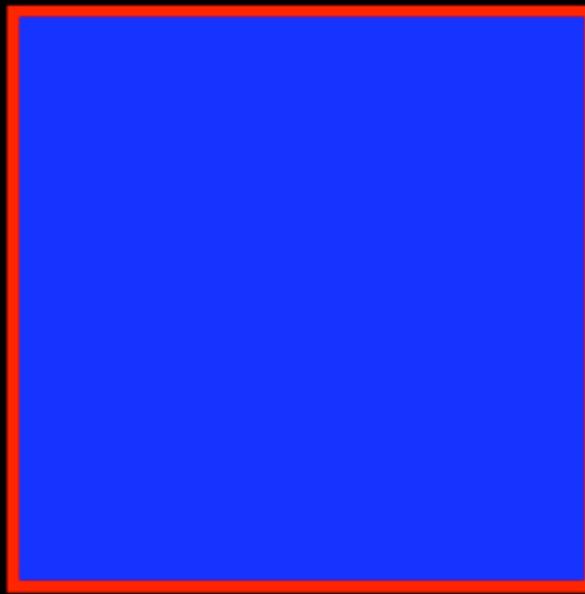
- We should know the potential at the boundary (Dirichlet)
- We could know the derivative of the potential at the boundary (Neumann)



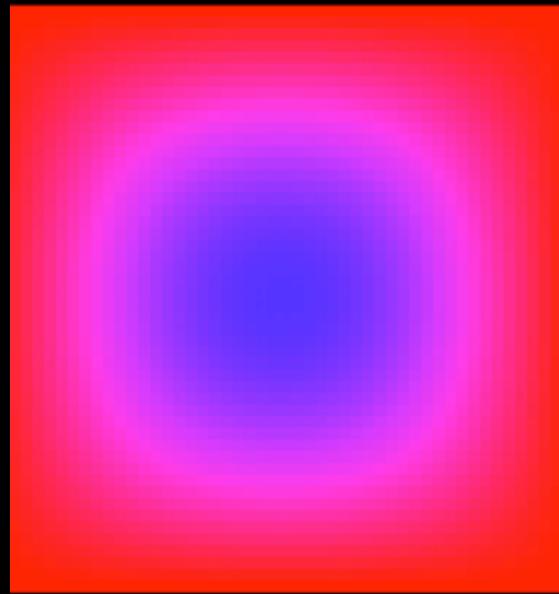
$V=10 \text{ VOLT}$

$V=0 \text{ VOLT}$

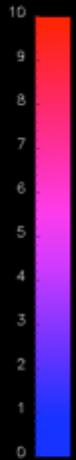
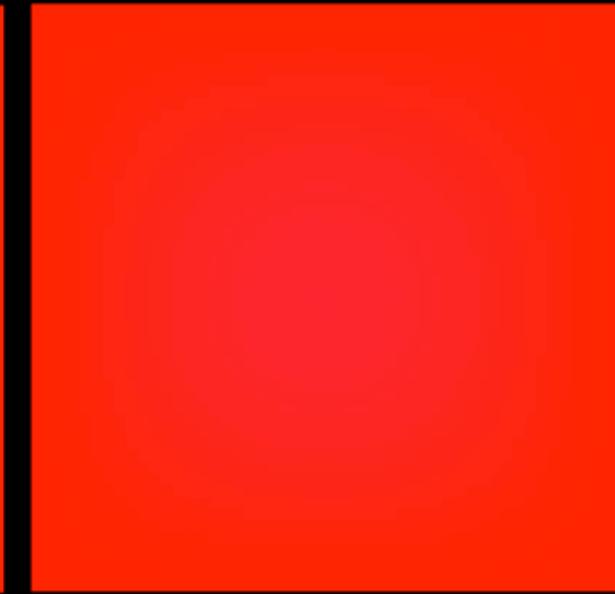
**INITIAL**



**AFTER FEW ITERATIONS**



**CONVERGED**



**$V=10$  V ON THE BOUNDARIES**

# HOW TO ACCELERATE CONVERGENCE: OVER- AND UNDER- RELAXATIONS

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$$U_{i,j}^{(\text{new})} = U_{i,j}^{(\text{old})} + r_{i,j}$$

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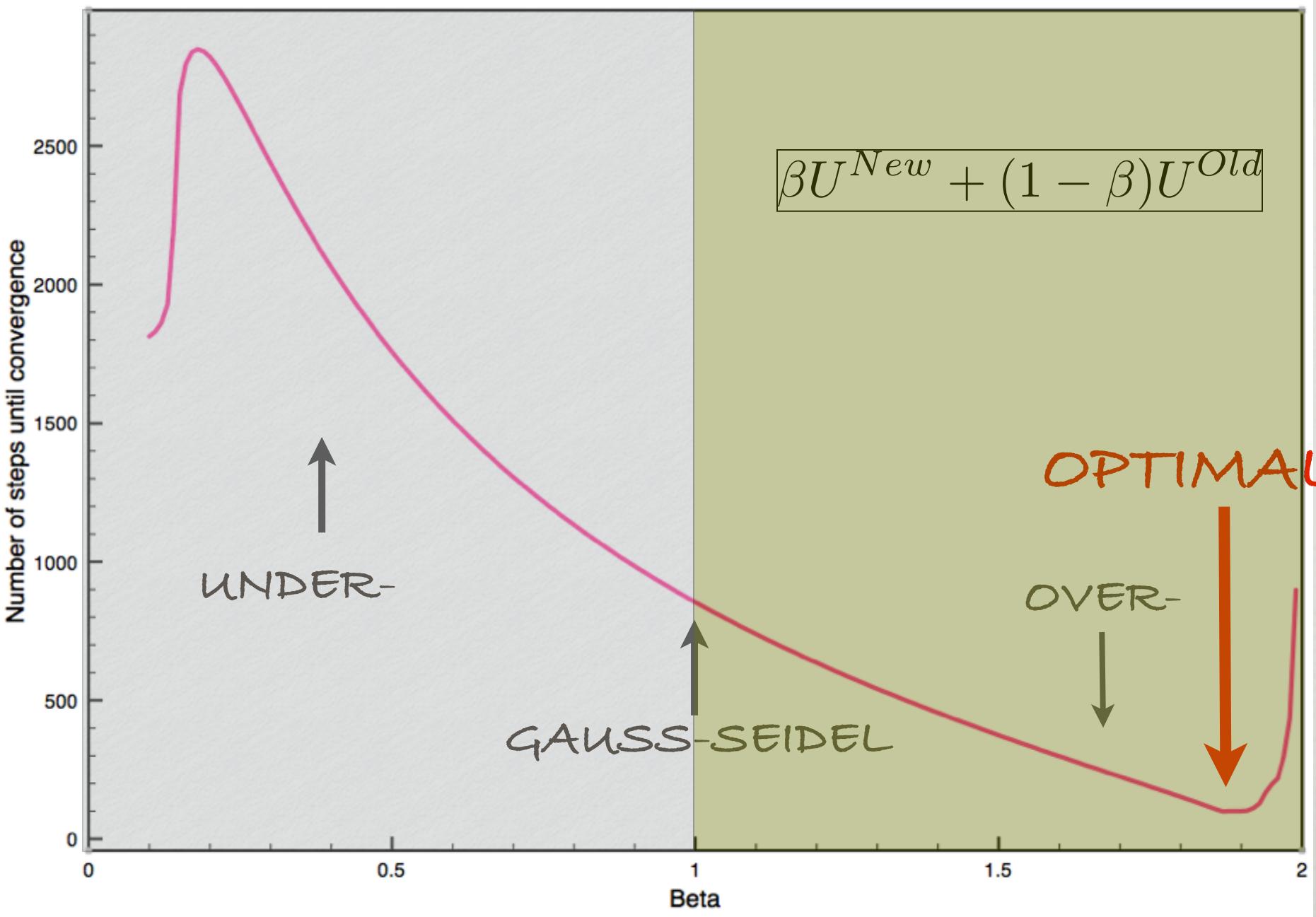
$$\begin{aligned} r_{i,j} &\equiv U_{i,j}^{(\text{new})} - U_{i,j}^{(\text{old})} \\ &= \frac{1}{4} \left[ U_{i+1,j}^{(\text{old})} + U_{i-1,j}^{(\text{new})} + U_{i,j+1}^{(\text{old})} + U_{i,j-1}^{(\text{new})} \right] - U_{i,j}^{(\text{old})} \end{aligned}$$

$$U_{i,j}^{(\text{new})} = U_{i,j}^{(\text{old})} + \omega r_{i,j}$$

$\omega < 1$ : under-relaxation

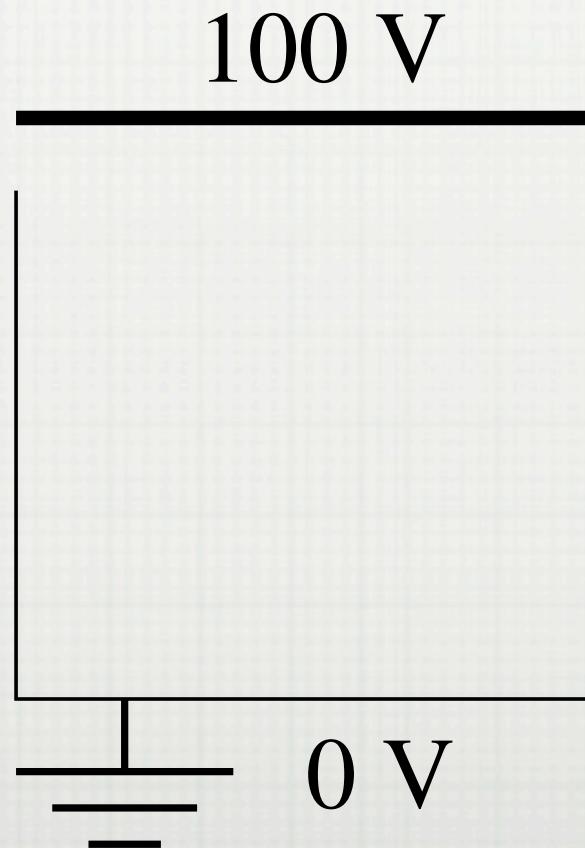
$\omega = 1$ : Gauss-Seidel

$\omega \geq 1$ : over-relaxation

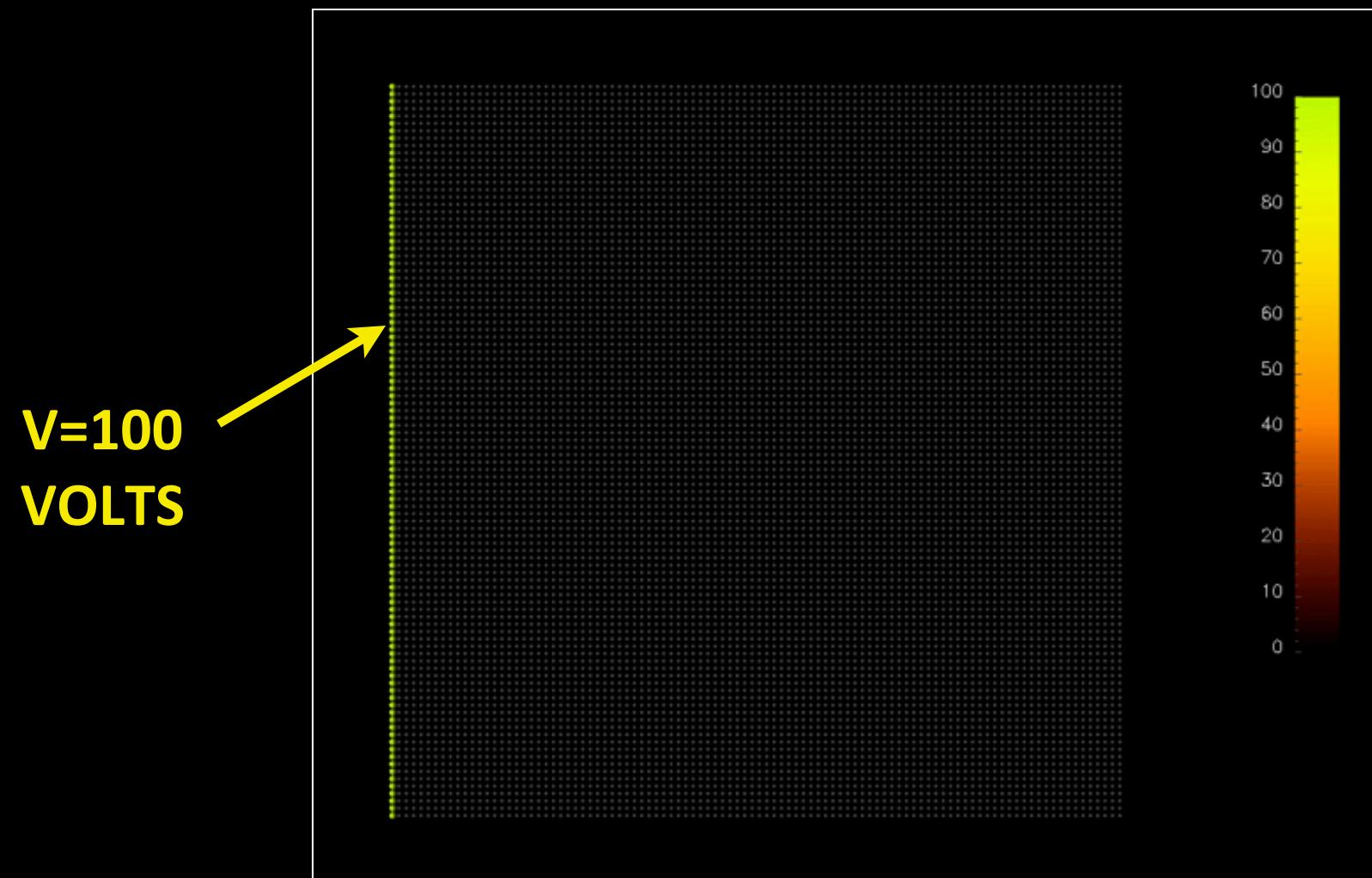


## OTHER EXAMPLE: SINGLE METALLIC PLATE CAPACITOR

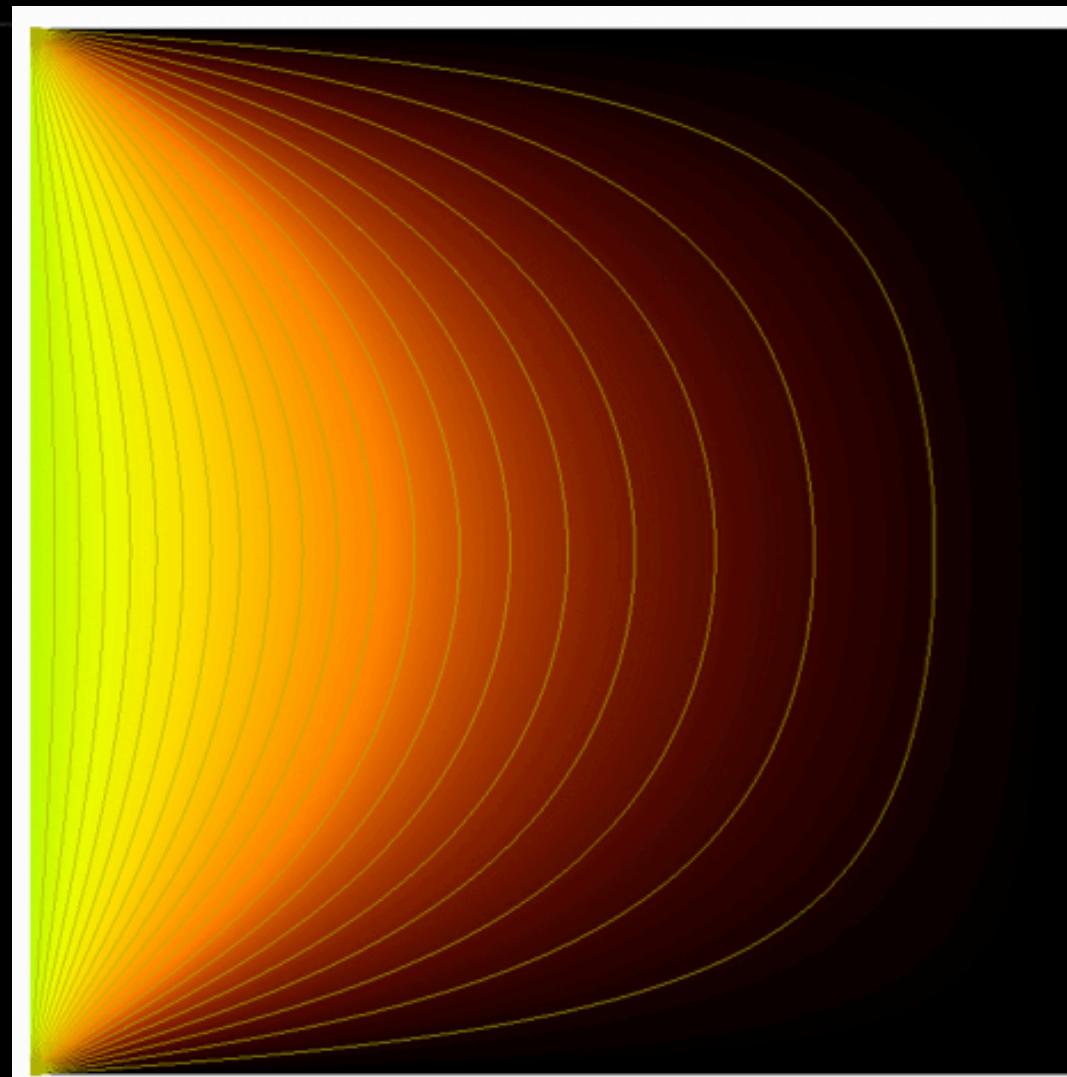
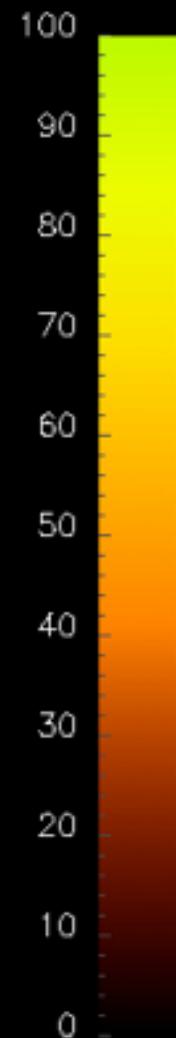
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# INITIAL CONDITIONS

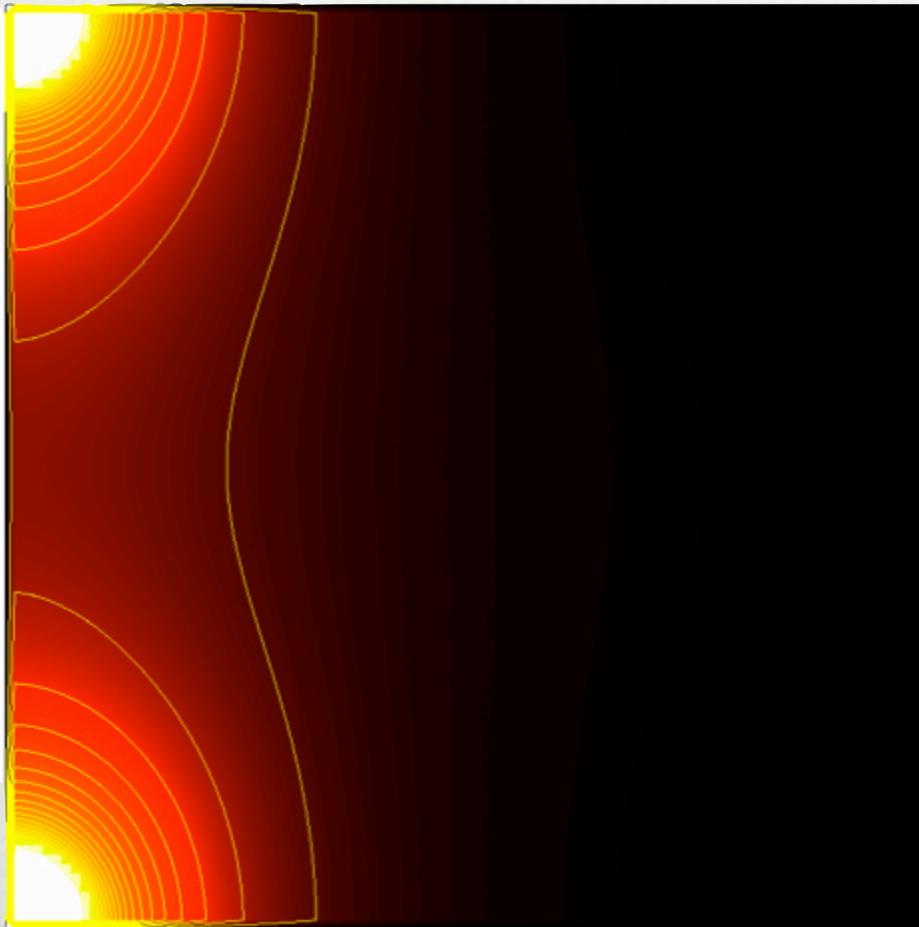


# AFTER CONVERGENCE

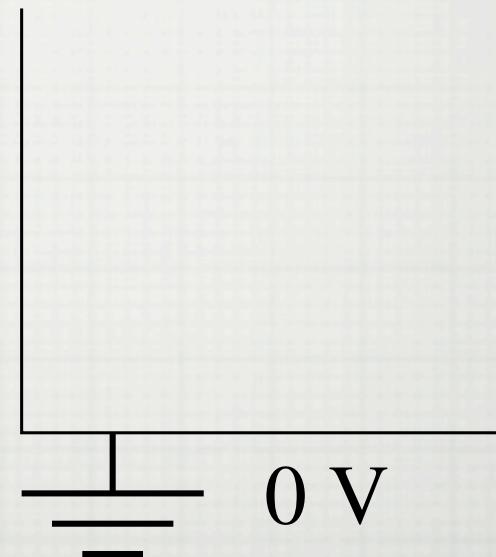


# ELECTRIC FIELD

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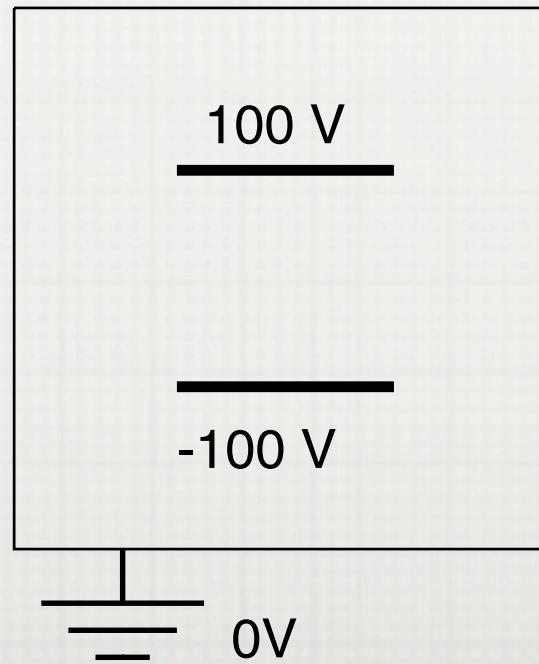


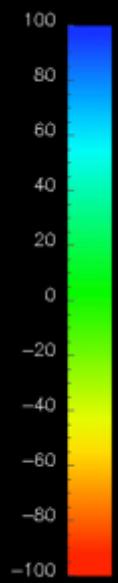
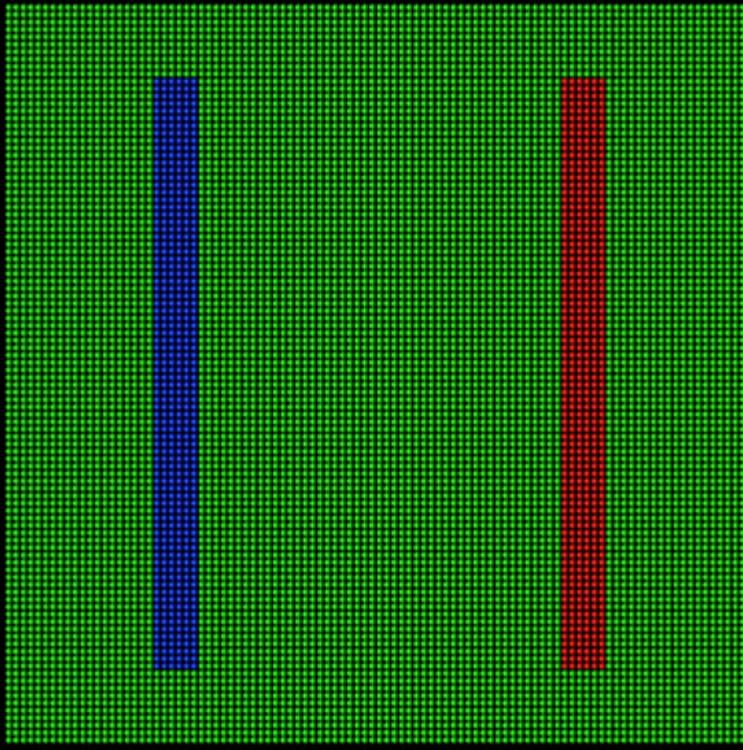
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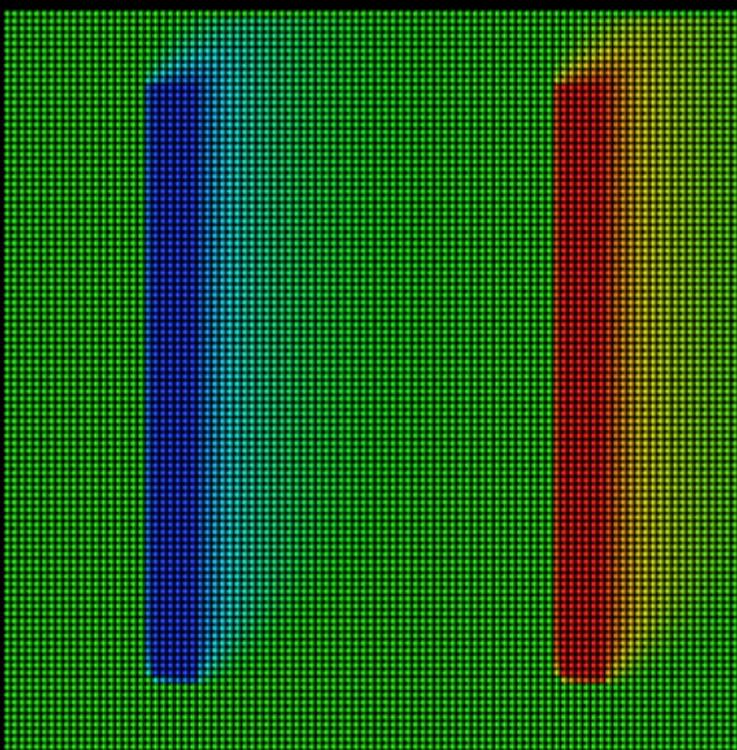
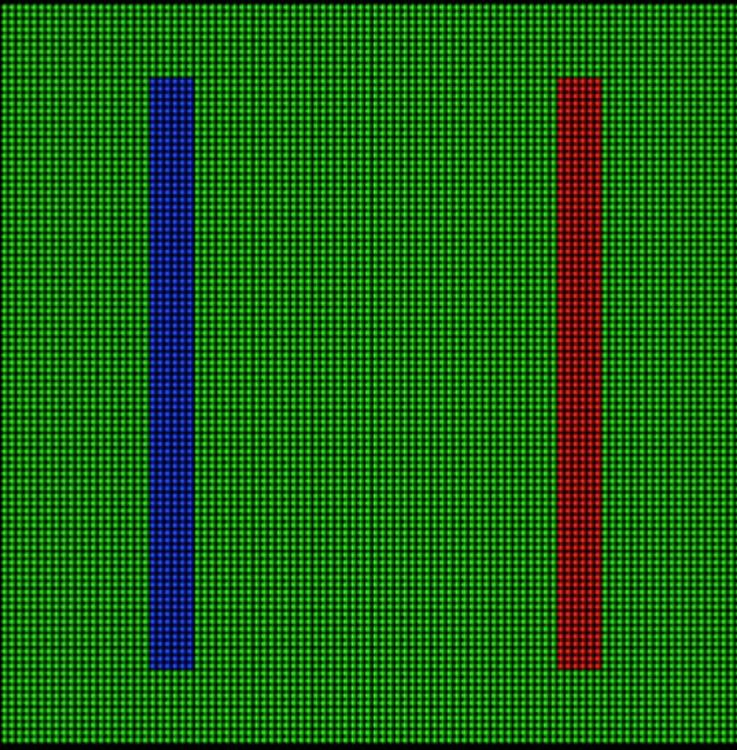


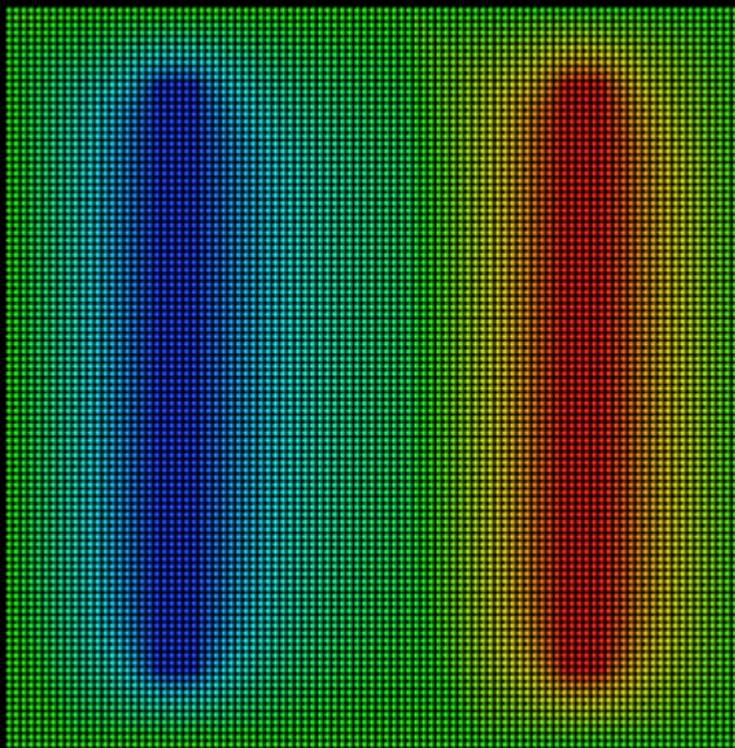
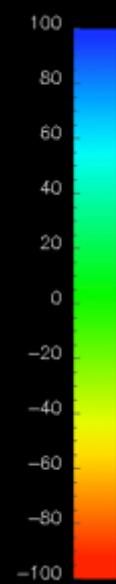
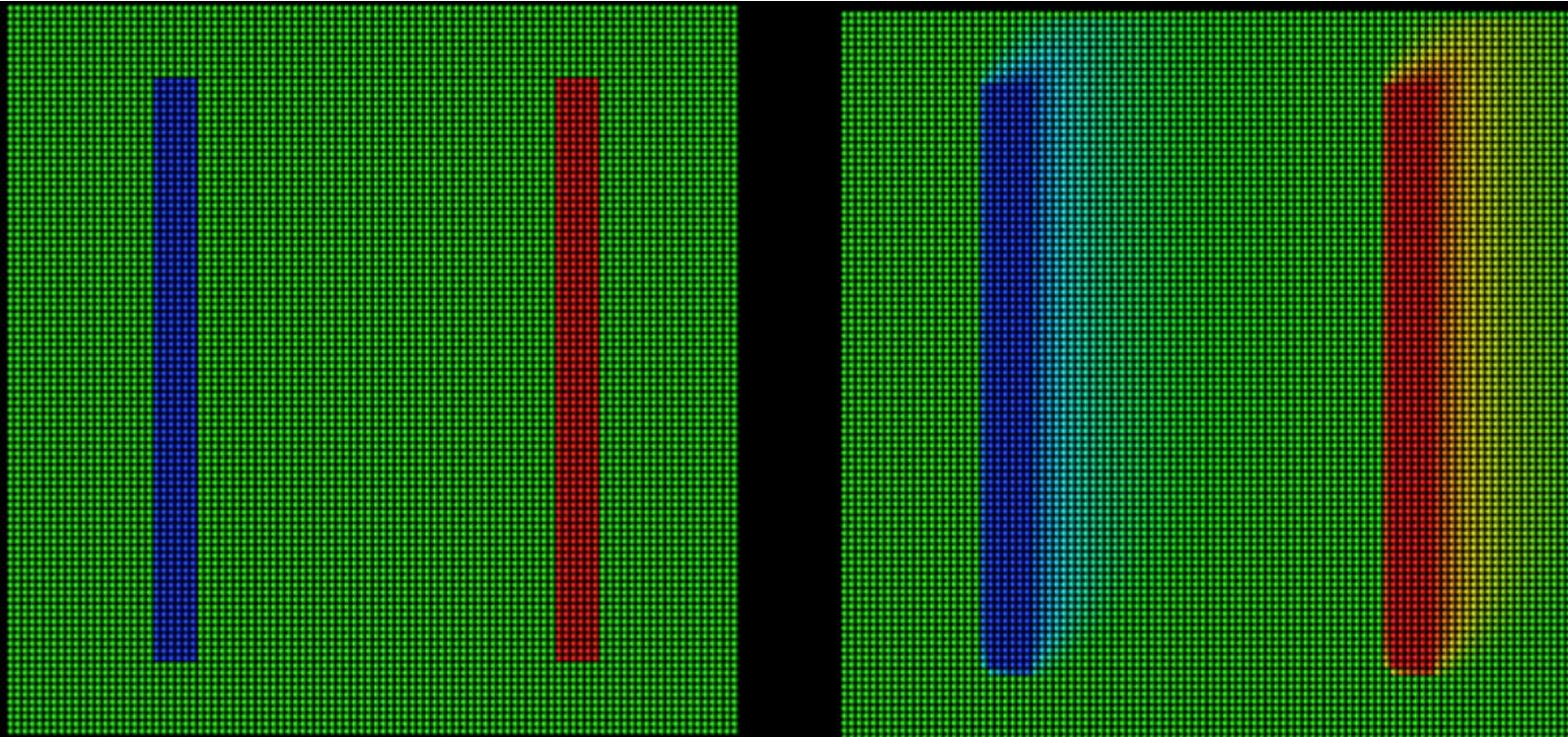
# PARALLEL PLATE CAPACITOR

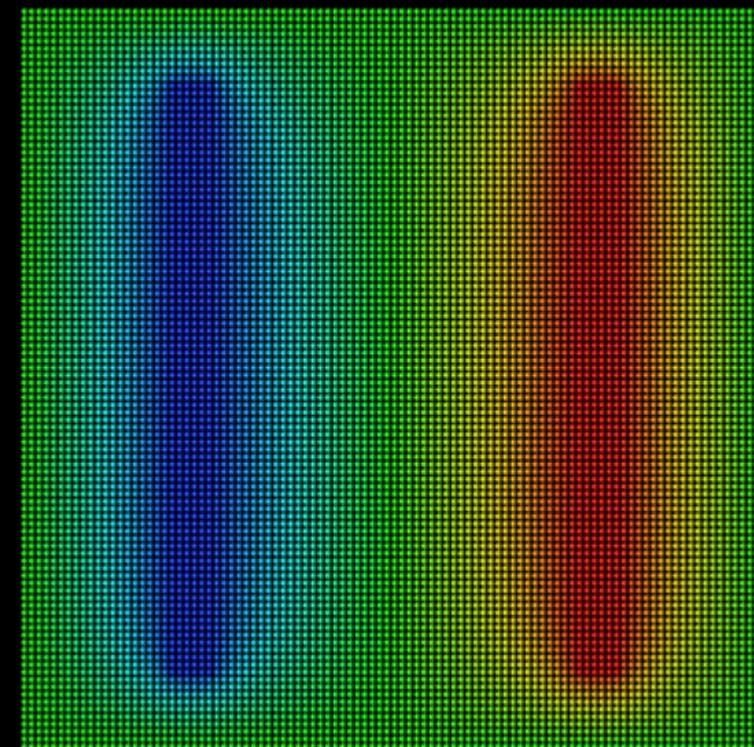
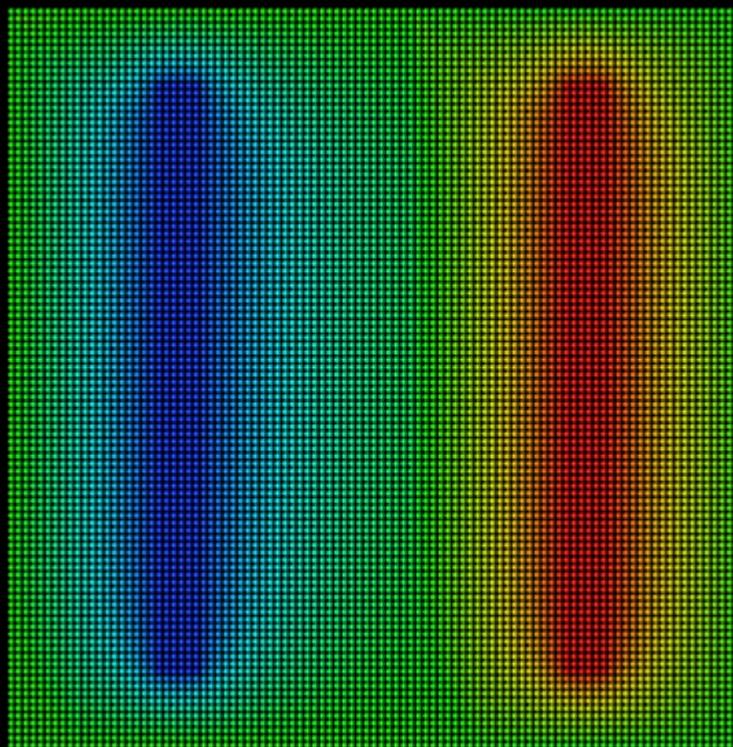
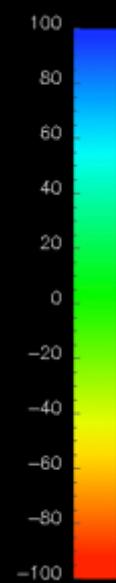
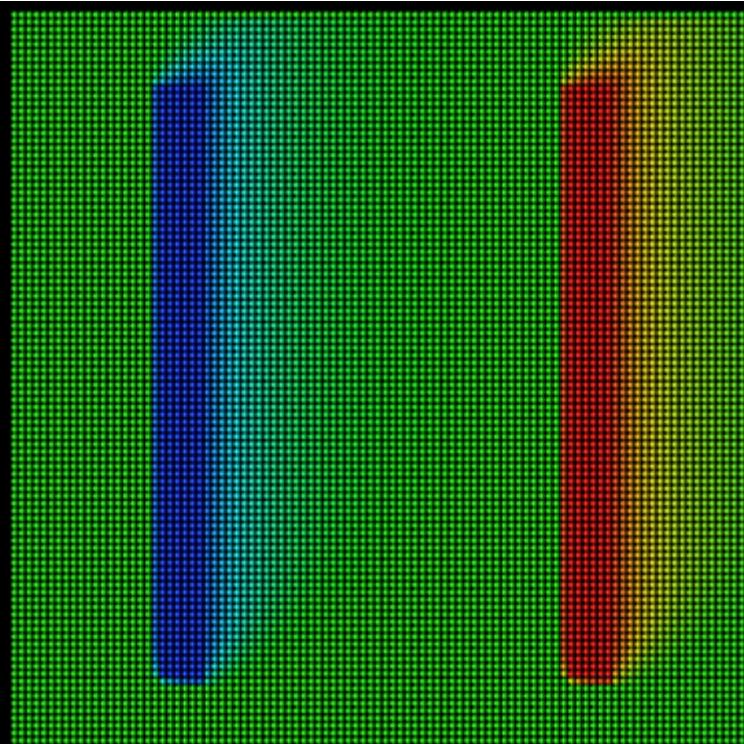
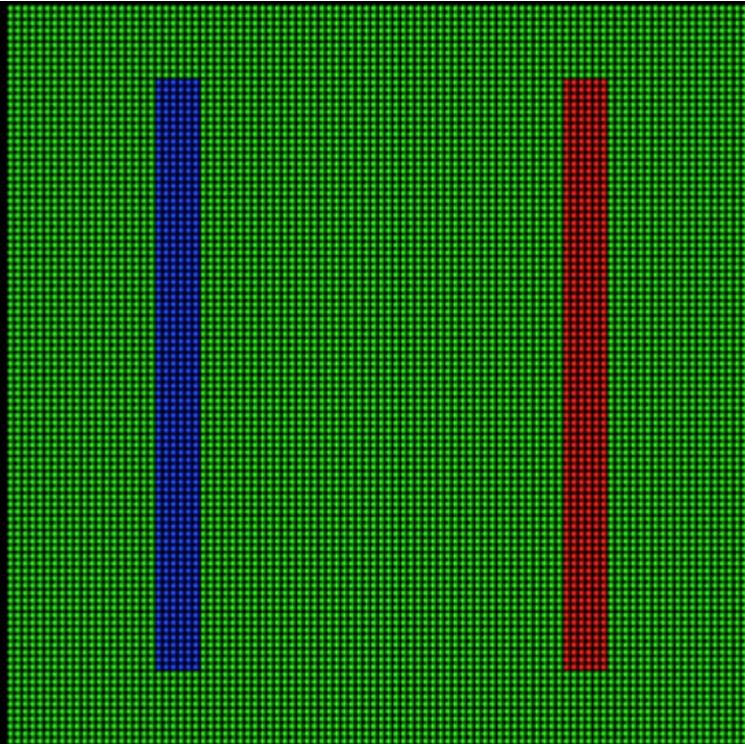
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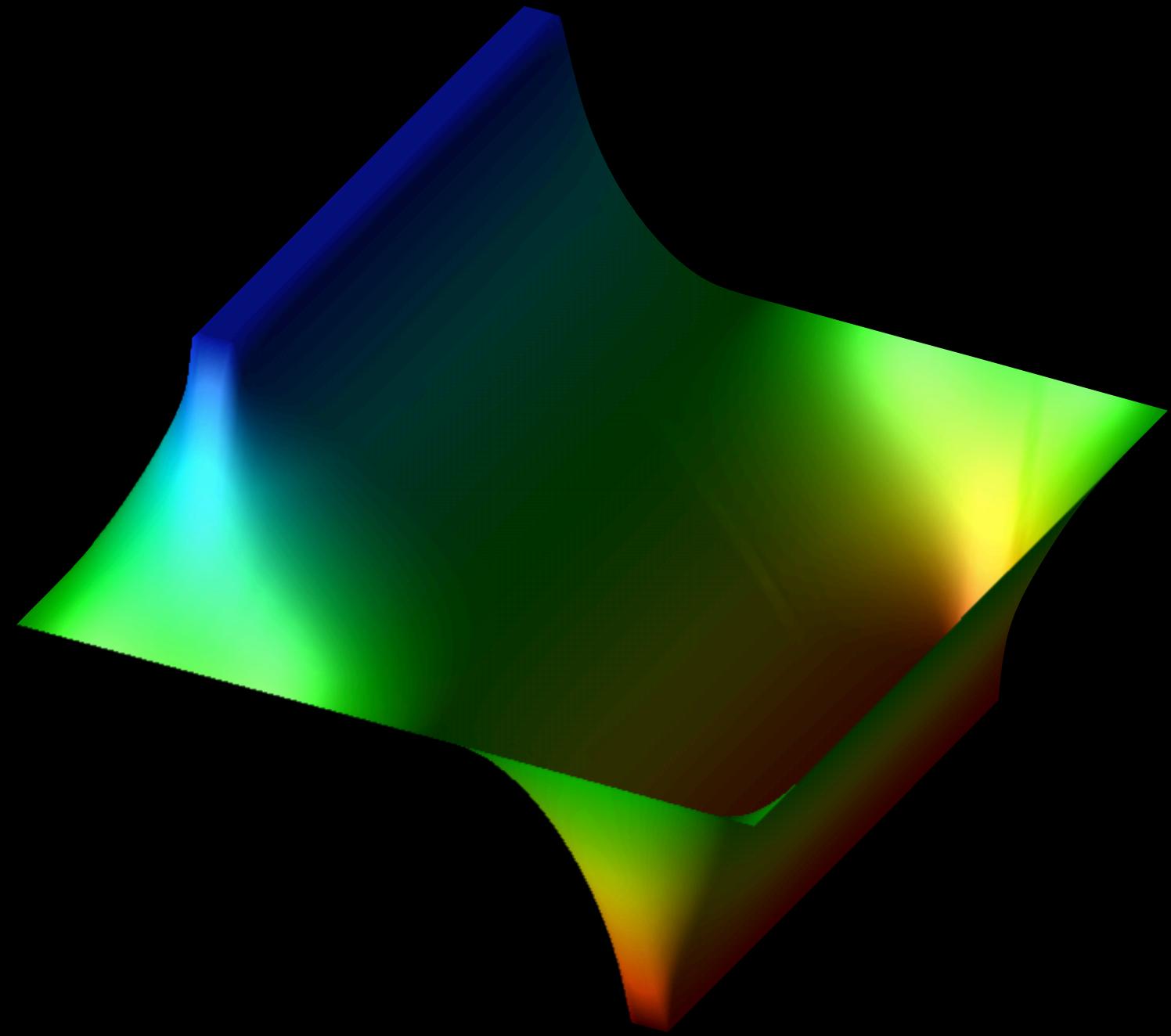






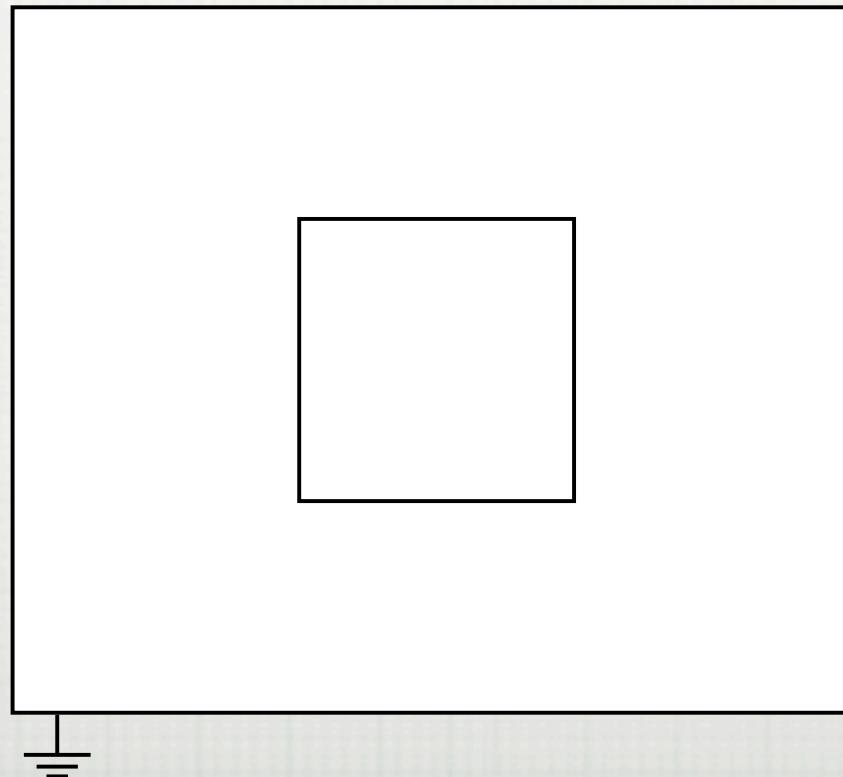




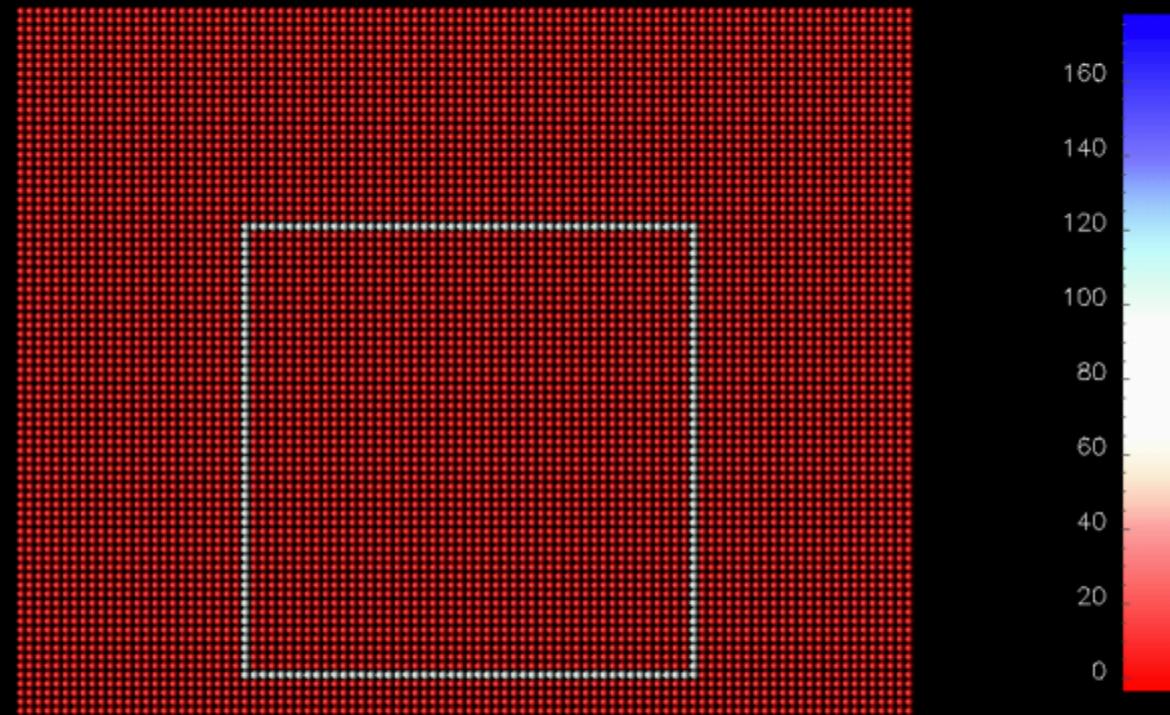


# SQUARE IN A SQUARE

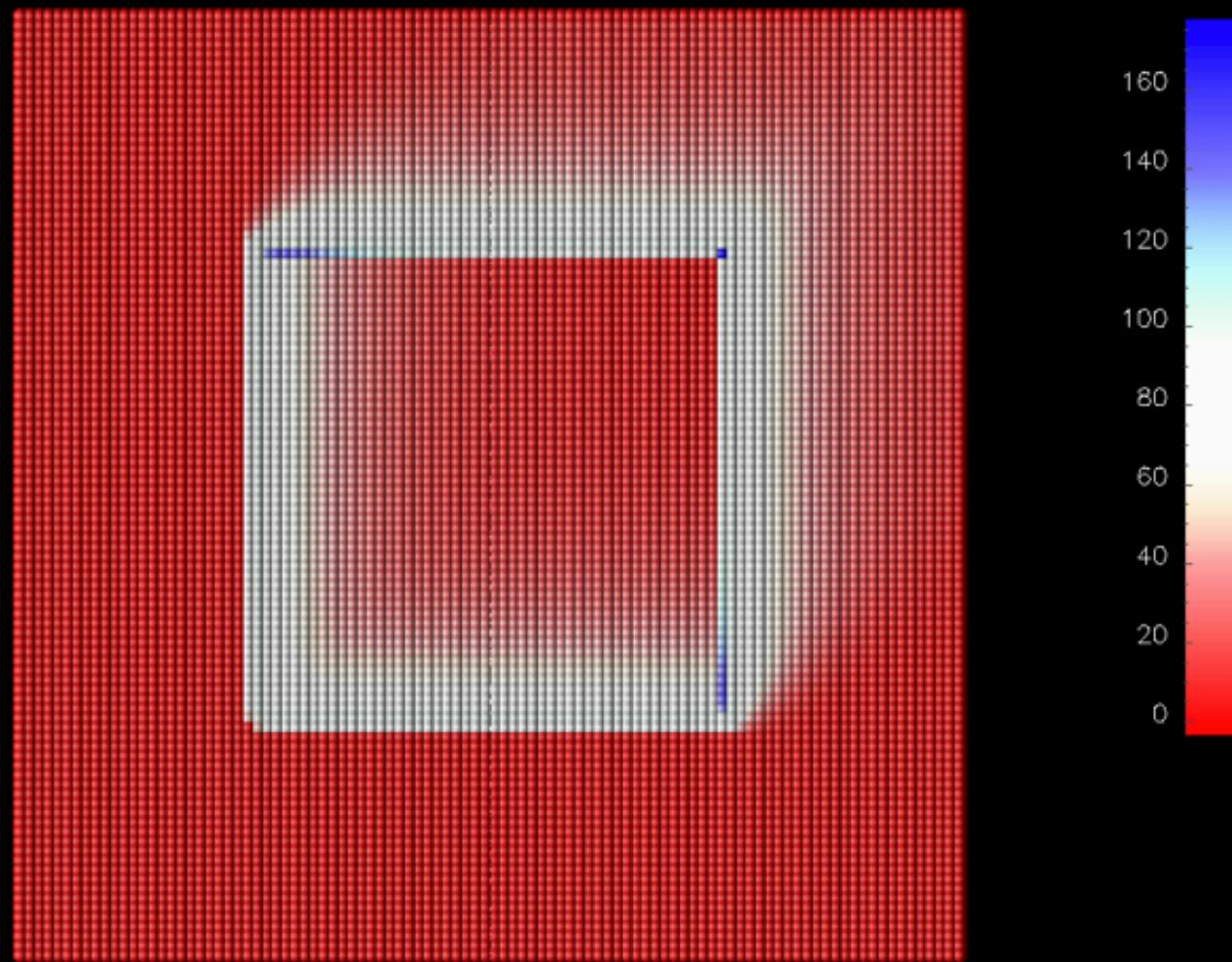
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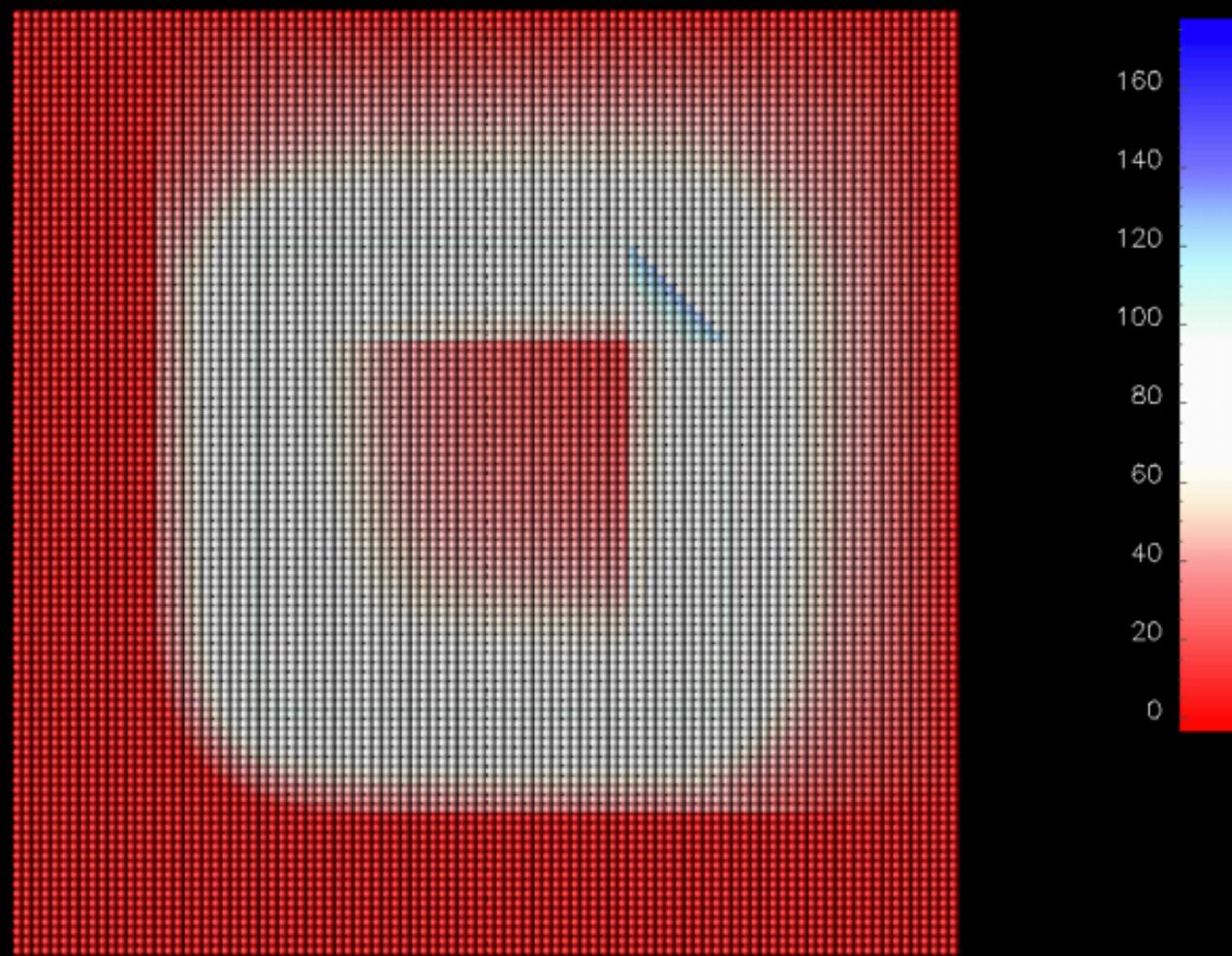
# INITIAL CONDITIONS



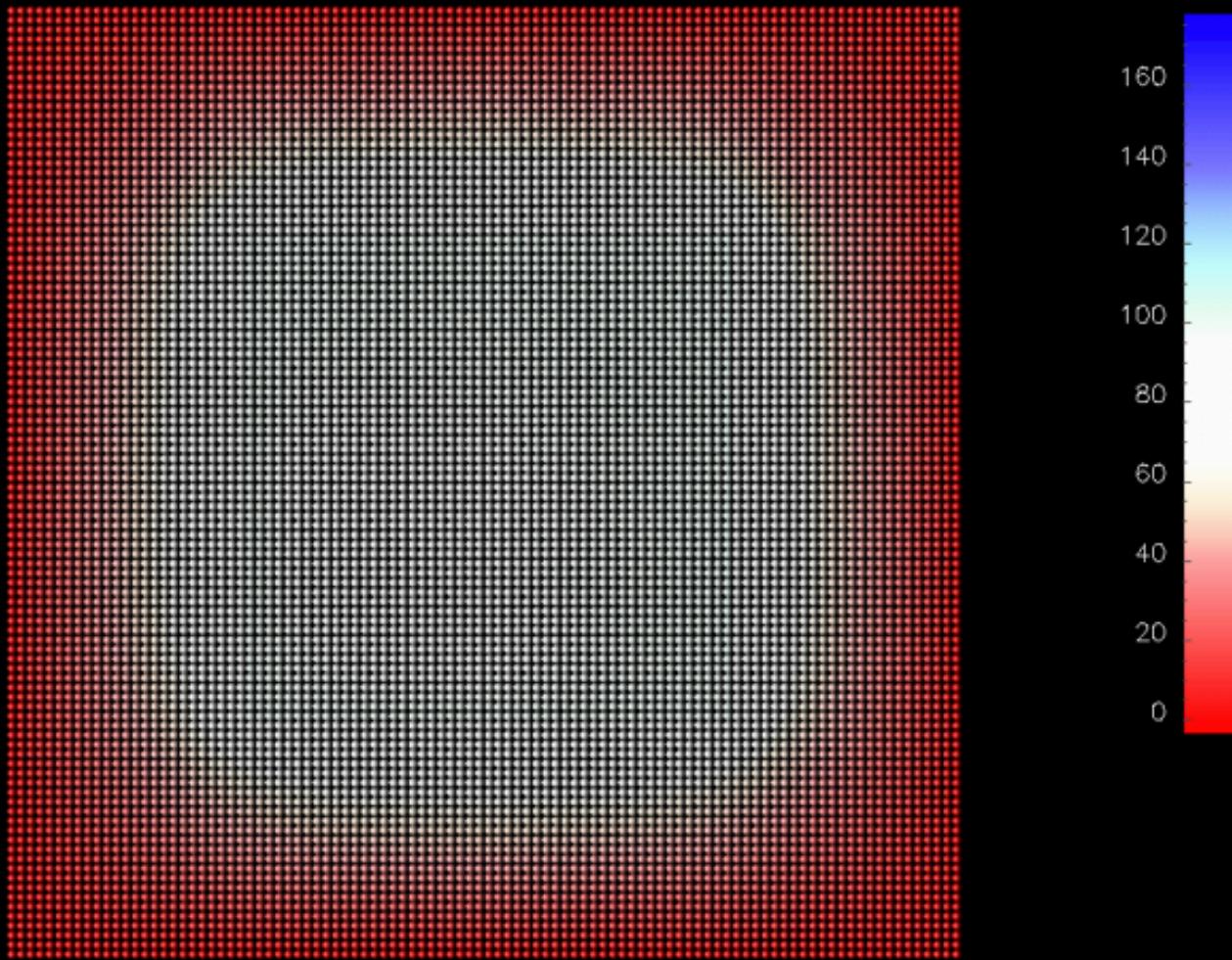
# STEP 1



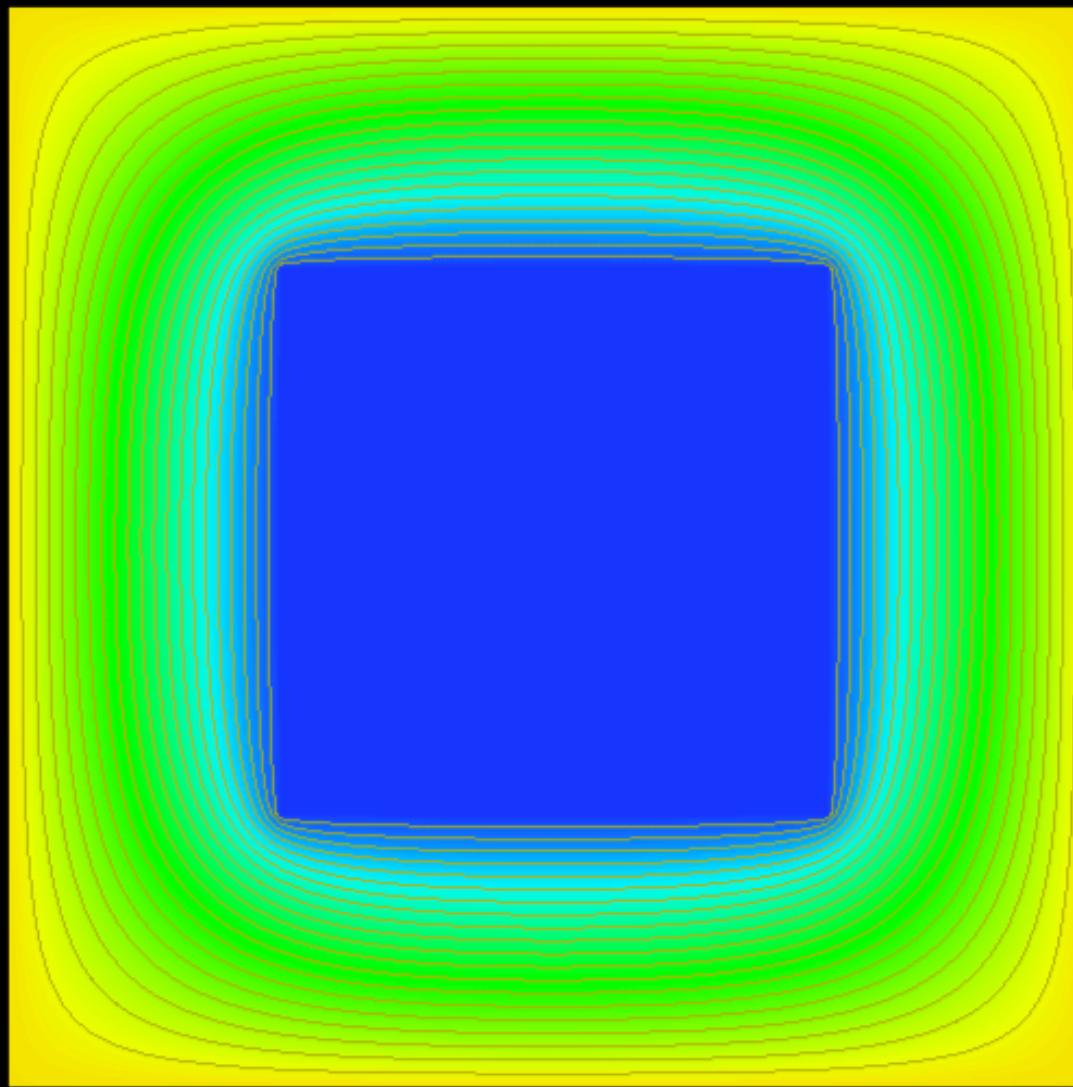
# STEP 50



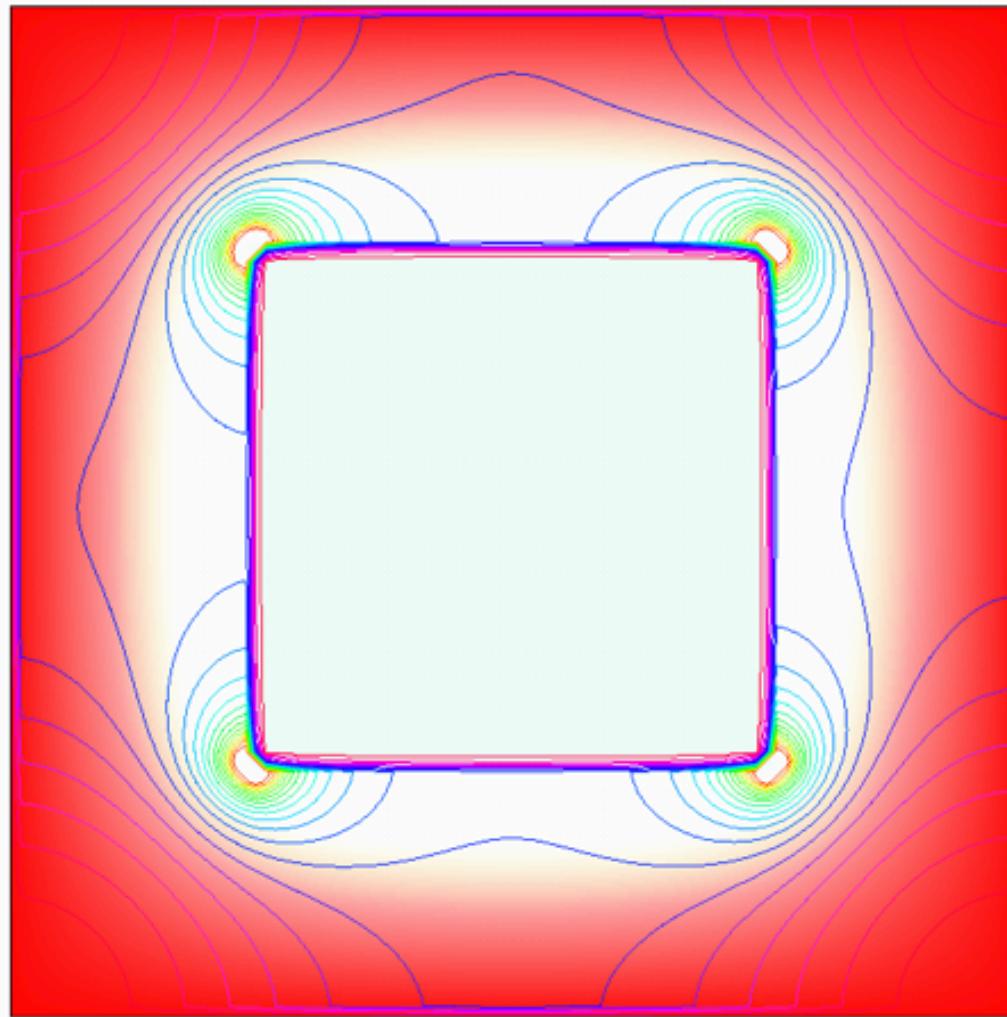
# CONVERGED

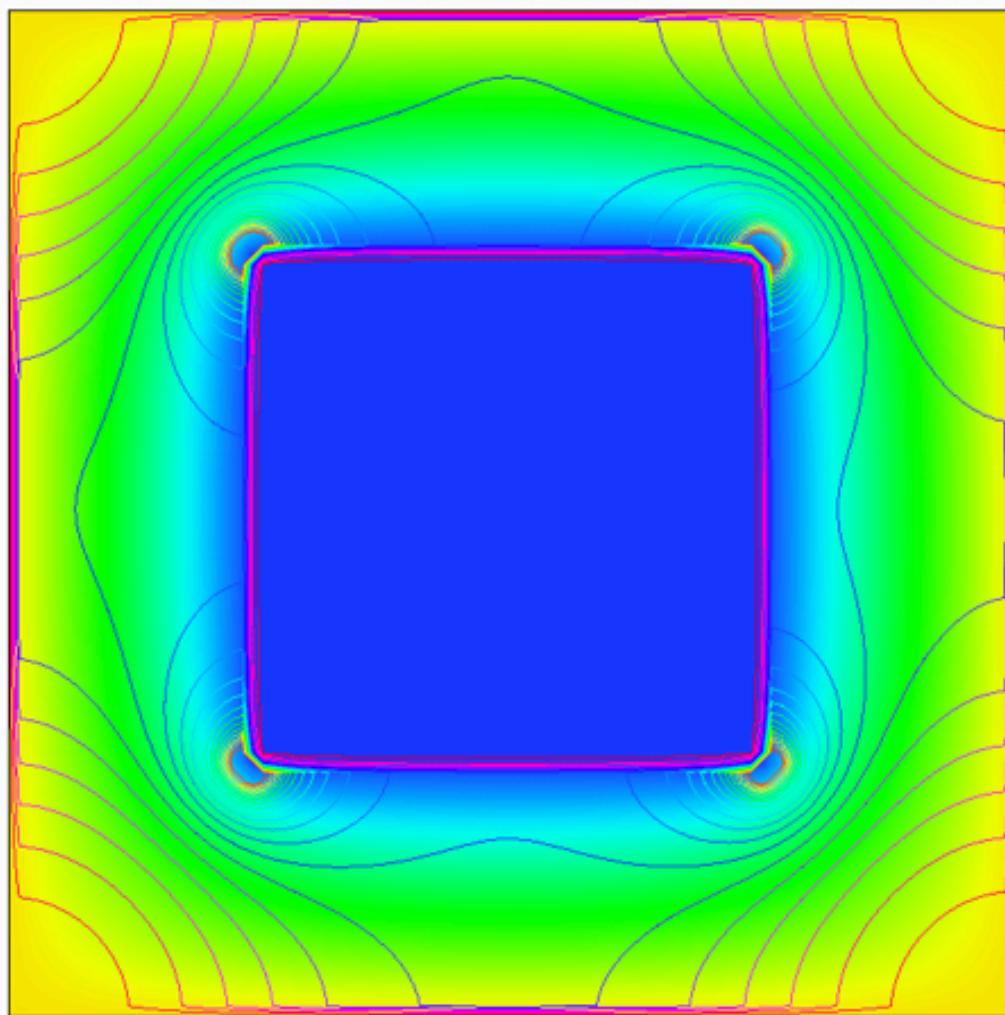


# POTENTIAL



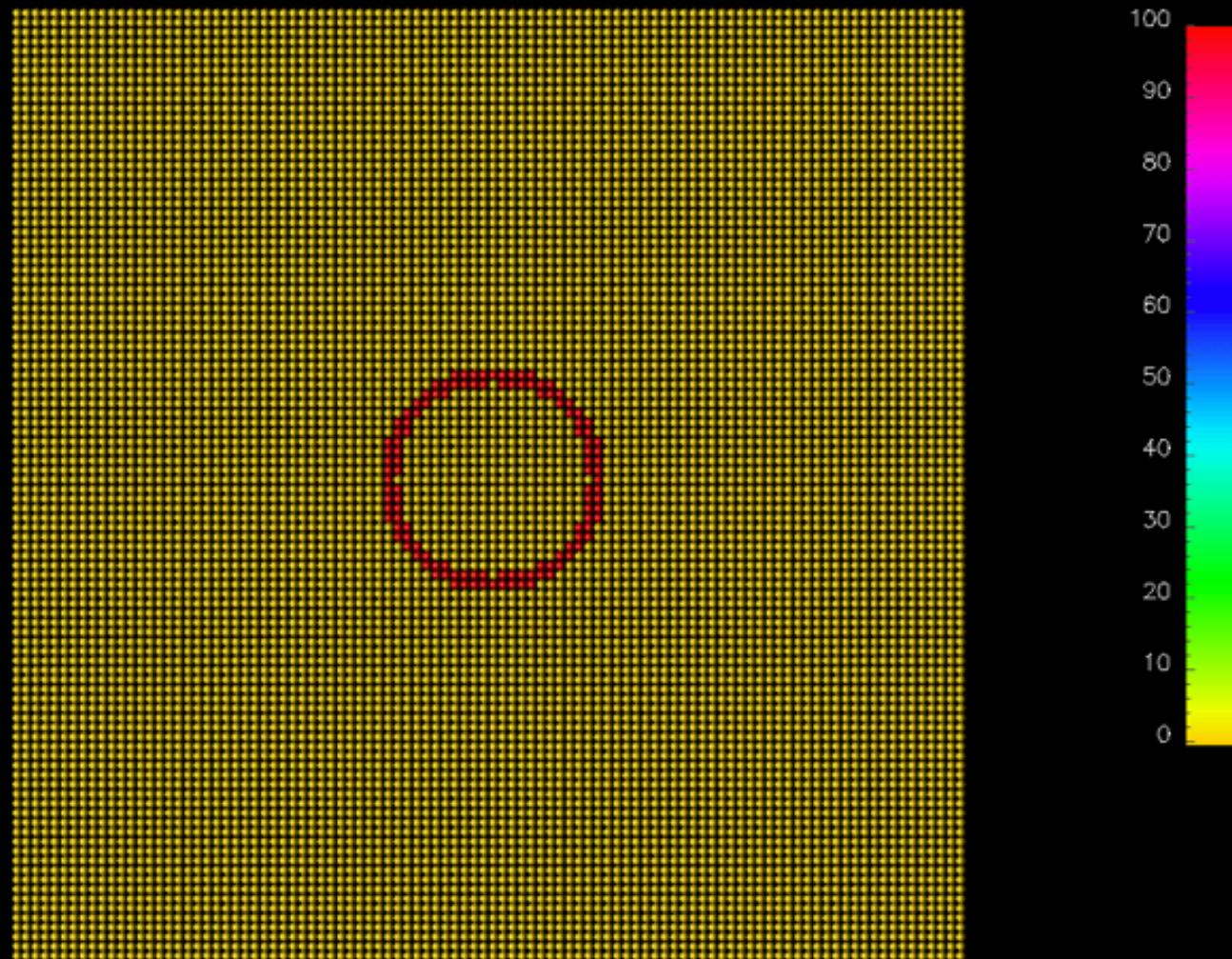
# ELECTRIC FIELD



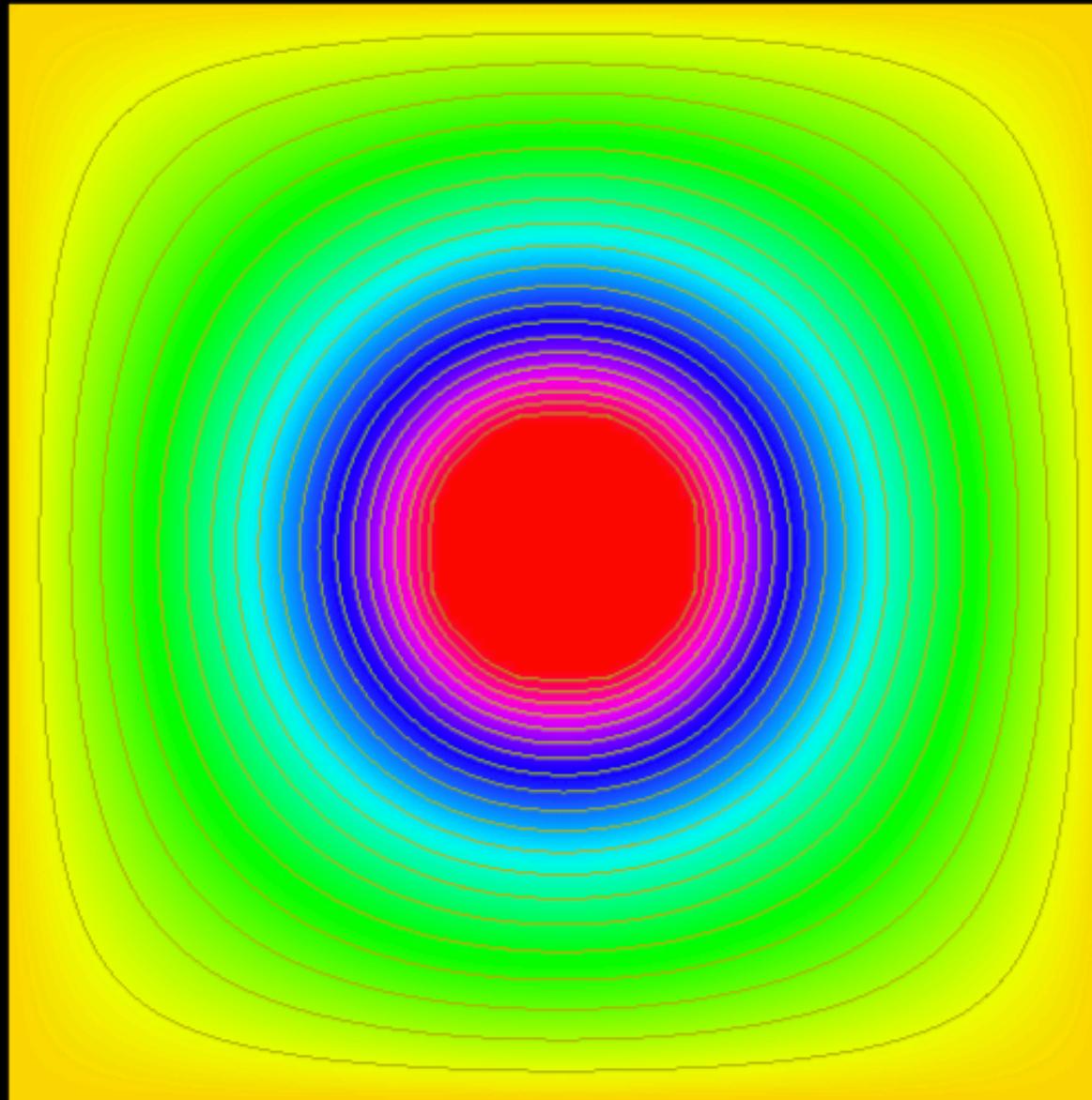


HOW CAN WE REDUCE  
THE ELECTRIC FIELD?

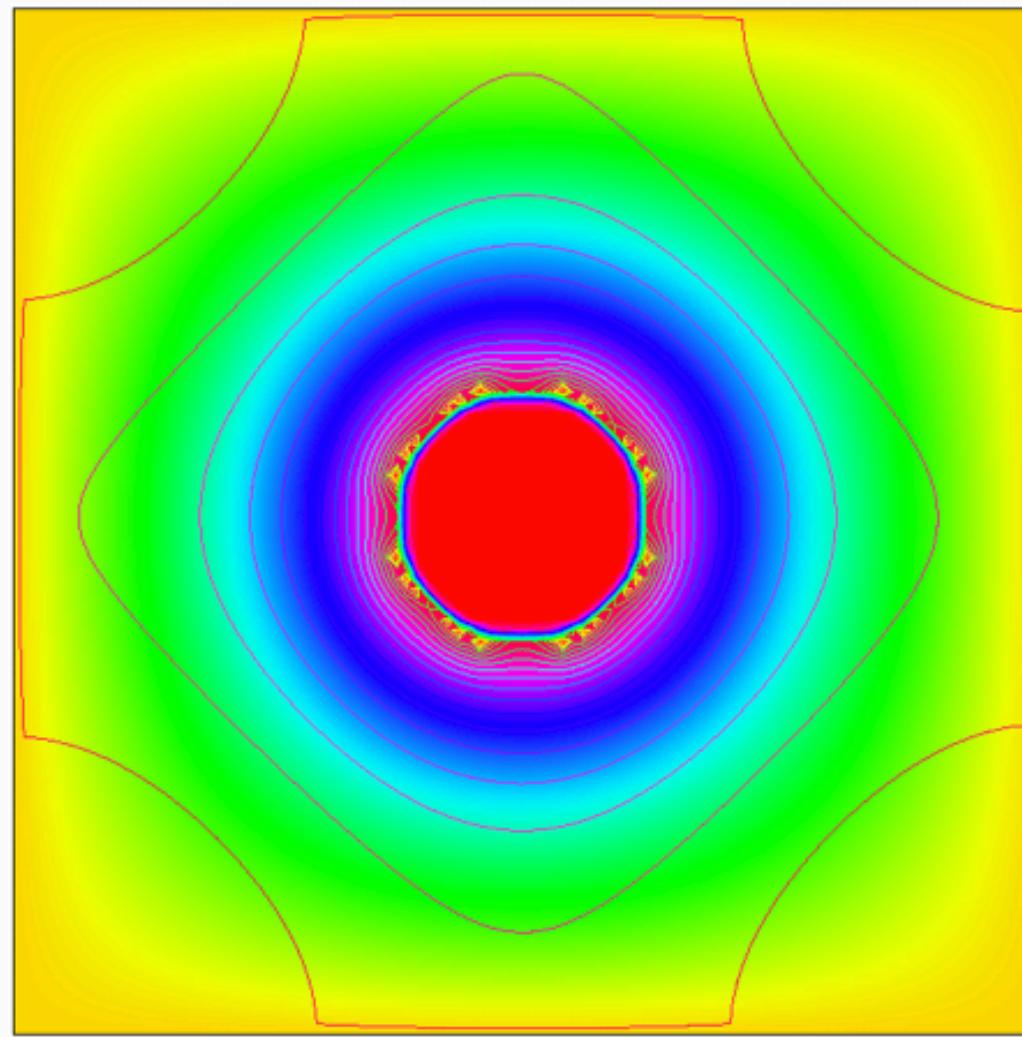
# INITIAL CONDITIONS



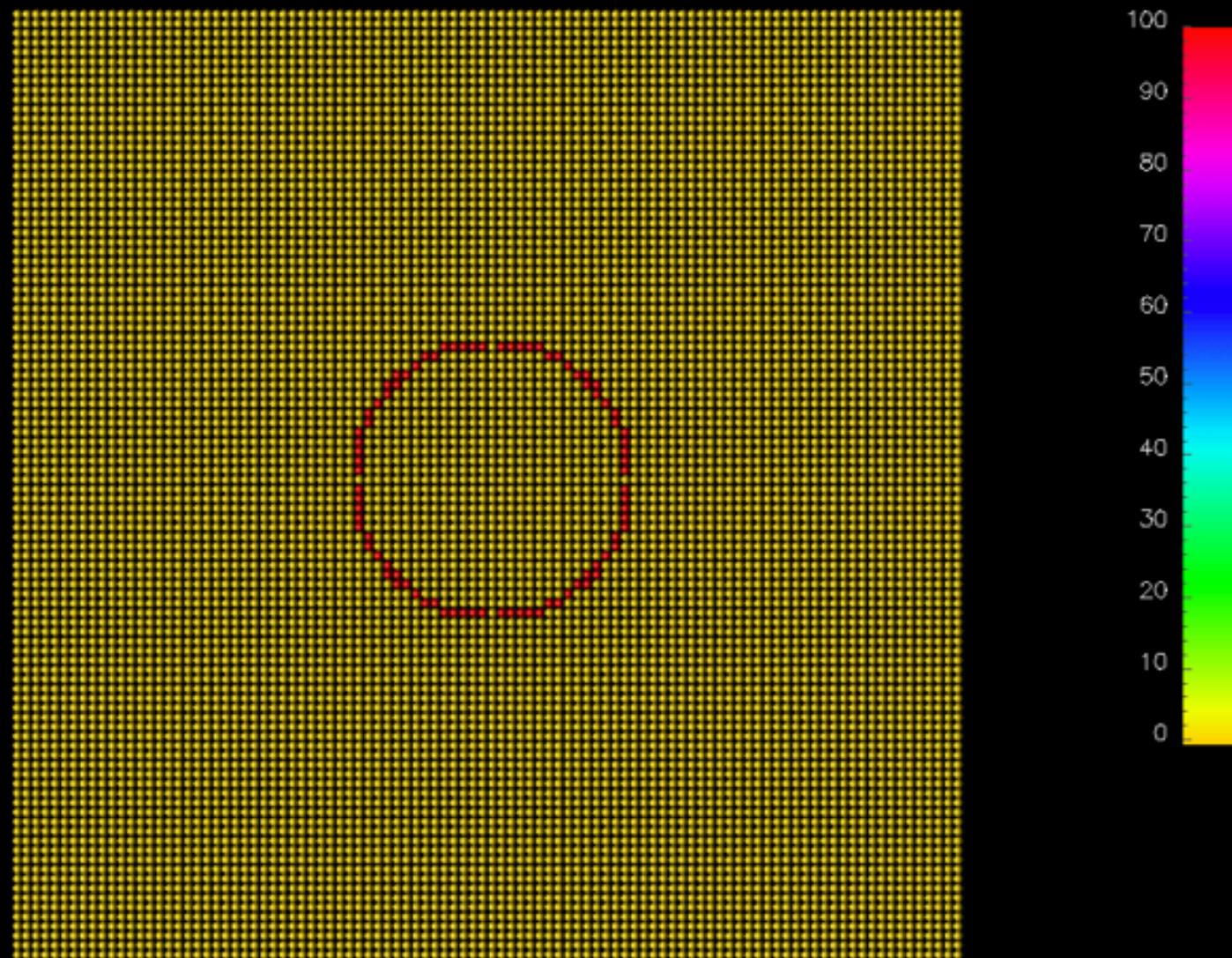
# CONVERGED POTENTIAL

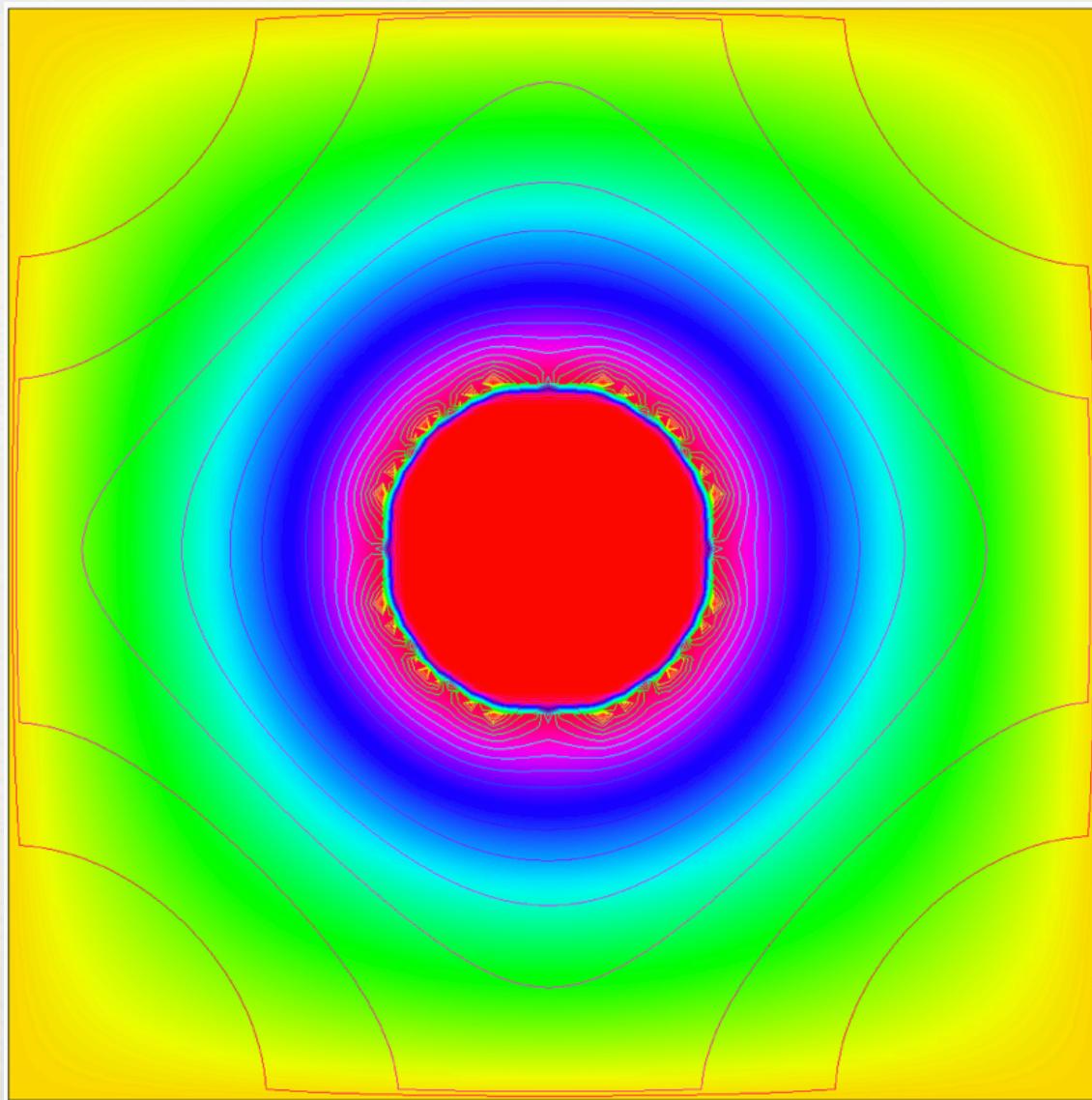


# FIELD

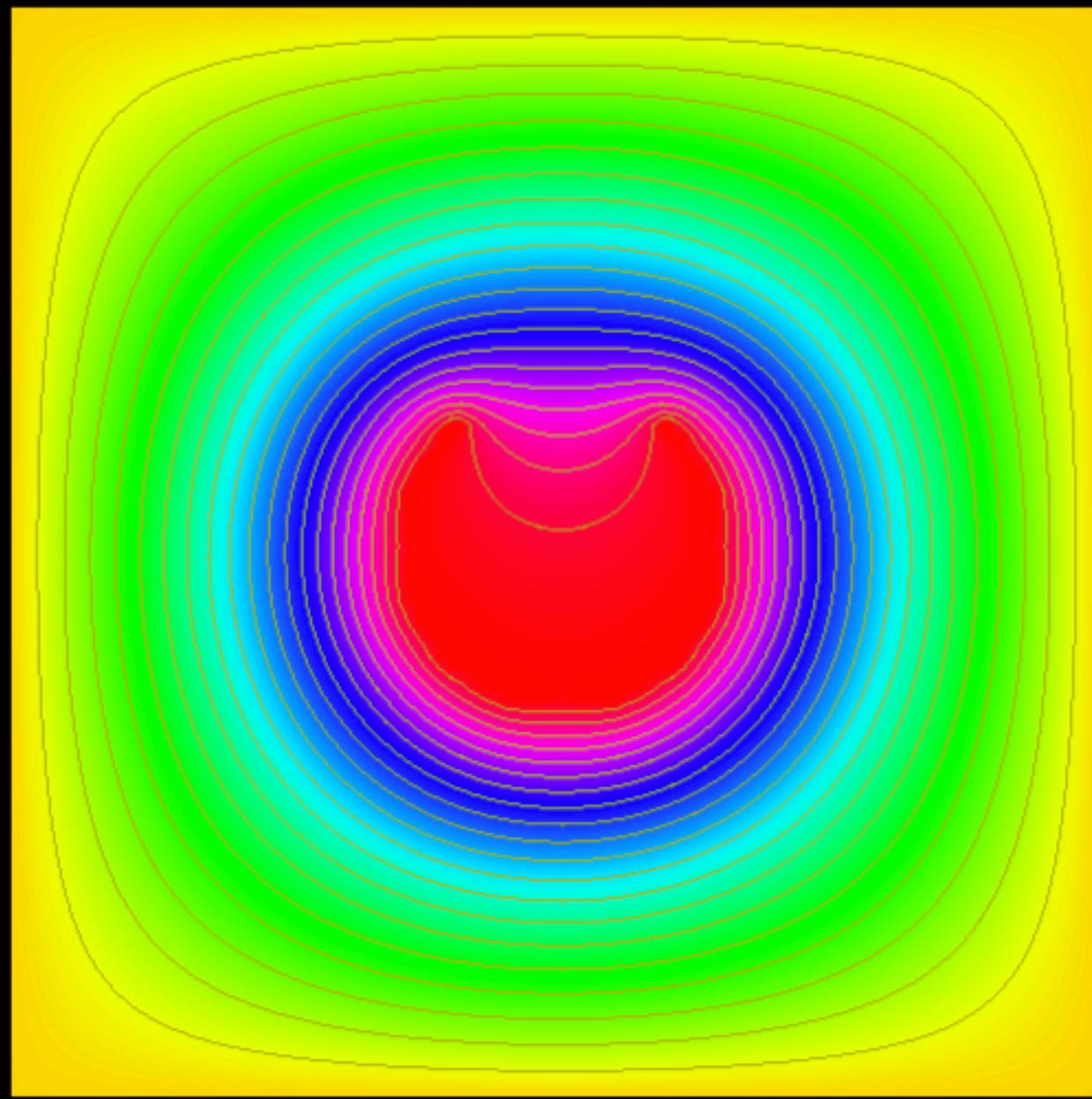


# INITIAL CONDITIONS (MODIFIED)

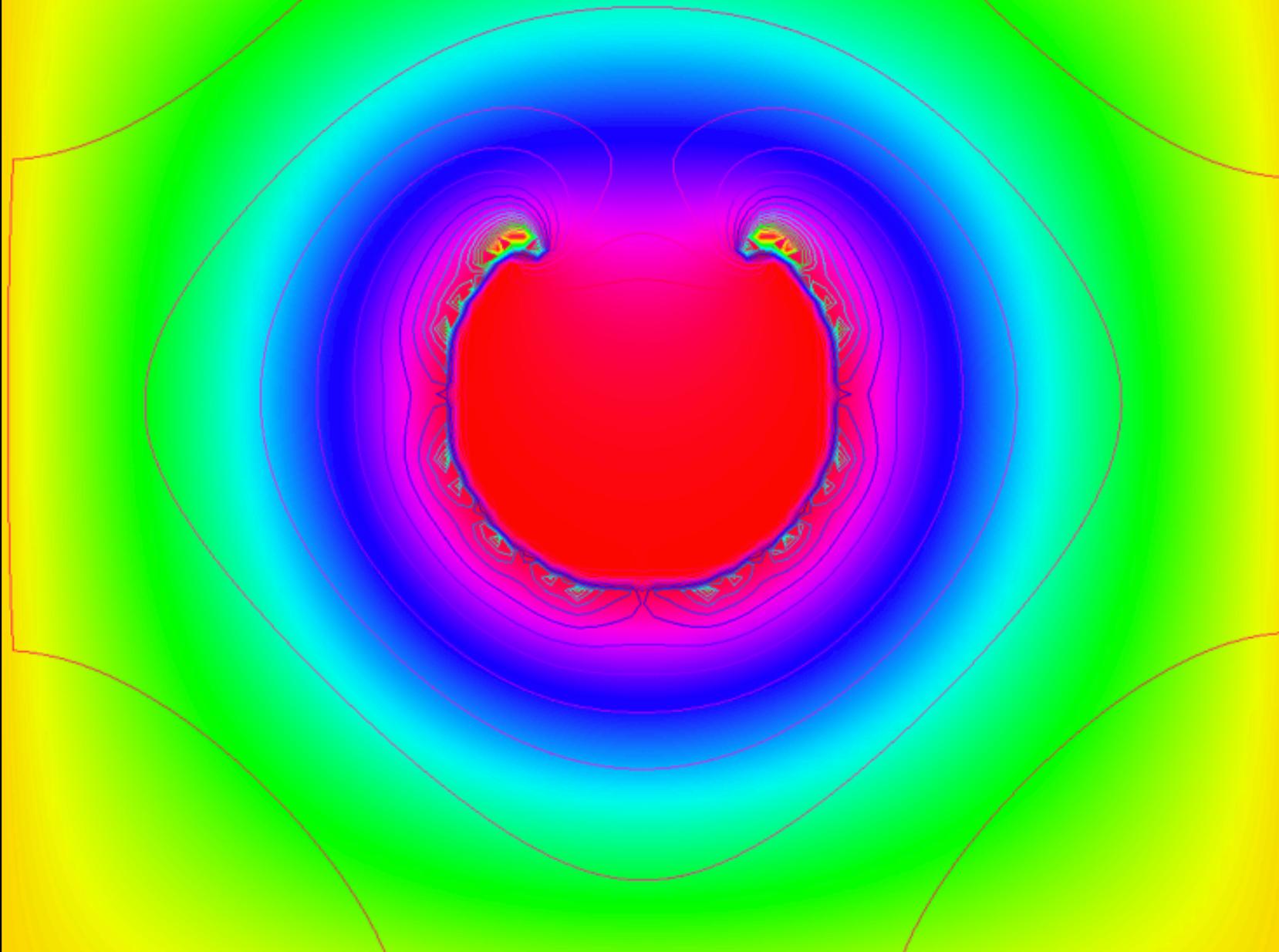


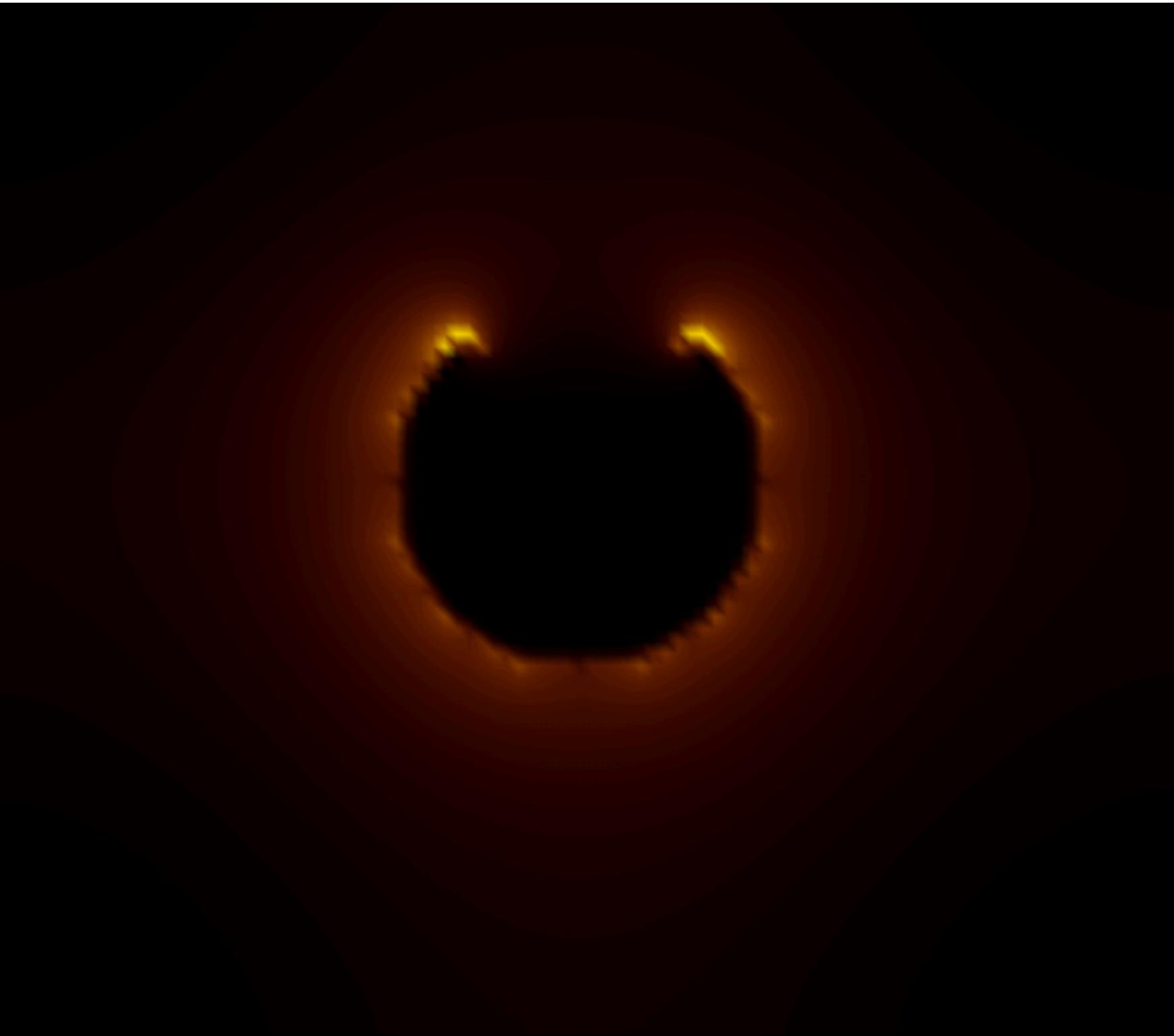


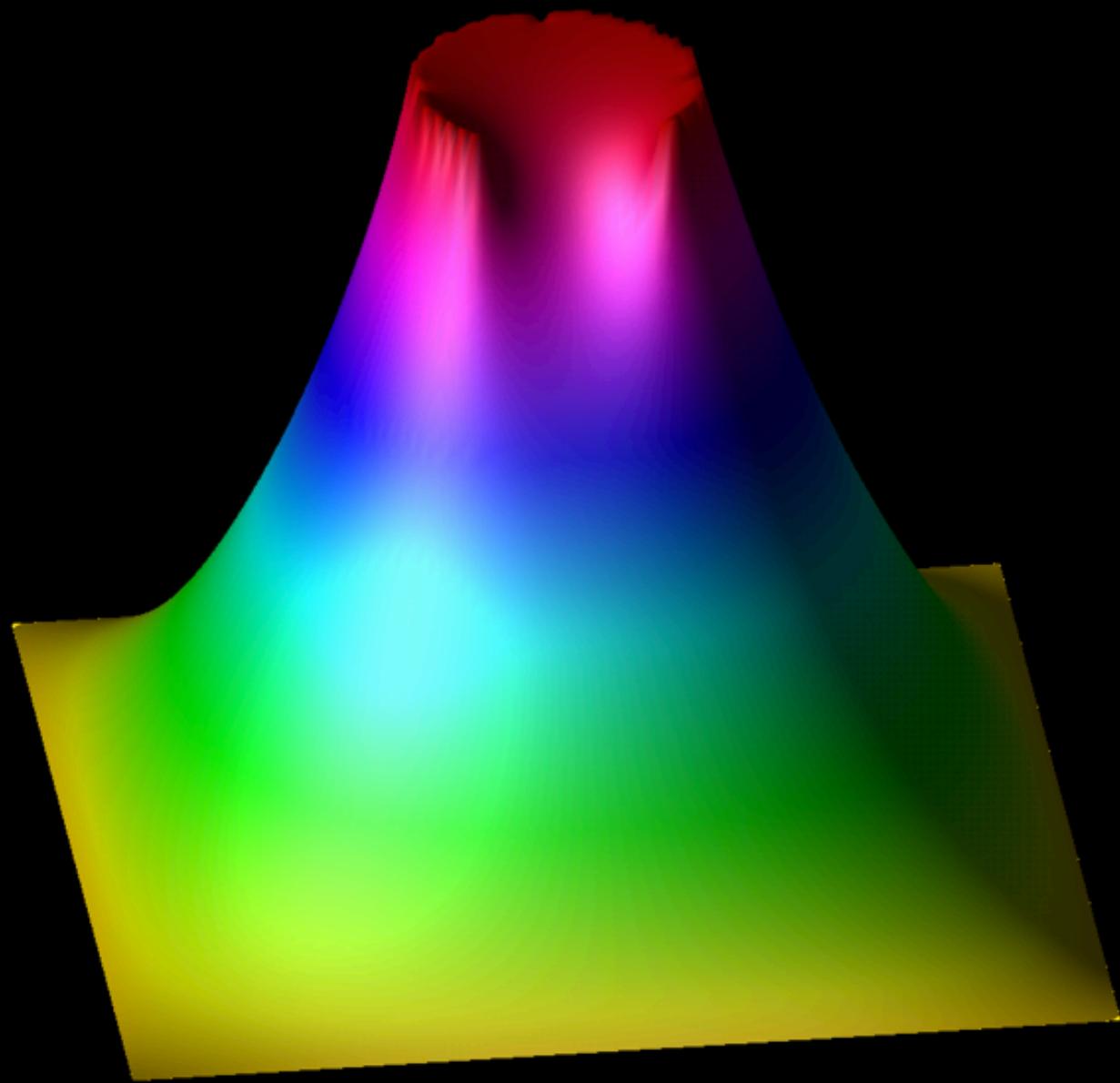
# WITH A HOLE



# FIELD PENETRATES THE HOLE



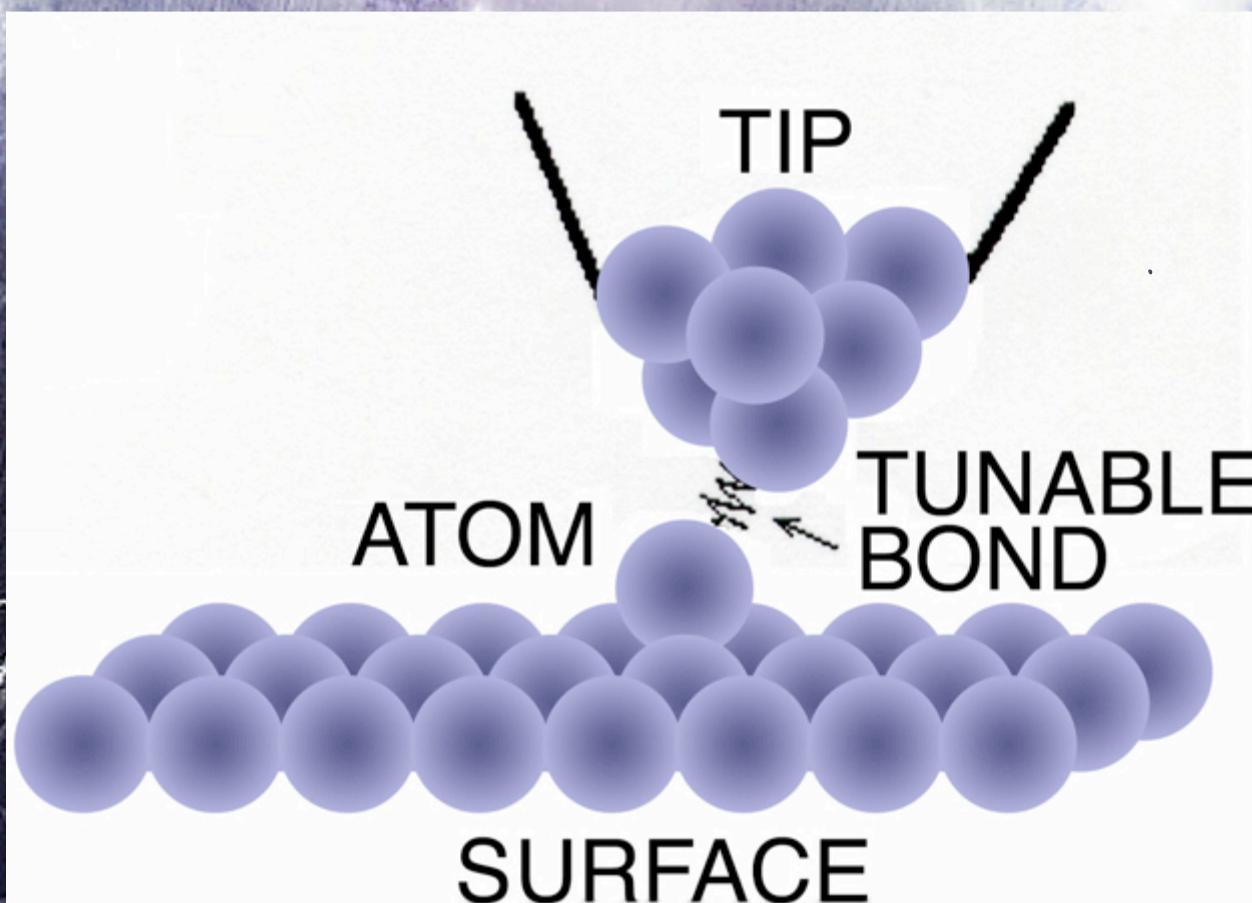




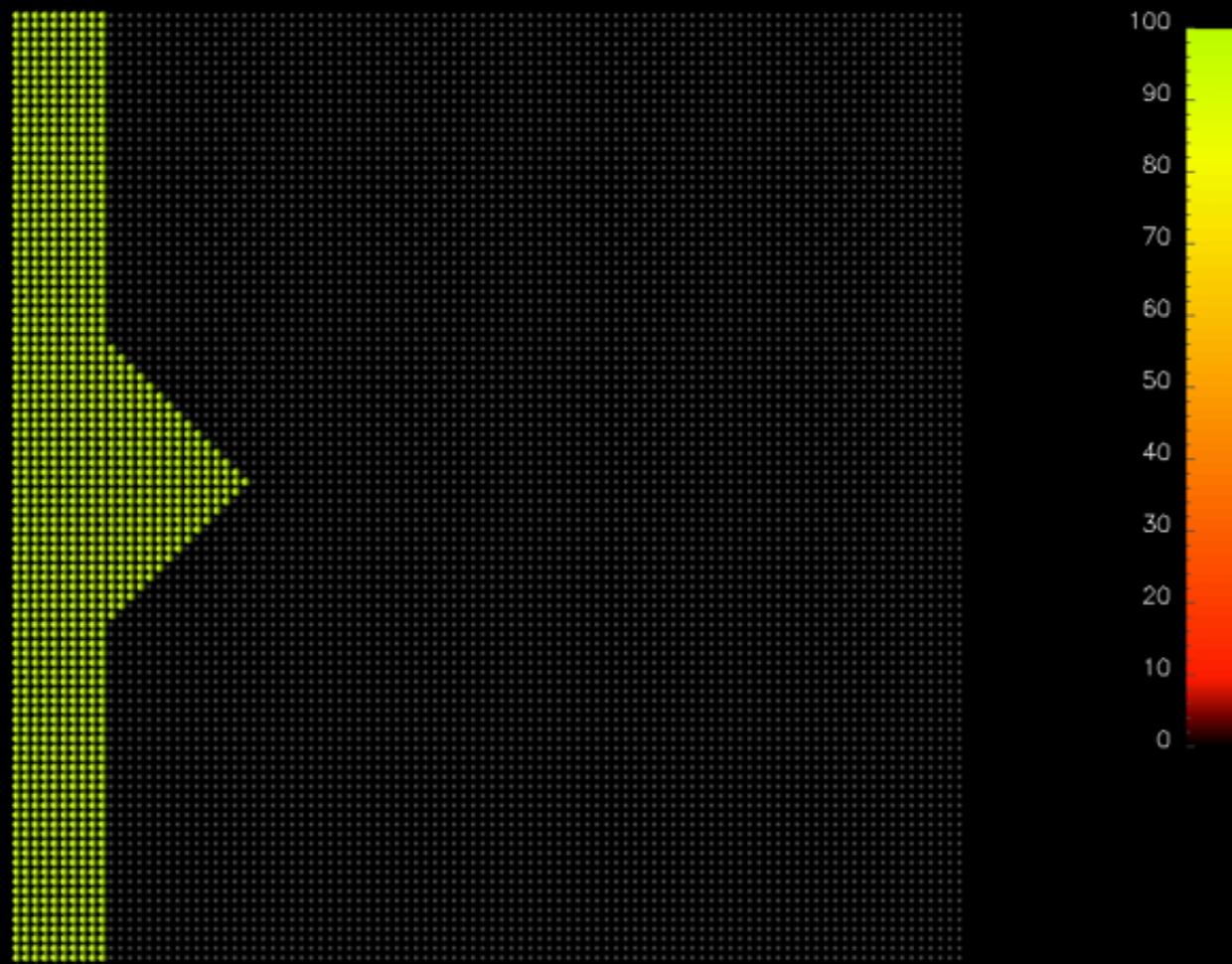
$\begin{matrix} z \\ \leftarrow \\ x \\ \downarrow \\ y \end{matrix}$

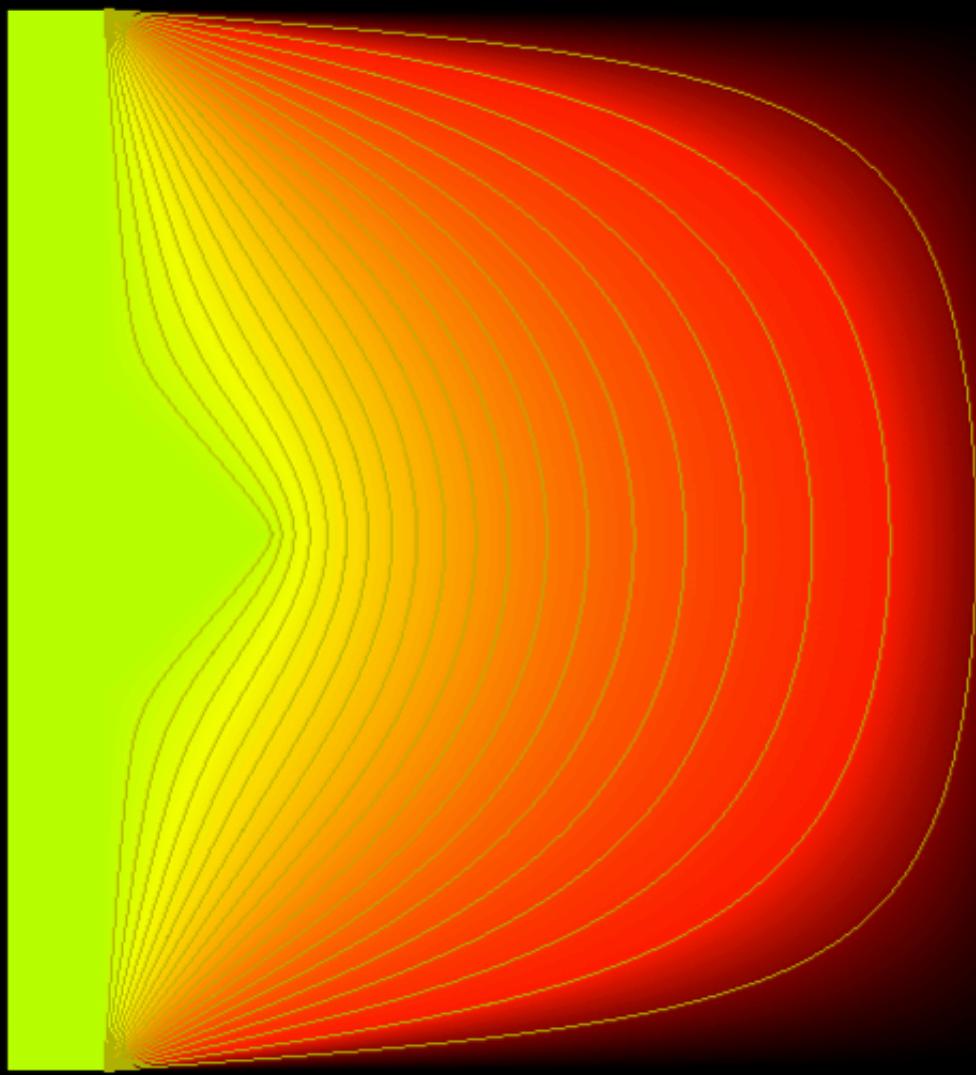


TIP EFFECT

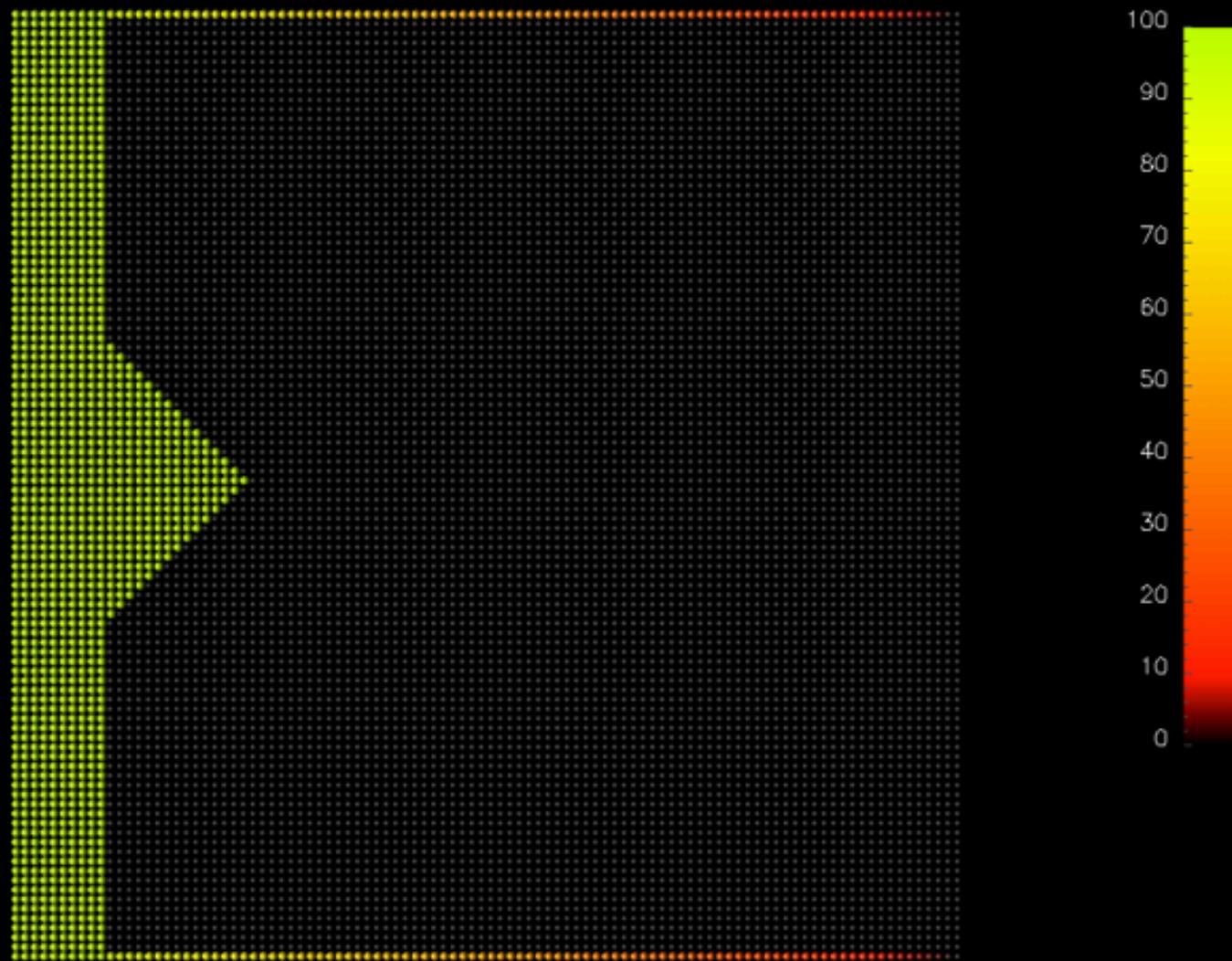


# INITIAL CONDITIONS

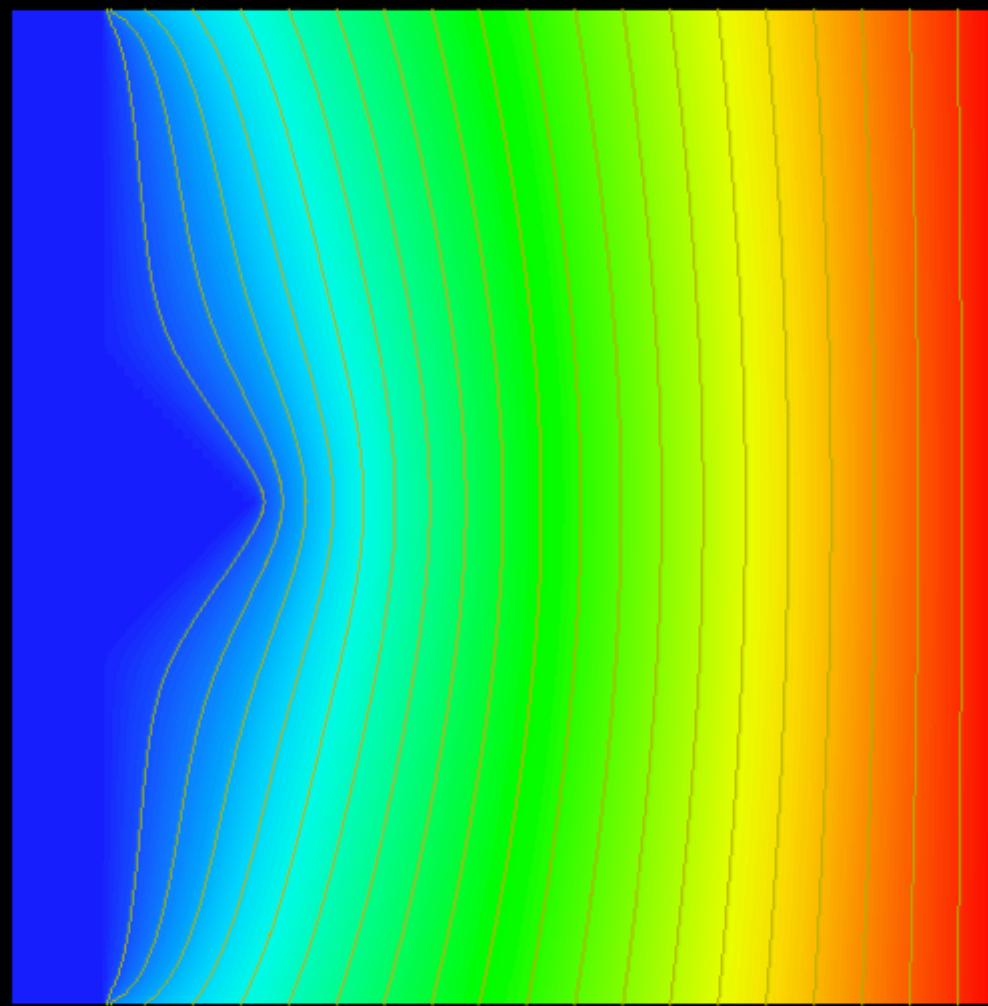




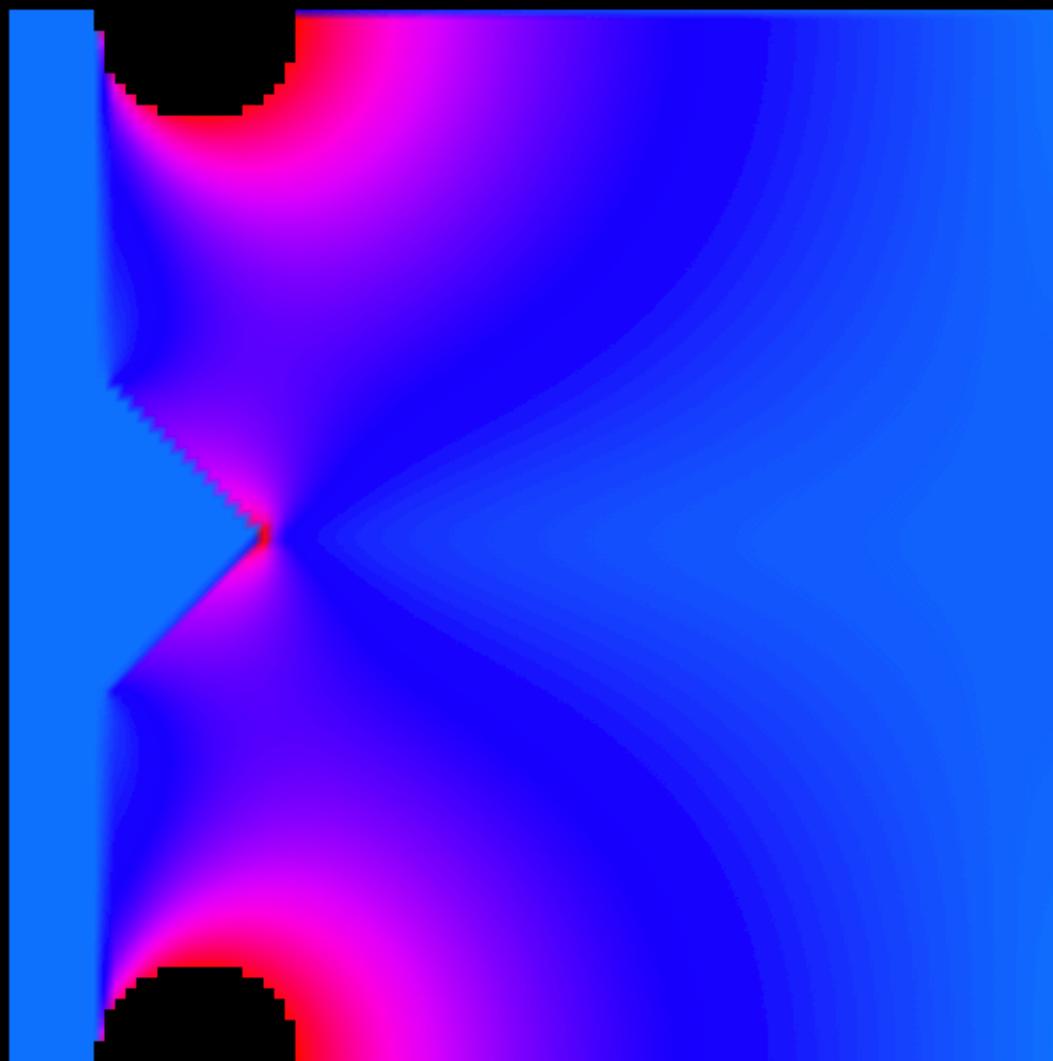
# MODIFIED INITIAL CONDITIONS



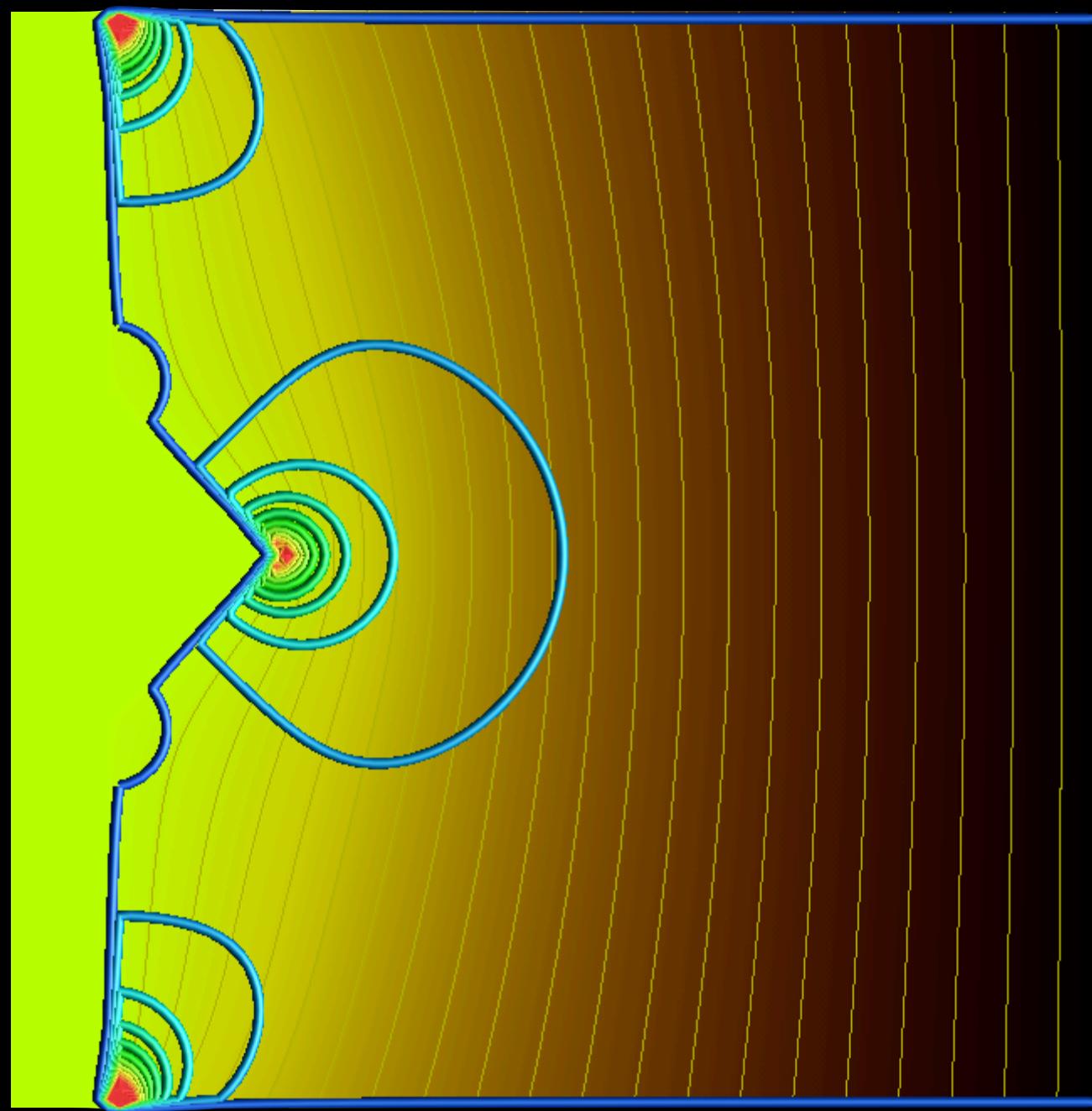
# CONVERGED POTENTIAL



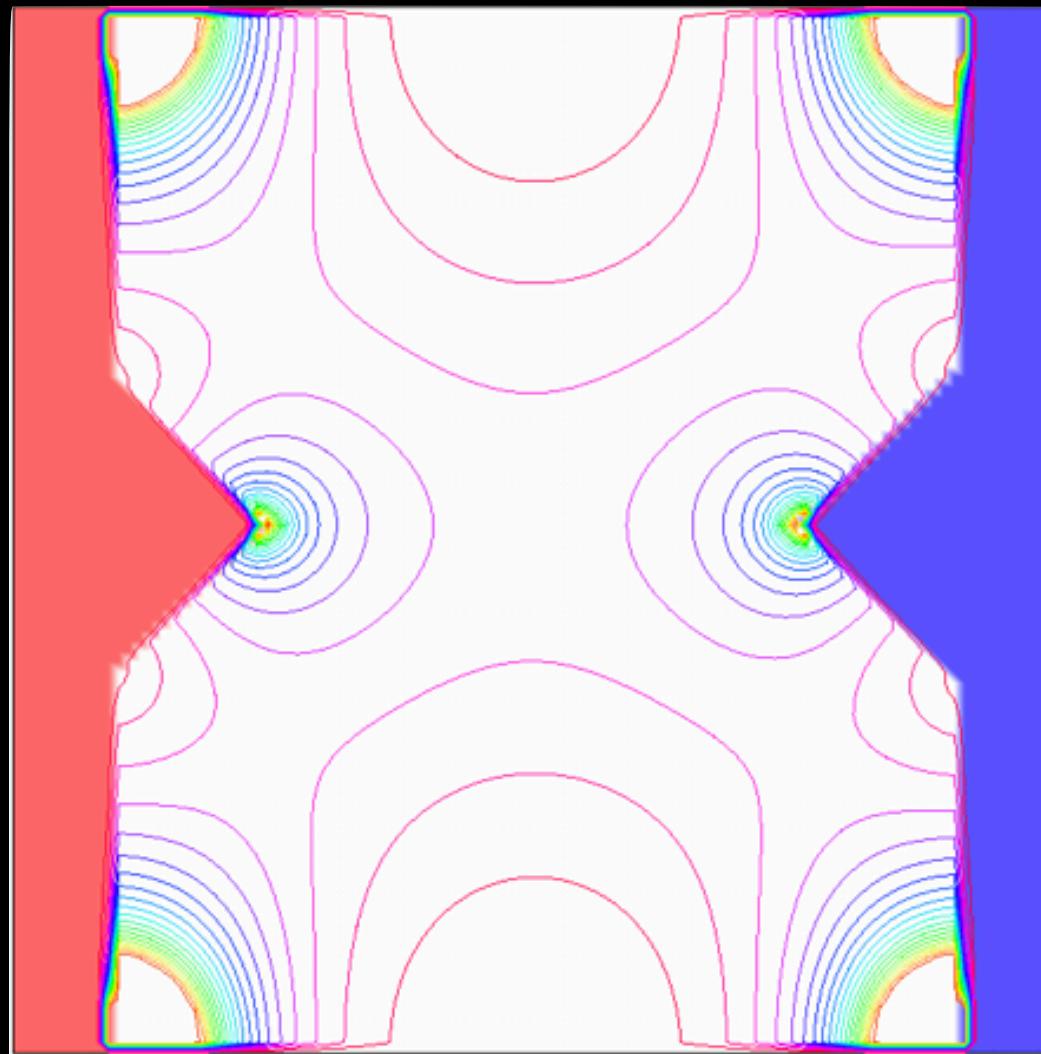
# ELECTRIC FIELD



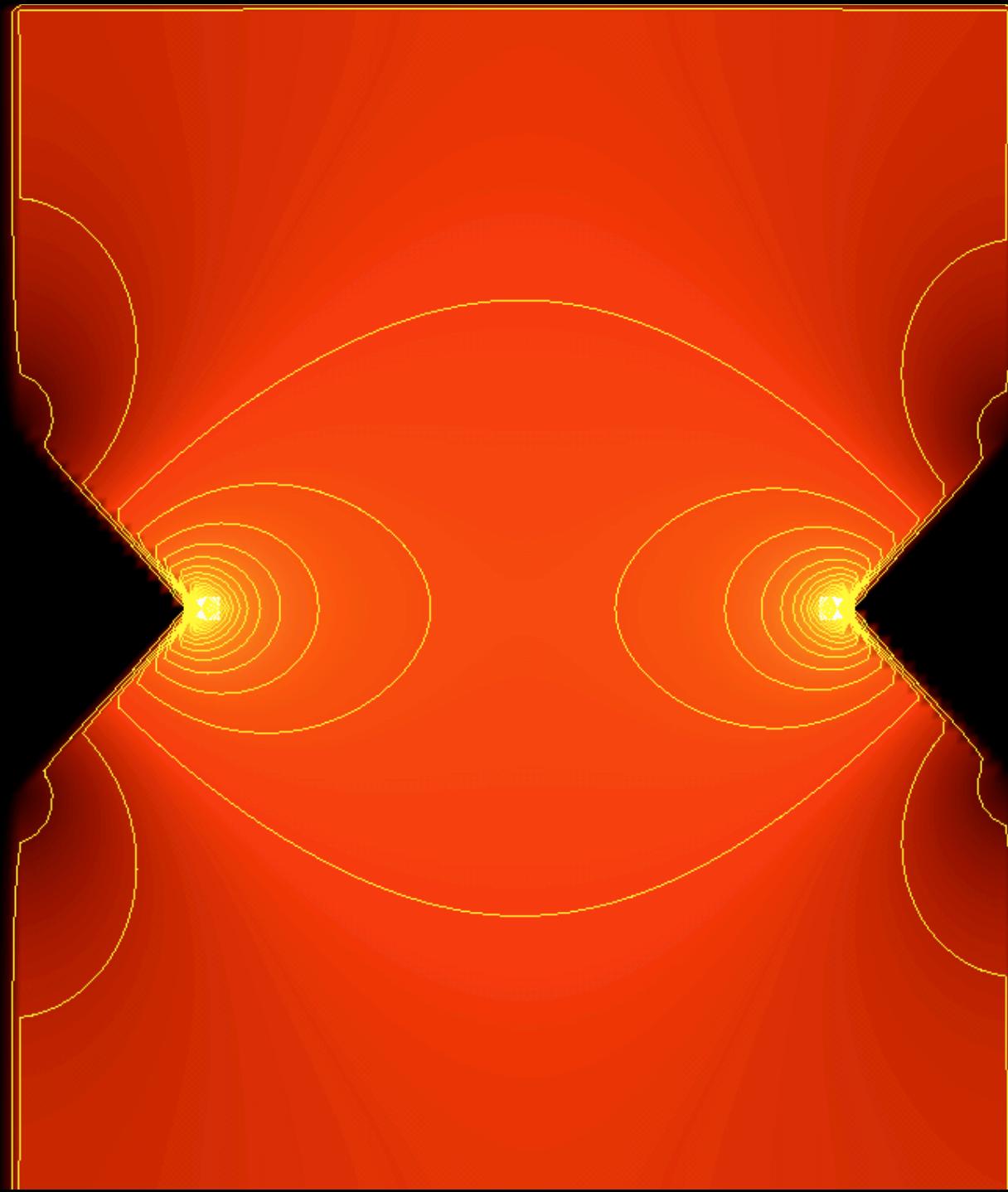
# ELECTRIC FIELD AND POTENTIAL

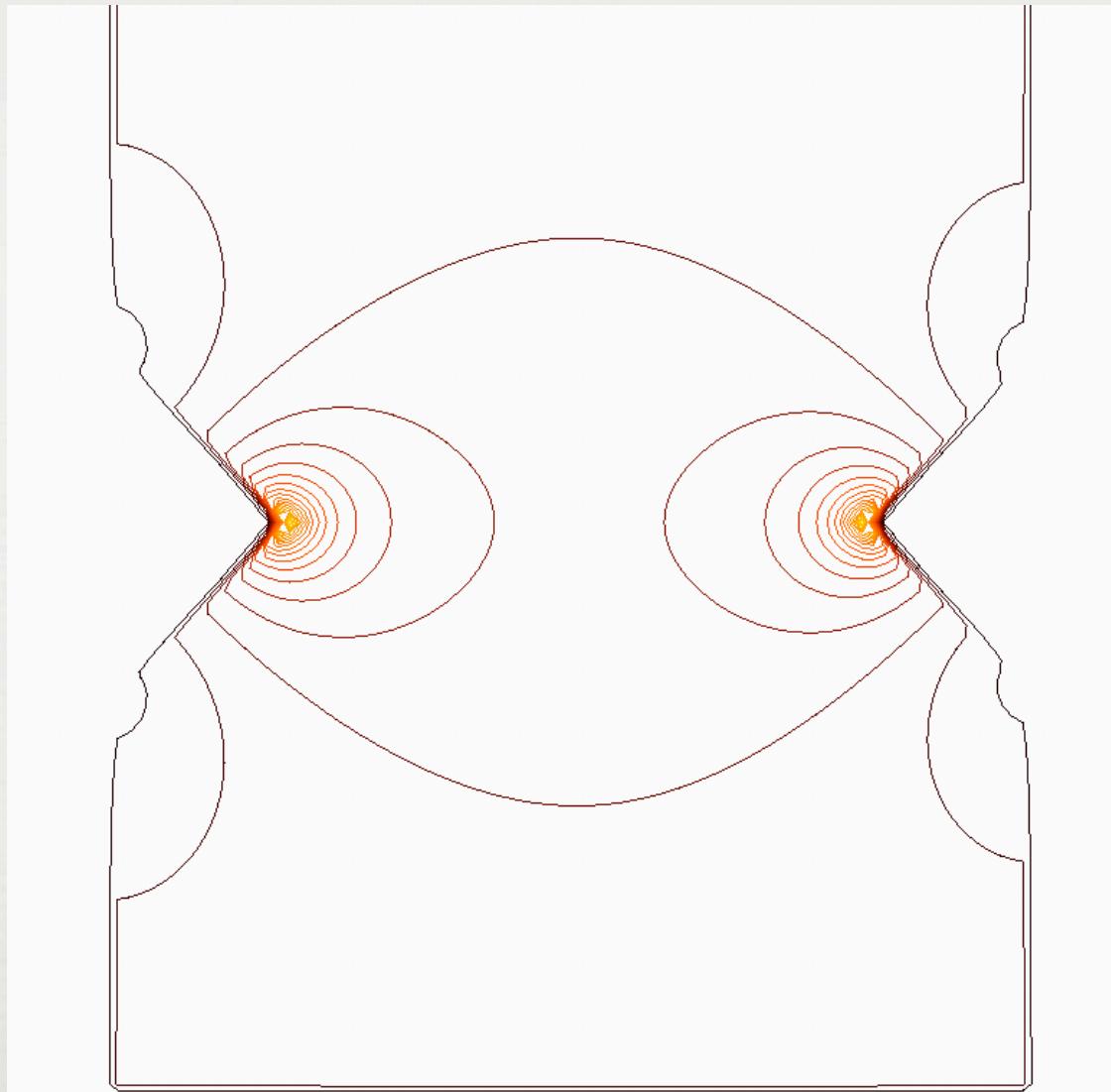


# TWO-TIPS AT OPPOSITE POTENTIALS



HOW CAN WE GET RID OF  
SIDE FIELD?



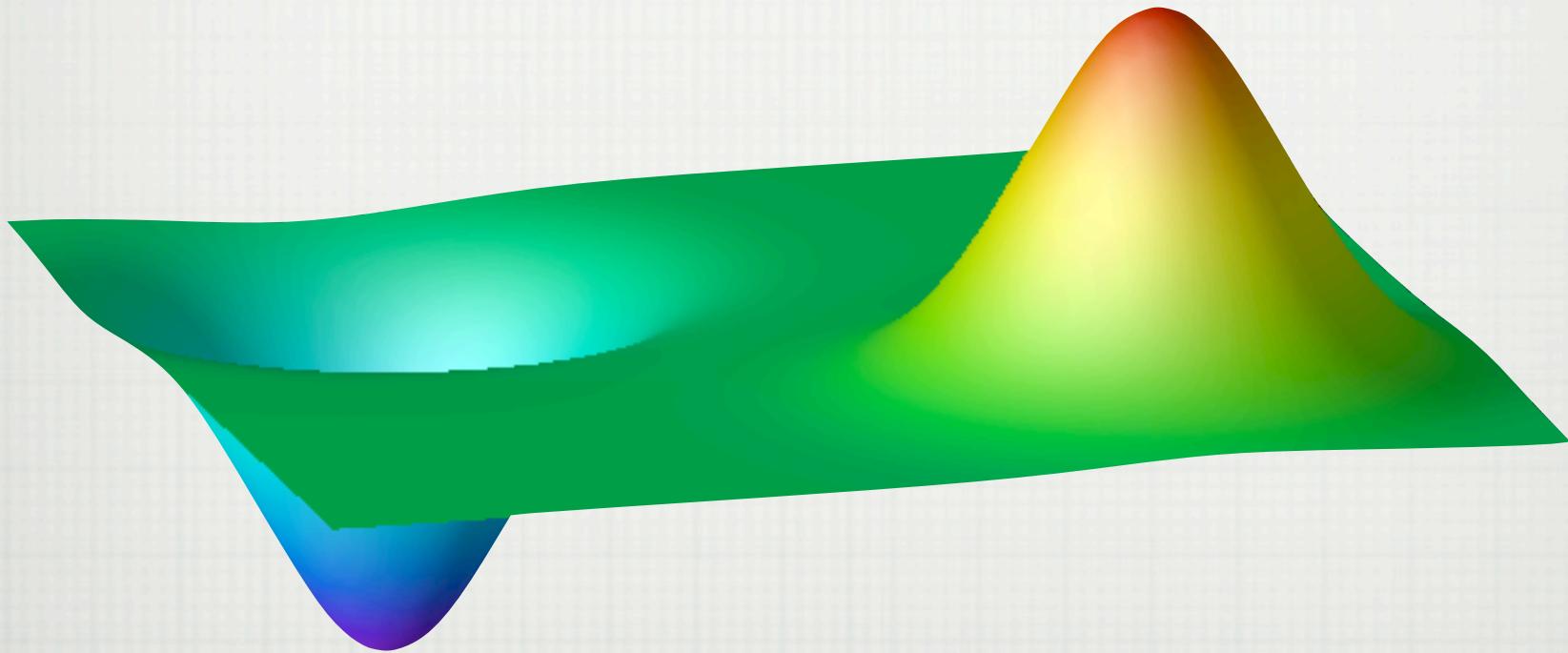


# POISSON EQUATION

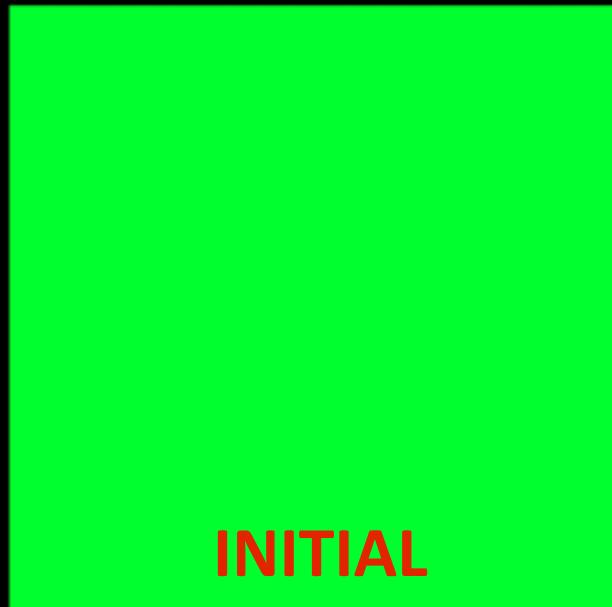
$$\frac{\partial^2 U(x, y)}{\partial x^2} + \frac{\partial^2 U(x, y)}{\partial y^2} = \begin{cases} 0 & \text{Laplace's equation} \\ -4\pi\rho(\mathbf{x}) & \text{Poisson's equation} \end{cases}$$

# POISSON: POTENTIAL OF A DIPOLE?

---



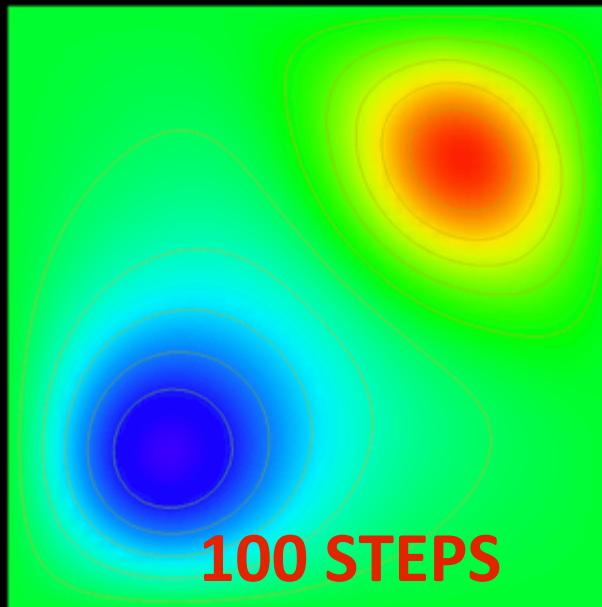
# POISSON EQUATION: DIPOLE



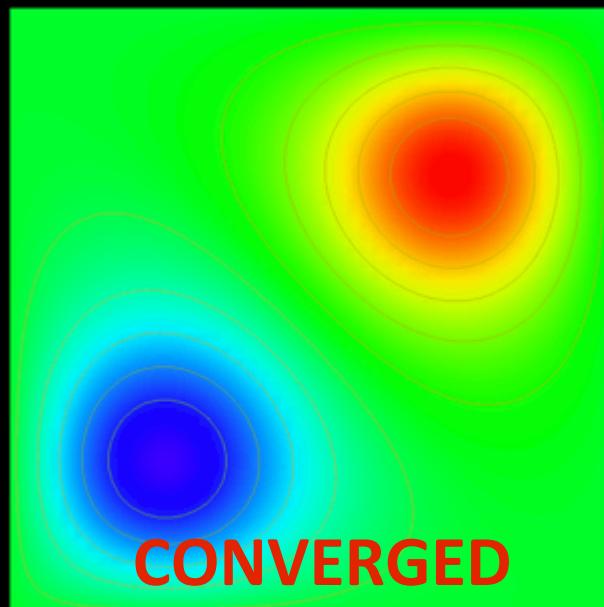
INITIAL



10 STEPS



100 STEPS



CONVERGED

# ACTUAL IMPLEMENTATION: C++ CODES

# BOUNDARY CONDITIONS

---

```
for (ix=0;ix<n;ix++){
    for (iy=0;iy<n;iy++){
        donotcompute[ix][iy]=false;
        func[ix][iy]=0.0;
    }
}
for (iy=0;iy<n;iy++){
    if(periodic==0){
        donotcompute[0][iy]=true;
        donotcompute[n-1][iy]=true;
    }
    func[0][iy]=0.0;
    func[n-1][iy]=0.0;

}
for (ix=0;ix<n;ix++){
    donotcompute[ix][0]=true;
    donotcompute[ix][n-1]=true;
    func[ix][0]=0.0;
    func[ix][n-1]=0.0;

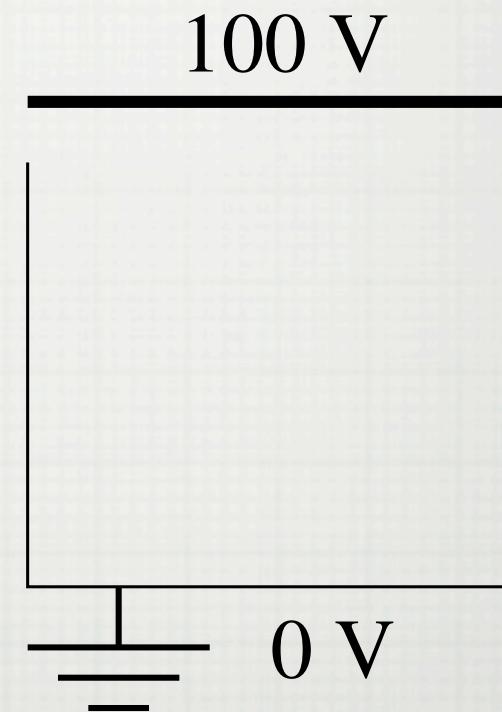
}
```

# SPECIAL CONDITIONS

---

```
for (ix=0;ix<n;ix++){
    for (iy=0;iy<n;iy++){
        donotcompute[ix][iy]=false;
        func[ix][iy]=0.0;
    }
}
for (iy=0;iy<n;iy++){
    if(periodic==0){
        donotcompute[0][iy]=true;
        donotcompute[n-1][iy]=true;
    }
    func[0][iy]=0.0;
    func[n-1][iy]=0.0;

}
for (ix=0;ix<n;ix++){
    donotcompute[ix][0]=true;
    donotcompute[ix][n-1]=true;
    func[ix][0]=100.0;
    func[ix][n-1]=0.0;
}
```



# SPECIAL CONDITIONS: DOUBLE-PLATE

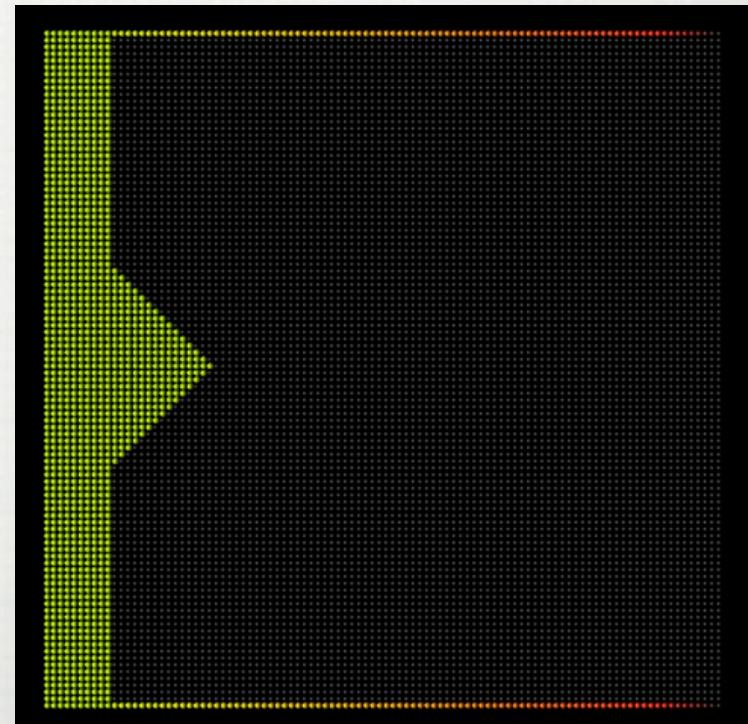
---

```
for (ix=10; ix<n-10; ix++){
    for(iy=20; iy<26; iy++){
        donotcompute[ix][iy]=true;
        donotcompute[ix][n-iy]=true;
        func[ix][iy]=result[ix][iy];
        func[ix][n-iy]=result[ix][n-iy];
    }
}
```

# SPECIAL CONDITIONS: TIP

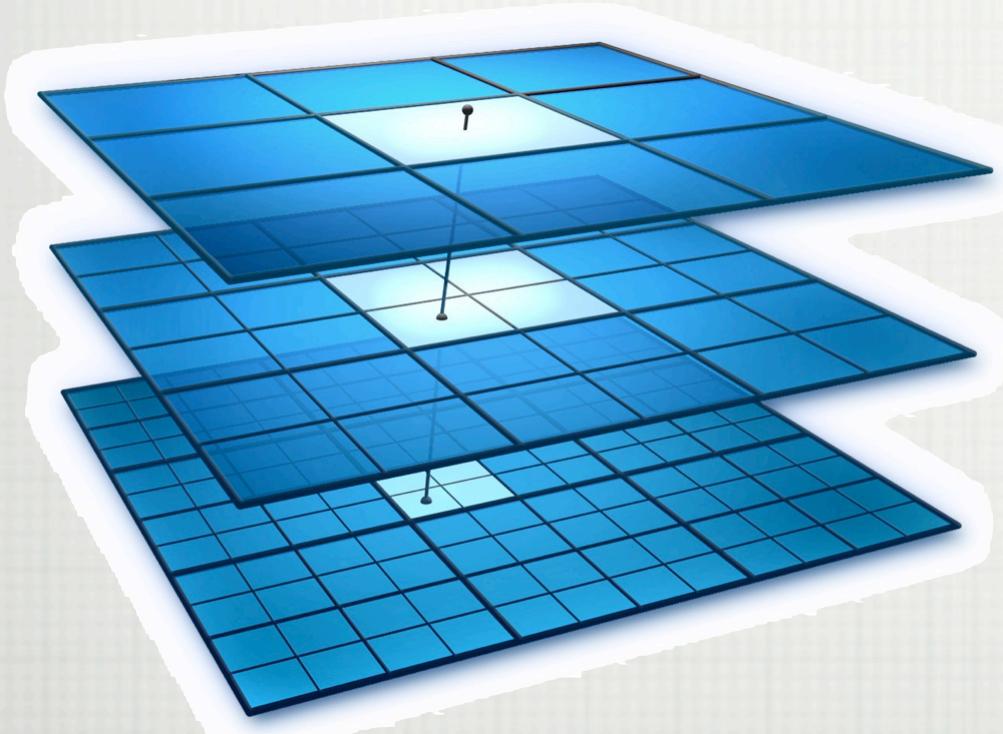
---

```
for (ix=0;ix<n;ix++){  
    for(iy=0; iy<n;iy++){  
        if(iy<ix-25 && iy < -ix+75)  
        {  
            donotcompute[ix][iy]=true;  
            func[ix][iy]=100;  
        }  
    }  
}  
for (ix=0;ix<n;ix++){  
    for (iy=0;iy<10;iy++){  
        donotcompute[ix][iy]=true;  
        func[ix][iy]=100.0;  
    }  
}
```



# HOW TO MAKE IT EVEN FASTER?

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- Solve on coarse grid and use the converged result on finer and finer grid.
- The preconditioning avoid spending too much time on the early stages of convergence

# SUMMARY

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- We covered relaxation methods for Laplace and Poisson's equation
- Over- and under- relaxation yield vast improvements to overall speed
- The “art” is to set up boundary conditions