## MiMes: Misalignment Mechanism Solver

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Manchester U

23/11/2021 El Journal Club más Sabroso

### Outline

- Axion Dark Matter Why particle dark matter
  - The dark matter particle
  - The axion (like) particle
- Calculating the Relic Abundance
  - The axion EOM
  - How hard can it be?
  - Initial conditions
  - (Bad) Analytical approximations
- Need for accuracy, speed, and automation
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### **Axion Dark Matter**

- Axion Dark Matter

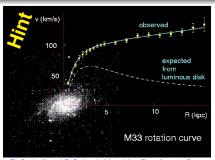
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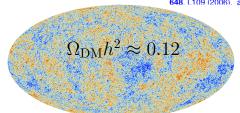
# Why particle dark matter



E. Corbelli and P. Salucci, Mon. Not. Roy. Astron. Soc. 311 441 (2000), arXiv:astro-ph/9909252.



M. Markevitch, ESA Spec. Publ. 604 (2006) 723, astro-ph/0511345.Clowe, Bradac, et. al. Astrophys. J. 648. L109 (2006). astro-ph/0608407



N. Aghanim et al. [Planck Collaboration], arXiv:1807.06209 [astro-ph.CO].

# The dark matter particle

"Έν οἴδα, ὅτι οὐδὲν οἴδα." "I know one thing, that I know nothing."

-Socrates

# The dark matter particle

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- Gravitational interactions.
- Mostly electrically neutral.
- Stable or very slow decay rate.
- Non-Baryonic.
- Cold/Warm and non-relativistic today.

## The axion (like) particle

Notably, the original Axion was originally introduced in order to solve the *strong-CP problem* of the SM. Axion-Like-Particles (ALPs) arise in a number of new physics models, beyond the SM.

#### Axions and ALPs generally:

- Have suppressed interactions with photons.
- Are (mostly) stable.
- Were non-relativistic around the epoch of structure formation.

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#### Maybe DM has Axionic nature!

# Calculating the Relic Abundance

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### The axion EOM

Axions and ALPs follow a similar equation of motion (EOM):

$$\left(\frac{d^2}{dt^2} + 3H(t) \; \frac{d}{dt}\right)\theta(t) + \tilde{m}_a{}^2(t) \; \sin\theta(t) = 0 \; , \label{eq:delta_total}$$

where  $\theta = A f_a$ , with A the axion filed, and  $f_a$  some energy scale that characterises the potential (Peccei-Quinn breaking scale).

## How hard can it be?

Hard (in general).

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The classical analogue is the dumped pendulum with both frequency (length) and friction being time-dependent:

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- There are no constants of motion (wait a minute).
- No package/library/program available!

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MiMeS simulates the evolution of the Axion/ALP, for (virtually) any cosmological scenario and Axion/ALP (thermal) mass.

### **Initial conditions**

Some time at the very early Universe,  $\tilde{m}_a \ll H(T)$ , <sup>1</sup> with

$$\ddot{\theta} + 3H \; \dot{\theta} \approx 0 \; .$$

The solution is

$$\theta = \theta_{\text{ini}} + C \int_0^t dt' \left( \frac{a(t'=0)}{a(t')} \right)^3.$$

So,  $\dot{\theta} \sim a^{-3}$ . Since we are interested in  $\theta$  at much later times (once the potential becomes relevant),  $\dot{\theta} \approx 0$ . Therefore, we can start integration at some point  $(t=t_{\rm ini})$  with  $3H \gg \tilde{m}_a$ , and set  $\theta(t=t_{\rm ini}) = \theta_{\rm ini}$  and  $\dot{\theta}(t=t_{\rm ini}) = 0$ .

<sup>1</sup> This is an assumption that MiMeS has to make, for the sake of generality.

<sup>&</sup>lt;sup>2</sup> Standard misalignment mechanism. For the kinetic one see R. T. Co, L. J. Hall and K. Harigaya, Phys. Rev. Lett. 124 (2020) no.25, 251802 [arXiv:1910.14152 [hep-ph]], C. F. Chang and Y. Cui, Phys. Rev. D 102 (2020) no.1, 015003 [arXiv:1911.11885 [hep-ph]], or B. Barman, N. Bernal, N. Ramberg and L. Visinelli, [arXiv:2111.03677 [hep-ph]].

Once we agree on the initial conditions, we move to the next important things:

<sup>&</sup>lt;sup>3</sup> Defined from  $s(T_0) = \gamma a_{\rm osc}^3 s_{\rm osc}$ .

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Then, we get the "WKB"-approximate solution

$$\theta(t) \approx \theta_{\rm ini} \left(\frac{3}{4}\right)^{1/4} \sqrt{\frac{\tilde{m}_a(T_{\rm osc})}{\tilde{m}_a(T)}} \left(\frac{a}{a_{\rm osc}}\right)^{-3/2} \, \cos\left(\int_{t_{\rm osc}}^t dt' \; \tilde{m}_a(t')\right) \; . \label{eq:theta}$$

The advantage of this approximation is that we get an easy formula for the axion/ALP energy density today:

$$\rho_{a,0} = \gamma^{-1} \frac{s_0}{s_{\text{occ}}} \frac{1}{2} f_a^2 m_a \tilde{m}_{a,\text{osc}} \theta_{\text{ini}}^2$$

where  $\gamma$  the amount of entropy injection between  $T_{\rm osc}$  and today. <sup>3</sup>

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## Need for accuracy, speed, and automation

#### Serious disadvantages of the approximate results:

- The approximations can be tested against numerical results in a case-by-case basis; there is no way to tell if they will work in new models and cosmological scenarios.
- There is no available tool that can help us reproduce published results obtained by numerical integration; people use their own private code.
- If someone wants to simply see if an ALP model is compatible with a cosmological scenario, they have to develop their own private code; the overall effort of the community increases.

#### MiMeS:

- Easy to use; anyone can run it and see if their model can work.
- ullet Reasonably fast; less than  $0.05\ s$  for the scenarios tested.
- Tools that can help determine if the algorithm is accurate enough.
- The user provides too much input. This helps the user determine whether convergence of the algorithm is consistent.

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### MiMeS: General

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### MiMes: Under the hood

MiMeS relies on Nabbodes <sup>4</sup> for the numerical integration, and SimpleSplines <sup>5</sup> for the various interpolations.

#### Advantages:

- You only need to have the standard C++ library.
- The two libraries are developed by myself, so their integration with MiMes is seamless.
- There is always going to be a compatible version of these libraries that works with MiMeS.

#### Disadvantages:

- These are not well tested libraries.
- No community of contributors; if it doesn't work, I have to fix it.
- Slow development.

<sup>4</sup> https://github.com/dkaramit/NaBBODES.

<sup>5</sup> https://github.com/dkaramit/SimpleSplines.

#### MiMes: Notation

MiMeS uses a notation suitable (any) underlying cosmology, since it is up to the user to define the cosmological evolution. First, we define

$$u \equiv \log \left( a/a_{\rm ini} \right) ,$$

with  $a_{\rm ini}$  some initial value of the scale factor. <sup>6</sup> in order to express the time derivatives as

$$\frac{d}{dt} \to H \frac{d}{du}$$
,  $\frac{d^2}{dt^2} \to H^2 \left( \frac{d^2}{du^2} + \frac{1}{2} \frac{d \log H^2}{du} \frac{d}{du} \right)$ .

Then, we express the EOM as a system of first order ordinary differential equations

$$\frac{d\zeta}{du} + \left[\frac{1}{2}\frac{d\log H^2}{du} + 3\right]\zeta + \left(\frac{\tilde{m}_a}{H}\right)^2 \sin\theta = 0.$$

$$\frac{d\theta}{du} - \zeta = 0.$$

Observe that, by definition,  $\zeta=d\theta/du$ . The initial conditions are  $\zeta(0)=0$  and  $\theta(0)=\theta_{\rm ini}$ .

 $<sup>^{6}</sup>$  Only the ratios  $a/a_{\mathrm{ini}}$  appear in the calculations. So, the choice does not matter as long as it is consistent.

## MiMes: When do you start and stop integrating?

The choice of a good starting point is important, as need to start at a temperature where  $\zeta=0$  is a good approximation. So you can start at some  $T_{\rm ini}$  with a given ratio  $3H(T_{\rm ini})/\tilde{m}_a(T_{\rm ini})\gg 1$ . This needs to be chosen carefully, as low values of  $3H(T_{\rm ini})/\tilde{m}_a(T_{\rm ini})$  result in inaccurate result, while high values may result in a slow calculation. *Advice*: Use various values of that ratio, and find where the relic abundance becomes  $T_{\rm ini}$ -independent.

The stopping condition is more difficult. You should stop at some point where the axion/ALP evolves "adiabatically". Find a quantity that becomes constant as the system relaxes, and use this to determine when adiabaticity is reached. Once  $\theta$  starts to evolve adiabatically, the amplitude of its oscillation is known at later times!

## MiMes: Some (classical) physics

If a system exhibits closed orbits, the quantity

$$J \equiv C \oint p \ d\theta \ ,$$

is the adiabatic invariant. In this case, it becomes

$$J = a^3 \tilde{m}_a \theta_{\text{peak}}^2 f(\theta_{\text{peak}})$$
,

with

$$f(\theta_{\rm peak}) = \frac{2\sqrt{2}}{\pi\theta_{\rm peak}^2} \int_{-\theta_{\rm peak}}^{\theta_{\rm peak}} d\theta \sqrt{\cos\theta - \cos\theta_{\rm peak}} ,$$

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**Important**:  $\theta_{\mathrm{peak}}$  is the peak of the oscillation. So, J can be used to determine how  $\theta_{\mathrm{peak}}$  changes with time. By definition, at  $\theta=\theta_{\mathrm{peak}}$ ,  $p\sim\dot{\theta}=0$ . This means that we can find  $\rho_{a,0}$  on the peak of today's  $\theta$ , as

$$\rho_{a,0} = \gamma^{-1} \frac{s_0}{s_a} m_a \tilde{m}_{a,*} \frac{1}{2} f_a^2 \theta_{\text{peak},*}^2 f(\theta_{\text{peak},*}),$$

where  $T_*$  the temperature at which adiabaticity was reached, and  $\gamma$  the entropy injection between  $T_*$  and today (i.e.  $s_0 = \gamma a_*^3 s_*$ ).

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## How to get MiMeS

There are several ways you can get a stable version of MiMes:

- git clone -b stable https://github.com/dkaramit/MiMeS.git.
  This is the preferred way, as it is guaranteed to be the latest stable version.
- Go to mimes.hepforge.org/downloads, and download it.
- Go to github.com/dkaramit/MiMeS/releases, and download a released version.

You can get the most up-to-date code — not always the most stable one — including the latest version of NaBBODES and SimpleSplines, by running

```
git clone https://github.com/dkaramit/MiMeS.git cd MiMeS git submodule init git submodule update --remote
```

## Configure (and make)

There is no need to install anything if you are going to use MiMeS in a C++ program. The only thing you *must* do is run

bash configure.sh

Alter that, you can include the header file MiMeS/MiMeS.hpp, and you are good to go.

However, you can also run

- make lib, in order to produce the (shared) libraries. This is needed in order to run the python interface.
- make examples, in order to compile the examples in MiMeS/UserSpace/Cpp.
- make exec, in order to produce some test executables (in MiMeS/exec). You just need to run then in order to see if you get any segfaults.

### Classes

There are three classes useful to the user. 7

 mimes::Cosmo<LD>: interpolation of relativistic degrees of freedom of the plasma. By default it uses the EOS2020 8 data. The user can choose another file easily.

<sup>&</sup>lt;sup>7</sup> There are various arguments that need to be passed to the constructors, and the are all listed and explained in the Appendix of the documentation.

<sup>&</sup>lt;sup>8</sup> K. Saikawa and S. Shirai, JCAP **08** (2020), 011 [arXiv:2005.03544 [hep-ph]].

<sup>&</sup>lt;sup>9</sup> S. Borsanyi, Z. Fodor, J. Guenther, K. H. Kampert, S. D. Katz, T. Kawanai, T. G. Kovacs, S. W. Mages, A. Pasztor and F. Pittler, *et al.* Nature **539** (2016) no.7627, 69-71 [arXiv:1606.07494 [hep-lat]].

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- mimes::Axion<LD, Solver, Method>: This is responsible for actually solving the EOM.

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MiMeS is designed to make as few assumption as possible. However, it still assumes that:

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- $\bullet$   $H/\tilde{m}_a$  increases monotonically with the temperature.
- $\zeta(0) = 0$ . This will be changed in the future.
- The energy density of the axion/ALP is always subdominant.

- $\mathbf{0} H/\tilde{m}_a$  increases monotonically with the temperature.
- The energy density of the axion/ALP is always subdominant.
- Only the EOM determines the energy density (no annihilations, no strings, etc.).

- The mass of the axion/ALP. A data file or an actual function.
- ② Data file with  $\log a/a_i$  ( $a_i$  is some arbitrary value; MiMeS rescales it appropriately), T, and  $\log H$  of the underlying cosmology.

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- **3** Value for  $3H/\tilde{m}_a \gg 1$ , which defines the point where integration begins.
- Relative difference of J between a given number of peaks at which we consider adiabaticity to have been reached.
- Other input, related to the algorithm, that might confuse you; *e.g.* temperature at which integration stops no *matter what*!

#### Template arguments

You need to choose what numeric type to use. This is done by the template argument LD which should be double (fast) or long double (accurate).  $^{10}$ 

You also need to tell MiMeS which integration strategy to use. This is done by choosing template arguments:

- Solver can be set to 1 for Rosenbrock (semi-implicit Runge-Kutta). The Method argument in this case can be:
  - RODASPR2<LD> (4th order).
  - ROS34PW2<LD> (3rd order).
  - ROS3W<LD> (2rd order, very bad).
- Solver can be set to 2 for explicit RK. The Method argument can be:
  - DormandPrinceRK45<LD> (7th order)
  - CashKarpRK45<LD> (5th order, very bad).
  - RK45<LD> (5th order, very bad).

<sup>10</sup> You could choose float, but we live in 2021.

### MiMeS from python

In order to call the python interface of MiMeS, we need to first call make lib in the root directory of MiMeS.

Before that, we can take some time to decide what the template arguments and compilation options should be. In the file MiMes/Definitions.mk, you can change the variables:

- LONG=long will compile the library with long double numeric types. LONG= will compile the library with double numeric types.
- SOLVER and METHOD, as in the template arguments.

Also, in the same file, you can change compilation options:

- Compiler:
  - CC=g++ in order to use the GNU C++ compiler.
  - CC=clang -lstdc++ in order to use the clang C++ compiler.
- Optimization level:
  - OPT=00: No optimization.
  - O=O1, O2, or O3 (marginally faster, default): some optimization occurs (read the compiler documentation for more information).
  - OPT=Ofast: full optimization (fast, but dangerous).

#### Examples

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  - MiMes: Notation MiMes: When do you start and stop integrating?
  - MiMes: Some (classical) physics

  - - How to get MiMeS Configure (and make)
    - Classes

    - Assumptions
    - What MiMeS expects from you
  - Template arguments MiMeS from python
  - Examples
    - C++
    - python
    - - Summing up
      - Future

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#### Define everything and solve in just a few lines of code!

```
// use chi PATH to interpolate the axion mass.
mimes::AxionMass<long double> axionMass(chi PATH.0.mimes::Cosmo<long double>::mP):
/*set \tilde{m}_a^2 for T \geq T_{\text{max}}*/
long double TMax=axionMass.getTMax(), chiMax=axionMass.getChiMax();
axionMass.set ma2 MAX(
    [&chiMax,&TMax](long double T, long double fa){
        return chiMax/fa/fa*std::pow(T/TMax.-8.16):
):
/*set \tilde{m}_a^2 for T \le T_{\min}*/ long double TMin=axionMass.getTMin(), chiMin=axionMass.getChiMin();
axionMass.set ma2 MIN(
    [&chiMin,&TMin](long double T, long double fa){
        return chiMin/fa/fa:
):
/*this path contains the cosmology*/
std::string inputFile = std::string(rootDir)+
    std::string("/UserSpace/InputExamples/MatterInput.dat"):
/*declare an instance of Axion*/
mimes::Axion<long double, 1, RODASPR2<long_double> > ax(
    0.1, 1e16, 500, 1e-4, 1e3, 15, 1e-2, inputFile, &axionMass,
    1e-1, 1e-8, 1e-1, 1e-11, 1e-11, 0.9, 1.2, 0.8, int(1e7)
/*solve the EQM!*/
ax.solveAxion();
```

## python

#### Outlook

- Why particle dark matter The dark matter particle
- The axion (like) particle
- - The axion EOM
  - How hard can it be?
  - Initial conditions
  - (Bad) Analytical approximations
  - Need for accuracy, speed, and automation
- - MiMes: General
  - MiMes: Under the hood
  - MiMes: Notation
  - MiMes: When do you start and stop integrating?
  - MiMes: Some (classical) physics

  - How to get MiMeS

    - Configure (and make)
    - Classes
    - Assumptions
    - What MiMeS expects from you
  - Template arguments
  - MiMeS from python
- - C++
  - python
- Outlook
  - Summing up
  - Future

## Summing up

Commend from D: blah blah, what you did

#### **Future**

Commend from D: Freeze-out/in, kinematic misalignment mechanism, automatically compare with searches.

# Thank you!

Language	files	blank	comment	code
C/C++ Header C++ Python	35 20 22	677 483 483	444 198 367	1595 1106 1000
SUM:	77	1643	1009	3701

# Backup

(equations, derivations, tables)







