Investigating Properties of Coherent States in Simple Harmonic Oscillator

In [1]: import vpython as vp

import numpy as np

import matplotlib.pyplot as pl

```
from vpython import rate
        vp.canvas()
In [2]: def SetArrowFromCN( cn, a):
            SetArrowWithCN takes a complex number on and an arrow object a .
            It sets the \, y \, and \, z \, components of the arrow s \, axis to the real
            and imaginary parts of the given complex number.
            Just like Computing Project 1, except y and z for real/imag.
            a.axis.y = cn.real
            a.axis.z = cn.imag
            a.axis.x = 0.0
In [5]: vp.canvas()
        gd = vp.graph(xtitle="t", ytitle="Expectation", width=400, height=300)
        gr = vp.gcurve(color = vp.color.black)
        NA=80
                                       # how many arrows?
        a=15.0
                                       # range of x is -a/2 to a/2 in units
                                       # of $\sqrt{\hbar/m\omega}$
        x = np.linspace(-a/2, a/2, NA) # NA locations from -a/2 to a/2
        NHs = 20
        hs = np.zeros((NHs,NA), np.double) # the hermite polynomials, an NHs x NA array
        coefs = np.zeros(NHs,np.double) # the coherent state coefficients, an NHs x 1 array
        psis = np.zeros((NHs,NA), np.double) # the stationary states, an NHs x NA array
        alpha = np.sqrt(10)
                                         \# < n > = 10.0
        hs[0] = 0*x + 1.0 # zeroth Hermite Polynomial, H0
        hs[1] = 2*x
                                      # first Hermite Polynomial, H1
        # Compute the first NHs Hermite Polynomials,
        # use recurrence relation to get the rest of the Hn(x)
        # (see e.g., http://en.wikipedia.org/wiki/Hermite_polynomials#Recursion_relation)
        for n in range(1,NHs-1):
           hs[n+1]=2*x*hs[n] - 2*n*hs[n-1]
        \# Use the coherent state coefficient relation to get the c[n]s.
        # avoid overflow by computing them in a loop. Don't worry about
        \# the overall factor of c[0] since we'll renormalize our discrete
        # psi array later anyway.
        coefs[0]=1.0 * np.exp(-alpha**2)
        for i in range(1,NHs):
           coefs[i]=coefs[i-1]*alpha/np.sqrt(i)
        #
        # Get the stationary states using the hs array and compute the
        # normalization factor in a loop to avoid overflow
        normFactor = 1.0/np.pi**0.25
        psis[0] = np.exp(-x**2/2)
        for i in range(1,NHs):
            normFactor = normFactor/np.sqrt(2.0*i)
            psis[i] = normFactor*hs[i]* np.exp(-x**2/2)
        # Now do the sum to compute the initial wavefunction
        psi = np.zeros(len(x),complex)
        for m in range(NHs):
           psi += coefs[m]*psis[m]
        # Normalize!
```

Coherent.png

alist = []

t = 0.0

rate(100)

psi=psi/np.sqrt((abs(psi)**2).sum())

SetArrowFromCN(3*psi[i],alist[i])

and computes the values for <x> and graphs them.

of 3 so they look nice.

vp.scene.autoscale = False

for i in range(NA):

dt = 2*np.pi/1000.0

build the arrows. Scale them on the screen by a factor

alist.append(vp.arrow(pos=vp.vec(x[i],0,0), color= vp.color.red))

After this.. create the "time loop" that animates the wavefunction

Expectation.png

```
In [6]: def timeloop(t):
            psi = np.zeros(len(x),complex)
            for m in range(NHs):
                psi += coefs[m] * psis[m]*np.exp(-1j*(0.5 + m)*t)
            psi = psi/np.sqrt((abs(psi)**2).sum())
            for i in range(NA):
                SetArrowFromCN(3*psi[i],alist[i])
            expect = (x*(abs(psi)**2)).sum()
            gr.plot(pos=(t, expect))
        vp.scene.autoscale = False
        t = 0.0
        dt = 2*np.pi/1000.0
        while t < 4*np.pi:</pre>
            rate(100)
            timeloop(t)
            t+=dt
```

Questions

1.

$$\psi_{lpha} = c_0 \sum_{n=0}^{\infty} rac{lpha}{\sqrt{n!}} \psi_n \ 1 = c_0^2 \sum_{n=0}^{\infty} rac{lpha^{2n}}{n!} \ 1 = c_0^2 e^{lpha^2} \ e^{-lpha^2} = c_0^2 \ c_0 = e^{rac{-lpha^2}{2}}$$

2.

$$egin{align} \langle n
angle &= e^{rac{-lpha^2}{2} 2} \sum_{n=0}^\infty rac{lpha^{2n}}{n!} n \ \langle n
angle &= e^{-lpha^2} |lpha^2| \sum_{n=1}^\infty rac{lpha^{2n-2}}{(n-1)!} rac{n}{n} \ \langle n
angle &= e^{-lpha^2} |lpha^2| \sum_{n=1}^\infty rac{lpha^{2n-2}}{(n-1)!} \ \langle n
angle &= e^{-lpha^2} |lpha^2| \sum_{n=0}^\infty rac{lpha^{2n}}{n!} \ \langle n
angle &= e^{-lpha^2} |lpha^2| e^{lpha^2} \ \langle n
angle &= |lpha^2| \end{aligned}$$

In [7]:

3.

$$E_n = \hbar\omega(n+rac{1}{2}) \ \langle E
angle = \sum_{n=1}^{\infty} P_n E_n \ \hbar\omega\langle n
angle + rac{\hbar\omega}{2} \ rac{kx^2}{2} = \hbar\omegalpha^2 \ x = \sqrt{rac{2\hbar\omega}{m}}lpha \ x = \sqrt{rac{2\hbar\omega}{m}}lpha \ x = \sqrt{rac{2\hbar\omega}{m}}lpha$$

```
alpha = np.linspace(1, 50, 100)
        x = np.sqrt((2*hbar)/m) * alpha
        pl.plot(alpha, x)
Out[7]: [<matplotlib.lines.Line2D at 0x1fe55c789d0>]
```

```
0.7
```

hbar = 1.0545718e-34m = 9.10938356e-31

```
0.6
0.5
0.4
0.2
0.1
0.0
                                          30
                 10
```