# Main notebook

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2017-11-25, Last update 2017-11-25

#### Drude model $\sigma$ and $\epsilon$ 1

The assumptions of the Drude model are simple: we have interaction-free electrons that occasionally undergo some scattering process during a time dt with probability  $\frac{dt}{\tau}$ , where  $\tau$  is some phenomenological parameter. This scattering will randomise electron momentum.

Our ultimate goal will be to find the conductivity  $\sigma$  and the dielectric constant  $\epsilon$  in the Drude model, with Drude relaxation time  $\tau$ , electron density n and electron mass m. We'll find

$$\sigma_{\rm DC} = \frac{ne^2\tau}{m}$$

$$\sigma_{\rm AC} = \frac{ne^2\tau}{m} \frac{1}{1 - i\omega\tau}$$
For SI and Gaussian (2)

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$$\epsilon_r = 1 + i \frac{4\pi\sigma}{\omega}$$
 Gaussian (3a)

$$\epsilon_r = 1 + i \frac{\widetilde{\sigma}}{\omega \epsilon_0}$$
 SI (3b)

Our dielectric constant can be rewritten to plug in for  $\sigma$ , giving us

$$\epsilon = 1 + i \frac{4\pi\sigma}{\omega} \tag{4}$$

$$= 1 + i \frac{4\pi\sigma_{\rm DC}}{\omega} \frac{1}{1 - i\omega\tau}$$

$$= 1 + i \frac{4\pi\sigma_{\rm DC}}{\omega} \frac{1}{1 - i\omega\tau} \frac{1 + i\omega\tau}{1 + i\omega\tau}$$

$$(5)$$

$$=1+i\frac{4\pi\sigma_{\rm DC}}{\omega}\frac{1}{1-i\omega\tau}\frac{1+i\omega\tau}{1+i\omega\tau}\tag{6}$$

$$=1+i\frac{4\pi\sigma_{\rm DC}}{\omega}\frac{1+i\omega\tau}{1+\omega^2\tau^2}\tag{7}$$

$$= \left(1 - \frac{4\pi\sigma_{\rm DC}\omega\tau}{\omega\left(1 + \omega^2\tau^2\right)}\right) + i\left(\frac{4\pi\sigma_{\rm DC}}{\omega\left(1 + \omega^2\tau^2\right)}\right) \tag{8}$$

$$= \left(1 - \frac{4\pi\sigma_{\rm DC}\tau}{1 + \omega^2\tau^2}\right) + i\left(\frac{4\pi\sigma_{\rm DC}}{\omega\left(1 + \omega^2\tau^2\right)}\right) \tag{9}$$

This lets us write down the explicit real and imaginary of the Drude dielectric function.

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#### 1.1 Limiting forms of the Drude model

We can look at the large and small  $\omega$  limits for the Drude dielectric function.

## 1.2 Derivations for Drude model

#### 1.2.1 Conductivity

We can start unit-system independently, with the expression  $\langle \mathbf{j} \rangle = \sigma \mathbf{E}$ . We can also relate our current to our average electron velocity:  $\mathbf{j} = ne\mathbf{v}$ . Imagine at time t = 0 our electron undergoes a Drude collision, and emerges with  $\mathbf{v}_{t=0} = \mathbf{v_0}$ . After a time t, the electron will accelerate with acceleration  $-\frac{e\mathbf{E}}{m}$  (which fortunately remains unit independent). Because it will only accelerate for a time  $\tau$  on average before a collision, it will end up with velocity  $\mathbf{v} = -\frac{e\mathbf{E}}{m}\tau + \mathbf{v_0}$ . The average velocity, and current, will be

$$\langle \mathbf{v} \rangle = -\frac{e\mathbf{E}}{m}\tau + \langle \mathbf{v_0} \rangle \tag{10}$$

$$= -\frac{e\mathbf{E}}{m}\tau\tag{11}$$

$$= -\frac{e\mathbf{E}}{m}\tau\tag{12}$$

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