

## Problem 2: Critical Field of a Thin Superconducting Film

a) Taking result:  $\mu_0 M(x) \approx -B_a \left( \frac{1}{8\lambda_c^2} \right) (\delta^2 - 4x^2)$  ;  $\delta \ll \lambda_c$   
 ;  $-\delta/2 < x < \delta/2$

Work density to bring a superconductor into  $B$  field

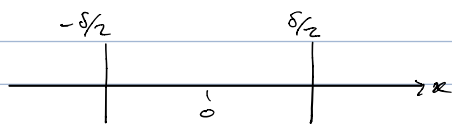
$$W = - \int_0^{B_a} \underline{M} \cdot d\underline{B}$$

$$= - \int_0^{B_a} \frac{-B_a}{\mu_0} \left( \frac{1}{8\lambda_c^2} \right) (\delta^2 - 4x^2) \cdot dB \quad \text{both } \underline{x}$$

$$= \frac{(\delta^2 - 4x^2)}{8\mu_0\lambda_c^2} \int_0^{B_a} B \, dB = \frac{B_a^2 (\delta^2 - 4x^2)}{16\mu_0\lambda_c^2}$$

The free energy density inside the superconductor at field strength  $B_a$  is the initial free energy density at 0 field strength;  $F_S(0)$  plus the additional amount gained from the work density to bring the superconductor into the field of strength  $B_a$

$$\therefore F_S(x, B_a) \approx F_S(B=0) + \frac{B_a^2 (\delta^2 - 4x^2)}{16\mu_0\lambda_c^2} \quad \text{when } \delta \ll \lambda_c \text{ from } M(x) \text{ sol.}$$



where  $x$  is the distance from the centre of the SC along  $\underline{x}$ .

b)  $F_{SM}(x, B_a) = \frac{(\delta^2 - 4x^2) B_a^2}{16\mu_0\lambda_c^2}$

$\frac{\int_{-\delta/2}^{\delta/2} F_{SM}(x, B_a) dx}{\delta}$  is the average <sup>(over  $x$ )</sup> magnetic contribution to the free energy density in the SC

$$= \frac{1}{\delta} \int_{-\delta/2}^{\delta/2} \frac{(\delta^2 - 4x^2) B_a^2}{16\mu_0\lambda_c^2} dx$$

$$= \frac{B_a^2}{16\mu_0\lambda_c^2} \left( \underbrace{\left[ \delta x \right]_{-\delta/2}^{\delta/2}}_{\frac{\delta^2}{2} + \frac{\delta^2}{2} = \delta^2} - 4 \underbrace{\left[ \frac{x^3}{3} \right]_{-\delta/2}^{\delta/2}}_{\frac{\delta^2}{24} + \frac{\delta^2}{24} = \frac{\delta^2}{12}} \right)$$

$$= \frac{B_a^2}{16\mu_0\lambda_c^2} \left( \delta^2 - \frac{\delta^2}{3} \right) = \frac{B_a^2}{24\mu_0} \left( \frac{\delta}{\lambda_c} \right)^2$$

c) The normal state has  $F_N(\mu, B_{ac}) \approx \text{const.}$

At the critical field  $B_{ac}$ ;  $F_S(B_{ac}) = F_N(B_{ac})$

We can write  $F_S$  using the average magnetic contribution, instead of the  $x$  dependent contribution, as:

$$F_S(B_{ac}) = F_S(0) + \frac{B_{ac}^2}{24\mu_0} \left( \frac{\delta}{\lambda_c} \right)^2$$

We can write  $F_N(B_{ac}) = F_N(0)$  as it is  $\sim \text{constant}$

And so the condensation energy is:

$$\begin{aligned} F_c &= F_S(B_{ac}) - F_S(0) = F_N(0) - F_S(0) \\ &= \frac{B_{ac}^2}{24\mu_0} \left( \frac{\delta}{\lambda_c} \right)^2 \end{aligned}$$

$$F_c = \frac{1}{2\mu_0} H_c^2$$

$$\therefore \frac{1}{2\mu_0} H_c^2 = \frac{B_{ac}^2}{24\mu_0} \left( \frac{\delta}{\lambda_c} \right)^2$$

$$B_{ac} = 2\sqrt{3} \left( \frac{\lambda_c}{\delta} \right) H_c$$

$F_c = \frac{1}{2\mu_0} H_c^2$  comes from considering a bulk type-I SC

$\hookrightarrow M \propto H$

$$\hookrightarrow W = - \int_0^{B_c} \underline{M} \cdot d\underline{B} \propto H^2$$

and then considering only the magnetic component of  $F \propto W \propto H^2$

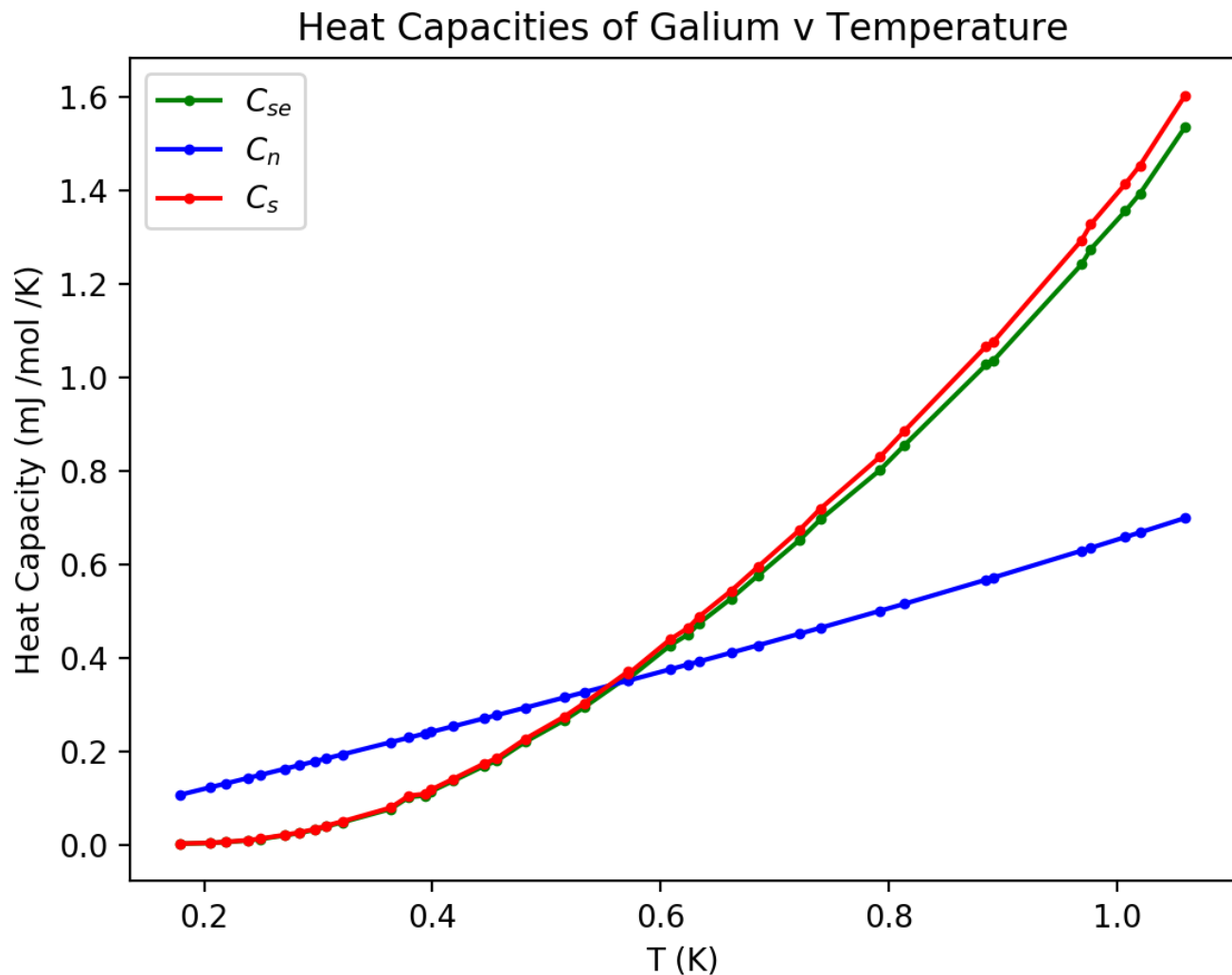
$\rightarrow$  the  $\frac{1}{2\mu_0}$  are constants of proportionality for a bulk type I SC

$$d) \delta = 10 \text{ nm} \quad \lambda_c = 37 \text{ nm} ; \text{Pb.} \Rightarrow B_{ac} = 2\sqrt{3} \left( \frac{37}{10} \right) H_c$$

$\therefore$  increase from bulk  $\rightarrow \frac{B_{ac}}{H_c} = 2\sqrt{3} \left( \frac{37}{10} \right) = 12.82$  factor increase.

## Problem 1: Excitation Gap in Superconducting Gallium

d) The calculated heat capacities from (a), (b) and (c) are plotted versus temperature below.



the data for these plots is included in the table titled 'Solid state Problem 1d Data'.

- $C_s$  was found by multiplying the given  $C$  data by  $T$
- $C_{se}$  was found by getting the phonon contribution,  $C_{ph}$ , by taking  $\beta = 0.0868$  (from figure 1) and multiplying by  $T^3$ , then subtracting  $C_{ph}$  from  $C_s$ .
- $C_n$  was found by  $\gamma T + C_{ph}$  where  $\gamma = 0.596$  from figure 1.

e)  $C_n(T_c) = \gamma T_c$  where  $\gamma = 0.596 \text{ mJ mol}^{-1} \text{ K}^{-2}$   
from the graph in Figure 1

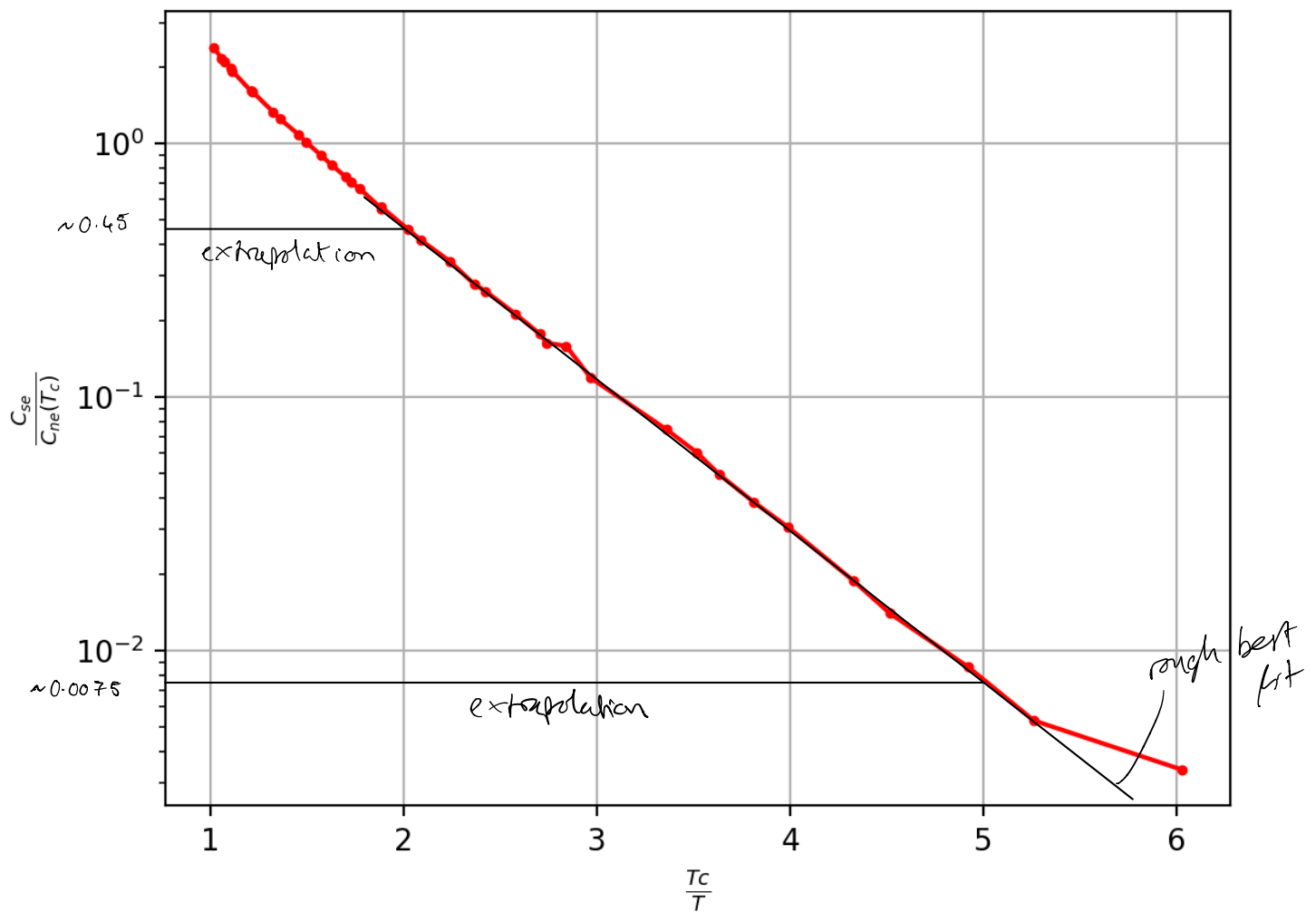
$$= 0.596 \cdot 1.078$$

$$C_n(T_c) = 0.642 \text{ mJ mol}^{-1} \text{ K}^{-1}$$

# Solid State Problem 1d Data

T (K)	C <sub>n</sub> (mJ /mol /K)	C <sub>s</sub> (mJ /mol /K)	C <sub>se</sub> (mJ /mol /K)
2.05E-01	1.23E-01	3.89E-03	3.40E-03
2.19E-01	1.31E-01	6.13E-03	5.54E-03
2.39E-01	1.43E-01	9.79E-03	9.02E-03
2.49E-01	1.49E-01	1.29E-02	1.21E-02
2.70E-01	1.62E-01	2.08E-02	1.97E-02
2.83E-01	1.70E-01	2.60E-02	2.47E-02
2.97E-01	1.78E-01	3.32E-02	3.17E-02
3.07E-01	1.84E-01	4.05E-02	3.88E-02
3.21E-01	1.93E-01	4.97E-02	4.79E-02
3.63E-01	2.19E-01	7.92E-02	7.65E-02
3.79E-01	2.29E-01	1.05E-01	1.02E-01
3.94E-01	2.38E-01	1.09E-01	1.05E-01
3.99E-01	2.41E-01	1.18E-01	1.14E-01
4.18E-01	2.53E-01	1.41E-01	1.37E-01
4.45E-01	2.70E-01	1.73E-01	1.68E-01
4.56E-01	2.77E-01	1.85E-01	1.79E-01
4.82E-01	2.93E-01	2.27E-01	2.21E-01
5.16E-01	3.15E-01	2.74E-01	2.67E-01
5.33E-01	3.26E-01	3.03E-01	2.94E-01
5.72E-01	3.51E-01	3.71E-01	3.61E-01
5.72E-01	3.51E-01	3.66E-01	3.55E-01
6.08E-01	3.75E-01	4.39E-01	4.26E-01
6.24E-01	3.86E-01	4.65E-01	4.51E-01
6.34E-01	3.92E-01	4.89E-01	4.74E-01
6.62E-01	4.11E-01	5.43E-01	5.27E-01
6.86E-01	4.27E-01	5.95E-01	5.77E-01
7.22E-01	4.52E-01	6.73E-01	6.52E-01
7.40E-01	4.64E-01	7.20E-01	6.97E-01
7.92E-01	5.01E-01	8.31E-01	8.02E-01
8.14E-01	5.16E-01	8.85E-01	8.55E-01
8.85E-01	5.67E-01	1.07E+00	1.03E+00
8.92E-01	5.72E-01	1.08E+00	1.04E+00
9.69E-01	6.29E-01	1.29E+00	1.24E+00
9.77E-01	6.35E-01	1.33E+00	1.28E+00
1.01E+00	6.59E-01	1.41E+00	1.36E+00
1.02E+00	6.68E-01	1.45E+00	1.39E+00
1.06E+00	6.99E-01	1.60E+00	1.54E+00

f)  $\frac{C_{se}(T)}{C_{se}(T_c)}$  versus  $\frac{T_c}{T}$  with a log scaled y axis.



g)  $C_{se}(T_c)$  and  $T_c$  provide a scaling factor. Looking at the proportionality:

$C_{se} \propto e^{-\frac{E_g}{k_B T}}$  against  $\frac{1}{T}$   $\therefore$  on a log graph, this will produce a straight line at gradient  $-\frac{E_g}{k_B}$

This means there is an energy difference between different  $T$  values which corresponds to an energy gap. This is due to the known quantisation of energy and so this is the gap between the states.

h)  $T < \frac{T_c}{2} \rightarrow \frac{T_c}{T} > 2$

From g) we see that the gradient between two points,  $m = -\frac{E_g}{k_B}$  where  $E_g$  is the energy gap.

From the plot in (f);  $E_g = -k_B m = 0.118 \text{ meV}$  using marked extrapolations

The energy gap determined from (f) is the energy to excite one electron from a Cooper pair;  $\Delta(0) = E_g$

BCS theory relates this.  $2\Delta(0) \approx 3.5 k_B T_c$ .

From BCS theory;  $\Delta(0) \approx \frac{7}{2} k_B T_c = 0.163 \text{ meV}$

therefore the result from BCS theory is roughly the same as derived here.

Note;  $E_g$  could be found using  $m = \frac{-E_g}{2k_B}$ , which then defines  $E_g$  as the energy to break and excite a Cooper pair, where  $E_g = 2\Delta(0)$ .  
→ I have used  $m = \frac{-E_g}{k_B T}$  as this was given on the lecture slides.

This method works when ignoring the temperature dependence as it introduces a very small difference when  $T < T_c/2$ . The gap also decreases as  $T$  increases and so this gives an estimate at the max gap.

(at  $T=0\text{K}$ )  
 $\Delta(0)$  is the energy per electron to break a Cooper pair and excite a single electron to a higher energy state