

# Transverse Momentum in Nucleons, From Raw Data to TMD Extraction

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# Chapter 1

## Introduction

The majority of this thesis explains how we performed our experimental measurement. The role of this introductory chapter is to provide an explanation of why we performed this measurement, and what exactly it is that was measured. First, a brief statement of the measurement is given. Then, in order to understand why we perform the current measurement, a historical look back at the parallel development of nucleon structure in theoretical and experimental aspects is presented. After obtaining a historical perspective, it is evident why the current study is needed, and the scope of the present measurement is re-stated. Finally, a high-level discussion of how this measurement was performed is presented, leaving the details to the remainder of the document.

### 1.1 Statement of Purpose

This work aims to contribute to the understanding of nucleon structure in the framework of transverse momentum dependent parton distribution functions (TMDs) by measuring structure functions in semi-inclusive deeply inelastic scattering (SIDIS). By measuring the cross section for charged pi-mesons  $\pi^\pm$ , we measure ( $F_{UU,T} + \epsilon F_{UU,L}$ ,  $F_{UU}^{\cos\phi}$ , and  $F_{UU}^{\cos(2\phi)}$ ). By analyzing the beam spin asymmetry (BSA) we measure the ratio

$$A_{LU}^{\sin\phi} = \sqrt{2\epsilon(1-\epsilon)} \frac{F_{LU}^{\sin\phi}}{F_{UU,T} + \epsilon F_{UU,L}} \quad (1.1)$$

for positively charged k-mesons. Finally, we are able to use these structure functions to estimate the model parameters in TMD models.

## 1.2 Nucleon Structure from Experiment and Theory

Nucleon structure is currently a very active field of research which aims to explain the structure of protons and neutrons in terms of the fundamental particles known as quarks. This sub-field was born in 1958 following Hofstadters demonstration that the electron-proton scattering cross sections are not consistent with the predictions based on a pointlike proton theory [1]. At the same time, an ever growing number of particles discovered by bubble and spark chamber experiments confused physicists.

Some theoretical progress was made during the 1960s as Murray Gell-Mann proposed the construction of this large number of observed particles from pairs or triplets of fundamental particles called quarks. The proposed quark came in 3 flavors called up, down, and strange. However, in order to correctly predict the charge of the observed particles the quarks had to be fractionally charged (up 2/3, down -1/3, and strange 1/3). Accordingly, the quarks had to be permanently confined somehow inside of bound states. This idea was troubling to many physicists, and support for the quark model only really established itself after the  $J/\psi$  ( $c\bar{c}$ ) particle was observed in 1974 and the inclusion of the charm quark into the existing set of 3 quark flavors accommodated its existence. This quark model picture enjoyed some success predicting the masses of hadrons, but gives no information about the dynamics of quarks in hadrons.

Almost all modern measurements that relate to quark dynamics are done by scattering electrons on protons or proton/neutron pairs in deuterium. Since this process proceeds via virtual photon exchange between the charged electron and the charged quarks, it is easier to understand using the tools of QED. If the virtual photon is of sufficiently high energy ( $\approx 2\text{GeV}$ ) its wavelength is an order of magnitude less than the diameter of the proton/neutron and likely interacts with just one part (quark or gluon) of the nucleon. This idea is known as the parton model, and was proposed by Feynman in 1969 [2].

For the 30 years following this development, the measurement of parton distribution functions (PDFs) was performed at experiments around the world (HERMES, COMPASS, and Jefferson Lab). The PDFs (at leading order) describe the probability to observe a quark with a given momentum fraction  $x$  in a hadron. There is one such function for each quark flavor in each hadron. The unpolarized PDF is now known quite well, and the polarized PDFs have also been measured. During the measurement of polarized PDFs in 1989 the European Muon Collaboration (EMC) observed that only 30% of the total spin of the proton appeared to be due to the spin of the quarks. This result came to be known as the proton spin crisis.

One possible resolution to this problem is that the quarks carry orbital angular momentum inside of the proton and this contributes to the total observed spin. In this case, measurements of the three-dimensional momentum structure of the quarks inside of hadrons is expected to be very useful. The transverse momentum dependent parton distribution functions (TMD PDFs) describe the quark mo-

menta in both the longitudinal direction  $x$  (defined by the hard momentum transfer direction) and the momentum in the plane transverse to that as well  $p_T$ .

### **1.3 Measurement of Semi-Inclusive Deeply Inelastic Scattering with CLAS**

Our purpose is re-stated clearly.

### **1.4 Overview of our Measurement**

Let's talk about detectors, accelerators, and software written in c++.



# Chapter 2

## Experiment

Jefferson National Lab houses the continuous electron beam accelerator facility (CEBAF) which is currently capable of providing 11 GeV electron beams to three experimental end stations, and 12 GeV to a fourth. The data that are analyzed in this study are from the run period E1-F, which occurred before the 12 GeV CEBAF upgrade, the details in this chapter describe the accelerator and CLAS detector at the time of the E1-F run period.

### 2.1 CEBAF

Proposed in 1982 and constructed between 1987-1997 the CEBAF accelerator at Jefferson lab was composed of a pair of linear accelerators and 9 recirculating arcs arranged in a racetrack shape [3, 4]. Originally designed to provide 4 GeV unpolarized electrons to three experimental halls, CEBAF was fitted with twin polarized electron guns, and upgraded to 6 GeV beam energy before the E1-F run period. CEBAF was built to provide an extremely high duty factor and an average beam current of up to  $200 \mu A$ .

#### 2.1.1 Electron Injection & Polarization

CEBAF's injector provided 45 MeV electrons with 70%-80% longitudinal polarization for the main accelerator. In order to provide an apparent continuous stream of events to the detectors, electron bunches were produced at a rate of  $f = 1497 \text{ Mhz}$ . To accommodate different requests for energy and beam current simultaneously, the injector produced three interspersed bunch trains at a frequency of  $499 \text{ Mhz}$ . The output injector energy  $E = 45 \text{ MeV}$  was chosen so that injected electrons were *sufficiently relativistic*. In other words, when bunches of electrons at different energies simultaneously passed through the linear accelerators (LINACs), the relative phase difference (between different energy passes) accumu-

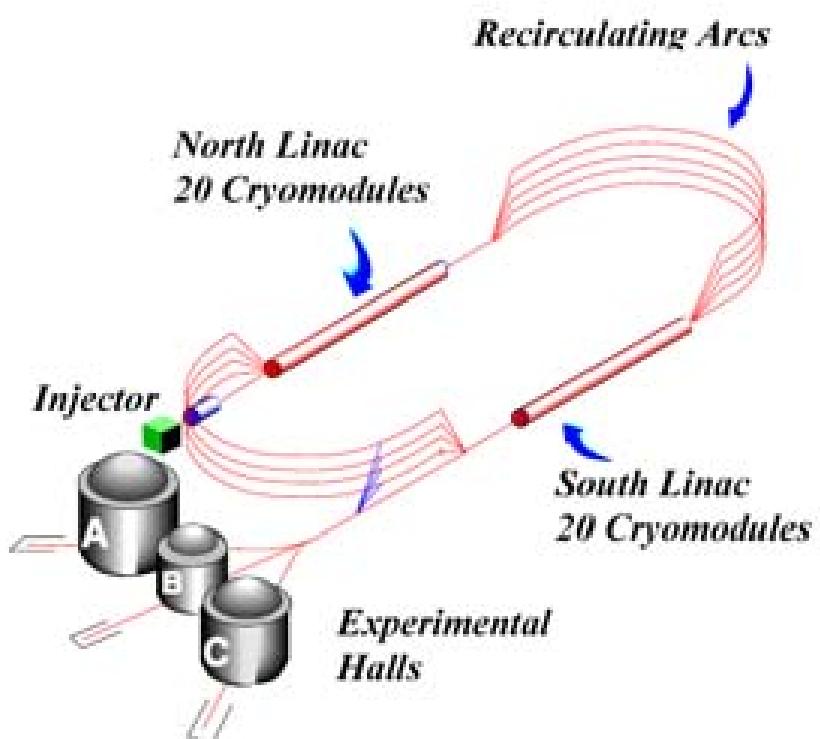


Figure 1: Layout of CEBAF.

Figure 2.1: Diagrammatic representation of CEBAF.

lated over the distance remained less than  $1^\circ$ .

During early accelerator construction, the need for polarized electrons became apparent and the final polarized electron gun was produced in 2000. Production of polarized electrons was achieved by using twin polarized electron guns mounted at  $15^\circ$  with respect to the injector beam axis. Inside of each gun, electrons were liberated from gallium arsenide photo-cathodes by three independent diode lasers operated at a repetition rate of  $499\text{ Mhz}$ . During polarized production, the diodes operated at a central wavelength of  $850\text{ nm}$ . By manipulating the laser polarization using a Pockels cell, the electron beam spin is flipped at a rate of  $60\text{ Hz}$ . Throughout the run period, an overall phase difference could be introduced by rotation of a wave-plate. As will be discussed in some detail later, changes in beam helicity due to wave-plate settings must be removed from the recorded data. After acceleration to  $5\text{ MeV}$  the beam polarization was measured in the injector facility using a Mott polarimeter.

### 2.1.2 Acceleration of Electrons

The north and south LINACs were responsible for increasing the energy of the electrons from  $45\text{ MeV}$  up to an impressive  $\approx 5.7\text{ GeV}$  before delivery to the experimental halls A, B, and C. In order to achieve this each bunch of electrons was accelerated through ten stages, five passes through each LINAC. The strong electric field needed to accelerate electron bunches was confined inside of superconducting 5-cell elliptical cavities. These cavities were machined from niobium, and operated at a temperture of  $2.2^\circ\text{ K}$ . Developed at Cornell University, the cavities were operated at  $1,500\text{ Mhz}$  with a gradient greater than  $5\text{ MV/m}$  and a quality factor  $Q_0 \geq 3 \cdot 10^9$  (the Q factor describes the monochromaticity of the cavity and is defined  $Q_0 = f_0/\Delta f$ ). Each cavity was sealed inside of a cryo-unit, four such units were connected together to form an 8.5 meter cryo-module. Each LINAC was composed of 20 such cryo-modules connected together to increase electron energies by more than  $500\text{ MeV}$  per pass.

The radio frequency (RF) that powers each cavity was sourced by a water-cooled 5 kilowatt Klystron located in groups of eight above each cryo-module. Phase locking of each cavity with a master oscillator ensured that the difference in phase between all cavities was less than one degree. The important super conductivity was maintained by circulation of liquid helium at  $2.2^\circ\text{ K}$ . Production of liquid helium ocurred on-site at the 5 kW helium liquefaction plant.

After re-circulation of fives passes through each LINAC, full energy beams were delivered to the halls. Bunches were separated using an RF separator before entering their respective experimental halls. The beamline leading to the halls was also equipped with the ability to separate bunches before



Figure 2.2: This aerial photograph contains annotations that show the accelerator path and the three experimental halls.

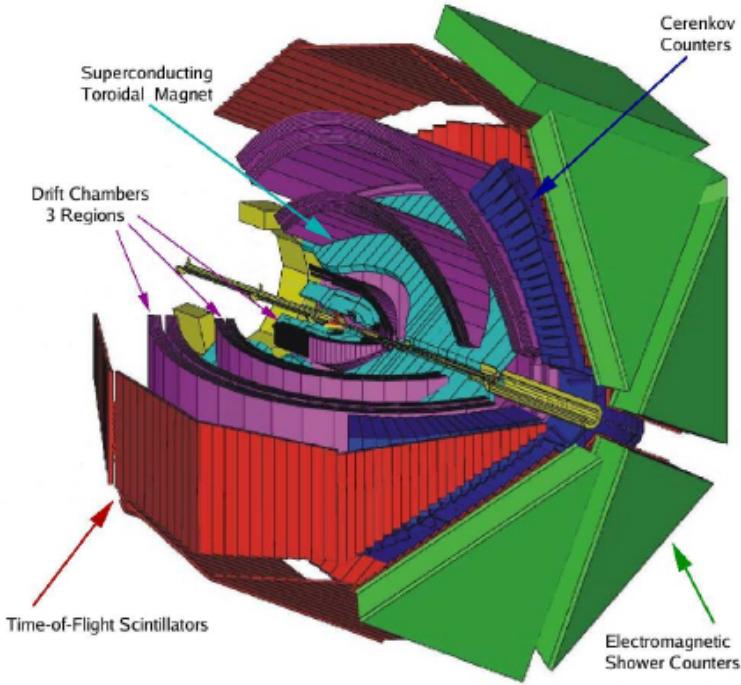


Figure 2.3: In this diagrammatic representation of CLAS, several important detector sub-systems are labeled.

they completed five full passes delivering either three full energy beams, two full energy beams and one lower, or one full energy beam and two of lower energy. This capability, combined with the flexible beam polarization and beam current provided by the injector ensured that each hall could experiment at its desired settings simultaneously.

## 2.2 CLAS in Jefferson Lab Experimental Hall B

The CEBAF large acceptance spectrometer (CLAS), housed in Jefferson lab's Hall B, was used to record the E1-F dataset that is used in this study. Capable of detecting particles over a very large angular range, the CLAS detector covered almost the full  $4\pi$  solid angle around the target region. The detector was also designed to perform efficiently for particles with wide range of momenta between 0.5-6.0 GeV. Overall detector design consisted of a large superconducting magnet that produced a toroidal field (this magnet was referred to as the torus), and six ideally identical *sectors*. Each sector of CLAS contained an identical set of sub-systems. After combining information from all sub-systems and running the reconstruction algorithm, complete events are measured. This capability made CLAS unique in comparison with the arm style spectrometers of halls A and C. The major components of CLAS are listed below.

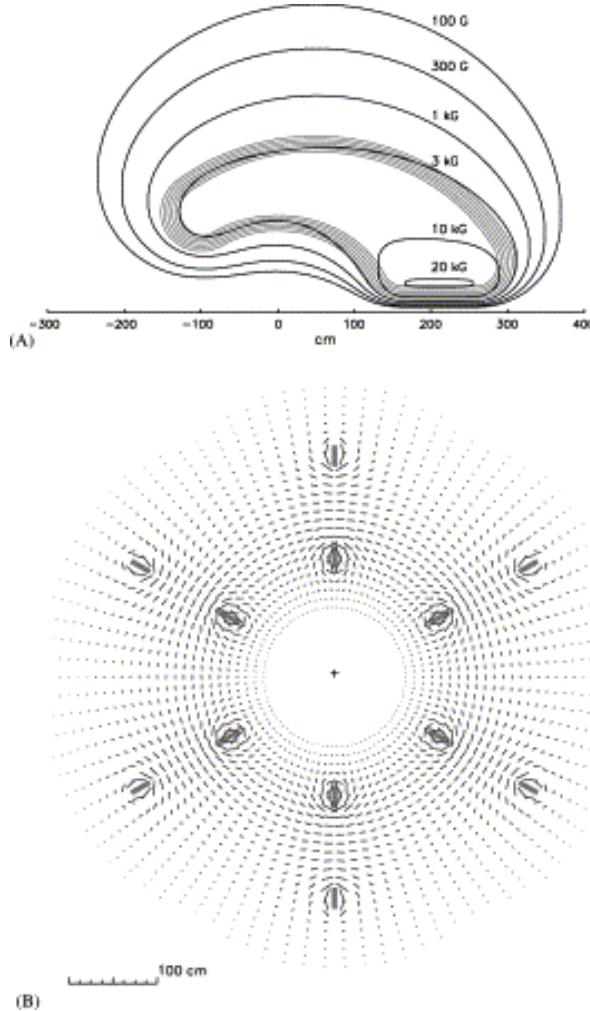


Figure 2.4: Diagrams of the CLAS torus field. This figure reproduced from [5]

### 2.2.1 CLAS Torus and Drift Chambers

Measuring the momentum of charged particles with  $p > 200 \text{ MeV}/c$  was accomplished by measuring the curvature of the particle trajectory as it passed through the CLAS toroidal magnetic field. Six superconducting coils were arranged  $60^\circ$  apart azimuthally around the beamline to create a 2 Tesla. The field produced, which varied from two tesla at lower angles to half a tesla at angles greater than  $90^\circ$ , curved charged particle trajectories in the  $\theta$  direction without altering the azimuthal  $\phi$  direction. The geometry of the torus magnet guided the development of the entire spectrometer.

Spatially, the 18 drift chambers were divided with six in each sector. In order to perform tracking before, inside, and after the torus three radially distinct drift chambers were constructed for each sector (these were called regions 1, 2, and 3). Each drift chamber consisted of 12 superlayers of hexagonal drift chamber cells. Angular measurements in the azimuthal direction were accomplished by offsetting the first six and last six superlayers by  $6^\circ$ . In total 35,148 individual drift chambers cells were used for tracking [6].

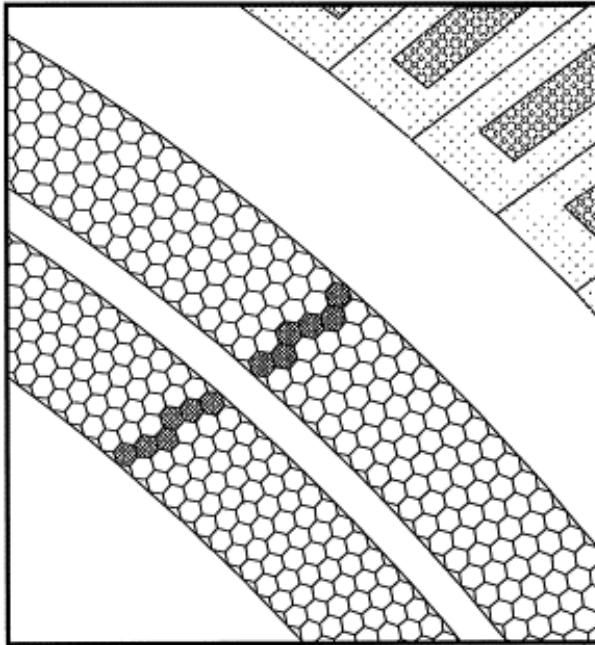


Figure 2.5: Illustration of a charged particle interacting with cells in the drift chamber. This figure was reproduced from [6]

### 2.2.2 CLAS Cherenkov Counter

The Cherenkov counters (CC), located radially outside of the region 3 drift chambers, greatly assisted in the separation of electrons and negatively charged pions for tracks with momentum less than the pion momentum threshold  $p < p_\pi \approx 2.5 \text{ GeV}/c$  [5]. The CLAS CCs were divided into 6 sectors, like most other detectors. Each sector was divided into 18 segments in the polar angle  $\theta$  away from the beamline. Furthermore, these segments were divided in half azimuthally to produce 12 half-sectors. Three mirrors, a light collecting Winston cone, a magnetic shield, and a 5 inch quartz face PMT were fitted to each of the 18 segments in all 12 half-sectors. During operation each CC was filled with  $6 \text{ m}^3$  of  $C_4F_{10}$  gas. The number of photo-electrons produced was recorded for tracks with polar angles between  $8^\circ < \theta < 45^\circ$ .

### 2.2.3 CLAS Time of Flight Scintillator

Measurements of average velocity can be made by simply knowing the distance some object has traveled in a given time period. Operating on this principle the CLAS time of flight (TOF) system allowed for the separation of  $\pi$  and  $K$  for momentum  $p \leq 2 \text{ GeV}/c$ . Constructed of 57 scintillating bars per sector, the TOF system covered an impressive area of  $206 \text{ m}^2$  and spanned the range of polar angles  $8^\circ \leq \theta \leq 142^\circ$ . Each of the scintillating bars measured 5.08 centimeters in thickness, 15 or 22 centimeters in width, and measured between 32 and 450 centimeters in length. Shorter bars, which covered the relatively higher rate low scattering angle, were built with an intrinsic timing resolution of 80 picoseconds.

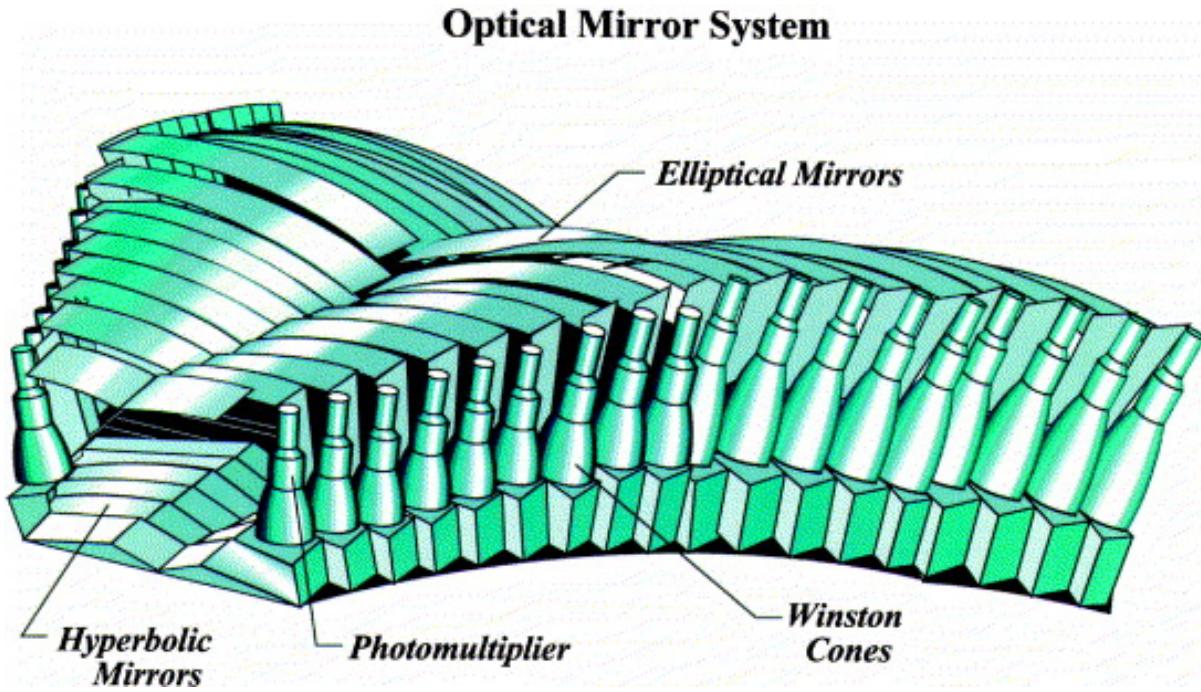


Figure 2.6: The CLAS Cherenkov Counter.

onds. The timing resolution for longer bars was designed to be 160 picoseconds, and measurements performed after the detector construction showed that the average timing resolution over the detector was 163 picoseconds [7].

During experiments where the electron beam was used, the start time for each event could be determined by assigning  $\beta = 1$  to the fastest electron measured in the final state. However, once the tagger magnet was powered on and photon beam was delivered on target, start time information came from the start counter. Originally designed as six separate counters, the start counter was constructed with three counters spanning the full angular range of the detector. The target was surrounded by these three thin scintillators which had sufficient resolution to determine the difference between two sequentially arriving bunches ( $\sigma \approx 350 \text{ ps}$ ).

## 2.2.4 CLAS Electromagnetic Calorimeter

The outermost layer of the CLAS detector was the electromagnetic calorimeter (EC). This sampling calorimeter was a main component of the CLAS trigger and had several important roles. Foremost, the EC detected and triggered on electrons with  $E > 0.5 \text{ GeV}$ . Detecting neutral particles such as photons and neutrons was a secondary role. By detecting photons with energies higher than 200 MeV, and having sufficient granular resolution, strongly decaying ( $\pi^0$  or  $\eta$ )  $\rightarrow \gamma\gamma$  particles were measured. Separation of neutrons and photons was achieved by combining information from the EC with timing information from the CLAS time of light system [8].

Structurally, the EC was composed in total of 1296 PMTs and 8424 scintillating strips in the six EC

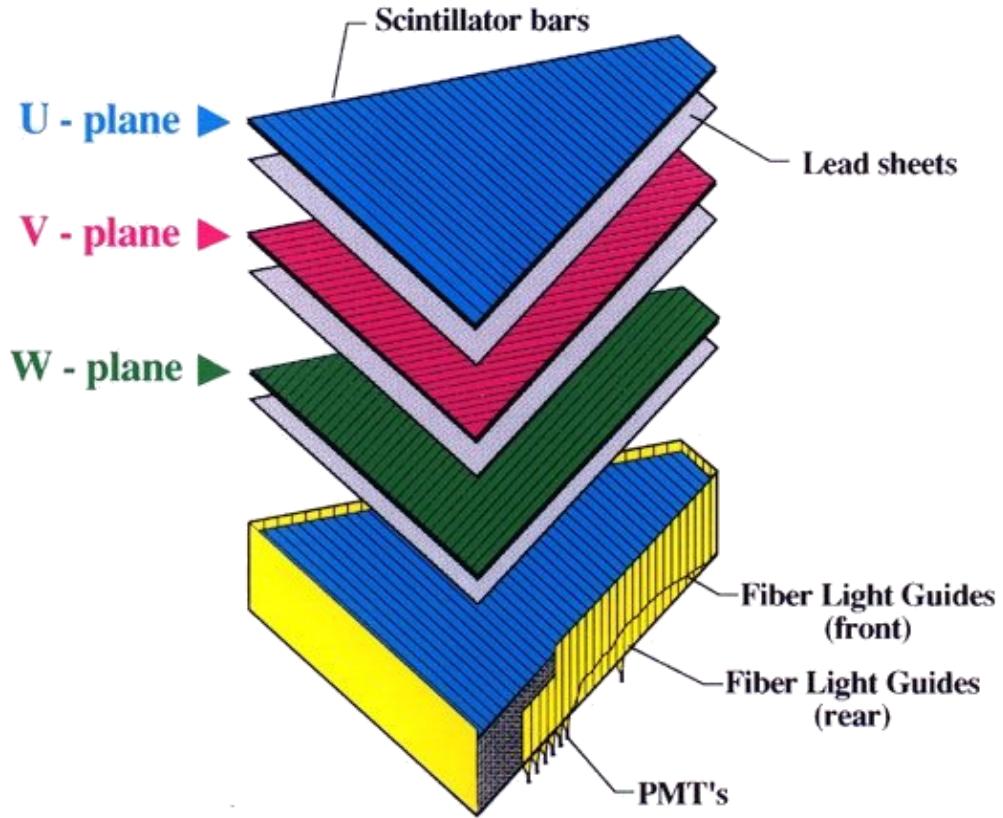


Figure 2.7: The U, V, W structure of CLAS electromagnetic calorimeter is shown above.

modules (one per sector). Alternating layers of lead and scintillator material were used to create a sampling fraction  $E_{dep}/p$  of approximately 0.3 for electrons. Measuring just 10 mm thick, and with a width of 10 cm (to balance cost of PMTs and granularity), the length of the scintillating strips depended on the angular location. Each EC module contained 3 sets of 13 layers offset by  $120^\circ$  to provide spatial information, these layers were referred to as U, V, and W.

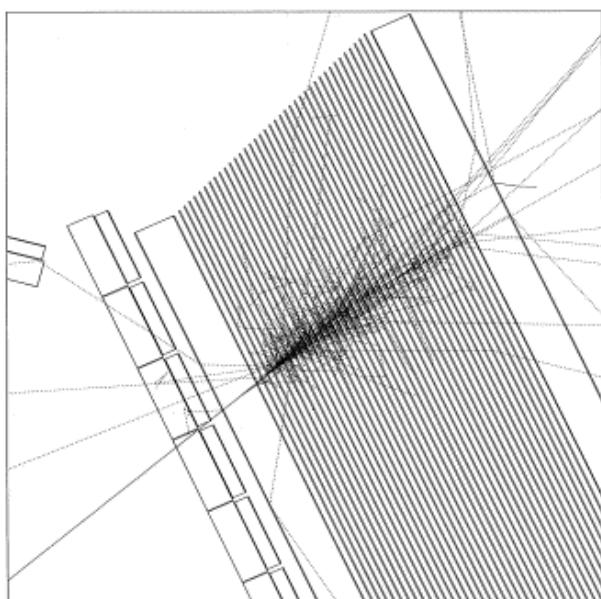


Figure 2.8: GEANT simulation of an electron showering in the EC. This figure originally appeared in [8]

# Chapter 3

## Basic Analysis & Corrections

### 3.1 Introduction

This chapter discusses analysis procedures that are common to the subsequent data analyses of pions and kaons. These procedures can be divided into two groups. The first type of basic analysis described is the aggregation or calculation of scalar values over the run-period (examples include luminosity and helicity). The second type of analysis procedure described is a correction to measured values. Vertex corrections, timing corrections, and kinematic corrections will be discussed.

### 3.2 Luminosity Calculation

A useful concept in accelerator/collider physics is the luminosity  $\mathcal{L}$ . Luminosity is defined as the number of collisions per unit area per unit time that could lead to some process of interest. Consider as an example elastic scattering of electrons from protons, the luminosity is the number of electron-proton collisions per unit time per unit area. The rate  $\frac{dN}{dt}$  of the occurrence of events for some process  $X$  can be written in terms of this luminosity and the cross section for the process.

$$\frac{dN_X}{dt} = \mathcal{L}\sigma_X \quad (3.1)$$

For the fixed target case, the luminosity has a simple expression.

$$\mathcal{L} = \frac{j_e \rho_p l_T}{e} \quad (3.2)$$

Here  $l_t$  is the target length,  $\rho_p$  is the proton number density in the target, and  $j_e$  is the beam current. To find the total number of events which accumulate in some time  $t_{exp}$  the event rate is integrated with respect to time.

$$N_X = \int_0^{t_{exp}} \frac{j_e \rho_p l_T}{e} \sigma_X dt = \frac{\rho_p l_T}{e} \sigma_X \int_0^{t_{exp}} j_e dt = \frac{\rho_p l_T}{e} \sigma_X \Delta Q \quad (3.3)$$

Thus the experimentally observed cross section for some process  $X$  is,

$$\sigma_X = \frac{N_X}{\mathcal{L}_{int}} \quad (3.4)$$

where the number of events  $N_X$  is corrected for all effects and  $\mathcal{L}_{int}$  is the integrated luminosity as shown above.

Experimentally, the factor  $\Delta Q$  can be calculated from charge deposition measurements by the Faraday cup. The Faraday cup charge is a scalar value written periodically into the output event stream, not with every recorded event. This information is stored in the output BOS files in a bank called TRGS, the variable is named FCUP\_G2.

For this data analysis, the final ntuple (root files) used did not contain the Faraday cup charge information. For this reason, the authors used the BOS files directly and recorded the value of FCUP\_G2 for every scalar reading, as well as the event number directly after each scalar entry (from the HEAD bank). This event number correlates directly to the event number stored in the root files used for analysis.

The total accumulated charge over a run is simply the sum over consecutive differences in the Faraday cup charge.

$$\Delta Q = \sum_{i=1}^{n-1} q[i+1] - q[i] \quad (3.5)$$

Here  $n$  denotes the number of scalar entries for a given file. Due to the periodic nature of the scalar bank writing events are also recorded after the last reading of the file, and before the first scalar reading of the next file in the run. To account for this the difference between consecutive files last and first readings is added to the total.

For the E1-F dataset, a run typically contains around 20 files, each representing a raw file size of 2 gigabytes. These files are named by run number, and given an index from 0 to  $n_{files} - 1$ . It is not uncommon that a run will contain missing files in the middle of the range. If this occurs, the charge difference between last/first reading is not added to the total.

Any charge which accumulates in a period of time where the number of events did not change is not

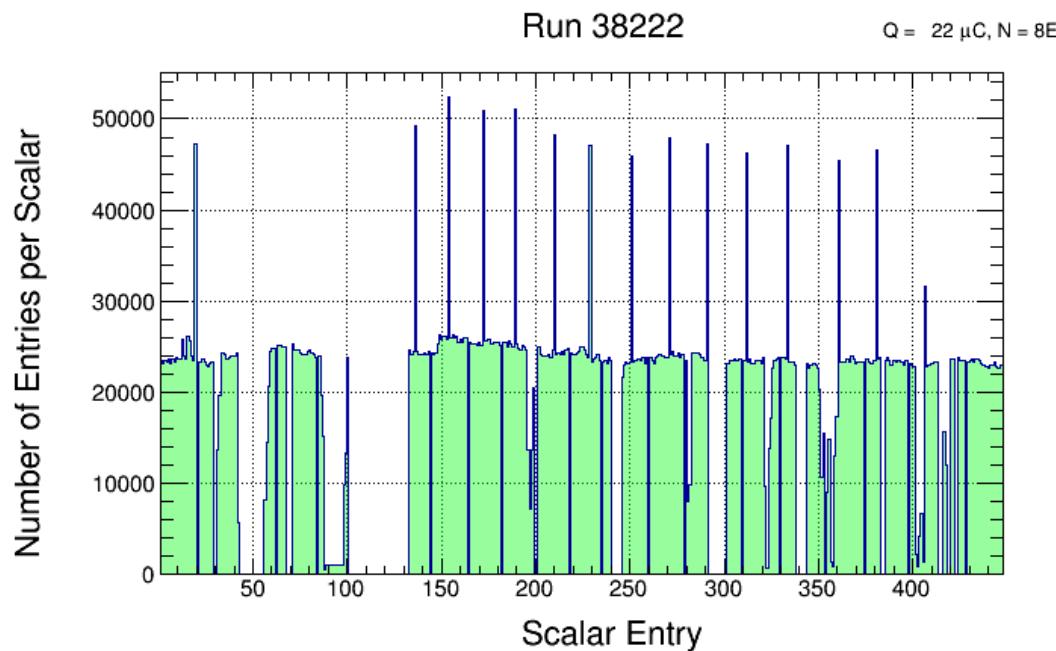


Figure 3.1: Entries per scalar reading for run 38222.

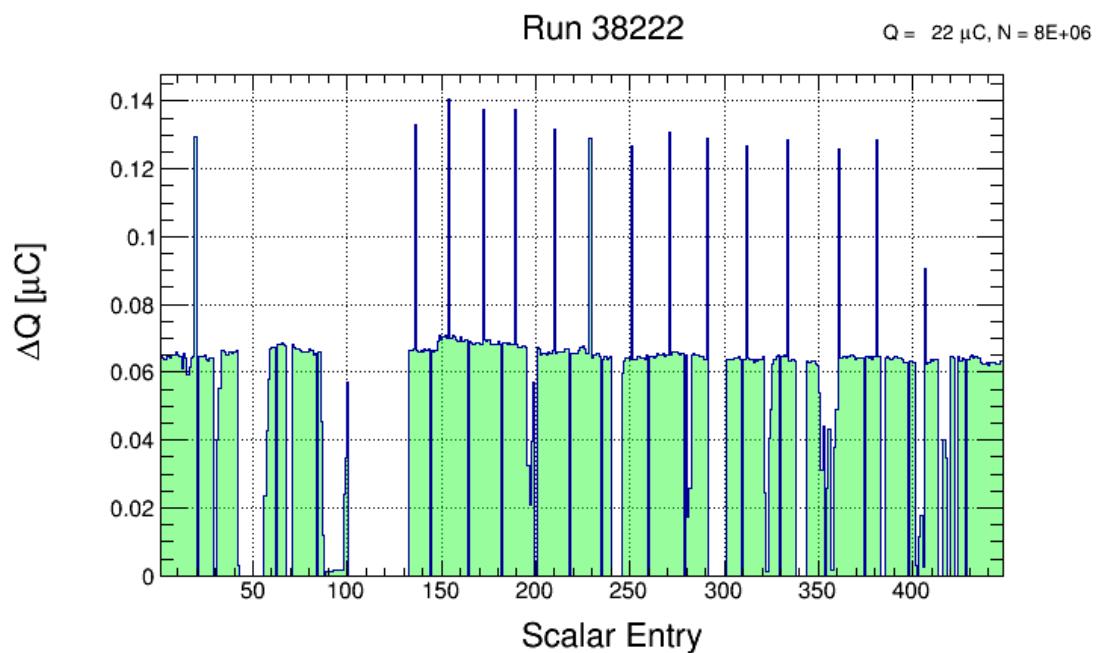


Figure 3.2: Charge per scalar reading for run 38222.

added to the total. Similarly, any events which occur within regions where no charge is recorded need to be discarded. This is accomplished by recording the bad event ranges for every file and removing these events from our analysis.

The result of this procedure is a numerical value of charge for each run. For practitioners, it is important to note that this value needs to be scaled by the DAQ scaling factor before it represents a value of charge. In our analysis, the numerical value of charge for a typical file is a few tens of micro-Coulombs.

### 3.3 Determination of Good Run List

The total dataset contains 831 runs. Due to the complexities of the CLAS experimental setup, it is not uncommon for run conditions to change during a small percentage of the runs such that the data collected are not of analysis quality. For this reason, a good run list is constructed.

Good runs are selected for the list by counting good electrons in each file and normalizing by the accumulated charge for the associated file. While the number of events collected varies from run to run the ratio defined above is a stable quantity – provided that the run conditions do not vary greatly. Runs which are within 3 standard deviations of the mean (calculated over the dataset) are used as good runs. The good run list used for this analysis contains 522 runs.

### 3.4 Helicity Determination

During the course of the E1-F run period the beam helicity convention was changed by the insertion of a half-wave plate at the injector. The definition of  $\pm$  helicity must change in accordance with these wave-plate insertions. To monitor these changes, the value of  $A_{LU}^{\sin \phi}$  for  $\pi^+$  is recorded for every run. Whenever the asymmetry (which has a magnitude of around 3%) changes sign, the sign convention has changed. These changes are taken into account in the data analysis.

### 3.5 Vertex Corrections

The track vertex position  $(v_x, v_y, v_z)$  is calculated based on the intersection of each track with the mid-plane (the plane which contains the beamline and bisects the sector at  $\phi_{rel} = 0$ ). If the beam is not centered at  $(x, y) = (0, 0)$ , the vertex position calculation needs to be corrected by shifting the mid-planes in accordance with the target offset. The offset  $(x, y)$  is identified by plotting events from the control foil placed near the target, which has a  $z$  position of -20 cm. For the E1-F run period, the beam

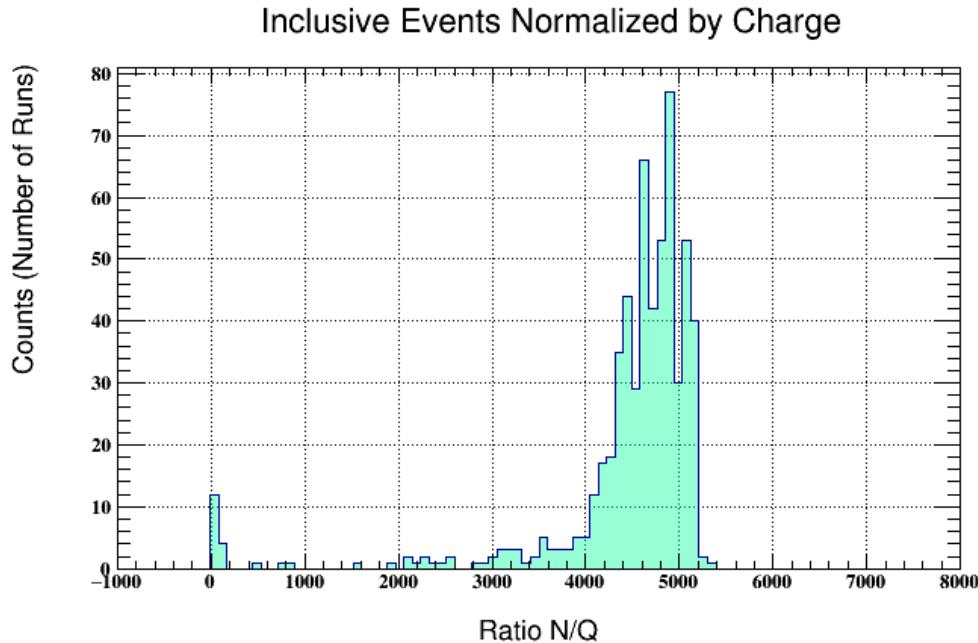


Figure 3.3: Inclusive electrons per file normalized by the total charge accumulated for the file. This quantity is used to make a good run list.

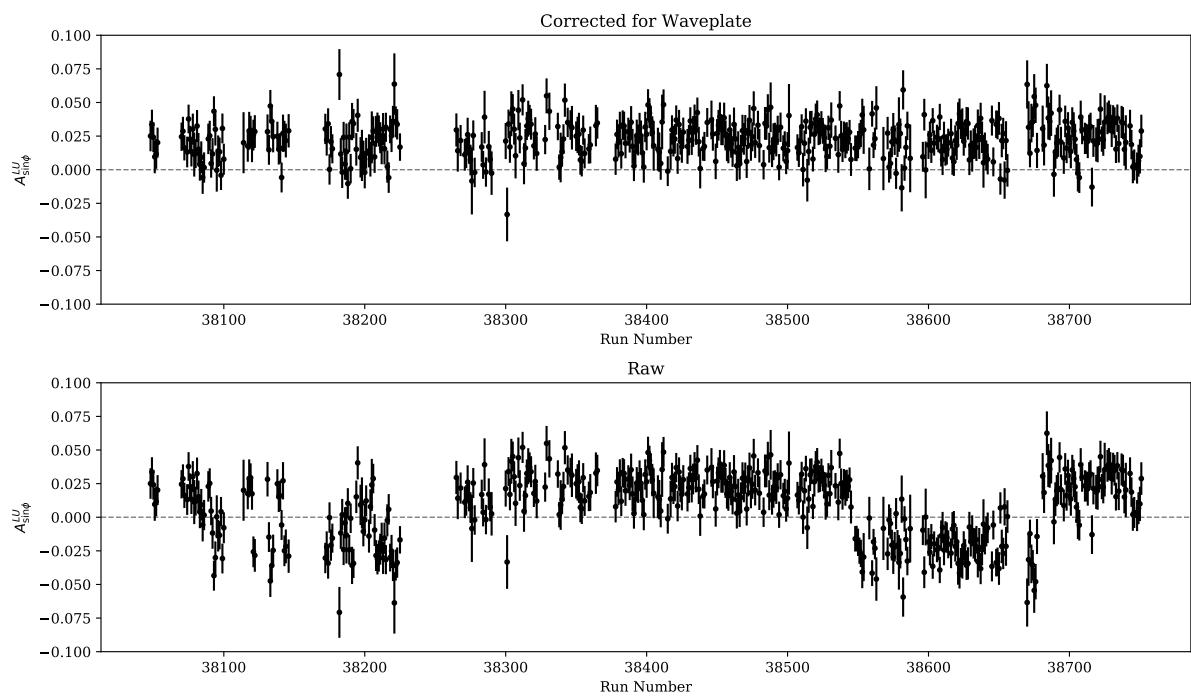


Figure 3.4: The waveplate position is determined and corrected by plotting the BSA for  $\pi^+$  mesons as a function of the run. The top panel shows the corrected results, the bottom shows the results before changing the helicity.

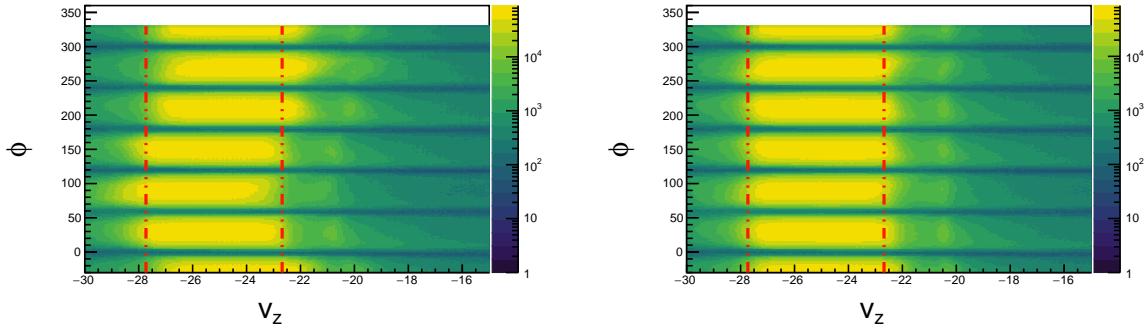


Figure 3.5: The z-vertex  $v_z$  position shown for different values of  $\phi$  the azimuthal angle in the hall. The left figure shows the distribution before corrections are applied, the right after. The vertical red lines bound the region which we define as acceptable for electrons in our analysis.

position was  $(0.15, -0.25)$  cm.

### 3.6 Timing Corrections

Timing information comes from the time-of-flight detector system. After calibration, small offsets in timing between time of flights paddles still exist for the E1-F dataset. These biases can be removed on a run-by-run and paddle-by-paddle basis by adding a small shift  $t_{corr}$ . In order to determine this shift  $t_{corr}$  for each paddle, charged pions are used.

Using momentum information from the drift chambers the value of  $\beta$  can be predicted and the difference  $\Delta\beta$  can be determined for each pion.

$$\Delta\beta = \beta_{obs} - \beta_{pred} = \frac{d}{ct_{obs}} - \sqrt{1 + (m/p)^2} \quad (3.6)$$

Here  $m$  is assumed to be  $m_\pi$ . The offset  $\Delta\beta$  from 0 is used to define the value of  $t_{corr}$  for each paddle. If this value is exceedingly small, no correction is applied. For some paddles with low statistics a reasonable value for  $t_{corr}$  cannot be obtained and these paddles are excluded from the analysis.

In the method described above, the calibrated paddle is the one which is struck by the pion. The electron paddle which was struck could also require calibration. In practice the magnitude of the correction term  $t_{corr}$  is small, and the paddle offset is (likely) randomly distributed about 0 when considering all paddles. By including events from many different (electron) paddles, miscalibration effects from the electron side cease to be important. This is demonstrated by the success of the technique in centering the  $\Delta\beta$  distributions. This work was first described in [9].

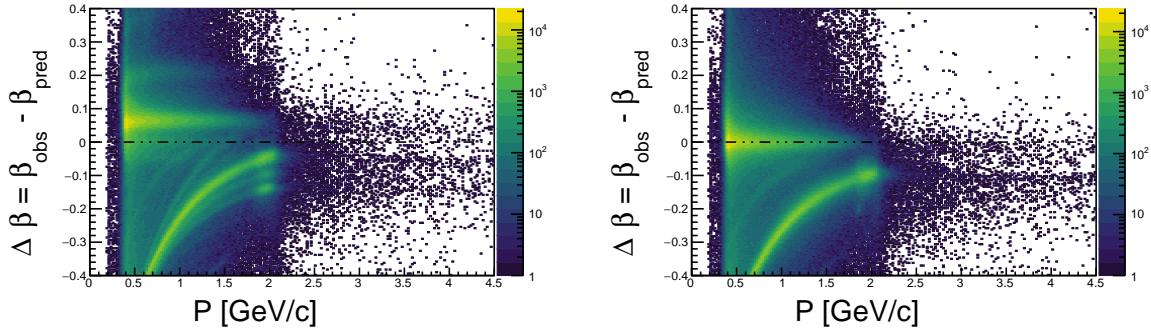


Figure 3.6: Timing corrections are shown for paddle 24 of sector 1. The left image shows the  $\Delta\beta$  distribution before corrections. On the right the same is shown after correction of the timing for this paddle. We assume the mass of the track to be the pion, these show up as the green band. Heavier protons are visible below the pion band.

### 3.7 Kinematic Corrections

The magnetic field map used in reconstruction to swim particle tracks cannot perfectly match the real magnetic field of the hall. As a result of this the reconstructed momentum of particles is often slightly off (of order 1%). Small misalignments in detector positions also contribute to this effect. In order to correct for these small differences, the momentum  $(p_x, p_y, p_z)$  and hence  $\theta$  of charged tracks is corrected.

Various procedures exist for the correction of kinematic variables of measured particles, and they all rely on energy and momentum conservation applied to standard processes (such as elastic scattering). The procedure used to derive corrections for the E1-F dataset was developed and described by Marco Mirazita in [10].

As mentioned previously, the need for correction to  $\theta$  (the polar angle measured from the beamline) arises from misalignments in the drift chambers. This implies that the correction will be the same for positives and negatives, and this assumption is used in the correction algorithm. First, elastic ( $ep \rightarrow ep$ ) events are selected by identifying events that contain at least one electron and one proton, then requiring that the missing mass  $M_X$  of the  $(ep \rightarrow epX)$  system is close to 0. The kinematics of the event are then calculated.

$$k^\mu = (k, 0, 0, k) \quad (3.7)$$

$$p^\mu = (M_p, 0, 0, 0) \quad (3.8)$$

$$k'^\mu = (k', k' \sin \theta, 0, k' \cos \theta) \quad (3.9)$$

$$p'^\mu = (E_p, -p' \sin \alpha, 0, p' \cos \alpha) \quad (3.10)$$

Applying energy and momentum conservation to the equations above yields 3 equations.

$$k + M_p = k' + \sqrt{M_p^2 + p'^2} \quad (3.11)$$

$$k' \sin \theta = p' \sin \alpha \quad (3.12)$$

$$k = k' \cos \theta + p' \cos \alpha \quad (3.13)$$

Using these equations, the electron angle  $\theta$  and the proton angle  $\alpha$  can be predicted by using the momenta  $(k', p')$ . These values are compared with measured values and iteratively corrected by tuning the parameters of a phi-dependent 2nd order polynomial.

$$\cos \theta = 1 - M_p \frac{k - k'}{kk'} \quad (3.14)$$

$$\tan \alpha = \frac{1}{p'} \frac{k' \sin \theta}{k - k' \cos \theta} \quad (3.15)$$

After  $\theta$  corrections are applied, the momentum of the electrons are corrected by using an analogous procedure for  $k'$  instead of  $\theta$  and  $\alpha$ . The momentum corrections are calculated as functions of  $\phi$  for each sector in one degree bins of  $\theta$ . Finally, the positively charged particles momenta are corrected by selecting the exclusive event ( $ep \rightarrow e\pi^+N$ ). In this reaction the scattered electron and pion are detected and the neutron is selected using a missing mass cut. Assuming the electron momentum, electron angle, and pion angle to be correct, the pion momentum correction is then calculated by iteratively improving the central position of the neutron mass peak to coincide with  $M_N$ . Marco Mirazita shows in his note that these corrections can be satisfactorily applied to all negative and positive particles.

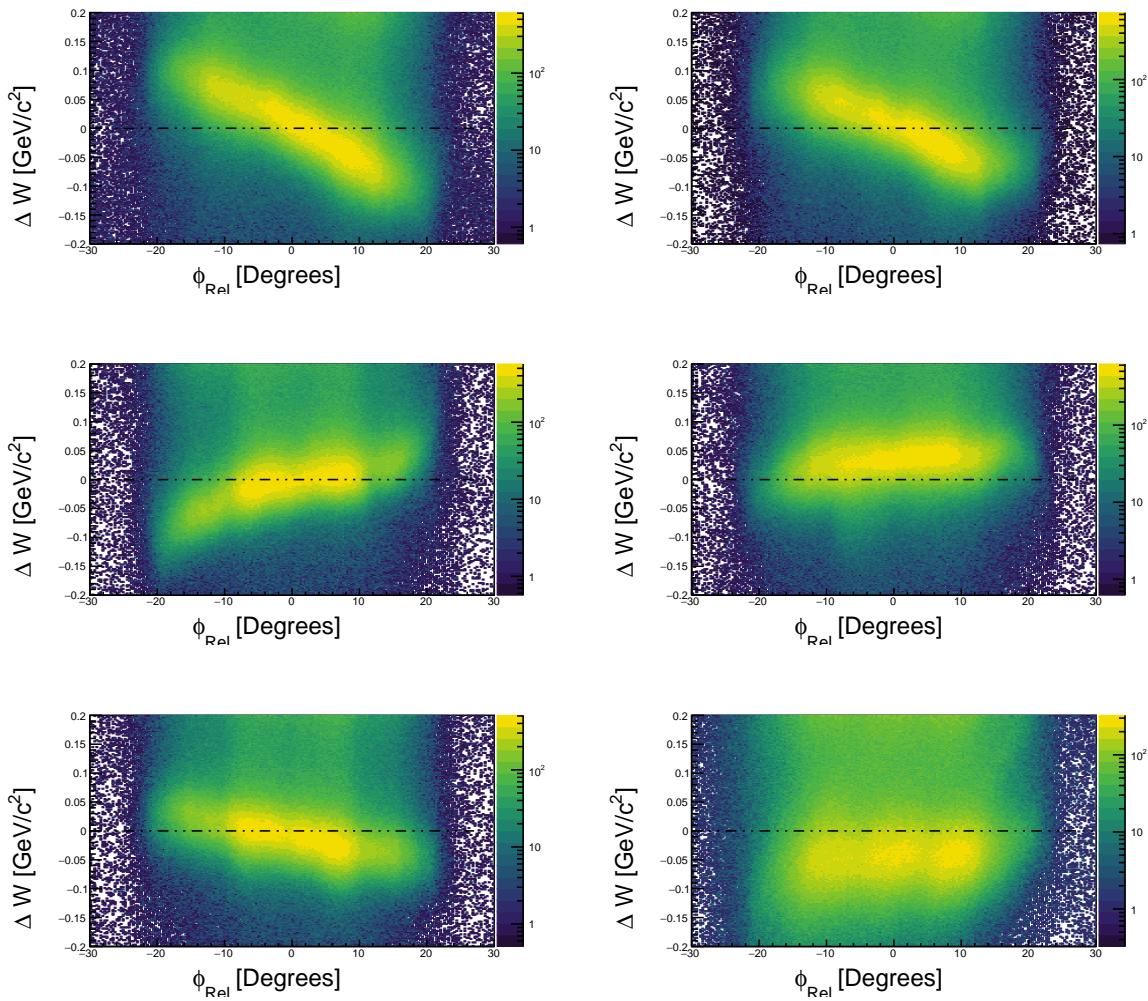


Figure 3.7: This figure shows the deviation from  $M_p$  of the  $W$  spectrum peak for elastic  $ep \rightarrow ep$  events (before corrections).

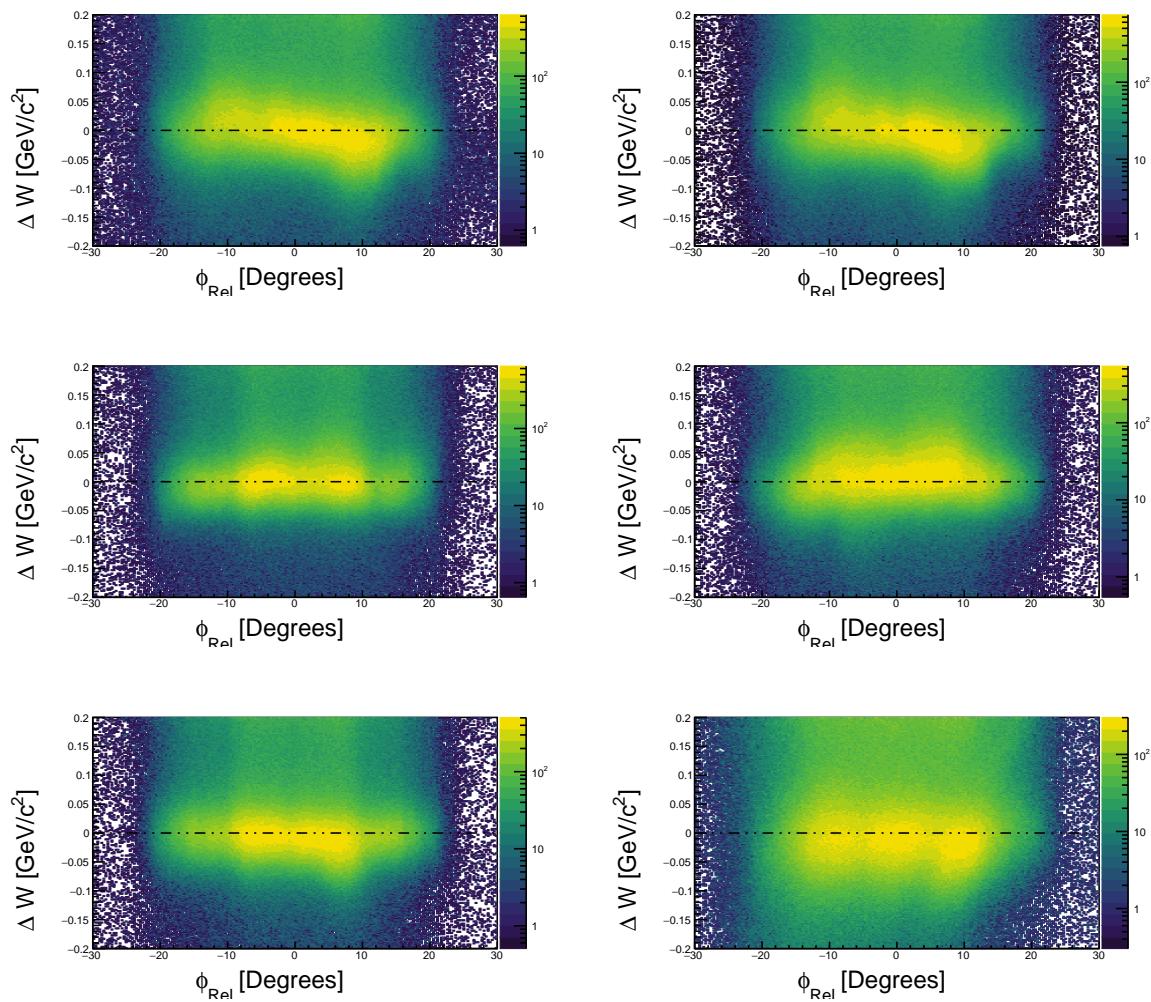


Figure 3.8: This figure shows the deviation from  $M_p$  of the  $W$  spectrum peak for elastic  $ep \rightarrow ep$  events (after  $\phi$ -dependent corrections).

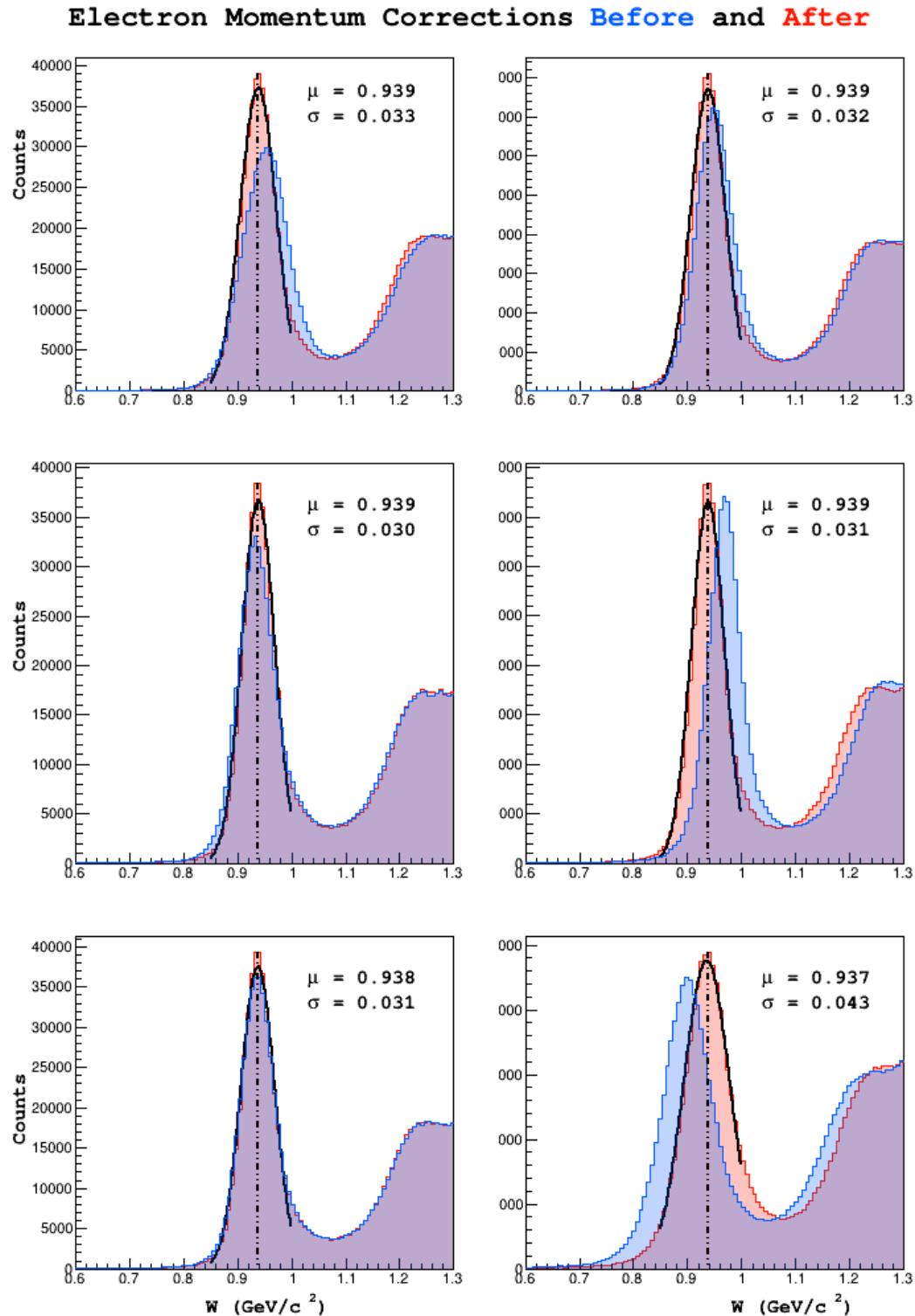


Figure 3.9: Elastic events shown in the spectrum of  $W$  before and after momentum corrections are applied.



# Chapter 4

## Particle Identification

### 4.1 Introduction

Particle identification (PID) is the process of classifying tracks as known particles. After reconstruction and matching of detector responses to each track, the reconstruction package `reccsis` assigns a preliminary particle identification based on loose selection criteria. In this analysis, tracks are classified based on a more stringent criteria. This chapter discusses the methodology used by the authors to classify particles.

### 4.2 Electron Identification

Electrons in CLAS are abundant, and the detection of an electron is a basic necessity for every event that will be analyzed. The most naive approach to performing electron identification would be to call all negatively charged tracks electrons. Doing this would provide an extremely efficient identification of electrons (none of them are missed), however the purity of the sample (the fraction of tracks identified as electrons that are actually electrons) would be low due to the vast quantity of negatively charged pions that are produced in during the experiment. Additionally, doing this would completely eliminate the possibility of identifying negatively charged pions or kaons, as all negative tracks would be called electrons. In practice then, the identification of electrons is concerned with removal of negative pions and kaons from the sample of negative tracks. This is accomplished by applying a series of cuts on measured variables that distinguish between electrons and pions (pions are the dominant background).

#### 4.2.1 Electron ID Cuts

The cuts used to select electrons are enumerated below.

- Negative charge

- Drift chamber region 1 fiducial
- Drift chamber region 3 fiducial
- Electromagnetic Calorimeter fiducial (UVW)
- EC minimum energy deposition
- Sampling Fraction (momentum dependent)
- z-vertex position
- Cherenkov counter  $\theta_{cc}$  matching to PMT number
- Cherenkov counter  $\phi_{rel}$  matching to PMT (left/right)

Each cut is now be described in more detail.

### Negativity Cut

Each track is assigned a charge based on the curvature of it's trajectory through the magnetic field of the torus. This is done during the track reconstruction phase. Tracks are eliminated as electron candidates if they are not negatively charged.

### Drift chamber fiducial

The fiducial region or volume is a term used to refer to the region of a sensitive detector which is unimpeded in it's acceptance of physics events. In practice, shadows from other detectors, poorly understood edge effects, or geometric obstacles may impede the flight of particles from the target, and render regions of sensitive detectors unreliable (to use the vocabulary presented above, these events fall outside of the fiducial region of the detector).

Negative tracks which pass geometrically close to the edges of the drift chamber are, from a tracking perspective, more difficult to understand. Additionally, tracks which fall outside of the fiducial region of the drift chambers are likely to fall outside of the fiducial region of the downstream detectors as well. For these reasons, it is common to remove tracks which are geometrically close to the boundaries of the drift chambers in region 1 as well as region 3 coordinate systems.

To implement this cut the  $(x, y)$  coordinates of the drift chambers are rotated into one sector. Then boundaries  $y_{left}, y_{right}$  are defined as linear functions of  $x$ . The boundary lines are parametrized by an offset  $h$  and an angle of the boundary line with respect to the center of the sector at  $x = 0$ . The slope of these lines is  $\pm \cot(\theta)$ .

Region	Height $h$ (cm)	Angle $\theta$ (degrees)
1	22	60
3	80	49

Table 4.1: Cut parameters used for the DC fiducial cut.

EC Coordinate	Min (cm)	Max (cm)
U	70	400
V	-	362
W	-	395

Table 4.2: Cut parameters used for the EC fiducial cut.

$$y_{right} = h + \cot(\theta)x \quad (4.1)$$

$$y_{left} = h - \cot(\theta)x \quad (4.2)$$

Tracks passing this criterion are those which have  $y > y_{left}(x)$  and  $y > y_{right}(x)$ .

### Electromagnetic Calorimeter fiducial (UVW)

As particles traverse the electromagnetic calorimeter they develop electromagnetic showers. If the track passes close to the edges of the detector, there is a chance that the shower will not be fully contained within the calorimeter volume (it spills out the edges). For this reason, it is standard to remove the hits which fall within the outer 10 centimeters of each layer of the EC (10 centimeters is the width of a scintillator bar). This cut is applied in the U, V, W coordinate system.

### EC minimum energy deposition

One way to differentiate between these electrons and pions is to exploit the difference in energy deposition between the two in the electromagnetic calorimeter. Electron typically develop a much larger and more energetic shower than  $\pi$  mesons, which minimally ionize the calorimeter material. The result is that the total energy deposition is typically larger for electrons than  $\pi$  mesons. In this analysis we require that at least 60 MeV was deposited in the inner calorimeter for electron candidates.

### Sampling Fraction (momentum dependent)

The electromagnetic calorimeter is designed such that electrons will deposit  $E_{dep}/p \approx 0.3$  approximately one-third of their energy, regardless of their momentum. In contrast to this, the ratio  $E_{dep}/p$  for  $\pi$  mesons decreases rapidly with momentum. To develop a momentum dependent cut for this distribution, all negative candidates are first filled into a two-dimensional histogram of  $E_{dep}/p$  vs.  $p$ . The histogram is then binned more coarsely in momentum, and projected into a series of 40 momentum

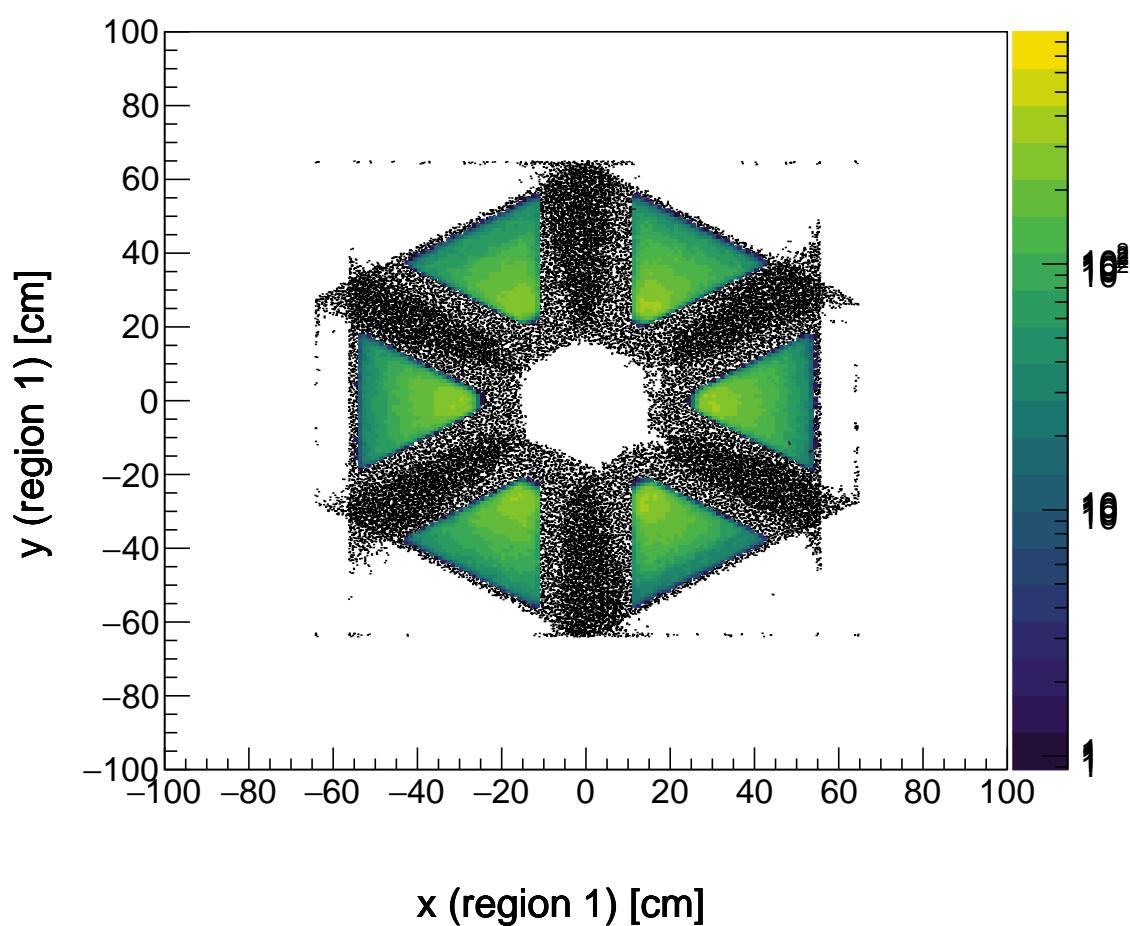


Figure 4.1: Tracks shown in color remain after the application of drift chamber region 1 fiducial cuts to all cuts, shown here as black points.

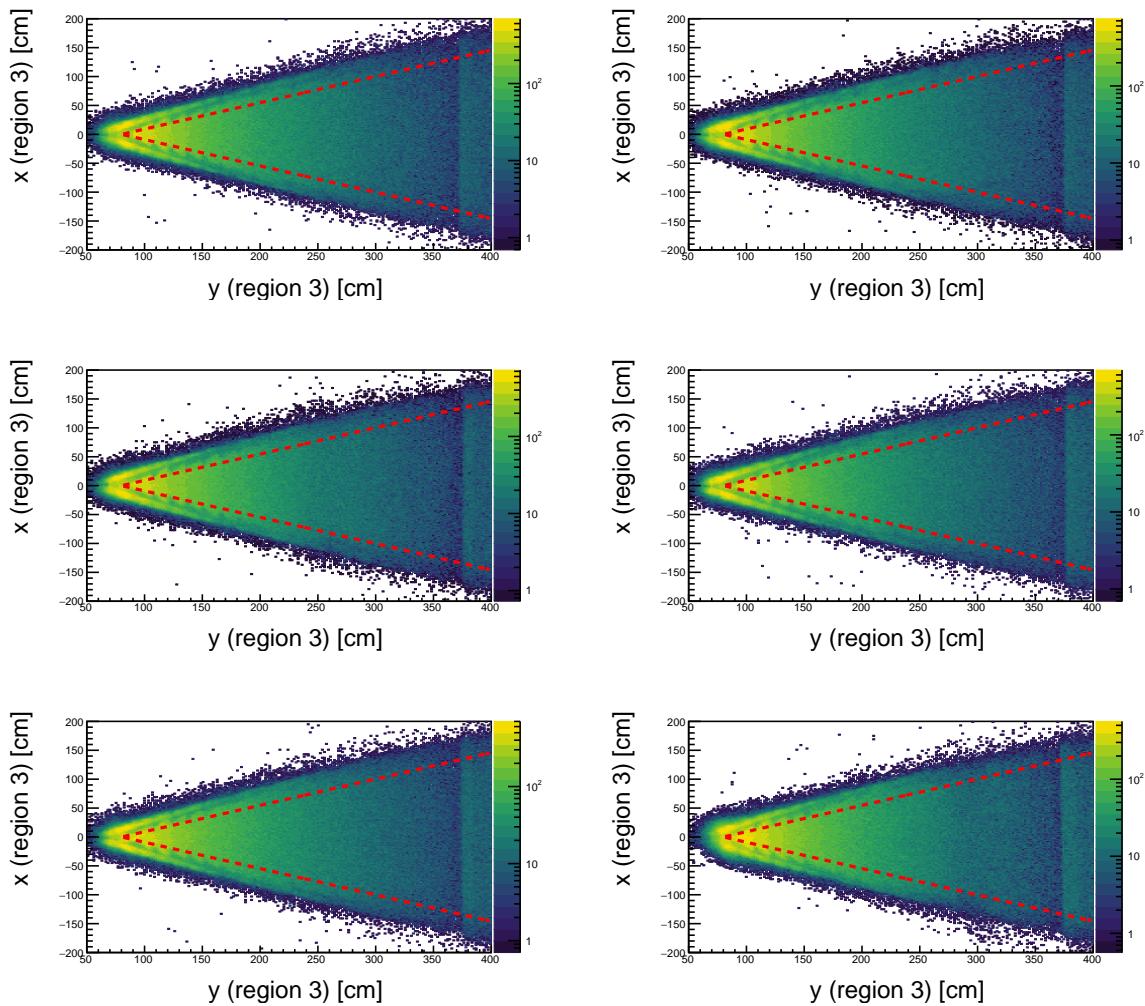


Figure 4.2: The selection criteria shown in red is applied to drift chamber region 3.

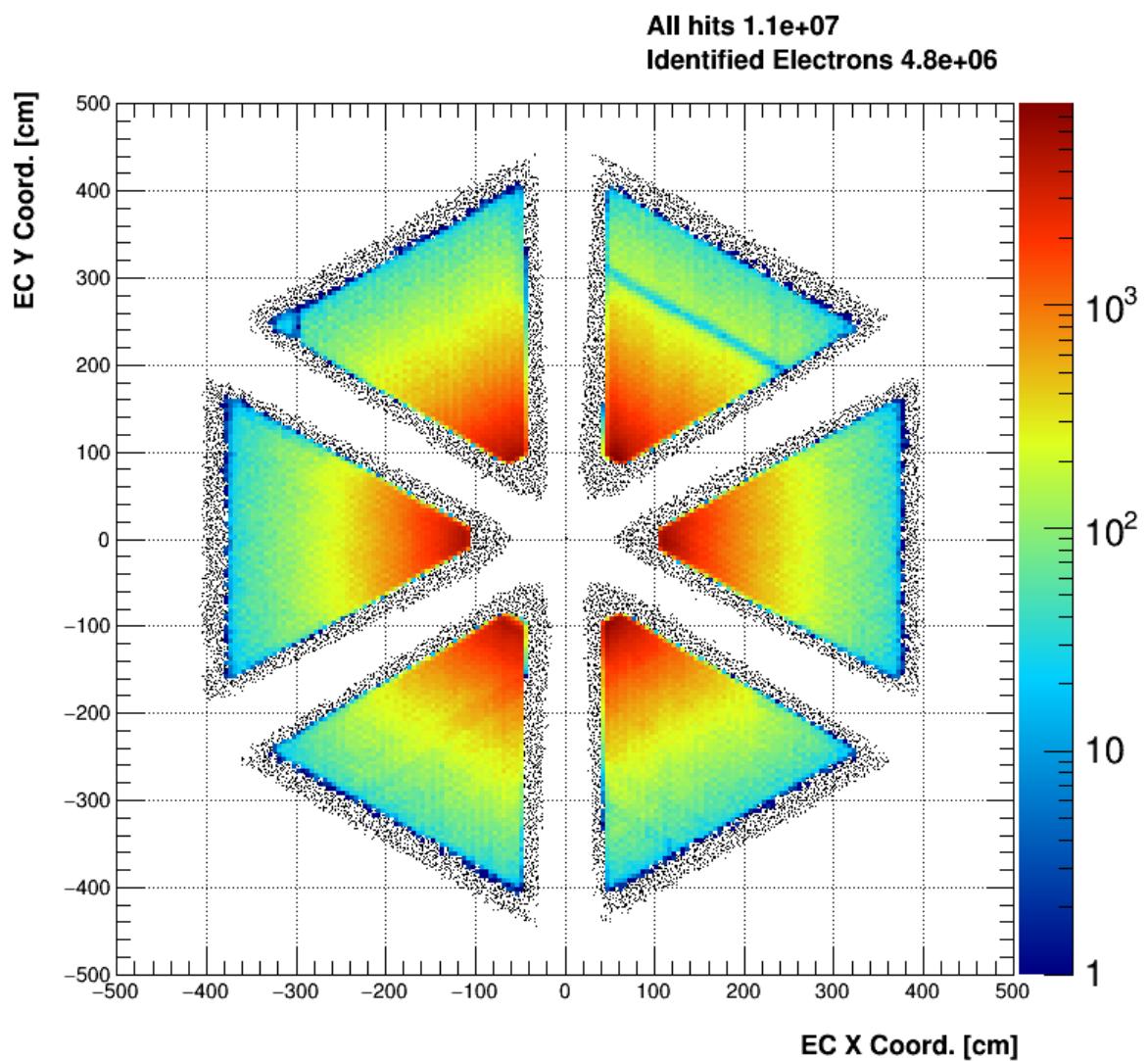


Figure 4.3: All negative tracks are shown here in black. In color, the tracks which pass the EC fiducial cut are shown.

slices. Each of these slices is fit with a Gaussian to extract the position  $\mu_i$  and width  $\sigma_i$  of the electron peak. Finally, a functional form for the mean and standard deviation of the distributions is chosen to be a third order polynomial in momentum.

$$\mu(p) = \mu_0 + \mu_1 p + \mu_2 p^2 + \mu_3 p^3 \quad (4.3)$$

$$\sigma(p) = \sigma_0 + \sigma_1 p + \sigma_2 p^2 + \sigma_3 p^3 \quad (4.4)$$

Boundaries are constructed from this information by adding (subtracting)  $n_\sigma$  from the mean. In the nominal case, we use  $n_\sigma = 2.5$ .

$$f_{max}(p) = \mu(p) + n_\sigma \sigma(p) = (\mu_0 + n_\sigma \sigma_0) + (\mu_1 + n_\sigma \sigma_1)p + (\mu_2 + n_\sigma \sigma_2)p^2 + (\mu_3 + n_\sigma \sigma_3)p^3 \quad (4.5)$$

$$f_{min}(p) = \mu(p) - n_\sigma \sigma(p) = (\mu_0 - n_\sigma \sigma_0) + (\mu_1 - n_\sigma \sigma_1)p + (\mu_2 - n_\sigma \sigma_2)p^2 + (\mu_3 - n_\sigma \sigma_3)p^3 \quad (4.6)$$

Due to slight differences between the 6 sectors of the CLAS detector, this cut is calibrated and applied for each sector individually. Results are shown in table 4.3.

Parameter	Sector 1	Sector 2	Sector 3	Sector 4	Sector 5	Sector 6
$\mu_3$	-8.68739e-05	0.000459313	9.94077e-05	-0.000244192	-7.65218e-05	-0.000392285
$\mu_2$	-0.000338957	-0.00621419	-0.00267522	-0.00103803	-0.00222768	-0.00105459
$\mu_1$	0.0191726	0.0393975	0.02881	0.0250629	0.0233171	0.0265662
$\mu_0$	0.2731	0.296993	0.285039	0.276795	0.266246	0.25919
$\sigma_3$	-0.000737136	0.000189105	-0.000472738	-0.000553545	-0.000646591	-0.000633567
$\sigma_2$	0.00676769	-0.000244009	0.00493599	0.00434321	0.00717978	0.00626044
$\sigma_1$	-0.0219814	-0.00681518	-0.0180929	-0.0140827	-0.0246181	-0.022029
$\sigma_0$	0.0474188	0.0475098	0.0461743	0.0492728	0.0546257	0.0517508

Table 4.3:  $\mu$  and  $\sigma$  values used to construct the momentum dependent sampling fraction cut.

### z-vertex position

Electrons can be produced as part of  $e^+e^-$  pairs, or by other processes. For this analysis, these are not of interest. For the purposes of this analysis it is then natural to accept only electron candidates which have a z-vertex  $v_z \in [-27.7302, -22.6864]$  within the expected target region. This cut is applied after the vertex position has been corrected (which is discussed in the basic analysis section).

### Cherenkov counter $\theta_{cc}$ and $\phi_{rel}$ matching to PMT

The angular arrangement of photo-multiplier tubes (PMTs) in the Cherenkov counter allows for additional consistency conditions to be applied. Each half-sector of the CC contains 18 PMTs increasing in polar angle away from the beamline, these divisions are known as segments. The polar angle measured

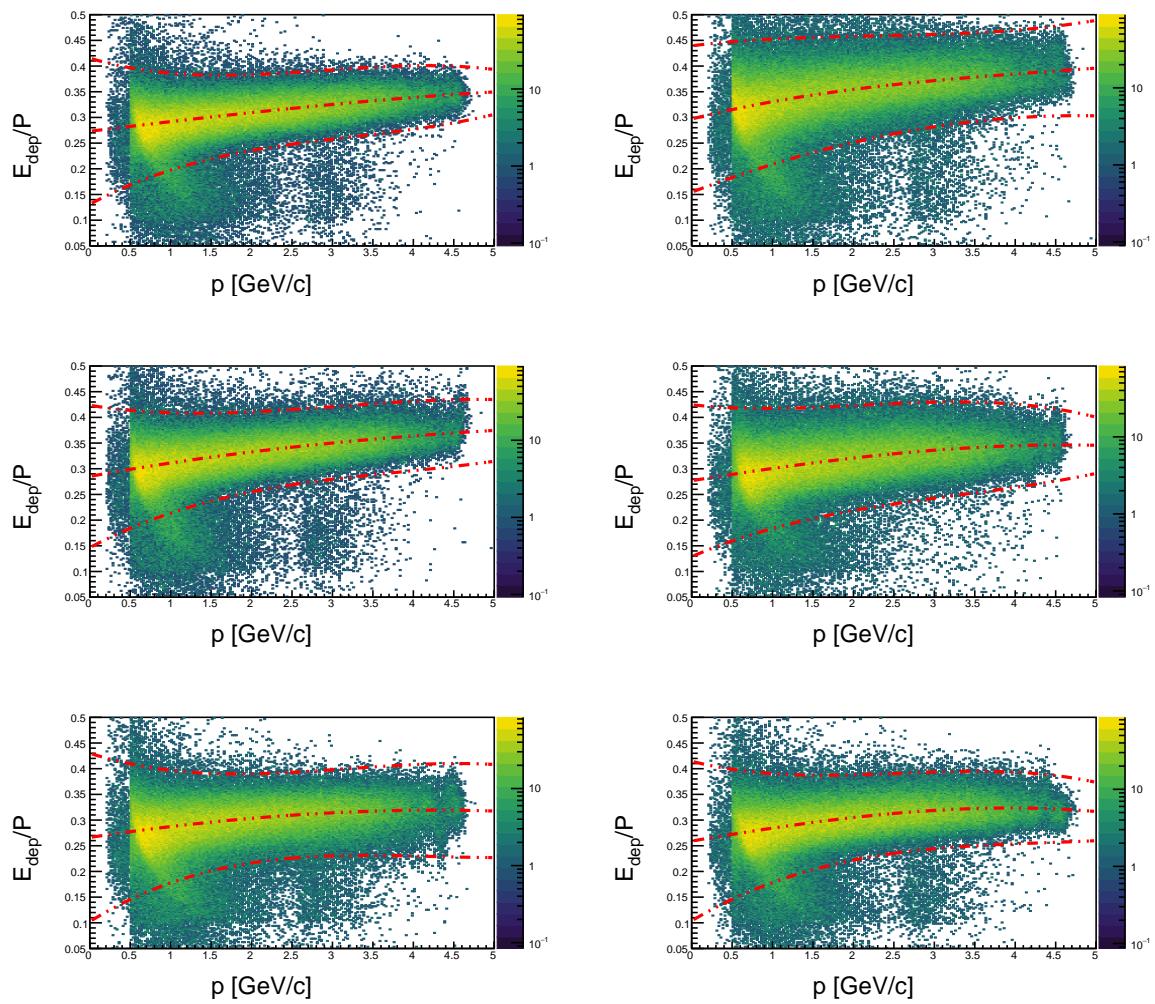


Figure 4.4: The sampling fraction selection boundary is shown here for the nominal value of  $N_{\text{sigma}} = 4$ .

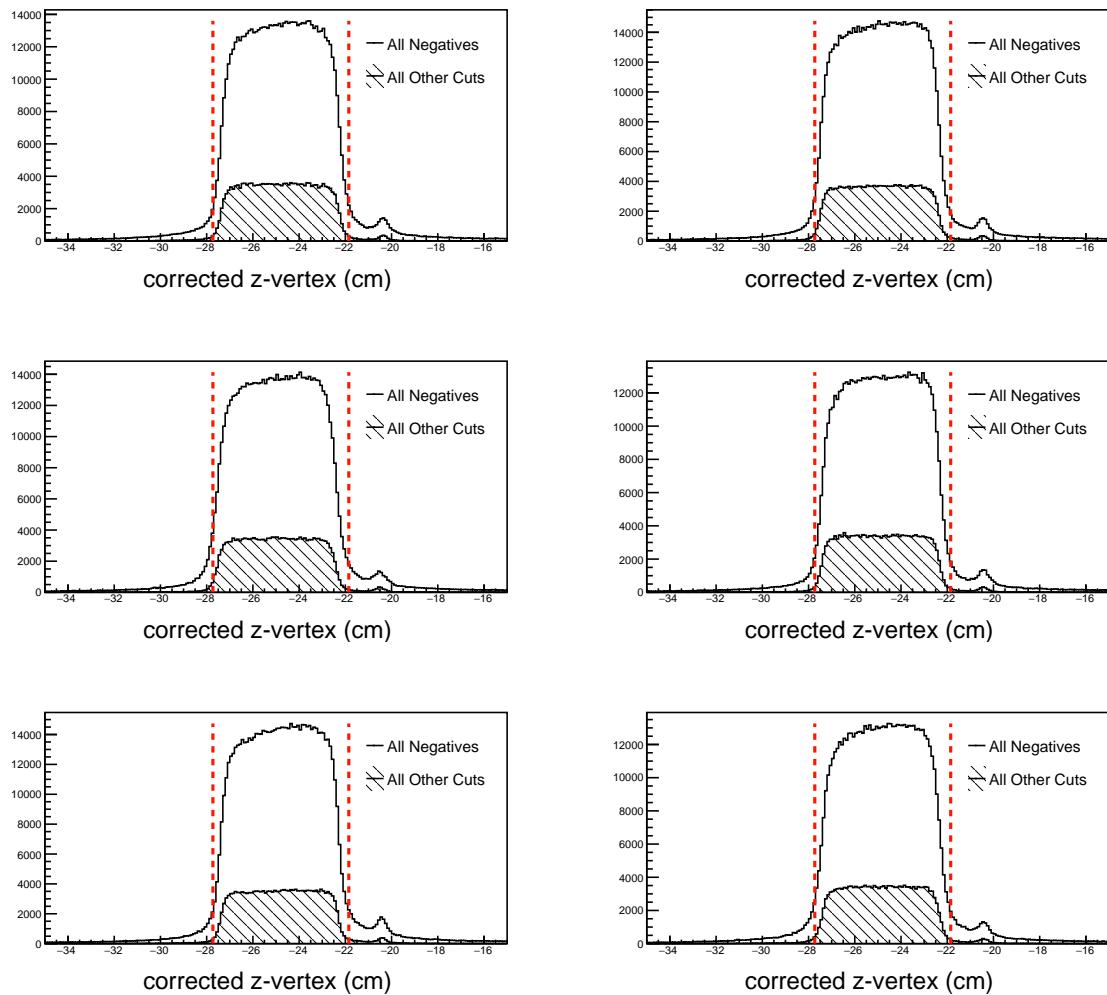


Figure 4.5: The track vertex cut is shown above. All negative tracks are shown in white, while the tracks passing all other criteria are shown in black hatch. The cut boundary is displayed as red lines. For E1-F the target center was located at -25 cm.

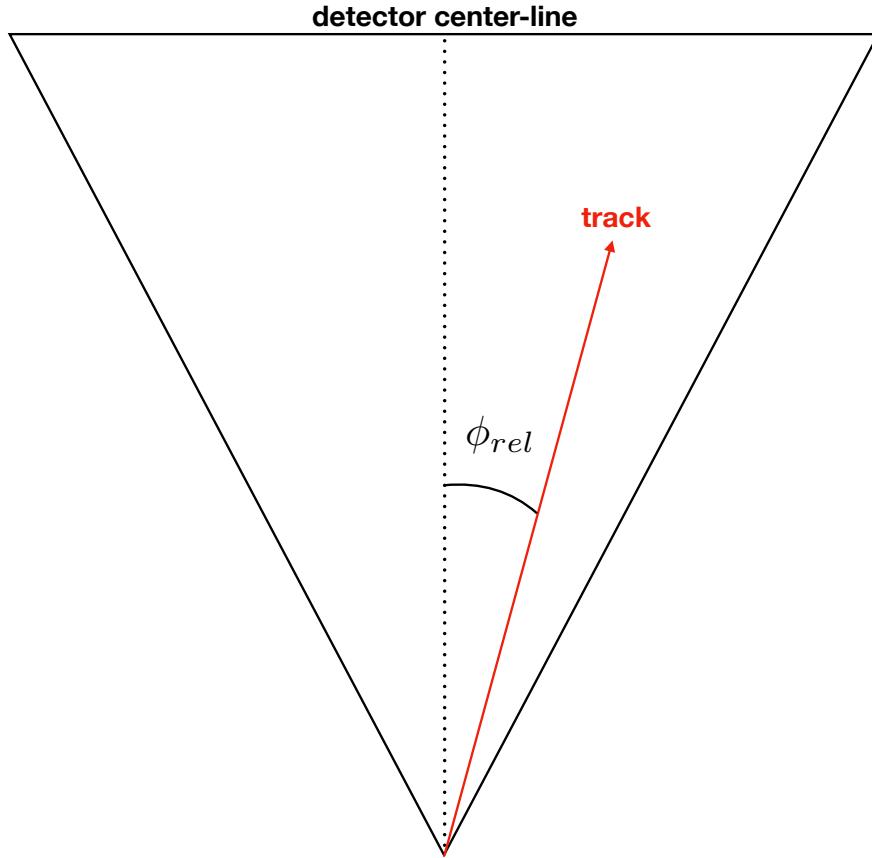


Figure 4.6: The angle  $\phi_{rel}$  is the azimuthal angle between the central line of the detector and the track.

at the Cherenkov counter  $\theta_{cc}$  is then correlated to the segment in which the track was detected. Additionally, PMTs that are placed on the left and right of the detector can be used to check consistency with the azimuthal angle the track forms with the central line of the detector (ie  $\phi_{rel} > 0$  means the track was in the right half of the sector,  $\phi_{rel} < 0$  means the track was in the left half of the sector). An integer is used to describe the PMT associated with the track. The left PMT is assigned value -1, the right 1, and a signal in both PMTs is assigned 0. If both PMTs have a signal, the track is allowed to pass. If the left PMT was the one that had a signal, only events with  $\phi_{rel} < 0$  passes. Similarly if the right PMT fired (code = 1), only events with  $\phi_{rel} > 0$  are allowed to pass. Technical note: the integers in question can be obtained from the ntuple22 format tree by doing the following.

---

```

for (int index = 0; index < event.gpart; index++) {
    int pmt = event.cc_segm[index]/1000 - 1;
    int segment = event.cc_segm[index]%1000/10;
}

```

---

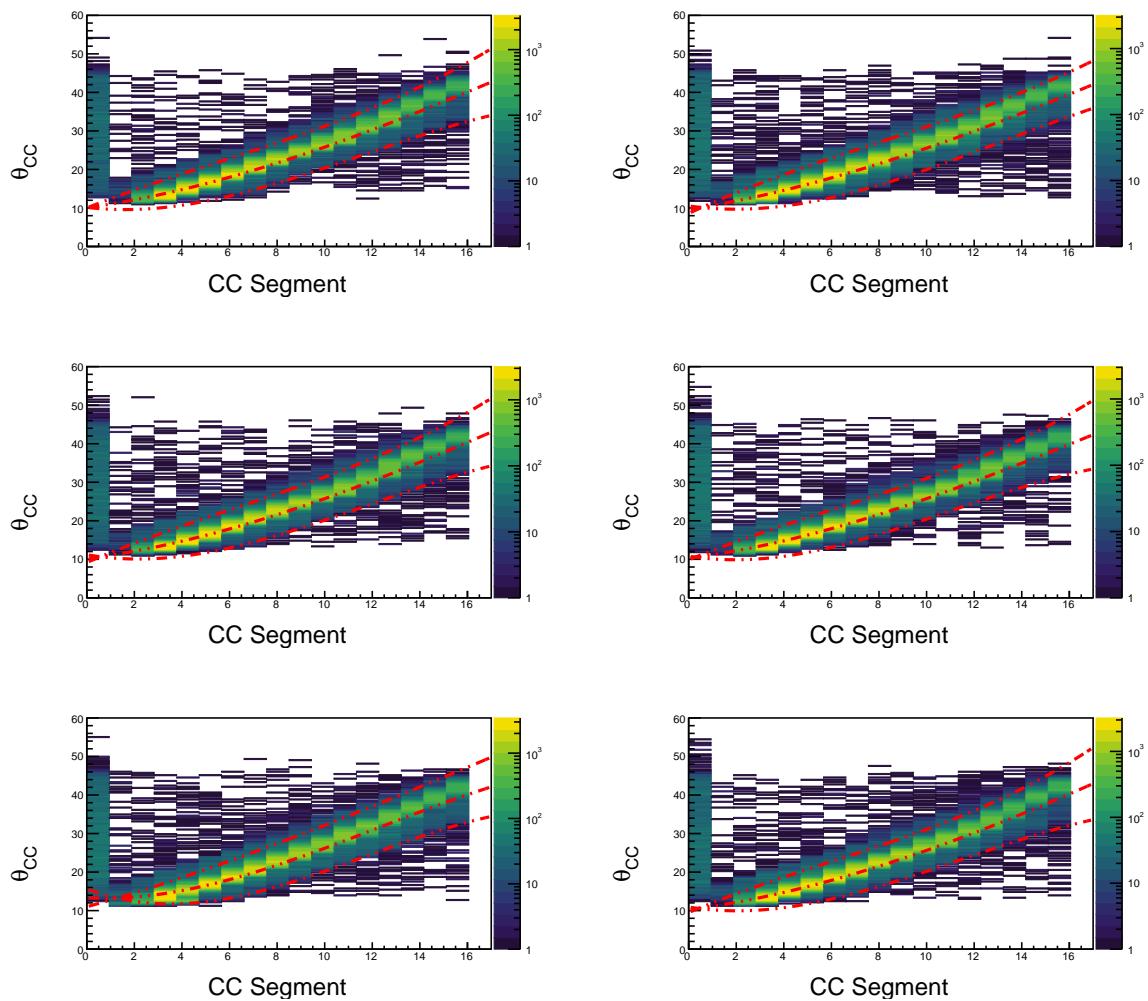


Figure 4.7: Correlation between  $\theta_{CC}$  and the CC segment is shown above, with our selection boundaries overlaid in red.

## 4.3 Hadron Identification

Hadron identification in CLAS is done by correlating particle momentum from the drift chambers with timing information supplied by the time of flight detector. In this analysis some quality assurance cuts are applied preliminarily, but they do not discriminate between different species of particle. The likelihood methodology described in this section is based on the discussion provided by the BES collaboration in [11].

### 4.3.1 Hadron ID Cuts

The cuts used for hadron classification are enumerated below.

- Drift chamber fiducial
- Hadron-electron vertex difference
- Likelihood maximization of  $\beta(p, h)$

#### Drift chamber fiducial

Drift chamber fiducial cuts are applied (only region 1) using the same procedure as described for electrons. The parameters are for negative hadrons are those which are used for the electron. The parameters used for positive tracks are  $h = 10, \theta = 60$ .

#### Hadron-electron vertex difference

The distance between the electron vertex and the hadron candidate track vertex is computed ( $\delta v_z = v_z^e - v_z^+$ ). This distance is constrained to be within the length of the target (5 cm) see figure 4.3.1. This cut would not be applicable to studies where a significantly detached vertex is expected.

#### Likelihood maximization of $\beta(p, h)$

In this section, positive hadrons are used as an example. The same method is applied to the negative hadrons. For each particle species considered, a normalized probability density function  $P(x; p, h)$  is constructed for each input into the likelihood analysis. Here,  $x$  corresponds to the feature being used to categorize different particles (in our case,  $x$  is the  $\beta$  value measured by CLAS time-of-flight),  $p$  is the particle momentum, and  $h$  is the hadron being hypothesized (eg: the possible values for positive hadrons are pion, kaon, proton). In general if one uses a set of  $N$  variables  $x = (x_1, x_2, \dots, x_N)$ , the likelihood for a hypothesis  $h$  is defined below.

$$\mathcal{L}_h = \prod_{i=1}^N P_i(x_i; p, h) \quad (4.7)$$

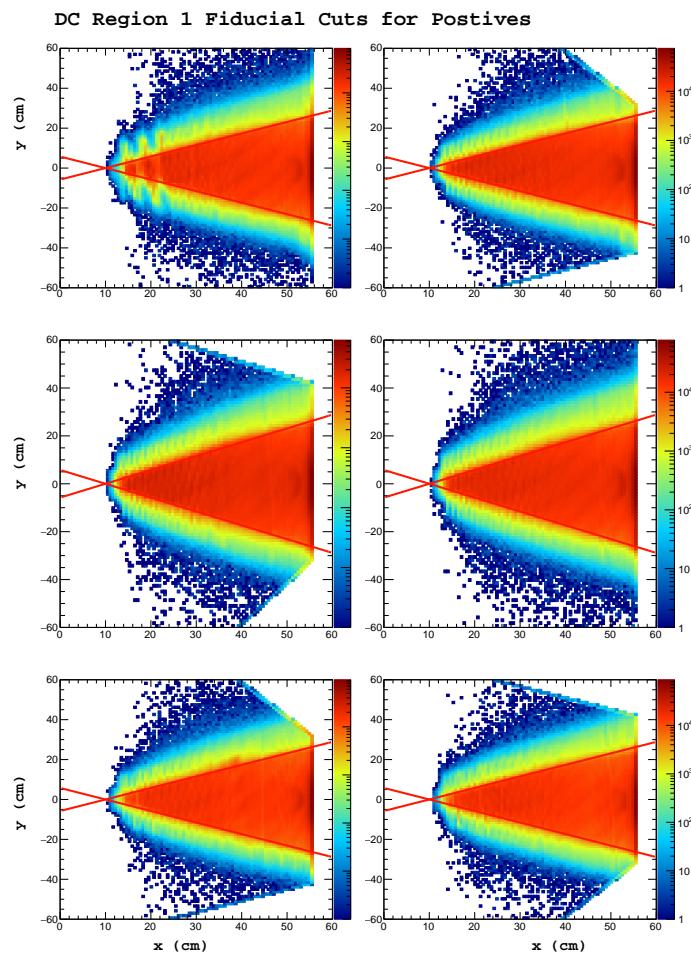


Figure 4.8: Shown above: Positive track hits on the region 1 drift chamber, events falling between the red lines are kept for analysis.

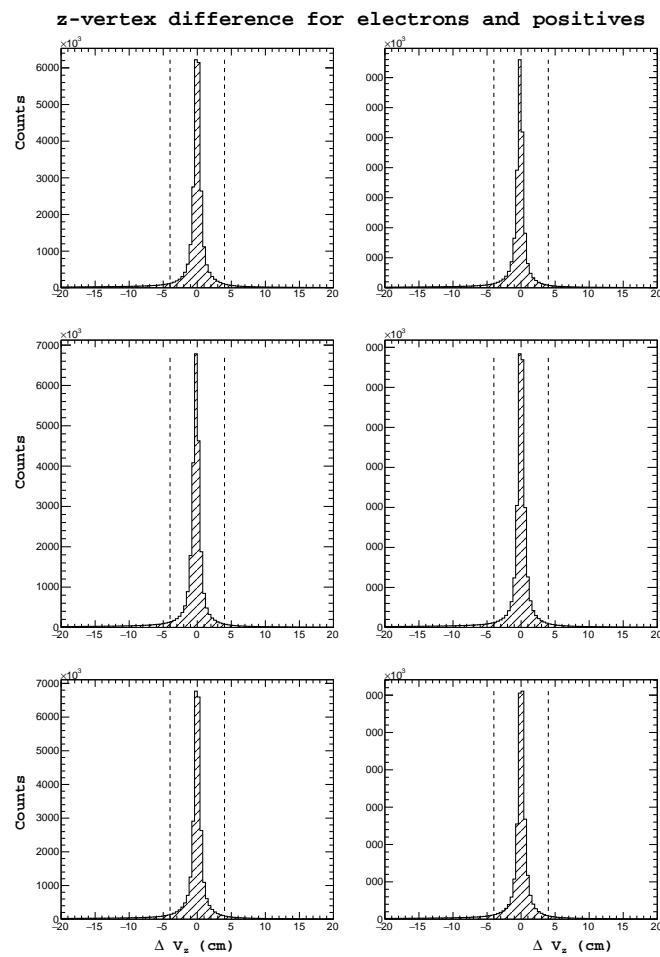


Figure 4.9: Shown above: The difference between the z-vertex position between detected electrons and positive tracks.

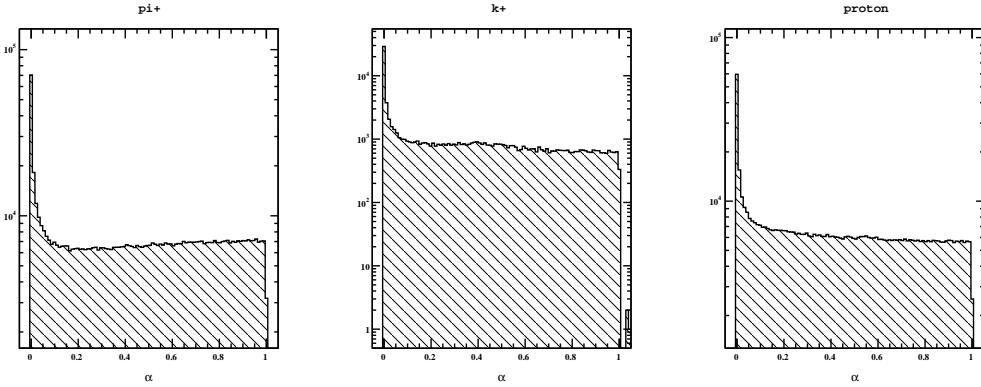


Figure 4.10: Shown above: The distribution of confidence level for all positive tracks after being classified by the likelihood ratio.

In our case, the only random variable we consider is  $\beta$ , and the likelihood is just the PDF. Here, and in many cases where the choice is statistically appropriate, it is possible to use a Gaussian PDF for the variable  $x_i$  (here  $\beta$ ).

$$P(\beta; p, h) = \frac{1}{\sqrt{2\pi}\sigma_\beta(p, h)} \exp \left\{ -\frac{1}{2} \left( \frac{\beta - \mu_\beta(p, h)}{\sigma_\beta(p, h)} \right)^2 \right\} \quad (4.8)$$

The identity is assigned by choosing the particle hypothesis  $h$  which maximizes the likelihood ratio.

$$\frac{\mathcal{L}_h}{\mathcal{L}_\pi + \mathcal{L}_K + \mathcal{L}_p} \quad (4.9)$$

Using this method, every positive track is assigned a particle identification. However, at times the likelihood value is quite small when compared with the maximum likelihood for that species. This is the case for positrons which are classified by this method as positive pions, because they are the closest particle for which a hypothesis has been provided. To avoid these situations, the confidence level  $\alpha$  of each track is calculated and a cut is applied on the minimum confidence. This cut can be easily varied to see how it changes the analysis result.

$$\alpha = 1 - \int_{\mu - \beta_{obs}}^{\mu + \beta_{obs}} P(\beta; p, h) d\beta \quad (4.10)$$

This quantity represents the probability to observe a value of  $\beta$  as far or farther from the mean as  $\beta_{obs}$ . Confidence levels close to zero correspond to tracks which are poorly identified as the class  $h$ . In the case that the PDF is Gaussian, the standard 1, 2, and 3  $\sigma$  cuts on  $\beta$  vs.  $p$  can be understood simply as confidence levels of approximately 0.32 = 1-0.68, 0.05 = 1-0.95, and 0.01 = 1-0.99.

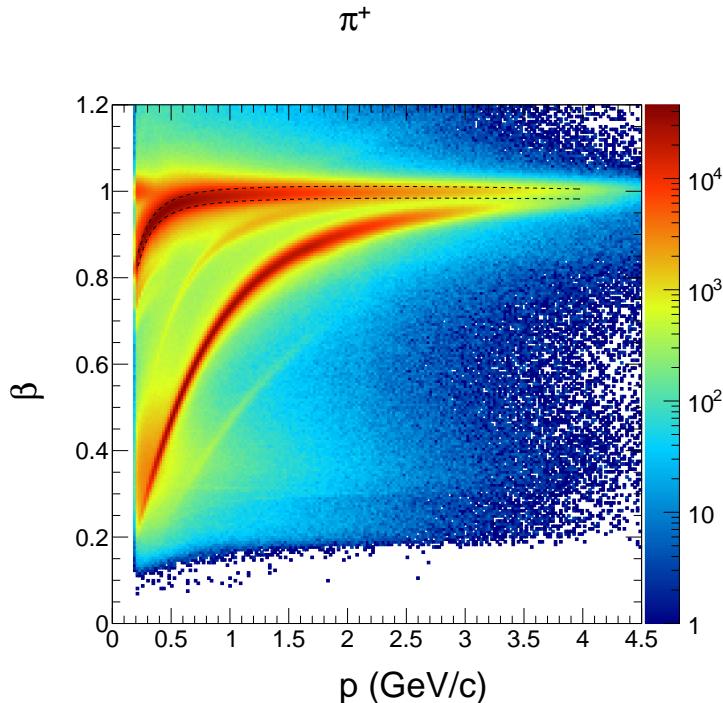


Figure 4.11: Shown above: All positive tracks overlaid with our determination of  $\mu(p) \pm \sigma(p)$  for  $\pi^+$

### Determination of probability density functions for likelihood method

The most important and most difficult part of constructing the likelihood ratio identification is the determination of the mean and standard deviation of the probability density function (which depends on momentum) for the different hypotheses. In the case where exceptionally accurate monte carlo (MC) simulations of the detector are available, one can use the truth information and track matching to construct the  $\beta$  vs.  $p$  2-dimensional histograms, and fit the  $\mu(p)$  and  $\sigma(p)$ . In the absence of high quality MC, analysts typically fit directly the spectrum of  $\beta$  vs.  $p$  and extract the mean and variance. In this work, an enhanced sample of candidates for each of the three positive particles in question is created before doing the fitting. In this way, we hope that our fit better represents the true  $\mu$  and  $\sigma$  for each particle. For fitting of pion and proton resolutions, positive tracks are assumed to be pions and the missing mass of the event is calculated. Then, a cut is placed around the neutron mass. In doing so, two main exclusive reactions are selected. The first is  $ep \rightarrow e\pi^+N$ , and the second is  $ep \rightarrow ep\pi^0$ . In this way most positrons, and positive kaons are removed from the sample prior to fitting. The mean and variance are fit using a third order polynomial in  $p$  (MINUIT  $\chi^2$  minimization is used). Negative pions and kaons are fit directly (as is normally done).

The parametrization used for the mean  $\mu(p, h)$  and resolutions  $\sigma(p, h)$  are shown below.

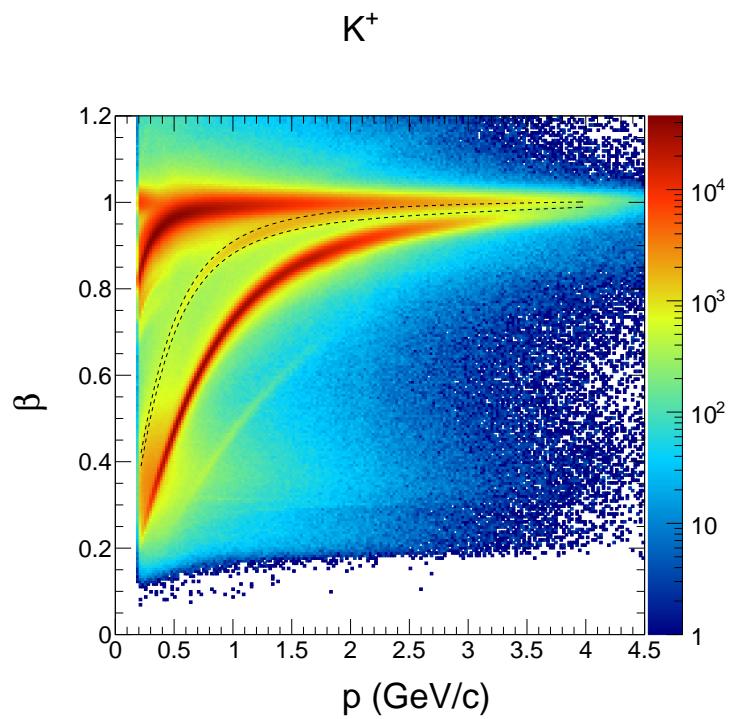


Figure 4.12: Shown above: All positive tracks overlaid with our determination of  $\mu(p) \pm \sigma(p)$  for  $K^+$

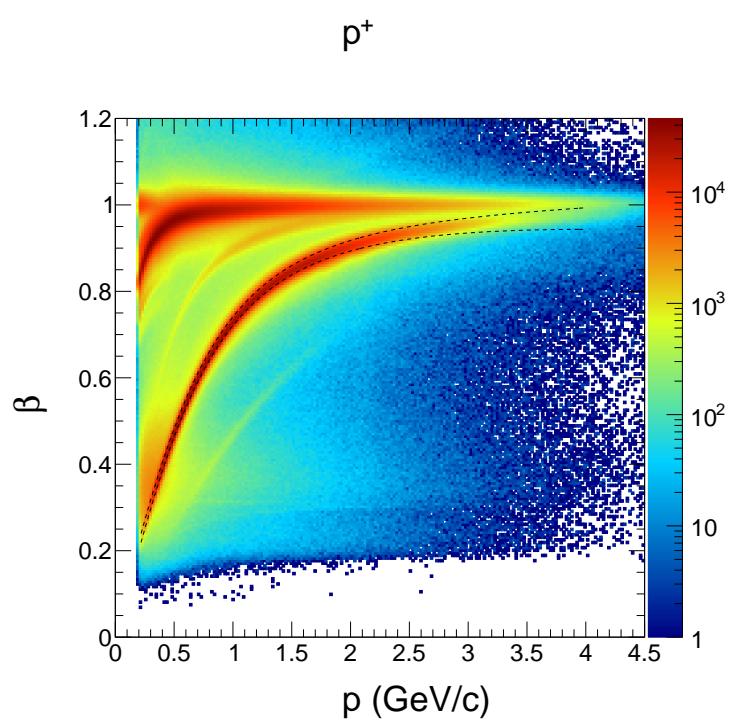


Figure 4.13: Shown above: All positive tracks overlaid with our determination of  $\mu(p) \pm \sigma(p)$  for  $p^+$

$$\mu(p, h) = \mu_{theory} + \Delta\mu \quad (4.11)$$

$$\mu_{theory} = \frac{1}{\sqrt{1 + (m_h/p)^2}} \quad (4.12)$$

$$\Delta\mu = \mu_0 + \mu_1 p + \mu_2 p^2 \quad (4.13)$$

$$\sigma(p, h) = \sigma_0 + \sigma_1 p + \sigma_2 p^2 \quad (4.14)$$

The values are displayed in the table below.

Hadron	Parameter	Sector 1	Sector 2	Sector 3	Sector 4	Sector 5	Sector 6
$K^+$	$\mu_2$	0.00111554	-8.97687e-05	4.78796e-05	0.000376425	-0.00204856	0.000652209
$K^+$	$\mu_1$	-0.00468038	6.19414e-05	-0.00081741	-0.00107931	0.00629181	-0.00264143
$K^+$	$\mu_0$	0.00361012	0.00134921	0.00299674	0.00220194	0.000117821	0.00162582
$K^+$	$\sigma_2$	-0.000331838	-0.00105807	-0.000712404	-0.000573934	-0.000259289	0.000508389
$K^+$	$\sigma_1$	-0.00105857	0.00236686	0.000509169	0.000163467	-0.00233617	-0.00461598
$K^+$	$\sigma_0$	0.0154964	0.0117702	0.0140748	0.0143761	0.0184055	0.0180945
$\pi^+$	$\mu_2$	-0.000962041	-0.000300602	-0.000306326	-3.2245e-05	-0.00226511	-0.000330818
$\pi^+$	$\mu_1$	0.00296349	0.0016512	0.0021962	0.00176045	0.00750862	0.00126443
$\pi^+$	$\mu_0$	-0.00225794	-0.00047045	0.000370406	0.000435526	-0.000449409	-0.00131045
$\pi^+$	$\sigma_2$	-0.000127659	0.000691895	-0.000289961	0.000315041	-0.000936521	-0.000131269
$\pi^+$	$\sigma_1$	-0.000489092	-0.0033948	0.00196853	-0.00197841	0.00212778	-0.000339411
$\pi^+$	$\sigma_0$	0.0155195	0.0167998	0.0124066	0.0157476	0.0145571	0.0141728
$p^+$	$\mu_2$	-0.00039358	-0.000701003	-0.000347651	0.0004854	-0.00121666	0.000563786
$p^+$	$\mu_1$	-0.000295423	0.00170899	0.000794901	-0.000744446	0.00376887	-0.00353545
$p^+$	$\mu_0$	0.00227353	0.00231676	0.00364672	0.00276859	0.00128827	0.00439605
$p^+$	$\sigma_2$	0.001429	0.00144256	0.00124456	0.00190709	0.00141039	0.0011516
$p^+$	$\sigma_1$	-0.0021472	-0.00262226	-0.00196308	-0.00385218	-0.00186708	-0.00186749
$p^+$	$\sigma_0$	0.0107541	0.0109091	0.0104381	0.0115449	0.0109969	0.0107759
$\pi^-$	$\mu_2$	3.28823666e-04	-1.30673670e-05	-2.32502052e-04	-9.75619848e-04	-5.89834444e-04	5.27496718e-04
$\pi^-$	$\mu_1$	-3.94924663e-03	-2.66028661e-03	-1.28565631e-03	9.09410075e-04	-2.01610684e-03	-4.42276918e-03
$\pi^-$	$\mu_0$	9.48011169e-04	1.55078786e-03	1.43431985e-03	1.35056935e-03	4.59833580e-03	2.30751866e-03
$\pi^-$	$\sigma_2$	4.37635504e-04	4.38306224e-04	5.32057510e-04	3.36999845e-04	7.74135462e-04	1.36515196e-04
$\pi^-$	$\sigma_1$	-3.28011836e-03	-3.28456104e-03	-3.82847286e-03	-3.11749323e-03	-4.63110728e-03	-2.21229710e-03
$\pi^-$	$\sigma_0$	1.63296567e-02	1.62229164e-02	1.59769911e-02	1.58803427e-02	1.74670064e-02	1.51753145e-02
$K^-$	$\mu_2$	-2.72020947e-03	-5.21081786e-03	-2.13868763e-02	-4.45600034e-03	-7.60703841e-03	-5.27074813e-03
$K^-$	$\mu_1$	1.78610401e-02	2.30787460e-02	9.49357818e-02	1.95764575e-02	3.63245785e-02	2.92417500e-02
$K^-$	$\mu_0$	-2.26190100e-02	-2.22562379e-02	-1.02704771e-01	-2.25931014e-02	-5.10484618e-02	-3.19918187e-02
$K^-$	$\sigma_2$	1.76905114e-02	1.62989708e-02	3.60928130e-02	1.51270521e-02	1.91308107e-02	2.38470033e-02
$K^-$	$\sigma_1$	-7.74901862e-02	-7.33041628e-02	-1.57454534e-01	-7.26870393e-02	-9.23654247e-02	-1.02397836e-01
$K^-$	$\sigma_0$	1.07082820e-01	1.00573410e-01	1.93148260e-01	1.00993689e-01	1.26963814e-01	1.30057621e-01

Table 4.4: Values used to calculate the mean and resolutions for hadron likelihood based identification.



# Chapter 5

# Beam Spin Asymmetry Analysis

## 5.1 Introduction

Measurement of the beam spin asymmetry is carried out for the positively charged k-meson. As discussed in the introduction, the beam spin asymmetry theoretically depends on  $F_{UU,L}$ ,  $F_{UU,T}$ ,  $F_{UU}^{\cos\phi}$ ,  $F_{UU}^{\cos 2\phi}$ , and  $F_{LU}^{\sin\phi}$ . By dividing the electron-kaon events into several bins of SIDIS kinematic variables, beam spin asymmetry measurements can be taken at different average values of the kinematic variables. Finally, the structure function ratios  $A_{LU}^{\sin\phi}$ ,  $A_{UU}^{\cos\phi}$ , and  $A_{UU}^{\cos 2\phi}$  can be extracted from each bin. In this chapter a discussion is provided of SIDIS event selection, the binning used in this analysis, measurement values with associated systematic uncertainties, and the extraction of structure function ratios using the  $\phi$  dependence in each kinematic bin.

## 5.2 Event Selection and Binning

### Event Selection

After particle identification, those events which have a trigger electron and a positive kaon are kept for analysis. Events are discarded that do not have  $W > 2$  and  $Q^2 > 1$ , because they are not considered part of the deeply inelastic region. Additionally, to avoid exclusive resonances in the  $ep \rightarrow eK^+X$  spectrum, a minimum value is imposed on the missing mass of the final state  $M_X$  ( $ep \rightarrow eK^+X$ ). Here, we use  $M_X(ep \rightarrow eK^+X) > 1.25$ . Finally, a cut is applied to exclude low values of  $z_h$  to constrain our kinematics to the current factorization region where TMD factorization has been demonstrated at leading order. It is additionally required that  $z_h < 0.75$  to avoid exclusive events. This restriction on  $z_h$  is not applied to the  $z_h$  axis, where we measure across the entire experimentally observed range.

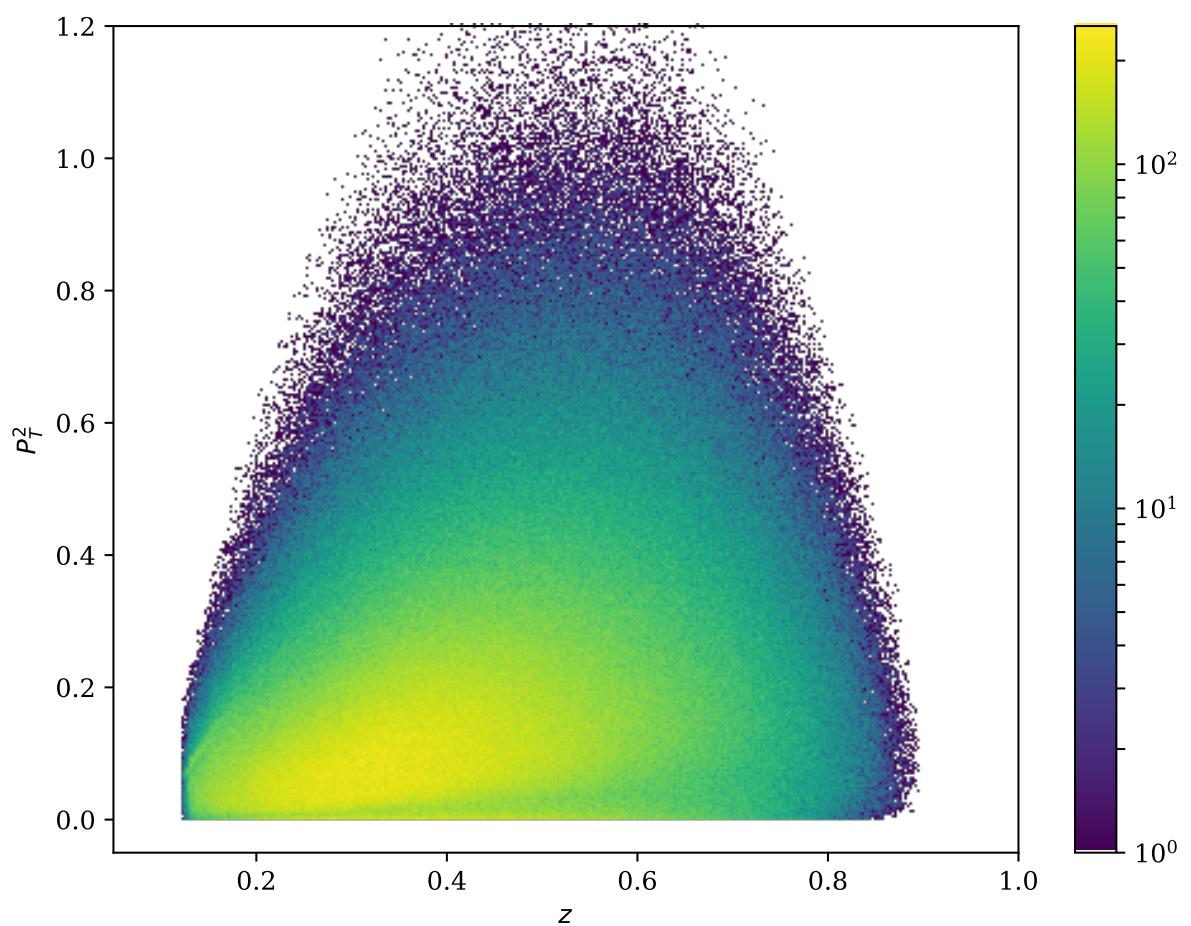


Figure 5.1: Correlation between  $z_h$  and  $P_T^2$  for each event in our analysis sample.

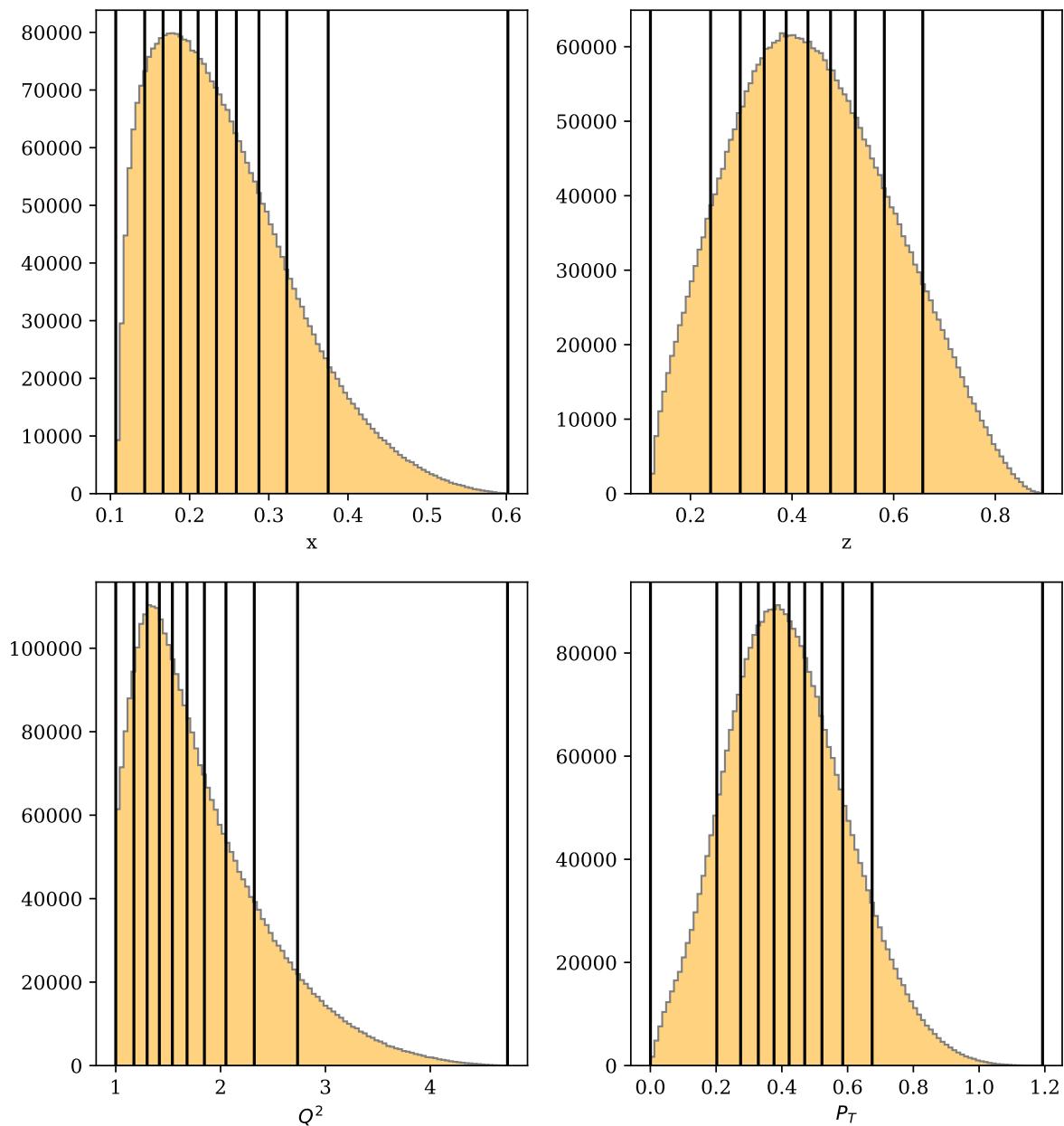


Figure 5.2: The binning used for each of the kinematic axes.

## Binning

The beam spin asymmetry measurement is performed for the kinematic variables  $x$ ,  $Q^2$ ,  $z_h$ , and  $P_T$ . For each variable 10 bins are chosen, as well as 12 bins in  $\phi$  for a total of 120 analysis bins.

Bins were chosen using a simple method to ensure equal statistics in each bin. The procedure is described using the axis  $x$  as an example. First, all events are sorted by their  $x$  value from smallest to largest. Then, the smallest and largest values are recorded, which are  $x_1$  and  $x_N$  if there are  $N$  events in the sample. Next, the target number of bins  $M$  is chosen (this choice depends on each analysis). Finally, the limits of each bin can be chosen by calculating the number of events per bin  $N/M$  and then using the value of  $x$  which corresponds to multiples of  $N/M$  in the sample.

$$\vec{b} = (x_1, x_{N/M}, x_{2N/M}, \dots, x_N) \quad (5.1)$$

Here, the symbol  $\vec{b}$  denotes a vector of  $(M+1)$   $x$  values which represent bin limits. The binning in  $\phi$  is chosen to be regularly spaced between -180 and 180 degrees.

## 5.3 $\phi_h$ Distributions

### Measured Asymmetry Values

In each bin  $i$  the beam spin asymmetry (here  $A_i$ ) is calculated according to,

$$A_i = \frac{1}{P_e} \frac{n_+^i - n_-^i}{n_+^i + n_-^i} \quad (5.2)$$

where  $P_e$  is the average beam polarization over the dataset (74.9%). The symbols  $n_\pm^i$  refer to the number of events counted in bin  $i$  with helicity  $\pm$ .

### Statistical Uncertainties

The uncertainty on the measured value of  $A_i$  can be attributed to statistical uncertainty on the counts  $n_i^\pm$ , and the uncertainty associated with the measurement of  $P_e$ . The statistical uncertainty reported on the measurement includes the contribution from counts, but not from the uncertainty in  $P_e$  which is included in the systematic errors. In general, the uncertainty in a measured observable  $\mathcal{O}$  depends on the uncertainty of the parameters (here denoted by  $\vec{\theta}$ ) used to construct it  $\vec{\theta}$  in the following way (see appendix for derivation).

$$\sigma_{\mathcal{O}}^2 = \sum_{i=1}^N \sum_{j=1}^N \frac{\partial \mathcal{O}}{\partial \theta_i} \frac{\partial \mathcal{O}}{\partial \theta_j} \rho_{ij} \sigma_i \sigma_j \quad (5.3)$$

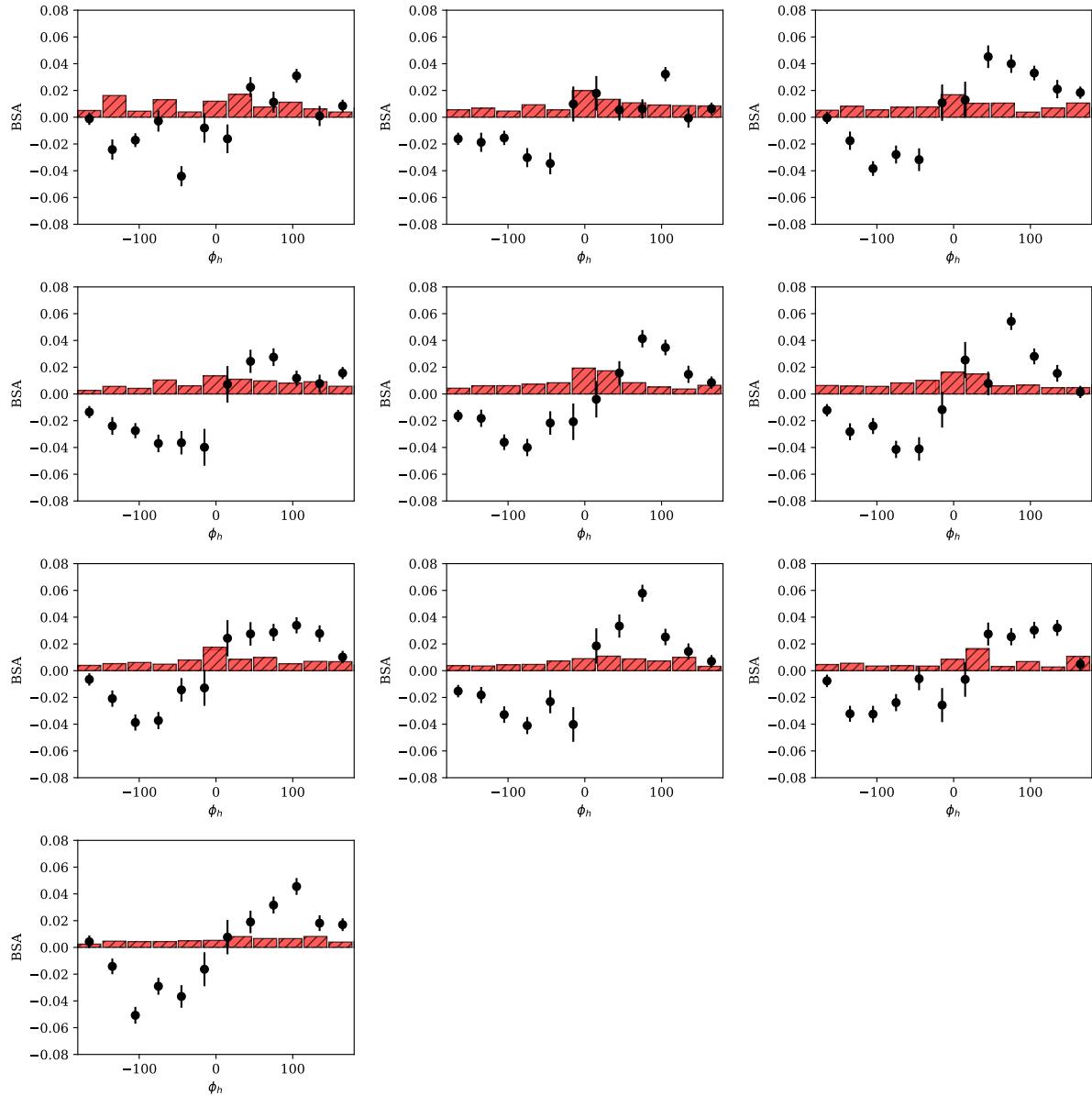


Figure 5.3: The  $\phi_h$  dependence is shown for each bin of  $x$ , increasing in value from the top left to the bottom right. The statistical uncertainty is shown as black error bars on each point. The total systematic uncertainty is shown as a red bar centered at zero.

For the beam spin asymmetry in the  $i^{th}$  bin  $A_i$  one finds that without correlations ( $\rho_{ij} = \delta_{ij}$ ) the error propagation proceeds as shown below.

$$\sigma_A^2 = \frac{A^2}{P_e^2} \sigma_{P_e}^2 + \frac{4(n_-^2 \sigma_+^2 + n_+^2 \sigma_-^2)}{P_e^2 (n_+ + n_-)^4} \quad (5.4)$$

The first term which is the contribution from the variance in the measurements of beam polarization will be included as a systematic error. The second term is used as the statistical error bars shown through the analysis. The counts  $n_{\pm}^i$  for the  $i^{th}$  bin are assumed to be Poisson in nature, and therefore have a variance equal to the expected number of counts  $\sigma_{\pm}^2 = n_{\pm}^i$ . With this expression for the statistical uncertainty on the counts, and dropping the beam polarization term for now, the expression becomes:

$$\sigma_A^2 = \frac{4n_+ n_-}{P_e^2 (n_+ + n_-)^3} \quad (5.5)$$

## Systematic Uncertainties

Systematic effects are shifts or biases in the measured result of some observable as a result of the procedure used in the measurement. Systematic effects can typically be identified and corrected for, or removed all together from the measurement. In the cases where an effect cannot be completely removed, the degree to which the correction for the effect is uncertain is included in the result of the measurement as a systematic uncertainty [12].

Sources of systematic effects can include background events from different processes which enter the sample, calibrations of different detector systems, misalignment in detector geometry, and biases in selection criteria. Each of the systematic sources mentioned here has at least one associated procedure for correcting it's effect on the analysis. As an example consider momentum corrections in CLAS. These corrections are performed to remove the effect of slight misalignment in detector geometry from what is in reconstruction, as well as slight differences between the true magnetic field and the field map used in reconstruction. These physical effects introduce a systematic effect, the particle 4-momenta reconstructed are shifted away from the true values. Standard reactions (elastic scattering) can be used to develop corrections for the 4-momenta of particles, and these corrections typically depend on a set of parameters  $\vec{\theta}$ , which have an associated parameter uncertainty described by a covariance matrix  $V_{ij}$ . It is these parameter uncertainties that propagate through to the final observables, and the assignment of the magnitude of such effects is then what is referred to as systematic uncertainty.

Systematic uncertainties are included using the standard equation for error propagation. In some cases it is possible to analytically find the derivatives needed to calculate the dependence of the ob-

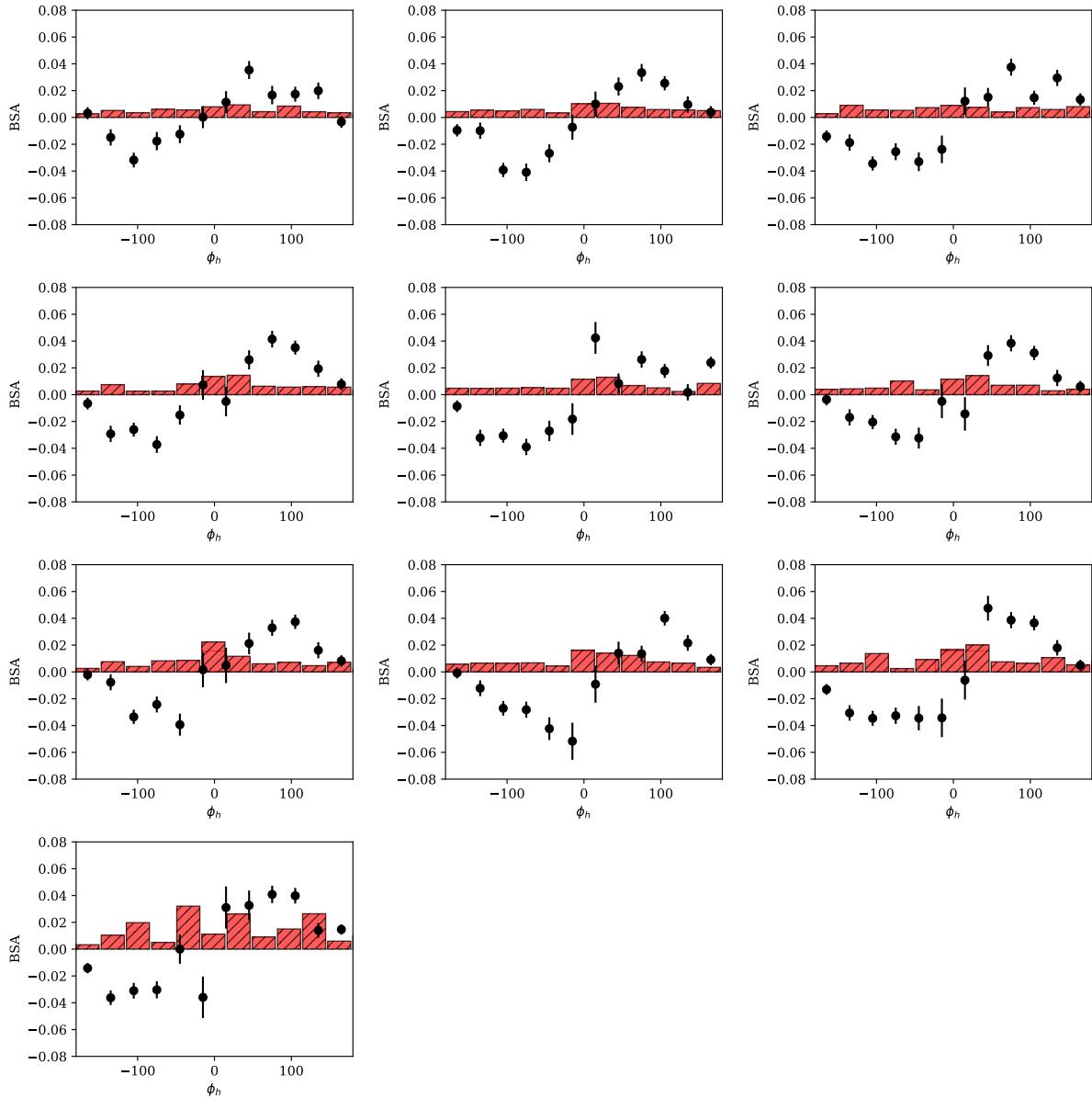


Figure 5.4: The  $\phi_h$  dependence is shown for each bin of  $z_h$ , increasing in value from the top left to the bottom right. The statistical uncertainty is shown as black error bars on each point. The total systematic uncertainty is shown as a red bar centered at zero.

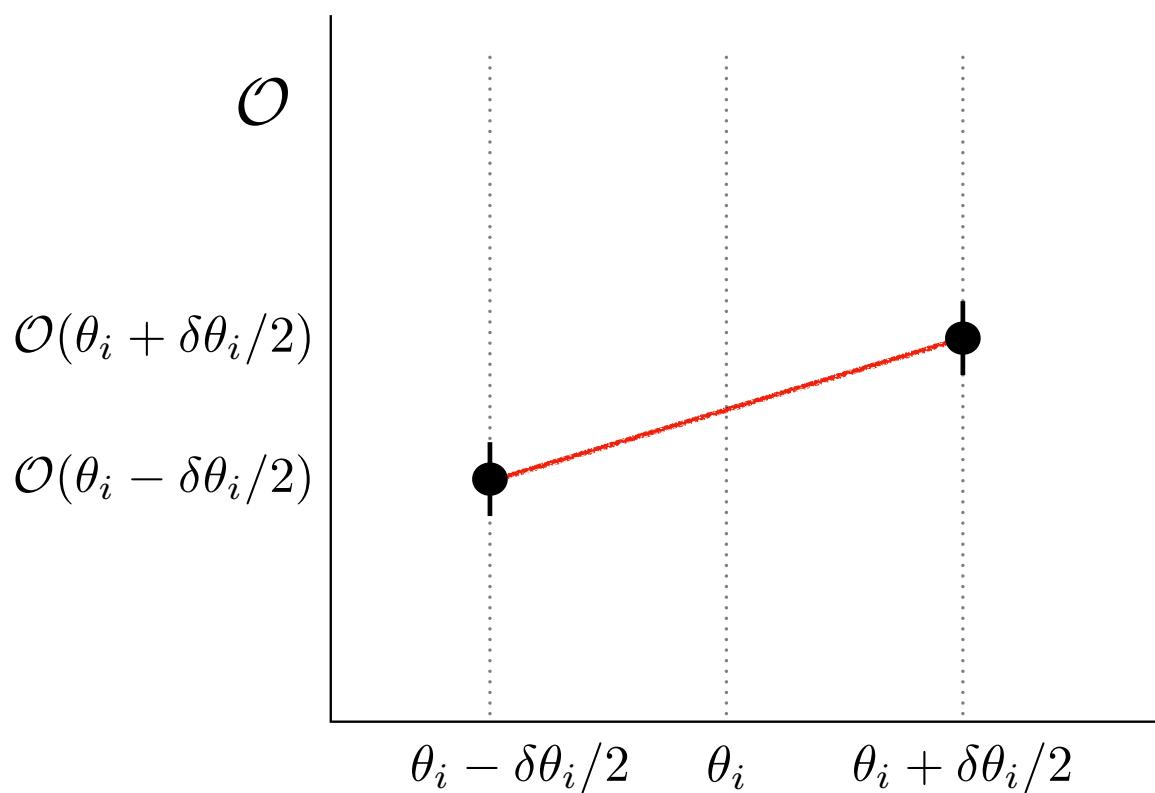


Figure 5.5: The analysis is run for variations in the input parameters  $\theta_i$  to calculate the dependence of the result  $\mathcal{O}$  on each parameter, as described in this section.

servable on a source of systematic uncertainty. This is the case for effect of the variance of the beam polarization on the beam spin asymmetry observable. However in many cases, it is not possible to analytically calculate the effect of an analysis parameter  $\theta_i$  on the observable  $\mathcal{O}$ . Since the observable is usually calculated using some computational chain which starts with the input parameters  $\vec{\theta}$ , it is possible to find the dependence of the observable  $\mathcal{O}$  on the inputs numerically.

$$\frac{\partial \mathcal{O}}{\partial \theta_i} \approx \frac{\mathcal{O}(\theta_i + \sigma_{\theta_i}/2) - \mathcal{O}(\theta_i - \sigma_{\theta_i}/2)}{\sigma_{\theta_i}} \quad (5.6)$$

After inserting the above into equation 5.3 one finds,

$$\sigma_{\mathcal{O}}^2 = \sum_{i=1}^n \sum_{j=1}^n \rho_{ij} (\mathcal{O}(\theta_i + \sigma_{\theta_i}/2) - \mathcal{O}(\theta_i - \sigma_{\theta_i}/2)) (\mathcal{O}(\theta_j + \sigma_{\theta_j}/2) - \mathcal{O}(\theta_j - \sigma_{\theta_j}/2)) \quad (5.7)$$

where  $\rho_{ij}$  is the correlation  $V_{ij}/\sigma_i \sigma_j$ . In most cases, these correlations are assumed to be zero. In some cases, when the parameters  $\theta_i, \theta_j$  come from a fit one may have a correlation provided by the covariance matrix and it should be used. In the case where correlations are assumed to be zero, the total systematic uncertainty is simply the quadrature sum of the shift in the observable within the uncertainty window on each parameter.

$$\sigma_{\mathcal{O}}^2 = \sum_{i=1}^n \left[ \mathcal{O}(\theta_i + \sigma_{\theta_i}/2) - \mathcal{O}(\theta_i - \sigma_{\theta_i}/2) \right]^2 \quad (5.8)$$

Another approach exists that takes into account possible correlations between the analysis parameters  $\theta_i$ . The approach consists of generating  $M$  sets of parameters  $\vec{\theta}$  and obtaining  $M$  results for the observable  $\mathcal{O}$ . The results are then interpreted statistically and the observable value and total systematic error are reported as the mean and standard deviation of the results. This Monte Carlo based technique has the advantage that it is conceptually simple, and allows the capture of correlations. However, it can be computationally intensive to perform the analysis hundreds of times. Additionally, only using this technique without first testing the individual size of each systematic source removes the possibility of correcting for or removing biases.

$$\langle \mathcal{O} \rangle = \frac{1}{M} \sum_{i=1}^M \mathcal{O}_i \quad (5.9)$$

$$\sigma_{\mathcal{O}}^2 = \frac{1}{M-1} \sum_{i=1}^M (\mathcal{O}_i - \langle \mathcal{O} \rangle)^2 \quad (5.10)$$

### Sources of Systematic Uncertainty

Table 5.1 below summarizes the sources of systematic uncertainty considered in this analysis.

Except for the beam polarization and the momentum of the kaon track, all parameters listed in the

Source	Variation	Magnitude
Beam polarization	0.024	0.000672
DC Region 1 Fid.	1 (cm)	0.001344
DC Region 3 Fid.	3 (cm)	0.001821
EC-W	12 (cm)	0.000948
EC-V	12 (cm)	0.000797
EC-U	12 (cm)	0.002487
Kaon Confidence ( $\alpha$ )	0.01-0.07	0.001827
$\theta_{cc}$ Matching	$\sigma$	0.001152
EC Energy Deposition	0.01 (GeV)	0.001644
$p_{K^+}$	$2.5-E_{beam}$	0.002360
EC Sampling Fraction	$0.5\sigma$	0.001240
Z-Vertex	0.5 (cm)	0.002581
Statistical	-	0.007494
MC Estimate	-	0.002917

Table 5.1: Different sources of systematic effect considered in this analysis. The magnitude of the effect is shown here averaged over all bins. The units of the shift are just the same units of the asymmetry value.

table are treated using the formalism outlined above. The beam polarization uncertainty quoted at 2.4% contains contributions from the standard deviation of the Moller polarimetry measurements (0.2%), residual target polarization effects (1.4%), and atomic motion/finite acceptance corrections (0.8%).

Because of the inability to distinguish kaon, pion, and proton tracks at higher momentum, the maximum kaon track momentum is varied between 5.5 (no maximum) and 2.5 GeV. The difference between these results is quoted as a systematic uncertainty and added in quadrature with the other sources. This source of systematic uncertainty has a larger effect on the  $z_h$ -axis, because  $z_{max}$  is limited by limiting  $p_{max}$ . While for the large  $z_h$ -bins this contribution is dominant, its size is comparable with other systematic sources throughout the remainder of the bins.

## 5.4 Extraction of Modulations

The motivation to measure the beam spin asymmetry in several kinematic bins as well as bins of  $\phi_h$  is to perform an estimate of the value of structure functions at the kinematic points (more precisely the average value of the structure functions over the range of values included in a point). To do this, the authors perform parameter estimation on the  $\phi_h$  distributions taking as a model the theoretical dependence of the beam spin asymmetry on  $\phi_h$ .

$$f(\phi_h, \vec{a}) = \frac{a_0 \sin \phi_h}{1 + a_1 \cos \phi_h + a_2 \cos(2\phi_h)} \quad (5.11)$$

The parameters  $\vec{a}$  are the structure function ratios to be extracted. The simplest way to extract these parameters is to use  $\chi^2$  minimization implemented in a standard fitting package. In these approaches,  $\chi^2$  is defined as the square difference between the observed data values and those predicted by the

model, normalized by the error. If the fluctuation between the data and theory predictions is on the order of the error, the  $\chi^2$  is simply on the order of the number of data points. The parameters  $\vec{a}$  which best describe the data are those which make the  $\chi^2(\vec{a})$  assume its minimum value. This minimization is done in practice with gradient descent or quasi-Newton's method based algorithms like those provided in `Minuit` or `scipy.optimize.minimize`, the details of such algorithms will not be discussed here. It is sufficient to say that these minimization methods produce the parameters  $\vec{a}$ , and an estimate of the covariance matrix  $V$ . The parameters and their errors become the extracted value and uncertainty of the structure function ratio in each bin.

Unfortunately, applying the standard single-fit procedure described above does not always produce stable results. In some cases, the resulting parameter sets are reasonable, in other cases however the parameters in the denominator become nonphysically large and oppose each other. This effect has motivated previous analysts to search for other means of extracting the dominant  $\sin \phi_h$  behavior from the distributions. One common technique is to assume that  $a_1$  and  $a_2$  of above are small compared to 1. The analyst can then fit the  $\phi_h$  distribution with just one linear parameter  $a_0$ . This produces a stable result, but has the disadvantage that one needs to introduce a systematic uncertainty associated with the difference observed between using the full model (with a restricted range for the parameters in the denominator) and the results obtained using the single parameter model. Additionally, the structure function decomposition of the SIDIS cross section relies on theoretically solid ground, therefore it should be used in its full form. If the data contain little information regarding the structure function ratios in the denominator, the authors believe it more valuable to demonstrate this by extracting those parameters with (large) errors, rather than ignore their contribution. In order to accomplish this, the method of replicas (or parametric bootstrapping) is used to perform the parameter extraction. The replica method consists of generating  $N_{rep}$  pseudo-data  $\phi_h$  distributions which have a normal distribution located at the observed value, and with a variance equal to the statistical errors on the associated data point.

$$\vec{A}_{rep} = \mathcal{N}(\vec{A}, \sigma_A) \quad (5.12)$$

Where here  $\vec{A}$  is a vector of length  $n_{phi}$  bins, representing the measured beam spin asymmetry for each value of  $\phi_h$  in a given kinematic bin. Each of these distributions is fit with the full model, and the resulting parameter values are saved. The final reported value for each fit parameter, as well as its uncertainty can be reported as the mean, and standard deviation of the fit results. This procedure which is similar to bootstrapping, can be seen as an attempt to fit the underlying distribution that generated the data while avoiding the statistical noise. This technique has been discussed in [13].

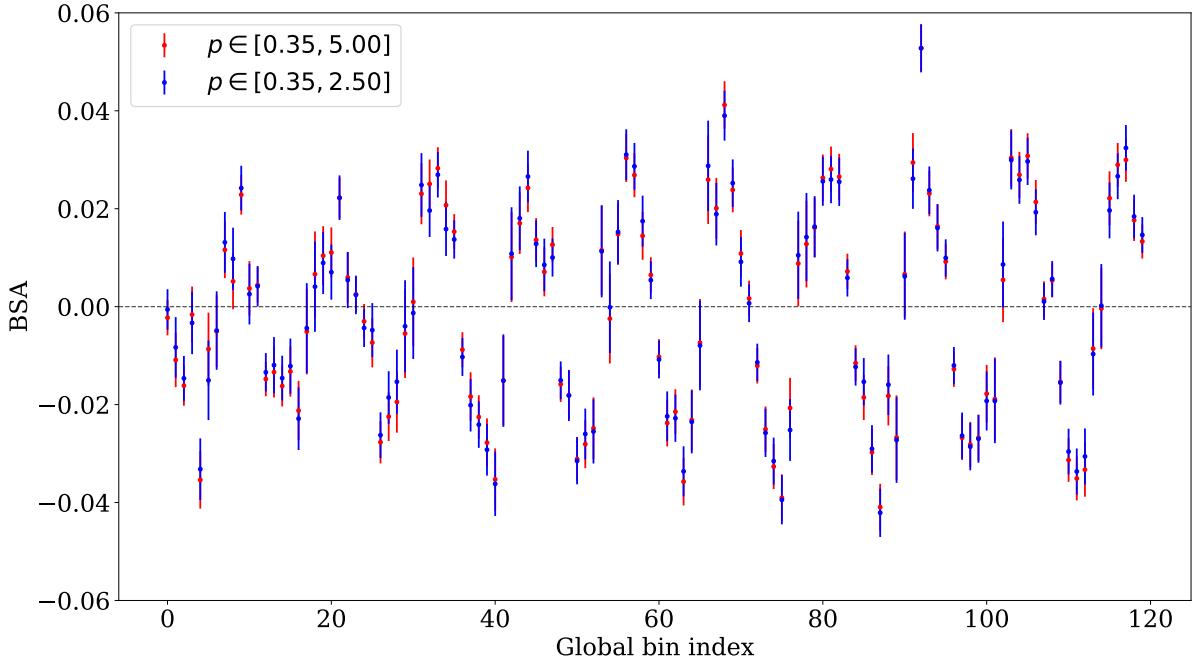


Figure 5.6: The BSA for each bin of  $x, \phi_h$  (all plotted together on the x-axis) is compared with and without a maximum momentum (no tracks exceed 5 GeV) for the kaon track. The global bin coordinate on the x-axis is  $i + n_\phi * j$  where  $i$  is the  $\phi_h$  bin,  $j$  is the  $x$  bin, and  $n_\phi$  is the number of  $\phi$  bins (12). Both  $i$  and  $j$  start at 0.

$$\langle a_j \rangle = \sum_{i=1}^{N_{rep}} a_j^{(i)} \quad (5.13)$$

$$\sigma_{a_j}^2 = \frac{1}{N_{rep} - 1} \sum_{i=1}^{N_{rep}} (a_j^{(i)} - \langle a_j \rangle)^2 \quad (5.14)$$

## Results

As is the case for positive pions, the observed structure function ratio  $A_{LU}^{\sin \phi}$  is positive for all kinematic points that were measured. In general, this extraction reveals that the  $\sin \phi_h$  moment has a magnitude around 3% for most kinematic points, and depends weakly on the kinematic variables used in this analysis. The relative asymmetry value to total error ratio is around 1.5 for most measured points.

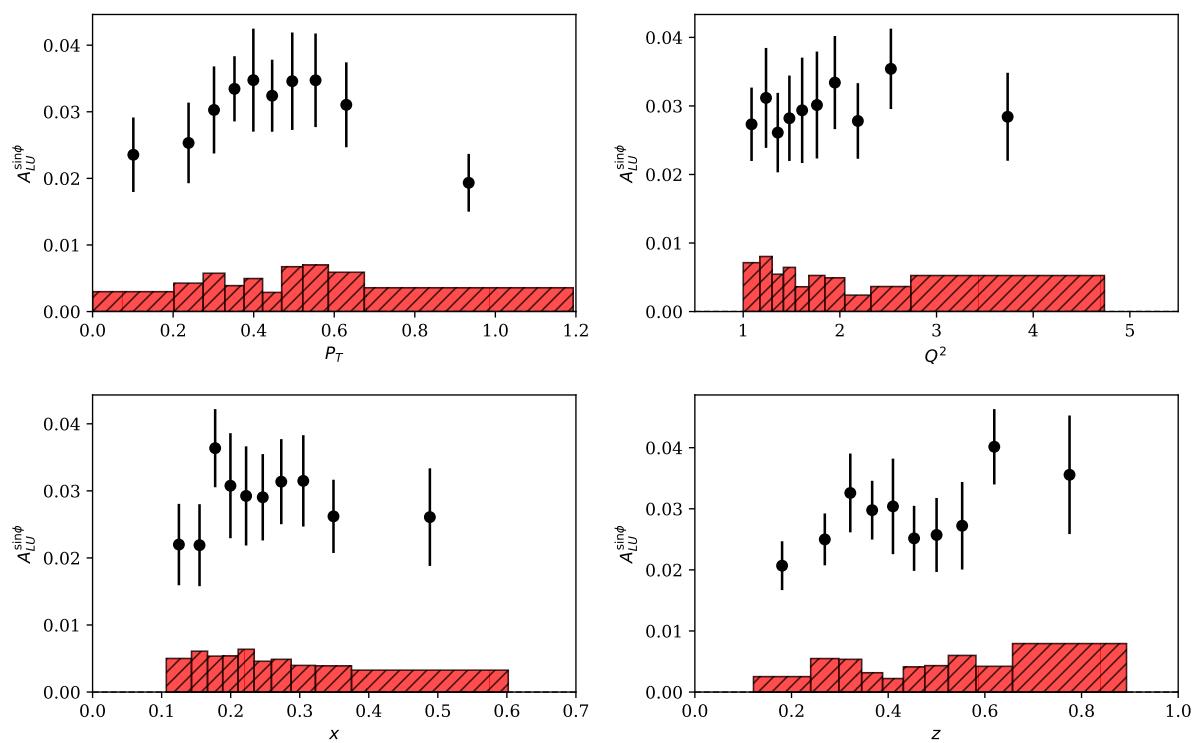


Figure 5.7: Our extraction of  $A_{LU}^{\sin\phi}$  for the kinematic bins described above. The black error bars represent uncertainty in the extraction of the parameter value. Red error bars are systematic uncertainties.



# Chapter 6

## Inclusive Cross Section

Inclusive electron scattering is the process  $ep \rightarrow eX$ , where only the final state electron is detected and the rest of the event is not (anything apart from the electron that is detected is not analyzed). As a function of  $W$  (the invariant mass of the final state ( $\gamma^* + p$ ) system) the region below 2 GeV contains resonances and is often referred to as the resonance region. Resonance structures are difficult to detect higher than about 2 GeV, and this region is typically called the *deeply inelastic* region. While the deeply inelastic region is used extensively for measurements in nuclear/particle physics, the goal of luminosity verification is more easily achieved in the resonance region. This fact is due principally to the excess of Bethe-Heitler events which collect in the  $2 < W < 3$  region for  $E_{beam} = 5.498$  (such events are difficult to remove when detecting only the final state electron).

It is true that the elastic scattering cross section for  $E_{beam} = 5.498$  is small compared to the inclusive cross section, but a significant number of electrons radiate photons before colliding with the target with  $E_i < E_{beam}$ . These lower energy electrons then have a significantly higher probability to scatter elastically and for our beam energy collect in the region of  $2 < W < 3$ .

In this chapter, the detailed procedure for selecting inclusive events and calculating the inclusive cross section in the resonance region is discussed.

### 6.0.1 Event Selection and Binning

A simple choice of 10 bins in  $Q^2$  and 35 bins in  $W$  is used. This choice is mainly driven by the desire to keep bin migration effects small. Events are generated and reconstructed in some bins  $R^{(j)}$  and  $G^{(i)}$  respectively. Due to finite detector resolution, it is not always the case that  $i = j$ . This effect is known as bin migration, and negatively impacts the acceptance calculation.

The only kinematic restriction that is imposed is applied to the *inelasticity*  $y = 1 - E'/E < 0.7$ . This restriction is applied because events with large- $y$  have a significantly higher probability to be Bethe-Heitler events. This cut is equivalent to enforcing a minimum energy for the scattered electron.

Variable	N	Min.	Max	Width
$W$	35	1.1	2.1	0.286
$Q^2$	10	1.7	4.2	0.25

Table 6.1: Summary of  $W$  and  $Q^2$  binning used for the inclusive cross section.

$$E_{min} = E_{beam}(1 - y_{max}) \approx 1.6 \text{ GeV} \quad (6.1)$$

### 6.0.2 Simulation

All processes that CLAS measures are observed through the combination of signals from several sub-detectors. During analysis all sub-systems are calibrated accurately, but such a complicated system still often produces distributions that do not look like the true physical distribution. This discrepancy arises from the combination of several effects.

1. Holes, barriers, obstructions, shadows of other detectors, and any other physical effects that prevent events from being measured in some range of  $\theta, \phi$  are known as geometrical acceptance effects. An important geometrical acceptance effect is the presence of the torus coils in between every sector. These represent a complete loss of information for a small range of  $\phi$  between each sector.
2. Inefficiencies due to the probabilistic nature of particle interaction in the detector subsystems also lower the overall acceptance.
3. Detectors have finite spatio-temporal resolutions.

In order to understand and limit the impact of these effects on the physics extracted from the experiment, a mock experiment is simulated. In the simulated experiment everything is modelled as realistically as possible. The simulation used for CLAS is called GSIM and is based on the CERN package GEANT3 (GEometry ANd Tracking).

In this controlled environment, control samples of events can be generated and fed into the simulation. The output of GSIM is a bos file that is similar to the raw data from the data acquisition system. This is then reconstructed using the same reconstruction algorithm that is applied to data (`userana`).

By retaining the truth information for all particles that are generated, the effect of the detector can be studied completely. These concepts can be stated more formally by considering the true  $t(x')$  and measured  $m(x)$  distributions of some observable. In the absence of background processes, the relationship between these distributions is expressed as a Fredholm integral equation of the first kind.

$$m(x) = \int_{\Omega} K(x, x') t(x') dx' \quad (6.2)$$

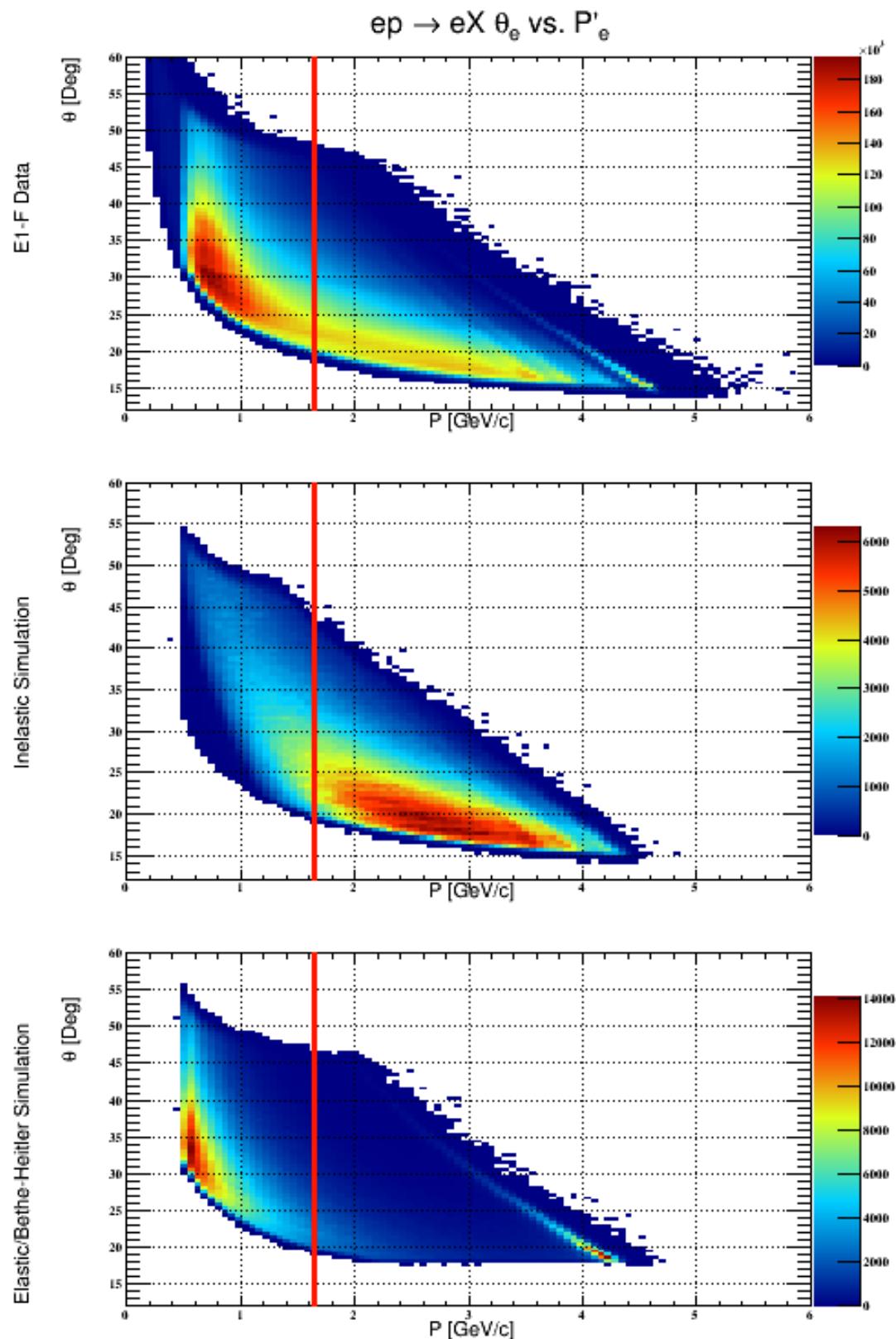


Figure 6.1: Event distributions ( $\theta_e$  vs  $p$ ) for data, simulated inelastic events, and simulated elastic events with radiation are shown. The red line indicates the momentum cut applied by restricting  $y < 0.7$ .

Here  $K(x, x')$  is a kernel which encodes information about detector acceptance due to the effects described above. The goal of the Monte Carlo simulation is then to *unfold* the measured distribution  $m(x)$  by providing an estimate of  $K(x, x')$  and finally corrected the data to get  $t(x)$ .

Observed events are usually aggregated into bins and the problem is naturally discretized and written in vector-matrix form.

$$\mathbf{A}\mathbf{x} = \mathbf{y} \quad (6.3)$$

In this notation  $\mathbf{A}$  represents the response matrix, a discretized version of the kernel function  $K$ , the vector  $\mathbf{y}$  represents the measured distribution in the bins, and the vector  $\mathbf{x}$  is the true distribution over the bins. The matrix elements  $A_{ij}$  can be estimated by using generating events, passing them through a monte-carlo detector simulation, and then counting the number of events that are reconstructed in bin  $i$  when generated in bin  $j$ . This quantity is then normalized by the total number of events generated in the  $j^{th}$  bin. In the absence of bin migration and with perfect acceptance this matrix is the identity matrix  $I^n$  where  $n$  is the total number of bins.

$$A_{ij} = \frac{n_{rec=i, gen=j}}{n_{gen=j}} \quad (6.4)$$

The binned true distribution can be recovered by inverting the response matrix and correcting the observed distribution.

$$\mathbf{x} = \mathbf{A}^{-1}\mathbf{y} \quad (6.5)$$

In the absence of bin migration, the matrix becomes diagonal with efficiency elements  $\epsilon_i$  that represent the fraction of events reconstructed in the bin  $i$ .

$$\mathbf{A} = \begin{pmatrix} \epsilon_0 & 0 & 0 \\ \vdots & \ddots & \\ 0 & & \epsilon_n \end{pmatrix} \quad (6.6)$$

The inverse is,

$$\mathbf{A}^{-1} = \begin{pmatrix} 1/\epsilon_0 & 0 & 0 \\ \vdots & \ddots & \\ 0 & & 1/\epsilon_n \end{pmatrix} \quad (6.7)$$

and the corrected observation for the  $i^{th}$  bin is simply given by the observation over the efficiency.

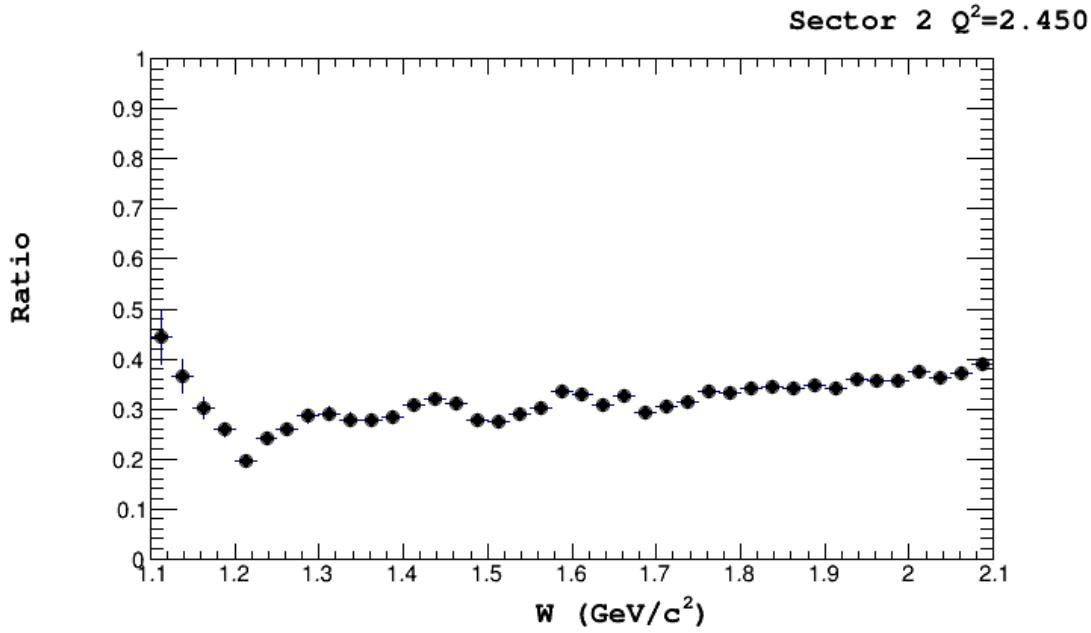


Figure 6.2: Acceptance for one bin of  $Q^2$  and one sector of CLAS.

$$t_i = \frac{m_i}{\epsilon_i} = m_i \frac{n_{gen=i}}{n_{rec=i}} \quad (6.8)$$

This is the simple *bin-by-bin* acceptance correction method, which is widely used and produces accurate results provided that bin migration is not significant. In this analysis the simple bin-by-bin acceptance correction is used.

### 6.0.3 Radiative Corrections

The removal of radiative effects from the measured distribution is a similar unfolding problem as described above for acceptance corrections. For this work we use two monte carlo event generators. The first generator produces inelastic events in the resonance region with radiative effects included that alter the kinematics. The second generator is the same parametrization but without radiative effects. The radiated generator is used to calculate the acceptance corrections, and both are used to try to remove the radiative effects on the cross section. The ratio  $R^{(i)}$  is defined for the  $i^{th}$  bin as shown below.

$$R^{(i)} = \frac{n_{unrad}^{(i)}}{n_{rad}^{(i)}} \quad (6.9)$$

This factor can be estimated without passing events through the simulation and we use the results directly from the output of the event generator to correct the cross section.

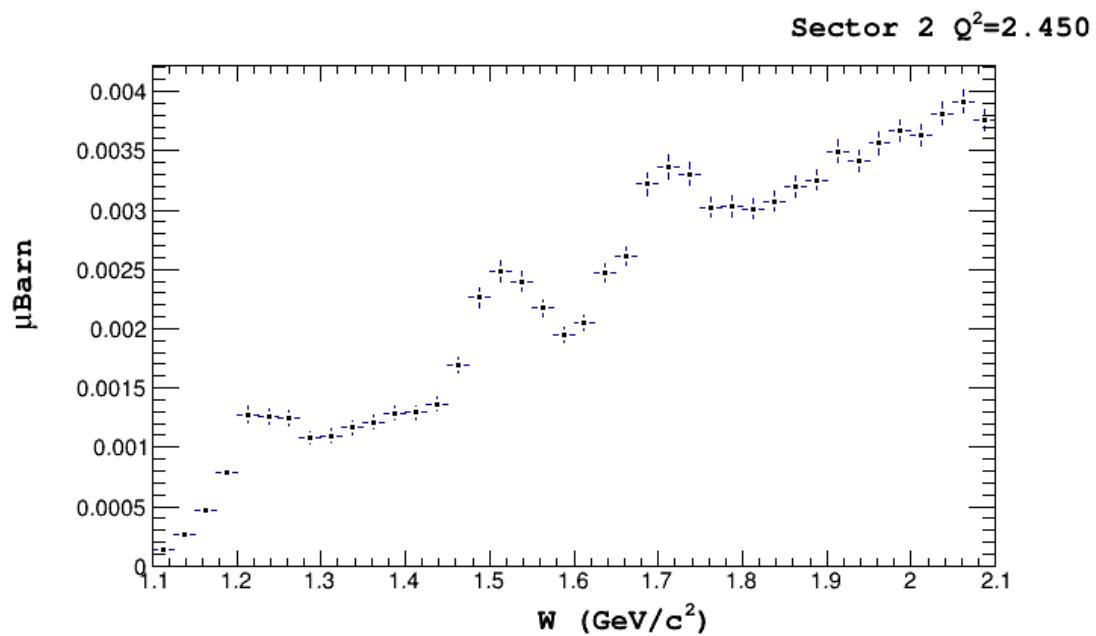


Figure 6.3: Inclusive cross section in the resonance region shown for one bin of  $Q^2$  in one sector of CLAS. Our calculation is compared with a trusted model.

#### 6.0.4 Model Comparison

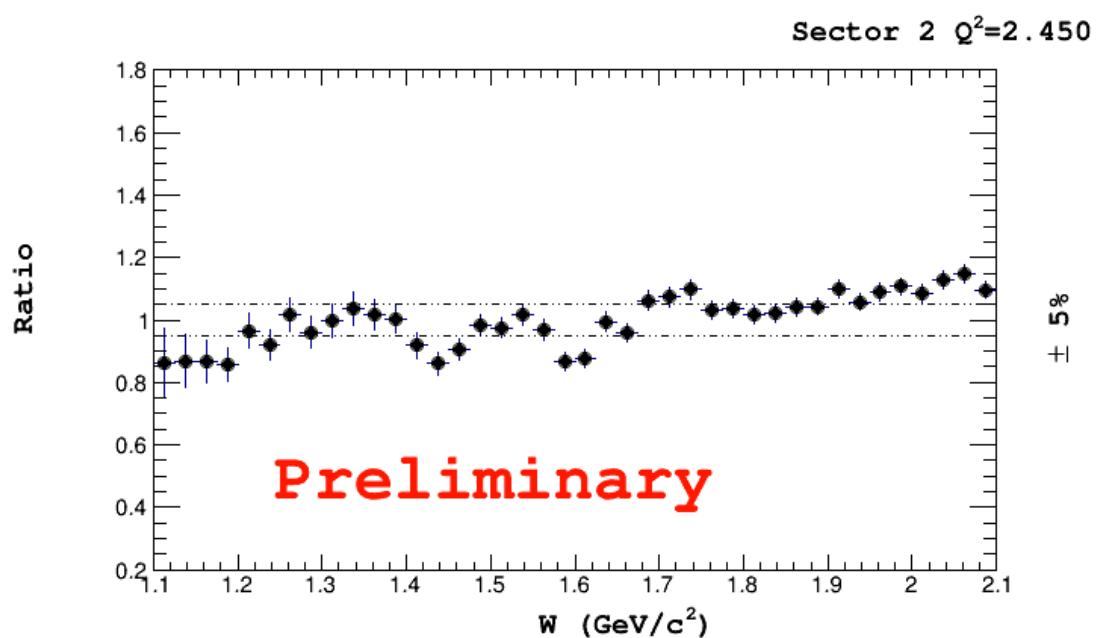


Figure 6.4: The ratio of cross section/model is shown here for one bin of  $Q^2$ .



# Chapter 7

## SIDIS Cross Section

This chapter discusses the analysis of semi-inclusive deeply inelastic events. After a brief motivation, the data analysis details are described.

### 7.1 Introduction

The primary goal of this work is to provide the SIDIS cross section for charged pions ( $\pi^\pm$ ) over a large kinematic range ( $0.1 < x < 0.6$ ,  $1.0 < Q^2 < 4.7$ ,  $0.0 < z < 0.9$ ,  $0.0 < P_T^2 < 1.0$ ,  $-180^\circ < \phi_h < 180^\circ$ ). Fortunately these channels were studied by Nathan Harrison and the CLAS collaboration [?] using the E1-F dataset (2010-2015). By writing the cross section as,

$$\frac{d^5\sigma}{dx dQ^2 dz dP_T^2 d\phi_h} = A_0 \left[ 1 + A_{UU}^{\cos \phi_h} \cos \phi_h + A_{UU}^{\cos(2\phi_h)} \cos(2\phi_h) \right] \quad (7.1)$$

Harrison et al. measured the un-normalized quantities  $A_0$ ,  $A_{UU}^{\cos \phi}$ , and  $A_{UU}^{\cos(2\phi)}$  which are defined below.

$$A_0 = \frac{\pi \alpha^2 y (1 + \gamma^2 / 2x)}{2EM_p x^2 Q^2 (1 - \varepsilon)} (F_{UU,T} + \varepsilon F_{UU,L}) \quad (7.2)$$

$$A_{UU}^{\cos \phi_h} = \sqrt{2\varepsilon(1 + \varepsilon)} \frac{F_{UU}^{\cos \phi_h}}{F_{UU,T} + \varepsilon F_{UU,L}} \quad (7.3)$$

$$A_{UU}^{\cos(2\phi_h)} = \varepsilon \frac{F_{UU}^{\cos(2\phi_h)}}{F_{UU,T} + \varepsilon F_{UU,L}} \quad (7.4)$$

In order to measure the structure functions  $F_{UU}^{\cos \phi_h}$  and  $F_{UU}^{\cos(2\phi_h)}$  directly, the integrated luminosity is needed. The calculation of this quantity for E1-F is described in detail during chapter 2 of this document. Experimentally, the cross section in the  $i^{th}$  bin is given as,

$$\frac{d\sigma}{dx dQ^2 dz dP_T^2 d\phi_h} = \frac{1}{\Delta x \Delta Q^2 \Delta z \Delta P_T^2 \Delta \phi_h} \frac{N_{obs}^{(i)}}{\mathcal{L} A^{(i)} R^{(i)}} \quad (7.5)$$

where the superscript  $(i)$  reminds the reader that these quantities are calculated for every bin. Throughout this chapter the symbols  $A^{(i)}$ ,  $R^{(i)}$ , and  $B^{(i)}$  refer to the acceptance correction, radiative correction, and bin centering correction respectively. These factors will be described in more details in this chapter. Finally, the  $\Delta$  factors here denote the width of each bin in 1 dimension of the 5-dimensional space (non-uniform sized bins may be used, in which case this factor also carries an index  $(i)$ ).

The integrated luminosity obtained in chapter 2 can be directly applied to the measurement of Harrison et al. to produce 5-dimensional differential cross sections. This procedure is carried out, but only after the luminosity factor is independently verified by calculating the cross section for inclusive inelastic electron scattering in the resonance region (here  $1.1 < W < 2.1 \text{ GeV}/c^2$ ). Accurate models for the inclusive cross section exist based on phenomenological fits to existing datasets. For verification, we compare the cross section from E1-F to a model created by Cynthia Keppel [? ].

The calculation of the inelastic scattering cross section as described here is non-trivial, and (together with the phenomenological analysis presented in chapter 6) constitutes the main original effort exerted by the authors.

## Appendix A: Derivation of formulas related to errors

### Propagation of errors

Let  $\vec{x}$  be a set of  $n$  random variables  $\vec{x} = (x_1, x_2, \dots, x_n)$  and known mean  $\mu_i = \langle x_i \rangle$  and covariance  $V_{ij} = \langle x_i x_j \rangle - \langle x_i \rangle \langle x_j \rangle$ . Suppose that we measure a function  $f(\vec{x})$  that depends on the variables  $\vec{x}$  and we want to understand how the covariances  $V_{ij}$  on  $\vec{x}$  will show up manifest themselves as errors on our measurement of  $f(\vec{x})$ . We can start by expanding our function around the expected value of  $x_i$ .

$$f(\vec{x}) \approx f(\vec{\mu}) + \sum_{i=1}^n \frac{\partial f}{\partial x_i} \Big|_{x_i=\mu_i} (x_i - \mu_i) \quad (6)$$

We can then take the expectation value of our function.

$$\langle f(\vec{x}) \rangle = \langle f(\vec{\mu}) \rangle + \sum_{i=1}^n \frac{\partial f}{\partial x_i} \Big|_{x_i=\mu_i} \langle x_i - \mu_i \rangle \quad (7)$$

Where here the term  $\langle x_i - \mu_i \rangle$  is zero.

$$\langle x_i - \mu_i \rangle = \langle x_i \rangle - \mu_i = \mu_i - \mu_i = 0 \quad (8)$$

It is apparent then that the expectation value of our function  $f$  evaluated close to the expected values of our variables  $\vec{x}$  is just the function evaluated at the expectation value of the random variables  $\vec{x}$ .

$$\langle f(\vec{x}) \rangle = \langle f(\mu) \rangle = f(\mu) \quad (9)$$

We can also consider the second moment  $\langle f^2(\vec{x}) \rangle$ ,

$$\langle f^2(\vec{x}) \rangle \approx \left\langle \left( f(\vec{\mu}) + \sum_{i=1}^n \frac{\partial f}{\partial x_i} \Big|_{x_i=\mu_i} (x_i - \mu_i) \right)^2 \right\rangle \quad (10)$$

which is,

$$= \langle f^2(\mu) \rangle + \sum_{i=1}^n \sum_{j=1}^n \frac{\partial f}{\partial x_i} \Big|_{x_i=\mu_i} \frac{\partial f}{\partial x_j} \Big|_{x_j=\mu_j} \langle (x_i - \mu_i)(x_j - \mu_j) \rangle + 2 \left\langle f(\mu) \sum_{i=1}^n \frac{\partial f}{\partial x_i} \Big|_{x_i=\mu_i} (x_i - \mu_i) \right\rangle \quad (11)$$

and the last expectation value vanishes due to the same logic used when calculating the first moment. We recognize the term  $\langle (x_i - \mu_i)(x_j - \mu_j) \rangle$  as the element of the covariance matrix  $V_{ij}$ . Our second moment is then complete as follows.

$$\langle f^2(\vec{x}) \rangle = f^2(\mu) + \sum_{i=1}^n \sum_{j=1}^n \frac{\partial f}{\partial x_i} \Big|_{x_i=\mu_i} \frac{\partial f}{\partial x_j} \Big|_{x_j=\mu_i} V_{ij} \quad (12)$$

We can then calculate the variance of the function.

$$\sigma_f^2 = \langle f^2(\vec{x}) \rangle - \langle f(\vec{x}) \rangle^2 \quad (13)$$

$$= (f^2(\vec{\mu}) - f^2(\vec{\mu})) + \sum_{i=1}^n \sum_{j=1}^n \frac{\partial f}{\partial x_i} \Big|_{x_i=\mu_i} \frac{\partial f}{\partial x_j} \Big|_{x_j=\mu_i} V_{ij} \quad (14)$$

$$= \sum_{i=1}^n \sum_{j=1}^n \frac{\partial f}{\partial x_i} \Big|_{x_i=\mu_i} \frac{\partial f}{\partial x_j} \Big|_{x_j=\mu_i} V_{ij} \quad (15)$$

This is the standard error propagation formula which is widely used. These correlations  $\sigma_{ij}$  can arise from several sources.

- Common measurement uncertainties.
- Correlations in  $x_i x_j$  leading to correlations in  $\sigma_i \sigma_j$ .

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