

Disturbance Torques^(Moments) on a Spacecraft (12 de Raita)

Gravity - Gradient Torque

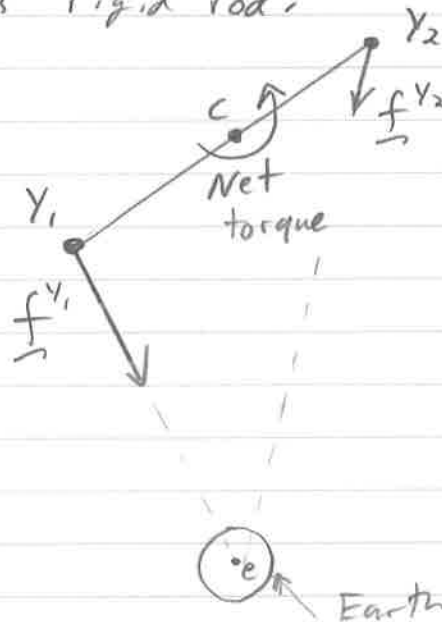
Recall that

$$\underline{f}^{ye} = - \frac{\mu m_y}{\|\underline{r}^{ye}\|^3} \underline{r}^{ye}$$

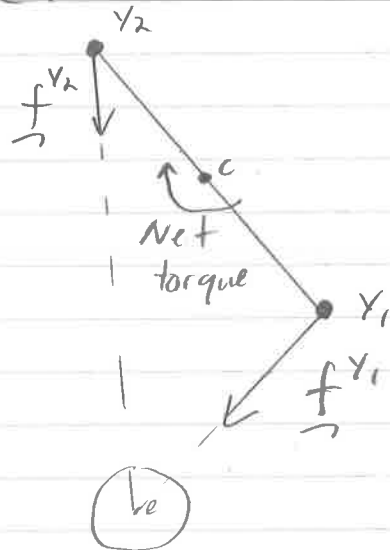
$$\|\underline{f}^{ye}\| \sim \|\underline{r}^{ye}\|^{-2}$$

Thus, gravitational force on a mass further from Earth is smaller than force on a mass closer to Earth.

Consider two particles of equal mass connected by a massless rigid rod.

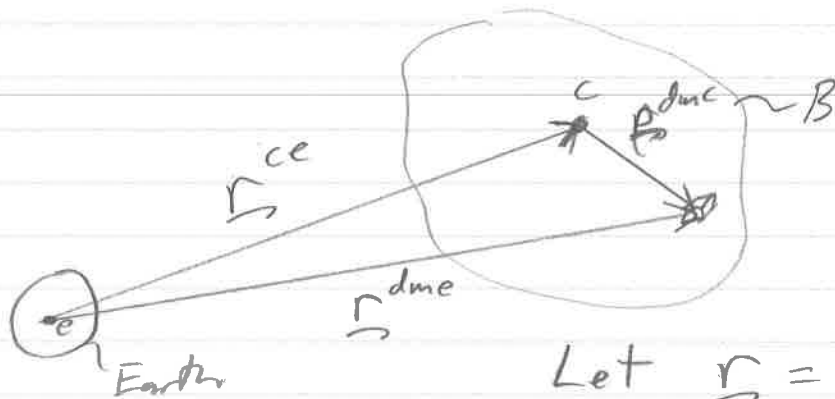


Force acting on Y_1 is larger than force acting on Y_2 , thus causing a net torque (moment) about pt. C.



Gravity-gradient torque will try to align body to point toward Earth.

Consider a continuous rigid body (CRB)



Let $\underline{r} = \underline{r}^{ce}$, $\underline{p} = \underline{p}^{dmc}$

Gravitational force acting on dm is

$$d\underline{f}^{dm} = -\frac{\mu dm}{\|\underline{r}^{dmc}\|_2^3} \underline{r}^{dmc}$$

$$= -\frac{\mu}{\|\underline{r}^{ce} + \underline{p}^{dmc}\|_2^3} (\underline{r}^{ce} + \underline{p}^{dmc}) dm$$

Note that

$$\begin{aligned} \|\underline{r}^{ce} + \underline{p}^{dmc}\|_2^3 &= [(\underline{r}^{ce} + \underline{p}^{dmc}) \cdot (\underline{r}^{ce} + \underline{p}^{dmc})]^{3/2} \\ &= [\underline{r}^{ce} \cdot \underline{r}^{ce} + \underline{p}^{dmc} \cdot \underline{p}^{dmc} + 2\underline{r}^{ce} \cdot \underline{p}^{dmc}]^{3/2} \\ &= (r^{ce^2} + p^{dmc^2} + 2r^{ce}p^{dmc}\cos\phi)^{3/2} \end{aligned}$$

where $r^{ce} = \|\underline{r}^{ce}\|_2$, $p^{dmc} = \|\underline{p}^{dmc}\|_2$, ϕ is the angle between \underline{r}^{ce} and \underline{p}^{dmc} .

$$\left(\| \underline{r}^{ce} + \underline{p}^{dmc} \|_2^3 = r^{ce3} \left(1 + \left(\frac{p^{dmc}}{r^{ce}} \right)^2 + 2 \left(\frac{p^{dmc}}{r^{ce}} \right) \cos \phi \right)^{3/2} \right)$$

Assume that $p^{dmc} \ll r^{ce}$

Define $\varepsilon = \frac{p^{dmc}}{r^{ce}} \ll 1$, therefore $\varepsilon^2 \approx 0$

Therefore,

$$\frac{1}{\| \underline{r}^{ce} + \underline{p}^{dmc} \|_2^3} \approx \frac{1}{r^{ce3} (1 + 2\varepsilon \cos \phi)^{3/2}} = f(\varepsilon)$$

Take Taylor series approximation

$$f(\varepsilon) \approx f(0) + \left. \frac{df}{d\varepsilon} \right|_{\varepsilon=0} \varepsilon$$

$$\left(\approx \frac{1}{r^{ce3}} + \left. \frac{d}{d\varepsilon} \left(\frac{1}{r^{ce3} (1 + 2\varepsilon \cos \phi)^{3/2}} \right) \right|_{\varepsilon=0} \varepsilon \right)$$

$$\left(= \frac{1}{r^{ce3}} - \frac{1}{r^{ce3}} (3 \cos \phi) \varepsilon \right)$$

$$\begin{aligned} \frac{1}{\| \underline{r}^{ce} + \underline{p}^{dmc} \|_2^3} &\approx \left(\frac{1}{r^{ce3}} - \frac{3}{r^{ce5}} (r^{ce} p^{dmc} \cos \phi) \right) \\ &= \frac{1}{r^{ce3}} - \frac{3 \underline{r}^{ce} \cdot \underline{p}^{dmc}}{r^{ce5}} \end{aligned}$$

Now,

$$d\vec{f}^{dm} \approx -\frac{\mu(\vec{r}^{ce} + \vec{r}^{dmc})}{r^{ce3}} \left(1 - \frac{3\vec{r}^{ce} \cdot \vec{r}^{dmc}}{r^{ce2}}\right) dm$$

Total torque on B relative to c is

$$\vec{\tau}_{Bc,gg} = \int_B \vec{r}^{dmc} \times d\vec{f}^{dm}$$

$$= \int_B \vec{r}^{dmc} \times \left(-\frac{\mu(\vec{r}^{ce} + \vec{r}^{dmc})}{r^{ce3}} \right) dm$$

$$+ \int_B \vec{r}^{dmc} \times \left(\frac{3\mu \vec{r}^{ce} \cdot \vec{r}^{dmc}}{r^{ce5}} (\vec{r}^{ce} + \vec{r}^{dmc}) \right) dm$$

$$= -\frac{\mu}{r^{ce3}} \left[\left(\int_B \vec{r}^{dmc} dm \right) \times \vec{r}^{ce} + \int_B \vec{r}^{dmc} \times \vec{r}^{dmc} dm \right]$$

$$+ \frac{3\mu}{r^{ce5}} \left(\int_B \vec{r}^{dmc} \times \vec{r}^{ce} (\vec{r}^{ce} \cdot \vec{r}^{dmc}) dm \right.$$

$$\left. + \int_B \vec{r}^{dmc} \times \vec{r}^{dmc} (\vec{r}^{ce} \cdot \vec{r}^{dmc}) dm \right)$$

$$= \left(\frac{3\mu}{r^{ce5}} \int_B \vec{r}^{dmc} \times \vec{r}^{ce} (\vec{r}^{ce} \cdot \vec{r}^{dmc}) dm \right)$$

Resolve all vectors in \mathcal{F}_b :

$$\underline{\tau}_b^{Bc,gg} = \frac{3\mu}{r^{ce5}} \int_B \underline{r}_b^{dmc^x} \underline{r}_b^{ce} \underline{r}_b^{dmc^T} \underline{r}_b^{ce} dm$$

$$= -\frac{3\mu}{r^{ce5}} \underline{r}_b^{ce^x} \int_B \underline{r}_b^{dmc} \underline{r}_b^{dmc^T} dm \underline{r}_b^{ce}$$

Recall that $\underline{r}_b^{dmc} \underline{r}_b^{dmc^T} = \underline{r}_b^{dmc^x} \underline{r}_b^{dmc^x} + \underline{r}_b^{dmc^T} \underline{r}_b^{dmc} \underline{1}$

$$\underline{\tau}_b^{Bc,gg} = -\frac{3\mu}{r^{ce5}} \underline{r}_b^{ce^x} \int_B \underline{r}_b^{dmc^x} \underline{r}_b^{dmc^x} dm \underline{r}_b^{ce}$$

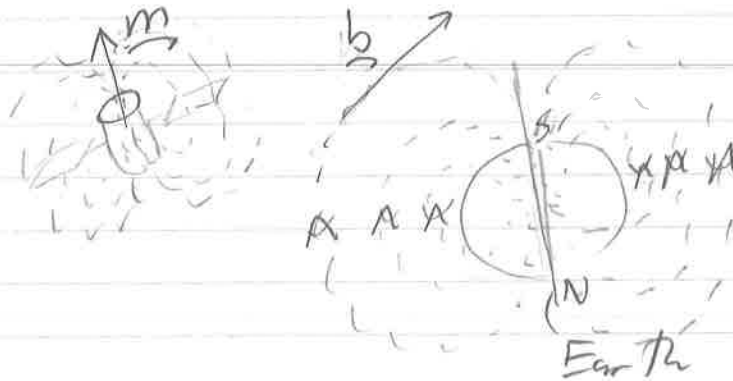
$$- \frac{3\mu}{r^{ce5}} \underline{r}_b^{ce^x} \int_B \underline{r}_b^{dmc^T} \underline{r}_b^{dmc} dm \underline{r}_b^{ce}$$

$$= \int_B \underline{r}_b^{dmc^T} \underline{r}_b^{dmc} dm \underline{r}_b^{ce^x} \underline{r}_b^{ce} \underline{0}$$

$$\boxed{\underline{\tau}_b^{Bc,gg} + \frac{3\mu}{r^{ce5}} \underline{r}_b^{ce^x} \underline{I}_b^B \underline{r}_b^{ce}}$$

↳ Gravity Gradient Torque

Magnetic Disturbance Torque



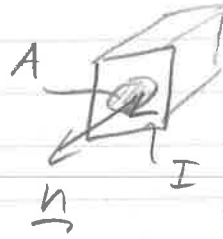
Earth's magnetic field, \underline{b} , interacts with magnetization of the spacecraft due to electronic components. The spacecraft has a residual magnetic dipole moment \underline{m} . This causes disturbance torque of

$$\underline{\tau}^{Bc} = \underline{m} \times \underline{b}$$

By adjusting magnetic dipole of spacecraft, we can take advantage of Earth's magnetic field to produce a desired magnetic torque. This is the concept used by magnetic torque rods (magnetorquers).

$$\underline{m} = N I A \underline{n}$$

↑ ↑ ↙
 number current area
 of turns



causes of
Other Relevant disturbance torques include

- Solar radiation pressure (SRP) ($\sim 10^{-5} \text{ N}\cdot\text{m}$)
(momentum transfer of photons)
- Aerodynamic forces ($\sim 10^{-7} - 10^{-1}$ depending on altitude)
(negligible above 1000 Km alt.)
- Meteoroidal impacts
- Structural dynamics

For comparison

Gravity gradient ($\sim 10^{-4} - 10^{-3} \text{ N}\cdot\text{m}$)

Magnetic torque ($\sim 10^{-5} - 10^{-4} \text{ N}\cdot\text{m}$)