

# 1 Introduction and Scientific Background

Galaxy–galaxy (g–g) lensing is a powerful probe of the relation between galaxies and dark matter (DM) haloes, but its theoretical interpretation requires a careful modelling of various contributions, such as the contributions from central and satellite galaxies. For this purpose, we developed a new method using the components of lensing shear  $\gamma_1$  and  $\gamma_2$  instead of tangential shear to extract information from the g–g weak lensing signal by comparing it to high-resolution dissipationless simulations that resolve subhaloes. We find that this method can better constrain subhalo masses and thus enabling us to further investigate halo–galaxy connections.

From the NFW density profile, we can derive the convergence field

$$\kappa(\theta) = \frac{\Sigma(D_l\theta)}{\Sigma_{cr}} \quad (1)$$

where

$$\Sigma(\xi) = \int dz \rho(\xi_1, \xi_2, z) \quad (2)$$

is the surface density,  $\xi = (\xi_1, \xi_2)$  is the impact two-vector in the lens plane,  $z$  is the perpendicular coordinate to the lens plane.

$$\Sigma_{cr} = \frac{1}{4\pi G} \frac{D_s}{D_l D_{ls}} \quad (3)$$

The deflection angle can be expressed in terms of convergence as

$$\alpha(\theta) = \frac{1}{\pi} \int d^2\theta' \kappa(\theta) \frac{\theta - \theta'}{|\theta - \theta'|^2} \quad (4)$$

and 2-dimensional deflection potential:

$$\Psi(\theta) = \frac{1}{\pi} \int d^2\theta' \kappa(\theta) \ln|\theta - \theta'| \quad (5)$$

such that:

$$\alpha(\theta) = \nabla\Psi(\theta) \quad (6)$$

$\Psi(\theta)$  then satisfies the 2-dimensional Poisson equation:

$$\kappa(\theta) = \frac{1}{2} \nabla^2 \Psi(\theta) \quad (7)$$

the lensing mapping Jacobian matrix is:

$$\mathcal{A}(\theta) = \delta_{ij} - \frac{\partial^2 \Psi(\theta)}{\partial \theta_i \partial \theta_j} = \begin{pmatrix} 1 - \kappa - \gamma_1 & \gamma_2 \\ \gamma_2 & 1 - \kappa + \gamma_1 \end{pmatrix} \quad (8)$$

So

$$\gamma_1 = \frac{A_{11} - A_{22}}{2} \quad (9)$$

$$\gamma_2 = A_{12} \quad (10)$$

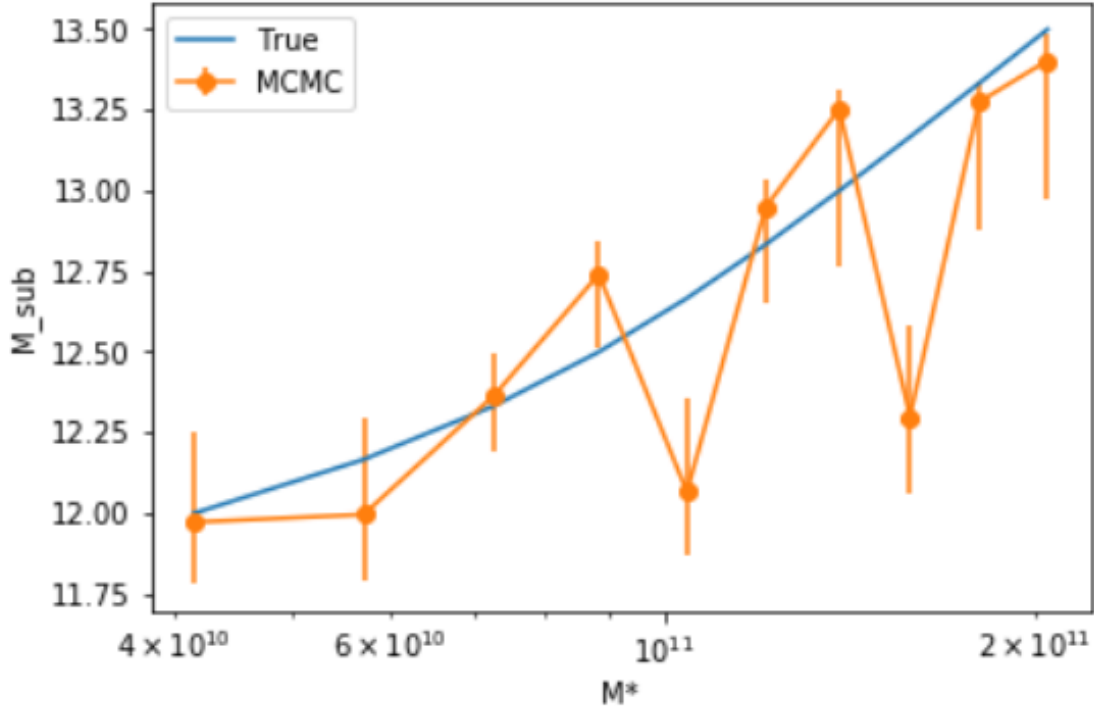


Figure 1: This plot shows the MCMC result based on  $\gamma_1$  and  $\gamma_2$  method comparing with true value.

## 2 Progress to Date and Motivation for Future Work

We constrain the subhalo masses using MCMC method, for a single sample case, the log-likelihood function is given by:

$$\log \mathcal{L}(M | \mathcal{D}) = \sum_k -\frac{1}{2}(\gamma_{1,k}^{theory}(M) - \gamma_{1,k}^{data})^T \Sigma^{-1}(\gamma_{1,k}^{theory}(M) - \gamma_{1,k}^{data}) - \frac{1}{2}(\gamma_{2,k}^{theory}(M) - \gamma_{2,k}^{data})^T \Sigma^{-1}(\gamma_{2,k}^{theory}(M) - \gamma_{2,k}^{data})$$

Where  $\gamma_{1,k}^{theory}(M)$  and  $\gamma_{2,k}^{theory}(M)$  are the theoretic prediction of  $\gamma_1$  and  $\gamma_2$  at pixel  $k$ , the covariance matrix is assumed to be diagonal with:

$$\Sigma = diag(\sigma_{\gamma_1}^2, \sigma_{\gamma_2}^2, \dots, \sigma_{\gamma_k}^2)$$

The total log-likelihood function is given by:

$$\log \mathcal{L} = \sum_{i=1}^n p_i \sum_{j=1}^m \log \mathcal{L}(M_i | \mathcal{D}_{i,j})$$

Where  $n = 10$  is the number of stellar mass bins,  $m = 50$  is the sampling number in each bin.  $p_i$  is the lenses number factor:

$$p_i = \frac{N_i}{m} \times \frac{S_{LSST}}{S_{cosmoDC2}}$$

Where  $N_i$  is the number of lenses in the  $i$ th mass bin between  $z = 0.3$  to  $z = 0.4$  in cosmoDC2 simulation,  $S_{LSST} = 20000\text{deg}^2$ ,  $S_{cosmoDC2} = 440\text{deg}^2$

### 3 Proposed Computational Methods

Given cosmological parameters in cosmoDC2, we first generate fake galaxies with subhalo masses and stellar masses. Then we fed the fake data to emcee, a module using affine invariant Markov chain Monte Carlo (MCMC) ensemble sampler (<https://github.com/dfm/emcee>). We ran a chain with 10 mass bins (10 parameters) with 10k steps.

### 4 Computational Resources

NERSC GPU

## REFERENCES