

# New constraints for QCD matter from improved Bayesian parameter estimation in heavy-ion collisions at LHC

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Based on PRC 104 (2021) 054904, arXiv:2111.08145

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Friday 1<sup>st</sup> July, 2022

Future Heavy-Ion Collision Projects: from LHC to FCC, July 1-2, 2022  
Yeosu, South Korea

Yeosu, 2015, D. J. Kim



UNIVERSITY OF JYVÄSKYLÄ

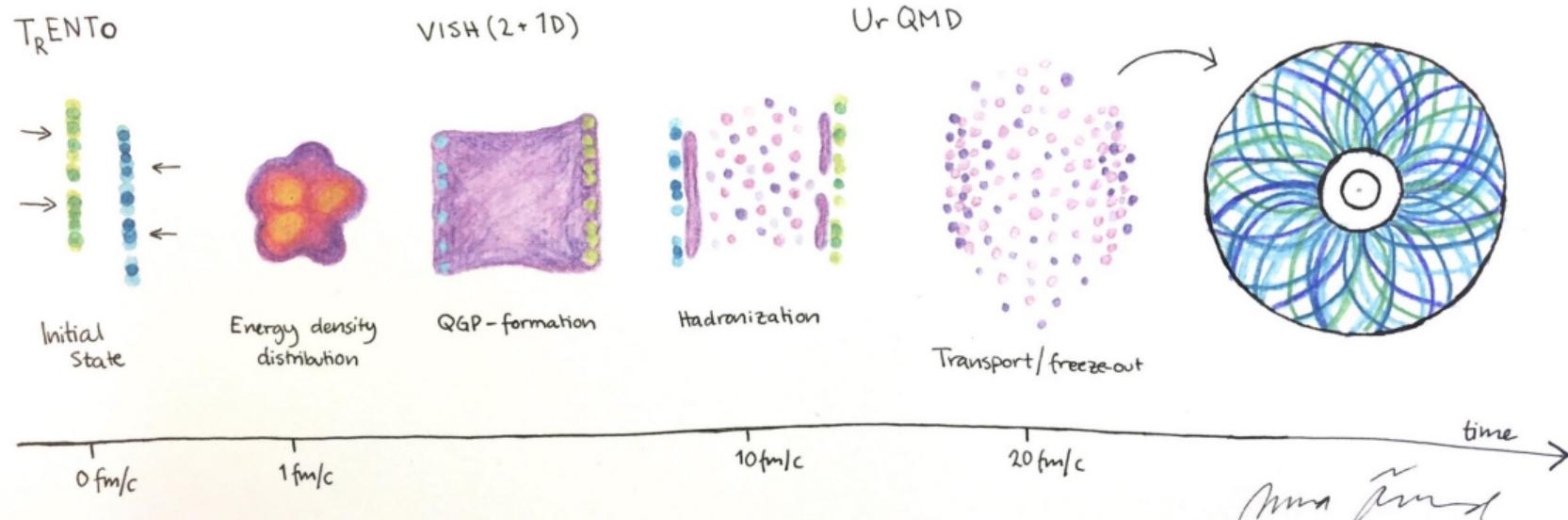


CoE  
QM



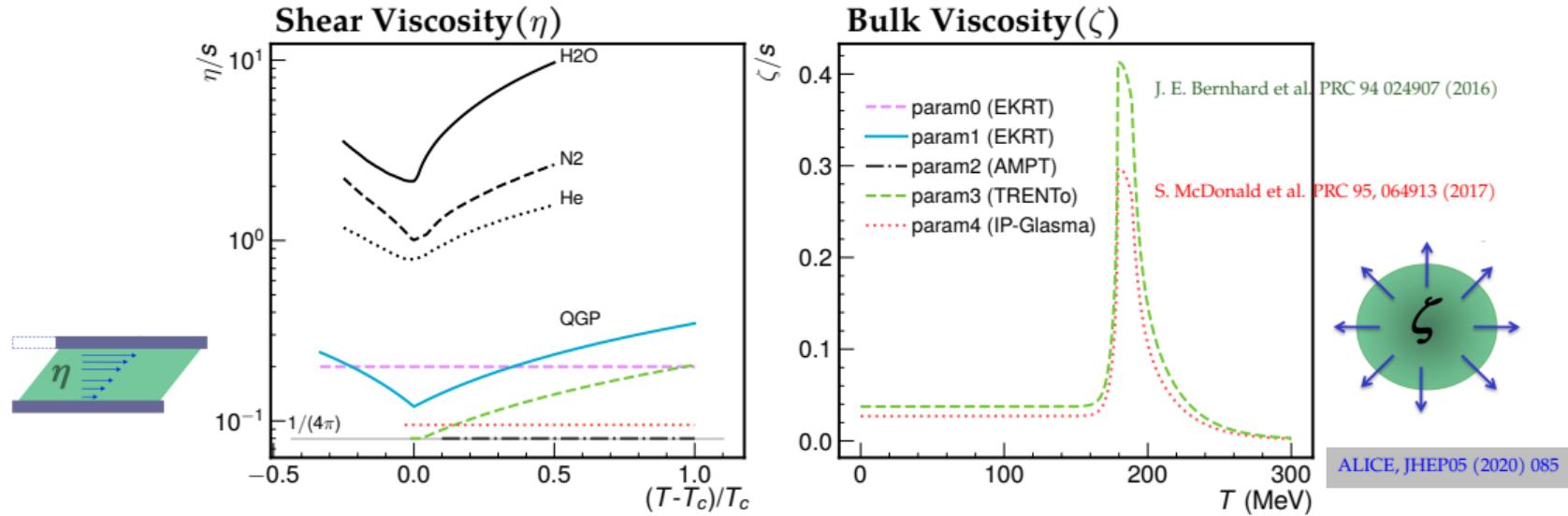
HELSINKI INSTITUTE OF PHYSICS

# THE DIFFERENT STAGES OF HEAVY-ION COLLISIONS



$$T^{\mu\nu} = eu^\mu u^\nu - (P + \Pi)\Delta_{\mu\nu} + \pi^{\mu\nu}, \quad \delta_\mu T^{\mu\nu} = 0$$

# TRANSPORT PROPERTIES IN HEAVY-ION COLLISIONS



$$(\eta/s)(T) = (\eta/s)(T_c) + (\eta/s)_{\text{slope}}(T - T_c) \left( \frac{T}{T_c} \right)^{(\eta/s)_{\text{curve}}}, (\zeta/s)(T) = \frac{(\zeta/s)_{\text{max}}}{1 + \left( \frac{T - (\zeta/s)_{T_{\text{peak}}}}{(\zeta/s)_{\text{width}}} \right)^2}$$

# LATEST MAP PARAMETERS, JYVASKYLA (2022)

Table: Input parameter ranges for the initial condition and hydrodynamic models.

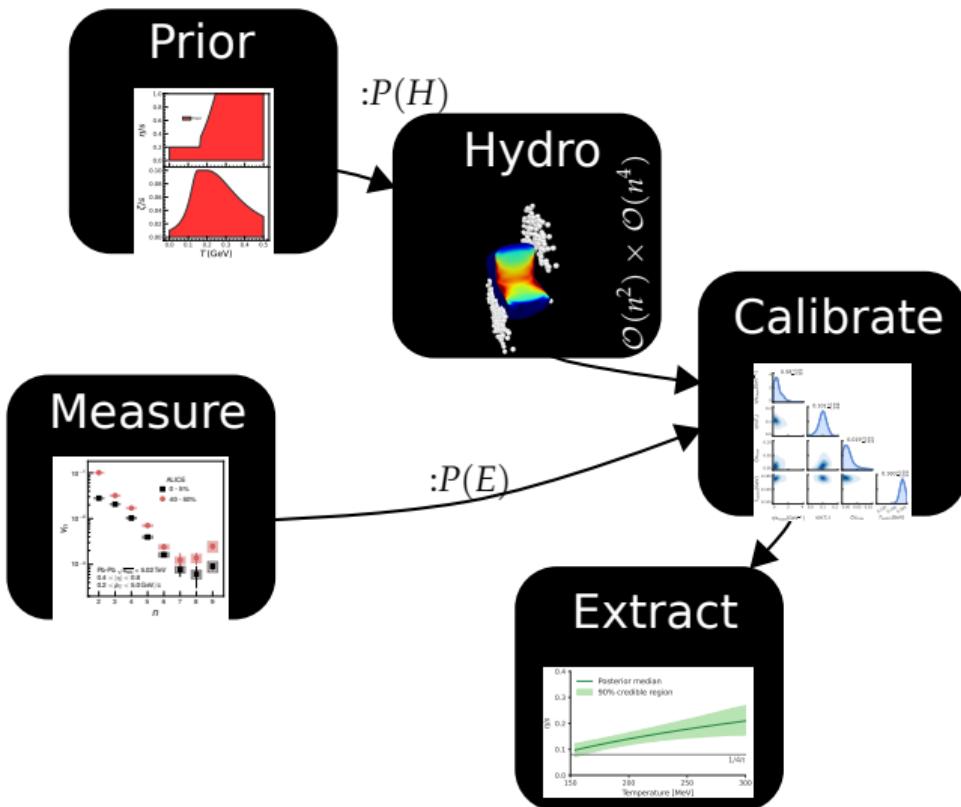
Parameter	Description	Range	MAP
$N(2.76 \text{ TeV})$	Overall normalization (2.76 TeV)	$[11.152, 18.960]$	14.373
$N(5.02 \text{ TeV})$	Overall normalization (5.02 TeV)	$[16.542, 25]$	21.044
$p$	Entropy deposition parameter	$[0.0042, 0.0098]$	0.0056
$\sigma_k$	Std. dev. of nucleon multiplicity fluctuations	$[0.5518, 1.2852]$	1.0468
$d_{\min}^3$	Minimum volume per nucleon	$[0.889^3, 1.524^3]$	$1.2367^3$
$\tau_{\text{fs}}$	Free-streaming time	$[0.03, 1.5]$	0.71
$T_c$	Temperature of const. $\eta/s(T)$ , $T < T_c$	$[0.135, 0.165]$	0.141
$\eta/s(T_c)$	Minimum $\eta/s(T)$	$[0, 0.2]$	0.093
$(\eta/s)_{\text{slope}}$	Slope of $\eta/s(T)$ above $T_c$	$[0, 4]$	0.8024
$(\eta/s)_{\text{curve}}$	Curvature of $\eta/s(T)$ above $T_c$	$[-1.3, 1]$	0.1568
$(\zeta/s)_{\text{peak}}$	Temperature of $\zeta/s(T)$ maximum	$[0.15, 0.2]$	0.1889
$(\zeta/s)_{\text{max}}$	Maximum $\zeta/s(T)$	$[0, 0.1]$	0.01844
$(\zeta/s)_{\text{width}}$	Width of $\zeta/s(T)$ peak	$[0, 0.1]$	0.04252
$T_{\text{switch}}$	Switching / particlization temperature	$[0.135, 0.165]$	0.1595

## BAYESIAN PARAMETER ESTIMATION

Bayes' theorem:

$$P(H|E) = \frac{P(E|H) \cdot P(H)}{P(E)}$$

$$, P(E) = \sum_{i=1}^n P(E|H_i)P(H_i)$$

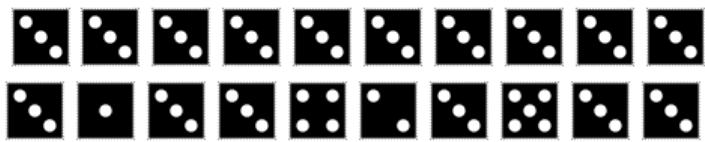


- Find optimal set of model parameters that best reproduce the experimental data
- Utilize constraints, such as flow observables, to help narrow down the  $\eta/s(T)$  and such.

Testing a single set of parameters requires  $\mathcal{O}(10^4)$  hydro events, and evaluating eight different parameters five times each requires  $5^8 \times 10^4 \approx 10^9$  hydro events. That's roughly  $10^5$  CPU years!

## BAYESIAN THEOREM - A SIMPLE QUIZ?

$$P(H|E) = \frac{P(E|H) \cdot P(H)}{P(E)}$$



1/6 1/6.....1/6, Prob(3) for 11th? It is 1/6

Up to 10th Prob(3) was 6/10, Prob for 11th?

Probability to have for 11th throw?

# BAYESIAN THEOREM

$$P(H|E) = \frac{P(E|H) \cdot P(H)}{P(E)}$$



$1/6$   $1/6$ ..... $1/6$ , Prob(3) for 11th? It is  $1/6$

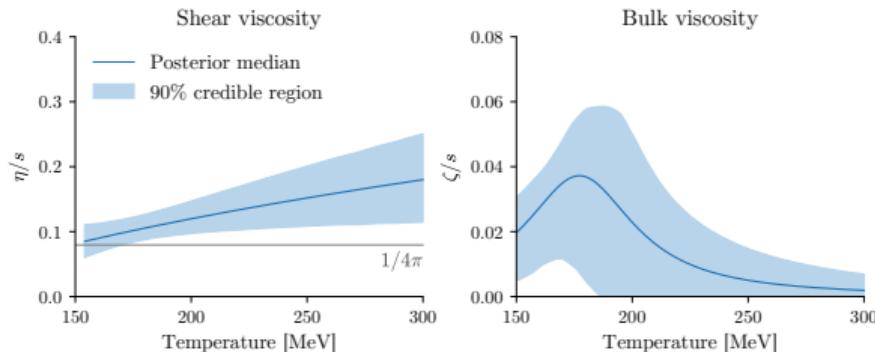
Up to 10th Prob(3) was  $6/10$ , Prob for 11th?

Probability to have for 11th throw is  $\frac{6+1}{10+1}$

# BAYESIAN PARAMETER ESTIMATION: PREVIOUS WORK

## JETSCAPE T<sub>R</sub>ENTo+MUSIC+SMASH

### Duke T<sub>R</sub>ENTo+VISH(2+1D)+UrQMD



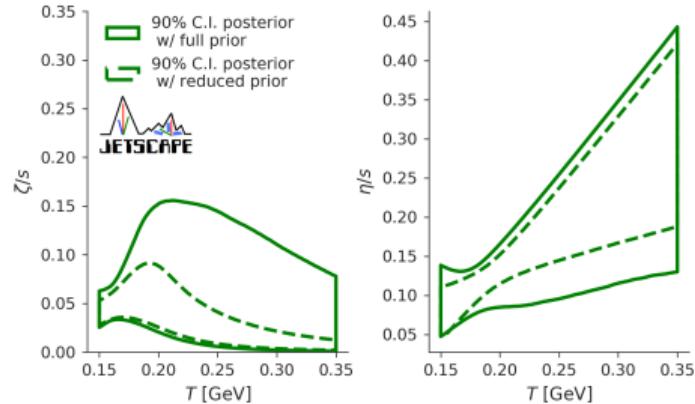
Steffen A. Bass *et. al.*, Nature Physics (2019)

- Low to moderate temperature dependence on  $\eta/s(T)$
- Moderate magnitude of  $\zeta/s(T)$  ( $\sim 0.1 \times$  w.r.t lattice QCD(PRL. **94**, 072305 (2005))
- Large uncertainty for both  $\eta/s(T)$  and  $\zeta/s(T)$ .
- Subsequent studies with still limited observables:
  - J. Auvilinen *et al.* PRC. **102**, 044911 (2020) ● G. Nijs *et al.* PRL. **126**, 202301 (2021)

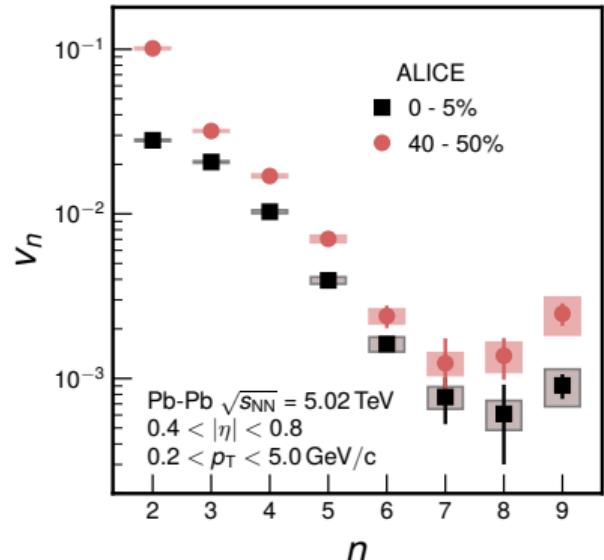
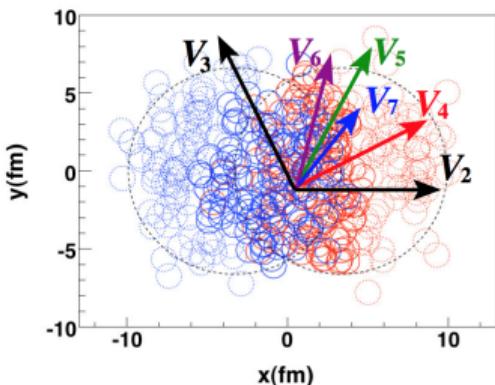
Uncertainties need to and can be further improved.

Only low-order harmonic  $v_n$  was used, including a *limited set of mostly 2.76 TeV* observables.

### P.B. Viscosity Posterior : Effect of Prior



# HIGHER FLOW HARMONICS AND FLOW FLUCTUATION

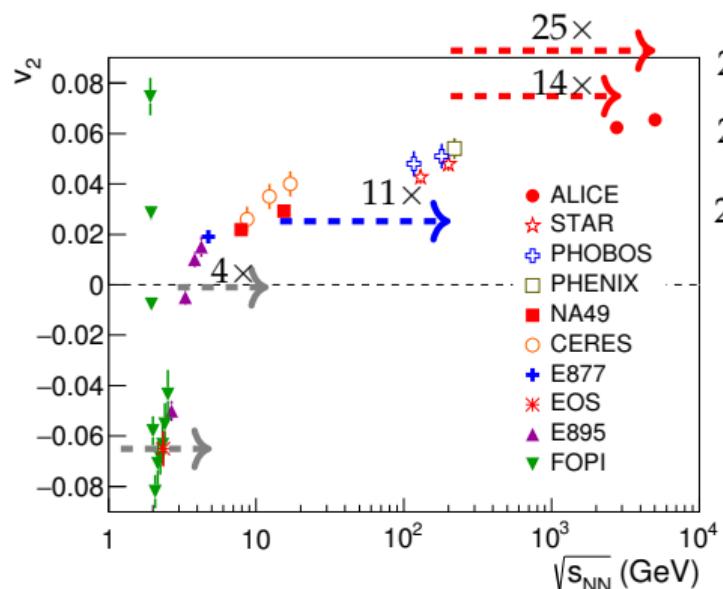


$$\text{blob} = \text{circle} + \text{blob} + \text{blob} + \text{blob}$$

$$P(\varphi) \propto \frac{1}{2\pi} \sum_{n=-\infty}^{+\infty} V_n e^{-in\varphi}$$

$$V_n \equiv v_n \{\psi_n\} e^{in(\psi_n - \phi)}$$

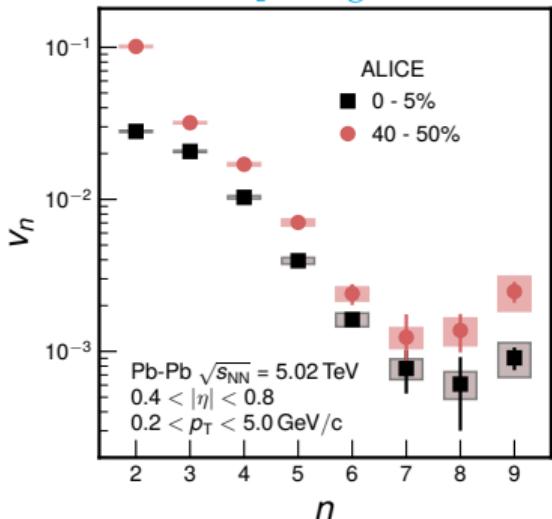
- Sensitive to initial state geometry and properties of the expanding QGP (viscosity( $\eta/s$ ), equation of state)

**$v_2$  VS  $\sqrt{s_{NN}}$  AND FLOW POWER SPECTRUM**

2015 LHC 5.02TeV CERN  
2010 LHC 2.76TeV CERN  
2000 RHIC 200GeV USA  
90s SPS 17GeV CERN  
80s AGS 4GeV USA

ALICE, Phys. Rev. Lett. 105 (2010) 252302

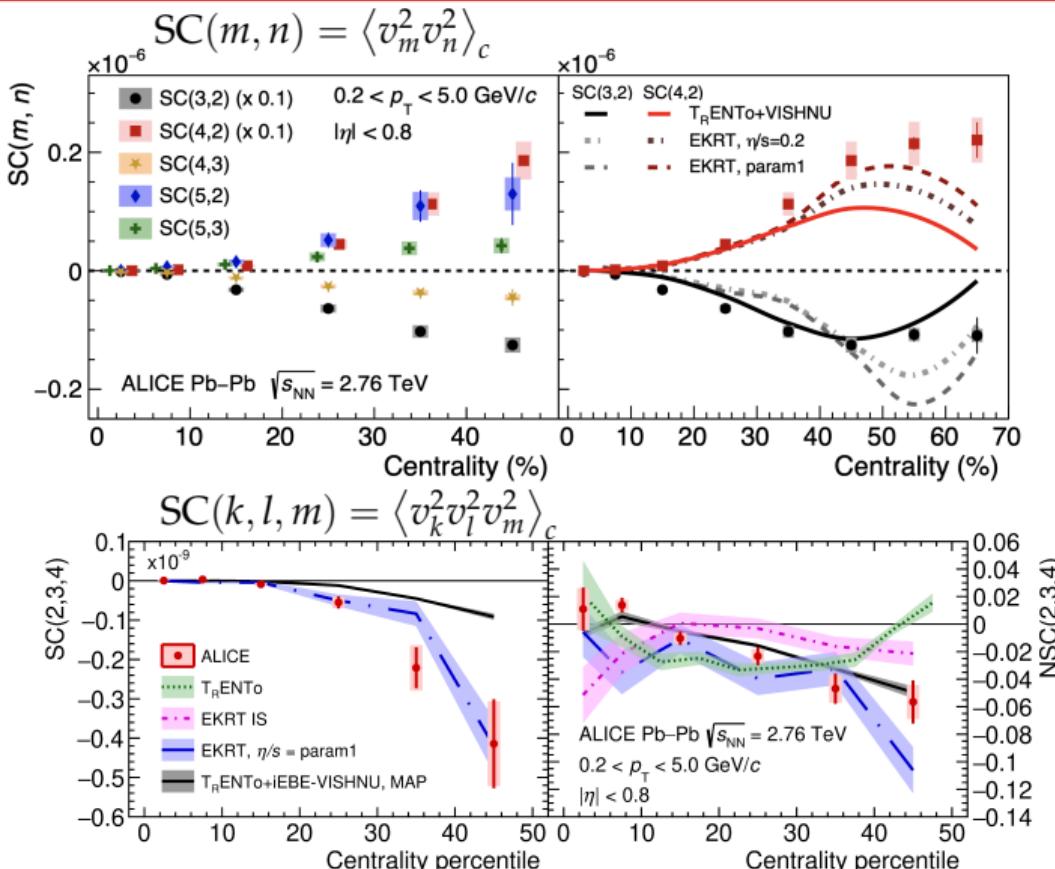
2020, cerncourier[Going with the flow]



ALICE, JHEP05 (2020) 085

Measured the largest flow  $v_2$  in 2010!  
Measured the largest harmonic order flow (up to  $v_9$ ) so far, 2020

# HIGH PRECISION FLOW RESULTS AND NEW DEVELOPMENTS- SYMMETRIC CUMULANTS



ALICE, Phys. Rev. Lett. 117 (2016) 182301

ALICE, Phys. Rev. C 97 no. 2, (2018) 024906

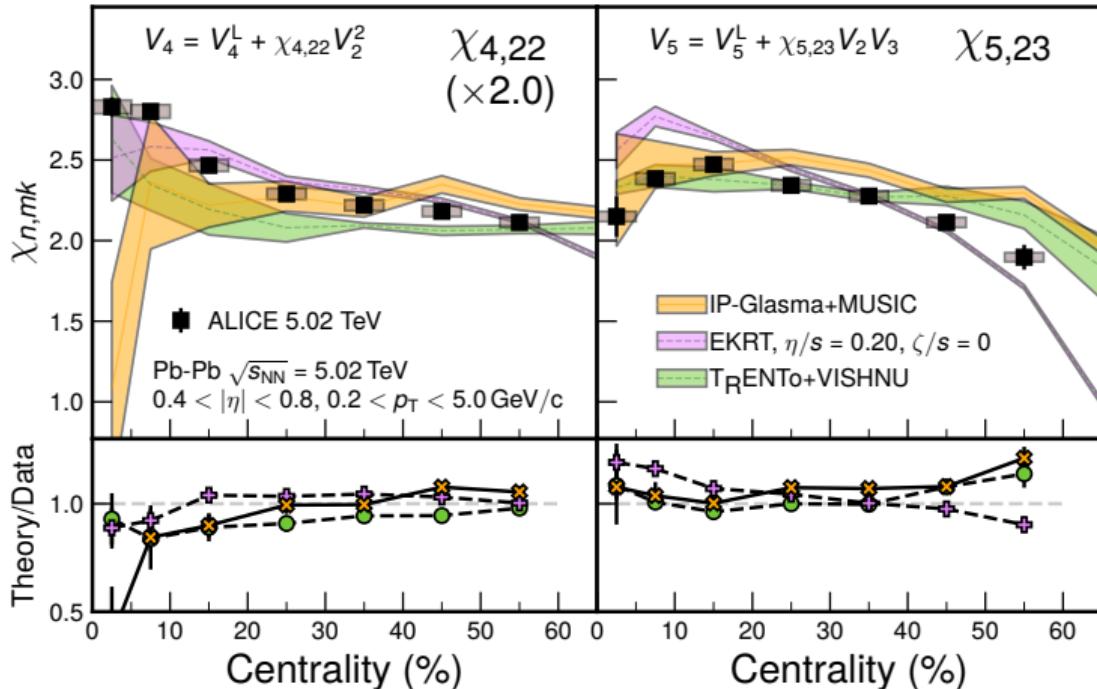
- Assessing the temperature dependence of  $\eta/s(T)$

ALICE, Phys. Rev. Lett. 127 (2021) 092302

- $\eta/s(T)$  and accessing  $\zeta/s(T)$

- Very challenging measurements because of their required high precisions (i.e  $10^{-6}$   $SC(m,n)$ ,  $10^{-12}$  for  $SC(k,l,m)$ ) and difficulties in correcting experimental biases.
- Symmetric Cumulants (Standard Candle)

# IMPROVING RESULTS WITH HIGHER HARMONICS AND MORE PRECISION - NON-LINEAR FLOW MODES



- Higher order  $v_n$ 's ( $n>3$ ) were studied → non-linear dependence on lower orders
- Characterised by the non-linear flow mode coefficients,  $\chi_{n,mk}$
- Better sensitivity to  $\eta/s(T)$ .

## OUR ARSENAL OF OBSERVABLES - STOCHASTIC APPROACH

- Together various flow observables cover the sensitivity for all components of transport properties.

Name	Symbol	Measure	Sensitivity-stochastic approach
Flow coefficients	$v_n$	System expansion and anisotropy of the flow	Average $\langle \eta/s \rangle$ and $\zeta/s(T)$ peak
(Normalized) Symmetric cumulants	$(N)SC(k, l, m)$	Correlation between magnitudes of flow harmonics	$\eta/s(T)$ temperature dependence
Linear and non-linear contributions	$v_{n,L}, v_{n,mk}$	Magnitude of the linear and non-linear contributions	$\eta/s(T)$ and initial conditions, not used
Non-linear flow mode coefficients	$\chi_{n,mk}$	Quantification of the non-linear response	$\eta/s(T)$ at the freeze-out
Symmetry-plane correlations	$\rho_{n,mk}$	Correlations between the directions of flow harmonics	$\eta/s(T)$

Thanks to excellent ALICE papers over years:

- Phys.Rev.Lett. 117 (2016) 182301, Phys.Lett. B773 (2017) 68, Phys.Rev. C 97 (2018) 024906, JHEP05 (2020) 085, Phys.Lett. B818 (2021) 136354, Phys.Rev.Lett. 127 (2021) 092302 - [flow](#)
- Phys.Rev.Lett. 106 (2011) 032301, Phys.Rev.C 88 (2013) 044910, Phys.Lett. B772 (2017) 567-577, Phys.Rev.C 101, 044907 (2020) - [N<sub>ch</sub>](#) and [⟨p<sub>T</sub>⟩](#)

# OUR ARSENAL OF OBSERVABLES

## Duke (2019)

**2.76 TeV**

- PID<sup>1</sup> mult. and  $N_{\text{ch}}$
- Transverse energy  $E_T$
- $\delta p_T / \langle p_T \rangle$
- $v_2$  to  $v_4$

**5.02 TeV**

- $N_{\text{ch}}$
- $v_2$  to  $v_4$

## [1] Jyvaskyla (2021)

**5.02 TeV**

- PID<sup>2</sup> mult. and  $N_{\text{ch}}$
- $v_2$  to  $v_7$
- NSC(3,2) to NSC(4,3)
- $\chi_{4,22}$  to  $\chi_{6,mk}$

## [2] Jyvaskyla (2022)

**2.76 TeV**

- $N_{\text{ch}}$
- NSC(3,2), NSC(4,2)
- NSC(2,3,4), NSC(2,3,5)
- PID<sup>1</sup>  $\langle p_T \rangle$
- $v_2$  to  $v_4$
- $\chi_{4,22}$  to  $\chi_{6,mk}$
- $\rho_{4,22}$  to  $\rho_{6,mk}$

**5.02 TeV**

- PID<sup>2</sup> mult. and  $N_{\text{ch}}$
- NSC(3,2) to NSC(4,3)
- PID  $\langle p_T \rangle$
- $v_2$  to  $v_7$
- $\chi_{4,22}$  to  $\chi_{6,mk}$
- $\rho_{4,22}$  to  $\rho_{6,mk}$

All reference data based on ALICE measurements.

Red: Missing from other group (Duke etc)

Blue: New since our PRC.

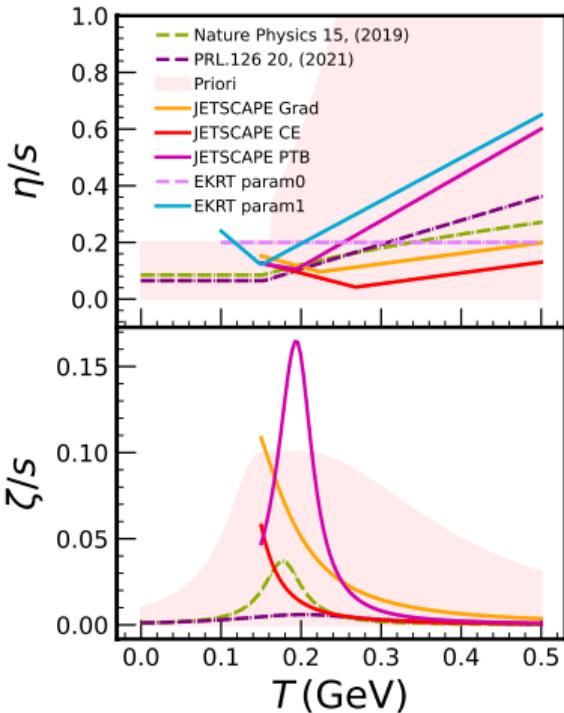
Orange: Not used in our studies.

$^1\pi^\pm, K^\pm$  and  $p^\pm$   
 $^2p^\pm$

[1]. J.E. Parkkila, A. Onnerstad, D.J. Kim, PRC **104** (2021) 054904

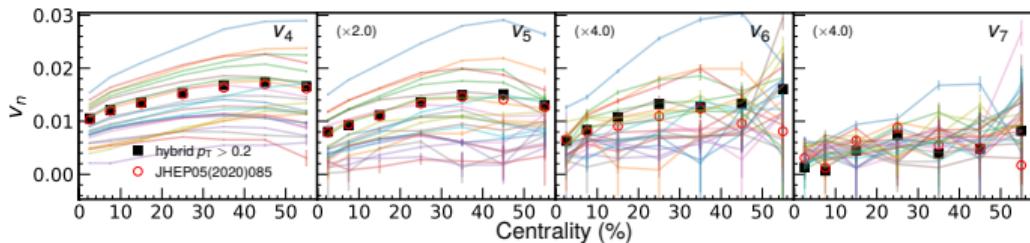
[2]. J.E. Parkkila, A. Onnerstad, S.F. Taghavi, C. Mordasini, A. Bilandzic, D.J. Kim, arXiv:2111.08145

# ANALYSIS STEPS AND PRIORI

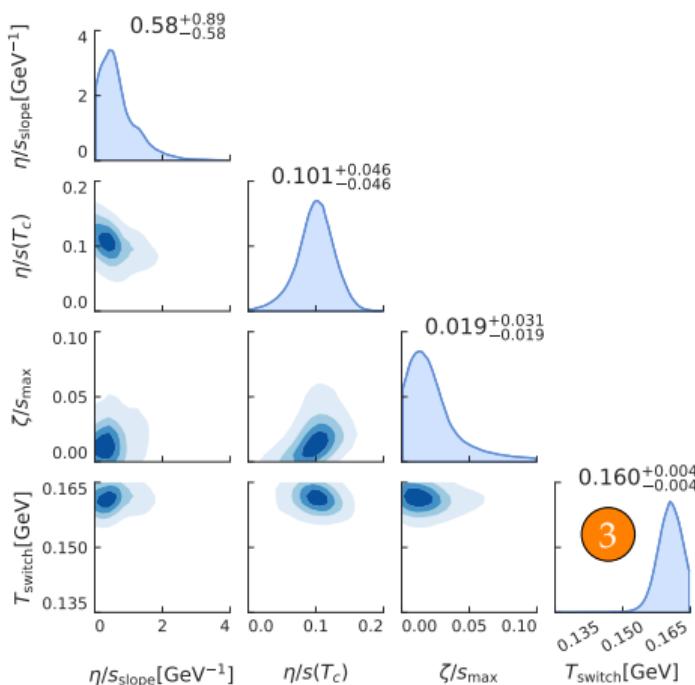
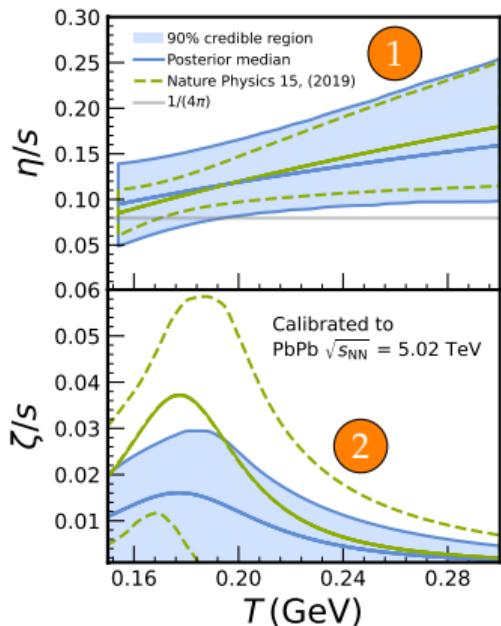


PRC 104 (2021) 054904

- 1 Choose prior parameter range based on results from 2019
- 2 Run hydro TRRENTO+VISH(2+1D)+UrQMD for 500 parameterizations, 3-5 million events ( $\times 100$  previous).
- 3 Calculate observables using our experimental framework
- 4 Train emulator and setup/run Bayesian analysis



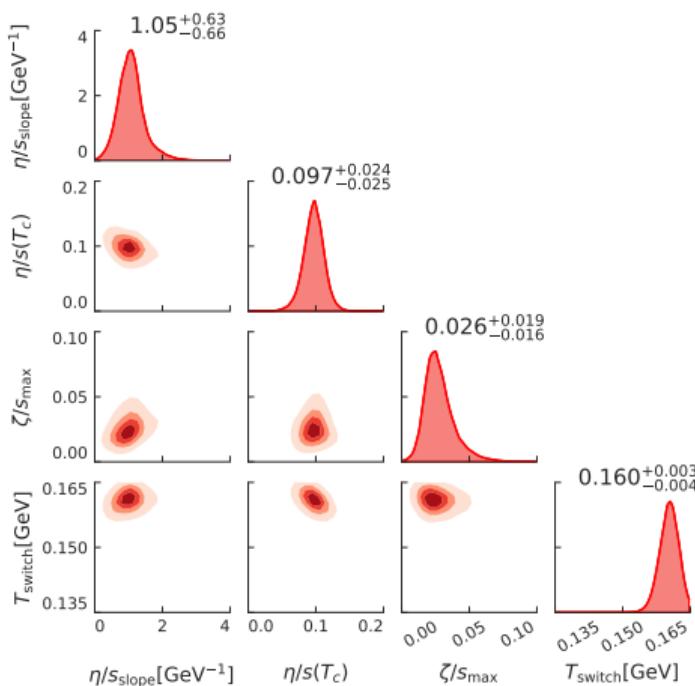
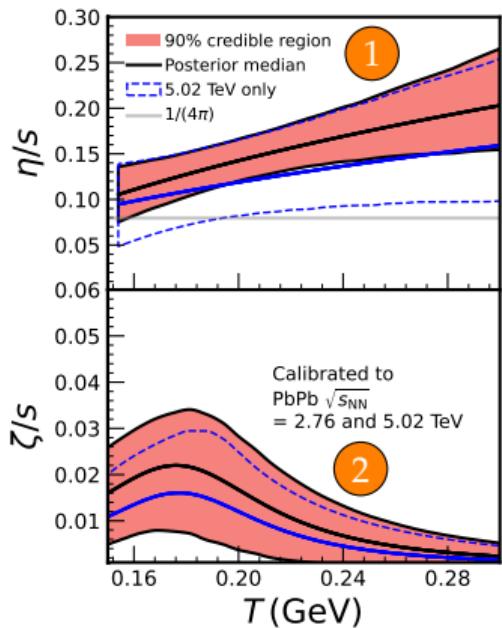
# RESULTS: JYVASKYLA (2021) (5.02 TeV ONLY) – HIGHER STATISTICS RUN



- 1 Similar  $\eta/s(T)$  to Duke (2019)
- 2 Lower  $\zeta/s(T)$  – much lower to previous calculations
- 3 Higher switching temperature  $T_{\text{switch}}$  (vs. Duke 152 MeV)

- Additional observables have reduced  $\zeta/s(T)$ .
- However, one collision energy only limits the potential of the additions.

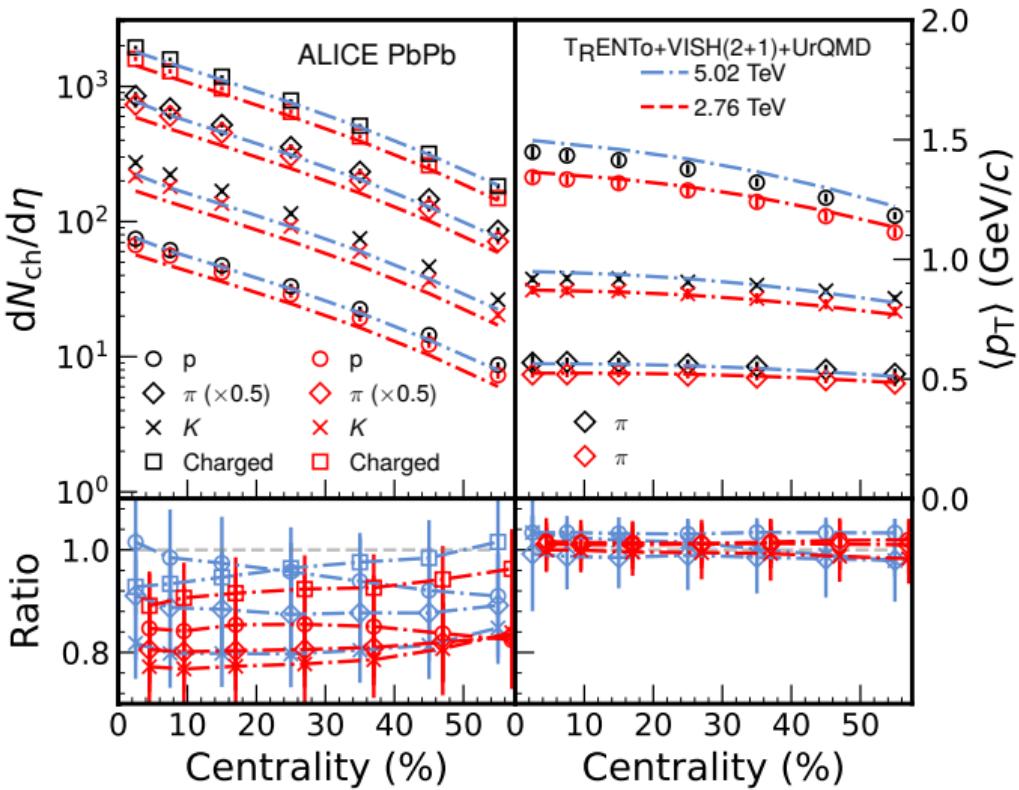
# RESULT: JYVASKYLA (2022) – COMBINED COLLISION ENERGY ANALYSIS (2.76 + 5.02 TeV)



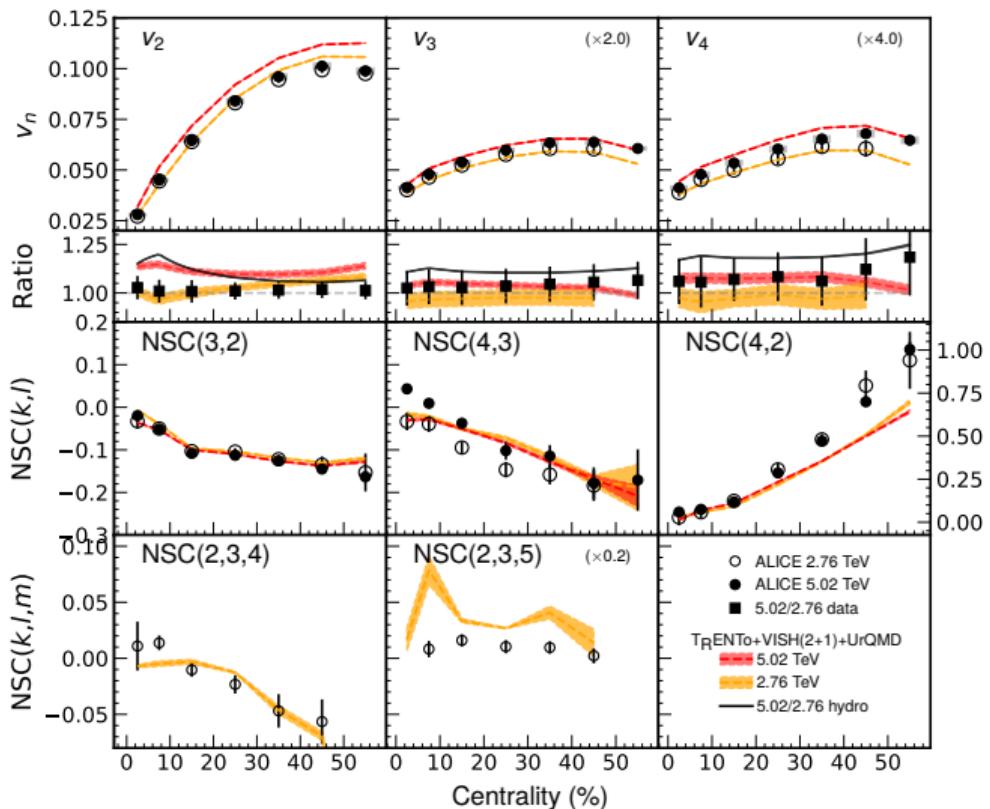
- ➊ Significantly improved  $\eta/s(T)$  uncertainty
- ➋ Non-zero  $\zeta/s(T)$
- ➌ Overall better convergence for parameter components

- Together with two collision energies and added observables, the uncertainty has reduced!

# PID MULTIPLICITY AND $\langle p_T \rangle$

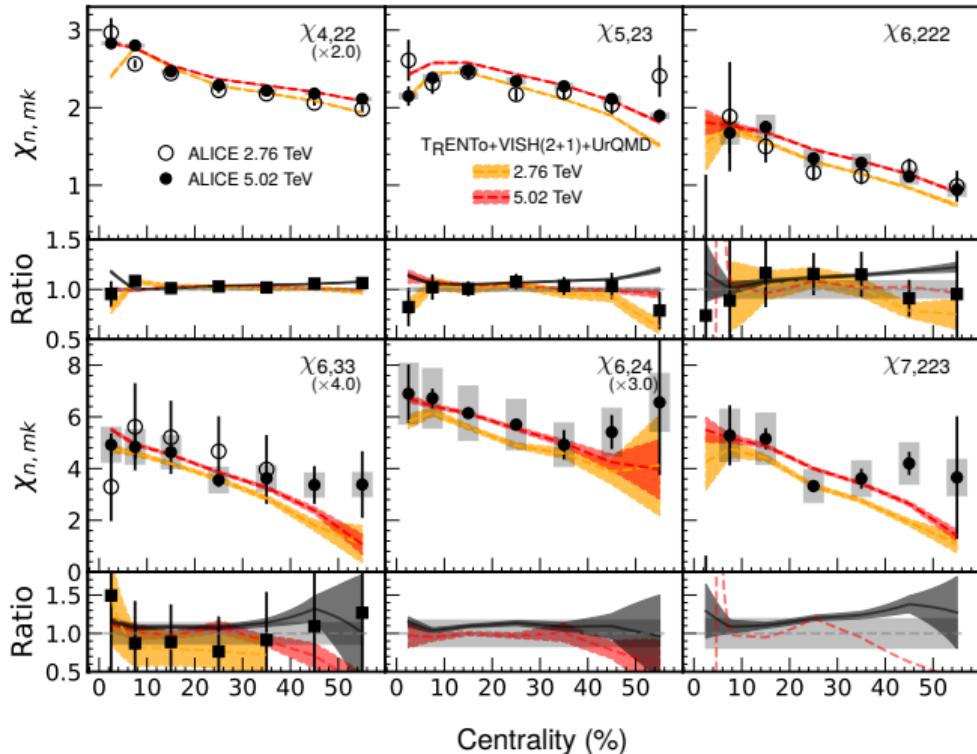


- Agreement for charged particle yield in 2.76 TeV and 5.02 TeV
- 10–20% difference for PID multiplicity
- Qualitative agreement for  $\langle p_T \rangle_{\pi,K}$

*v<sub>n</sub>* AND SYMMETRIC CUMULANTS

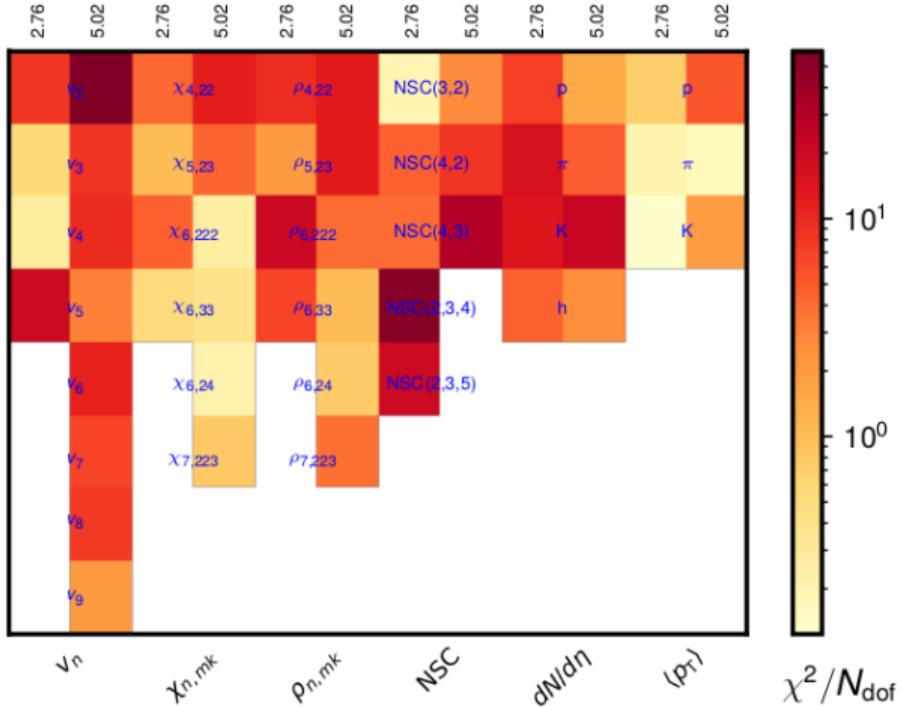
- Good agreement for 2.76 TeV  $v_n$ , overestimated  $v_2$  for 5.02 TeV by  $\sim 10\%$
- Magnitude and centrality dependence of NSC well captured. Further improved estimate for NSC(4,2).
- Good agreement for NSC(2,3,4). NSC(2,3,5) overestimated.

# NON-LINEAR FLOW COEFFICIENTS



- Qualitative agreement in both beam energies for all mode coupling coefficients.
- See arXiv:2111.08145 for all graphs.

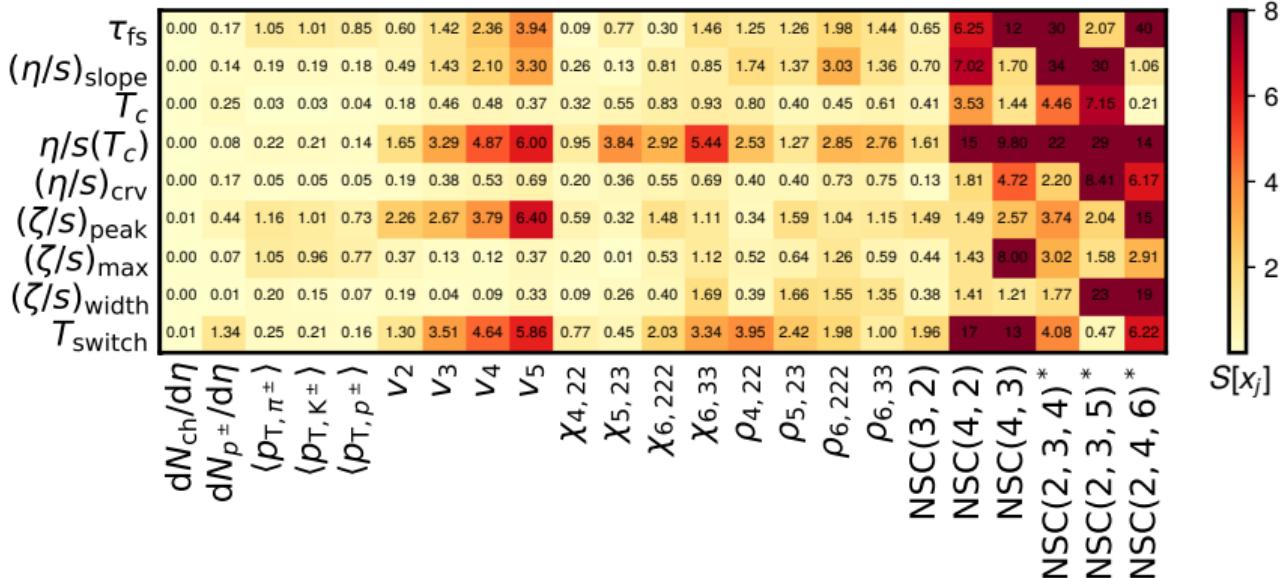
# REMAINING CONCERNs?



- Higher energy description worse for all observables except for:
  - $v_5$
  - $\chi_{6,222}$
  - charged particle multiplicity
- Concerns
  - overestimated  $v_n$  for 5.02 TeV by  $\sim 10\%$
  - still underestimated NSC(4,2)
  - overestimated NSC(2,3,5)
  - PID multiplicity (especially  $\pi^\pm$ )
- Why?
  - Reduction of the uncertainties is understood?

## SENSITIVITY OF THE OBSERVABLES TO PARAMETERS

Sensitivity of the observables:  $S[x_j] = \Delta/\delta.$ , where  $\Delta = \frac{|\hat{O}(\vec{x}') - \hat{O}(\vec{x})|}{|\hat{O}(\vec{x})|}$ .



- NSC( $m,n$ ) and NSC( $k,l,m$ ) are among the most sensitive observables followed by  $v_n$  and  $\chi_{n,mk}$ .
- The precision measurements of observables, reflecting mostly non-linear responses, are crucial.

## SUMMARY

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### Success:

- Higher harmonic orders and non-linear flow observables → better constraints.
- Improved the overall uncertainty by  $\times 2$  by combining two beam energy data.
- As a bonus, sensitivities of the observables are now quantified  
→ precision measurements of observables, reflecting non-linear hydrodynamic responses.

### Challenges:

- 10% difference for  $v_2$  (5.02 TeV)
- NSC description improved except for NSC(4,2)
- Remaining discrepancy for PID multiplicity (especially  $\pi^\pm$ )
- Improving the initial state model, with dynamical collision model or subnucleon structure à la IP-Glasma, might help us to improve the results.

# OUTLOOK

## Experiments

- RHIC data (AuAu collisions) - Energy and system size dependence
- LHC pPb and pp data - System size dependence
- Use new observables
  - Higher order ( $n > 5$ ) Symmetric cumulants
  - Improved Symmetric Plane Correlation (SPC) : independent from flow magnitude correlations
  - Asymmetric Cumulants (AC)

## Theory

- Improving the initial conditions with
  - EKRT
  - IP+Glasma
- Testing hydro limit of small systems?
- Role of the small system.

## THANKS

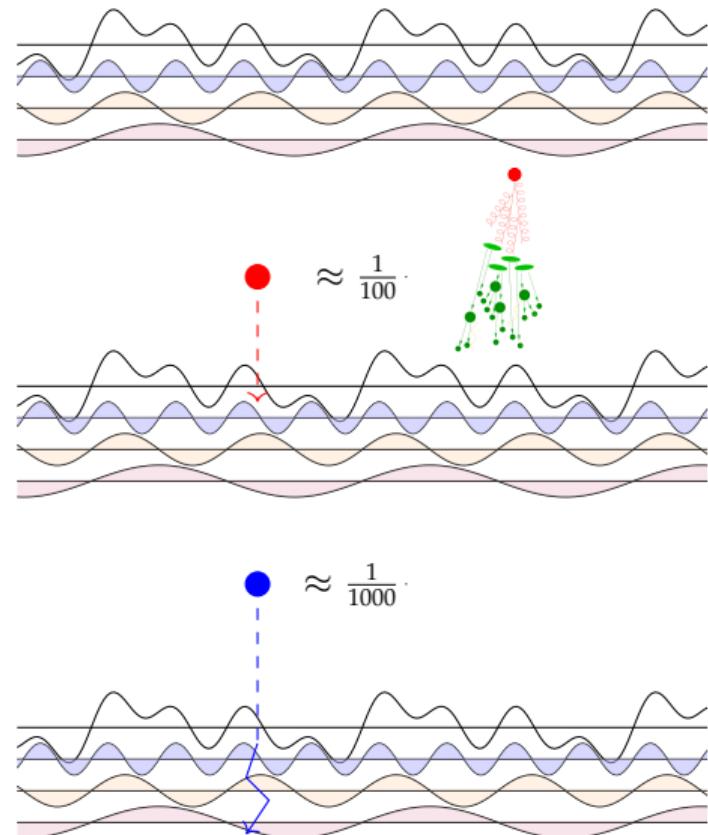
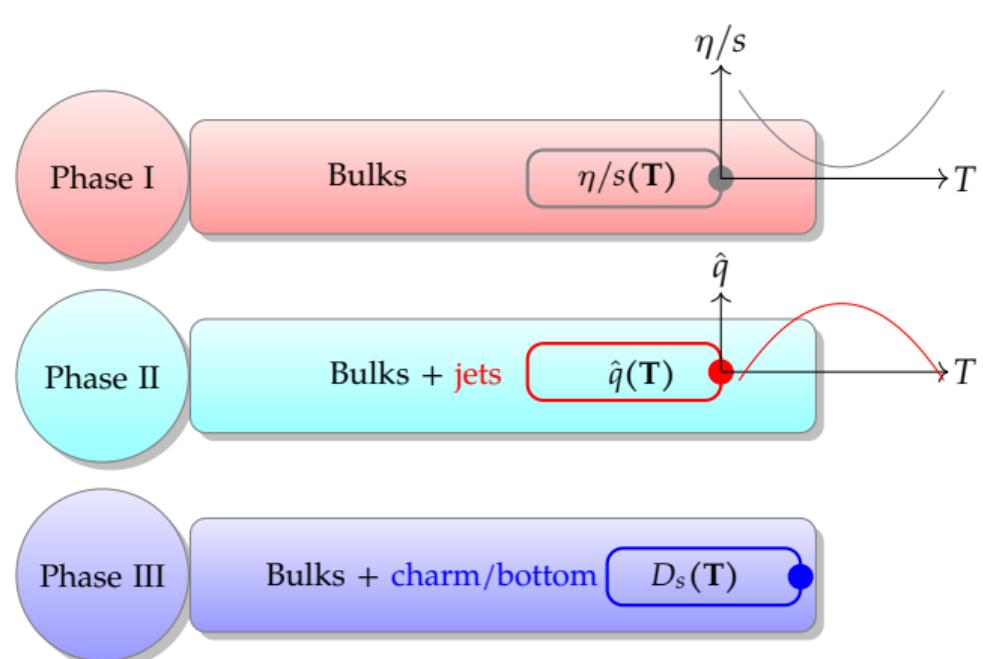
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Thank you for your attention!

Acknowledgments:

- CSC for providing the ~24 million CPU hours
- Harri Niemi, Kari Eskola, Jonah E. Bernhard, J. Scott Moreland and Steffen A. Bass for their useful comments

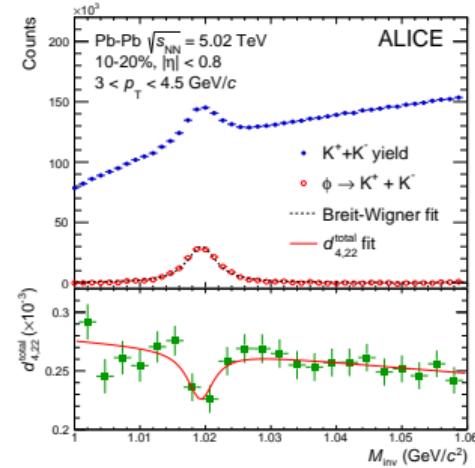
# TRANSPORT PROPERTIES



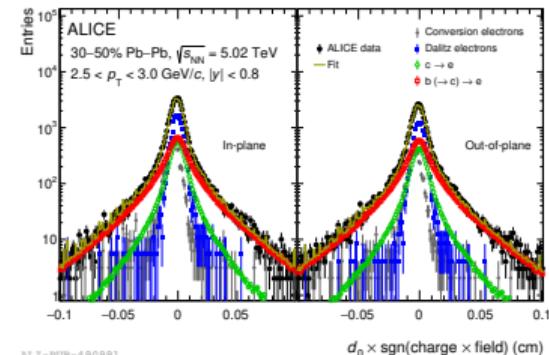
# HEAVY FLAVOUR OBSRVABLES

$$\begin{aligned}
 \text{SC}(m, n) &\equiv \langle\langle \cos(m\varphi_1 + n\varphi_2 - m\varphi_3 - n\varphi_4) \rangle\rangle_c \\
 &= \langle\langle \cos(m\varphi_1 + n\varphi_2 - m\varphi_3 - n\varphi_4) \rangle\rangle \\
 &- \langle\langle \cos[m(\varphi_1 - \varphi_3)] \rangle\rangle \langle\langle \cos[n(\varphi_2 - \varphi_4)] \rangle\rangle \\
 &= \left\langle v_m^2 v_n^2 \right\rangle - \left\langle v_m^2 \right\rangle \left\langle v_n^2 \right\rangle
 \end{aligned}$$

- Replace  $\varphi_1$  or  $\varphi_2$  with HF candidates
- Invariant mass or DCA approach
- compared to all tracks

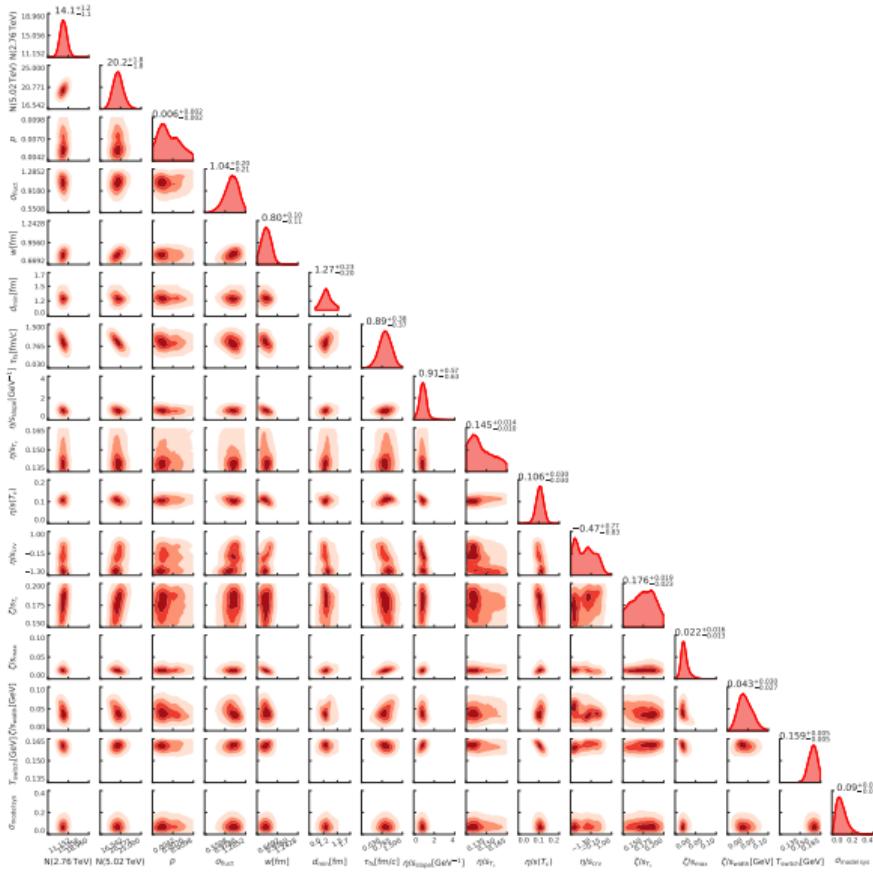


ALICE, JHEP06(2020) 147

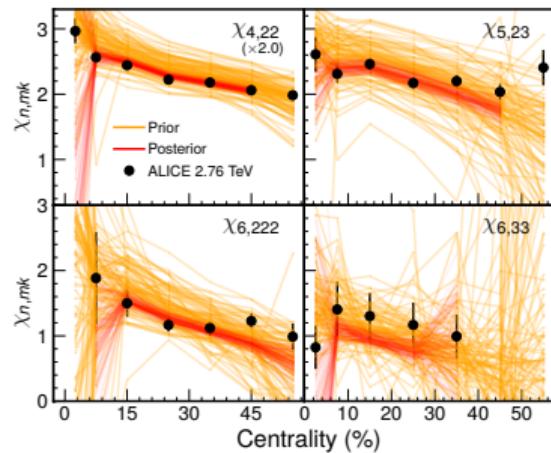
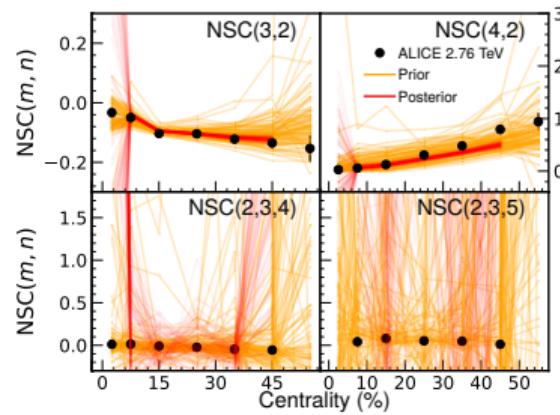
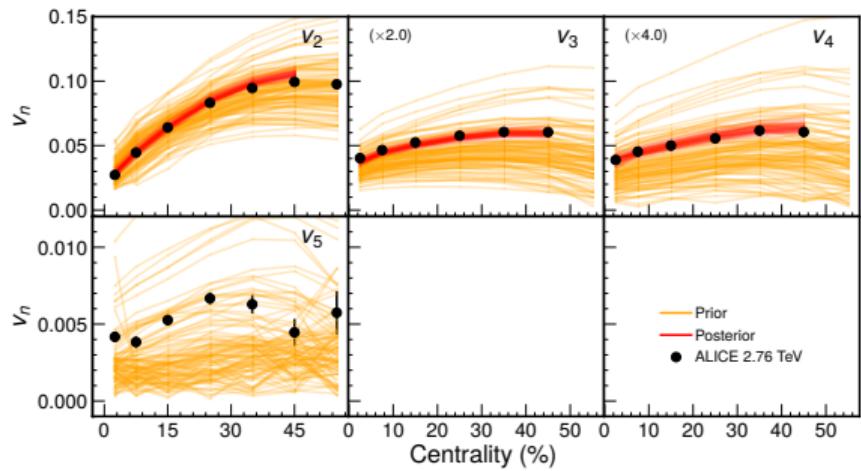


ALICE, Phys. Rev. Lett. 126 (2021) 162001

## MARGINAL AND JOINT MARGINAL PARTS OF THE POSTER DISTRIBUTION



# DESIGN PARAMETRIZATIONS



# CONCERNS ON THE HIGHER ORDER HARMONICS

