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Systematic studies of correlations between different order flow harmonics in Pb–Pb collisions at $\sqrt{s_{\text{NN}}} = 2.76$ TeV

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Abstract

The correlations between event-by-event fluctuations of anisotropic flow harmonic amplitudes have been measured in Pb–Pb collisions at $\sqrt{s_{\text{NN}}} = 2.76$ TeV with the ALICE detector at the Large Hadron Collider. The results are reported in terms of multiparticle correlation observables dubbed Symmetric Cumulants. These observables are robust against biases originating from nonflow effects. The centrality dependence of correlations between the higher order harmonics (the quadrangular v_4 and pentagonal v_5 flow) and the lower order harmonics (the elliptic v_2 and triangular v_3 flow) is presented. The transverse momentum dependence of correlations between v_3 and v_2 and between v_4 and v_2 is also reported. The results are compared to calculations from viscous hydrodynamics and A Multi-Phase Transport (AMPT) model calculations. The comparisons to viscous hydrodynamic models demonstrate that the different order harmonic correlations respond differently to the initial conditions and the temperature dependence of the ratio of shear viscosity to entropy density (η/s). A small average value of η/s is favored independent of the specific choice of initial conditions in the models. The calculations with the AMPT initial conditions yield results closest to the measurements. Correlations between the magnitudes of v_2 , v_3 and v_4 show moderate p_T dependence in mid-central collisions. This might be an indication of possible viscous corrections to the equilibrium distribution at hadronic freeze-out, which might help to understand the possible contribution of bulk viscosity in the hadronic phase of the system. Together with existing measurements of individual flow harmonics, the presented results provide further constraints on initial conditions and the transport properties of the system produced in heavy-ion collisions.

1 Introduction

The main emphasis of the ultra-relativistic heavy-ion collision programs at the Relativistic Heavy Ion Collider (RHIC) and the Large Hadron Collider (LHC) is to study the deconfined phase of strongly interacting QCD matter, the Quark-Gluon Plasma (QGP). The matter produced in a heavy-ion collision exhibits strong collective radial expansion [1, 2]. Difference in pressure gradients and the interactions among matter constituents produced in the spatially anisotropic overlap region of the two colliding nuclei result in anisotropic transverse flow in the momentum space. The large elliptic flow discovered at RHIC energies [3–7] is also observed at LHC energies [8–18]. The measurements are well described by calculations utilizing viscous hydrodynamics [19–24]. These calculations also demonstrated that the shear viscosity to the entropy density ratio (η/s) of the QGP in heavy-ion collisions at RHIC and LHC energies is close to a universal lower bound $1/4\pi$ [25].

The temperature dependence of η/s has some generic features typical to the most known fluids. This ratio reaches its minimum value close to the phase transition region [25, 26]. It was shown, using kinetic theory and quantum mechanical considerations [27], that $\eta/s \sim 0.1$ would be the correct order of magnitude for the lowest possible shear viscosity to entropy density ratio value found in nature. Later it was demonstrated that an exact lower bound $(\eta/s)_{\min} = 1/4\pi \approx 0.08$ can be conjectured using AdS/CFT correspondence [25]. Hydrodynamical simulations constrained by data support the view that η/s of the QGP is close to that limit [23]. It is argued that such a low value might imply that thermodynamic trajectories for the expanding matter would lie close to the quantum chromodynamics (QCD) critical end point, which is another subject of intensive experimental study [26, 28].

Anisotropic flow [29] is quantified with n^{th} -order flow harmonics v_n and corresponding symmetry plane angles Ψ_n in a Fourier decomposition of the particle azimuthal distribution in the plane transverse to the beam direction [30, 31]:

$$E \frac{d^3N}{dp^3} = \frac{1}{2\pi} \frac{d^2N}{p_T dp_T d\eta} \left\{ 1 + 2 \sum_{n=1}^{\infty} v_n(p_T, \eta) \cos[n(\varphi - \Psi_n)] \right\}, \quad (1)$$

where E , p , p_T , φ and η are the particle's energy, momentum, transverse momentum, azimuthal angle and pseudorapidity, respectively, and Ψ_n is the azimuthal angle of the symmetry plane of the n^{th} -order harmonic. Harmonic v_n can be calculated as $v_n = \langle \cos[n(\varphi - \Psi_n)] \rangle$, where the angular brackets denote an average over all particles in all events. The anisotropic flow in heavy-ion collisions is typically understood as the hydrodynamic response of the produced matter to spatial deformations of the initial energy density profile [32]. This profile fluctuates event-by-event due to fluctuating positions of the constituents inside the colliding nuclei, which implies that v_n also fluctuates [33, 34]. The recognition of the importance of flow fluctuations led to the discovery of triangular and higher flow harmonics [9, 35] as well as to the correlations between different v_n harmonics [36, 37]. The higher order harmonics are expected to be sensitive to fluctuations in the initial conditions and to the magnitude of η/s [38, 39], while v_n correlations have the potential to discriminate between these two respective contributions [36].

Difficulties in extracting η/s in heavy-ion collisions can be attributed mostly to the fact that it strongly depends on the specific choice of the initial conditions in the models used for comparison [19, 39, 40]. Viscous effects reduce the magnitude of the anisotropic flow. Furthermore, the magnitude of η/s used in hydrodynamic calculations should be considered as an average over the temperature evolution of the expanding fireball as it is known that η/s depends on temperature. In addition, part of the anisotropic flow can also originate from the hadronic phase [41–43]. Therefore, both the temperature dependence of η/s and the relative contributions from the partonic and hadronic phases should be understood better to quantify the η/s of the QGP.

An important input to the hydrodynamic model simulations is the initial distribution of energy density

in the transverse plane (the initial density profile), which is usually estimated from the probability distribution of nucleons in the incoming nuclei. This initial energy density profile can be quantified by calculating the distribution of the spatial eccentricities ε_n [35],

$$\varepsilon_n e^{in\Phi_n} = -\{r^n e^{in\phi}\} / \{r^n\}, \quad (2)$$

where the curly brackets denote the average over the transverse plane, i.e. $\{\dots\} = \int dx dy e(x, y, \tau_0) (\dots)$, r is the distance to the system's center of mass, $e(x, y, \tau_0)$ is the energy density at the initial time τ_0 , and Φ_n is the participant plane angle (see Refs. [44, 45]). There is experimental and theoretical evidence [9, 35, 46] that the lower order harmonics, v_2 and v_3 , to a good approximation, are linearly proportional to the deformations in the initial energy density in the transverse plane (e.g. $v_n \propto \varepsilon_n$ for $n = 2$ or 3). Higher order ($n > 3$) flow harmonics can arise from initial anisotropies in the same harmonic [35, 44, 47, 48] (linear response) or can be induced by lower order harmonics [49, 50] (nonlinear response). For instance, v_4 can develop both as a linear response to ε_4 and/or as a nonlinear response to ε_2^2 [51]. Therefore, the higher harmonics ($n > 3$) can be understood as superpositions of linear and nonlinear responses, through which they are correlated with lower order harmonics [47, 48, 50, 52, 53]. When the order of the harmonic is large, the nonlinear response contribution in viscous hydrodynamics is dominant and increases in more peripheral collisions [50, 52]. The magnitudes of the viscous corrections as a function of p_T for v_4 and v_5 are sensitive to the ansatz used for the viscous distribution function, a correction for the equilibrium distribution at hadronic freeze-out [52, 54]. Hence, studies of the correlations between higher order ($n > 3$) and lower order (v_2 or v_3) harmonics and their p_T dependence can help to understand the viscous correction to the momentum distribution at hadronic freeze-out which is among the least understood parts of hydrodynamic calculations [45, 52, 55, 56].

The first results for new multiparticle observables which quantify the relationship between event-by-event fluctuations of two different flow harmonics, the *Symmetric Cumulants* (SC), were recently reported by the ALICE Collaboration [57]. The new observables are particularly robust against few-particle nonflow correlations [8] and they provide independent, complementary information to recently analyzed symmetry plane correlators [37]. It was demonstrated that they are sensitive to the temperature dependence of η/s of the expanding medium and therefore simultaneous descriptions of correlations between different order harmonics would constrain both the initial conditions and the medium properties [57, 58]. In this article, we have extended the analysis of SC observables to higher order harmonics (up to 5th order) as well as to the measurement of the p_T dependence of correlations for the lower order harmonics (v_3 - v_2 and v_4 - v_2). We also present a systematic comparison to hydrodynamic and AMPT model calculations. In Sec. 2 we summarize our findings from the previous work [57] and present the analysis methods. The experimental setup and measurements are described in Sec. 3. The sources of systematic uncertainties are explained in Sec. 4. The results of the measurements are presented in Sec. 5. In Sec. 6 we present comparisons to model calculations. Finally, Sec. 7 summarizes our new results.

2 Experimental Observables

Existing measurements for anisotropic flow observables provide an estimate of the average value of η/s of the QGP, both at RHIC and LHC energies. What remains uncertain is how the η/s of the QGP depends on temperature (T). The temperature dependence of η/s of the QGP was discussed in [28]. The effects on hadron spectra and elliptic flow were studied in [59] for different parameterizations of $\eta/s(T)$. A more systematic study with event-by-event EKRT+viscous hydrodynamic calculations was recently initiated in [45], where the first (and only rather qualitative) possibilities were investigated (see Fig. 1 therein). The emerging picture is that the study of individual flow harmonics v_n alone is unlikely to reveal the details of the temperature dependence of η/s . It was already demonstrated in [45] that different $\eta/s(T)$ parameterizations can lead to the same centrality dependence of individual flow harmonics. In Ref. [36] new flow observables were introduced which quantify the degree of correlation between amplitudes of two different harmonics v_m and v_n . These new observables have the potential to discriminate

between the contributions to anisotropic flow development from initial conditions and from the transport properties of the QGP [36]. Therefore their measurement would provide experimental constraints on theoretical parameters used to describe the individual stages of the heavy-ion system evolution. In addition, it turned out that correlations of different flow harmonics are sensitive to the temperature dependence of η/s [57], to which individual flow harmonics are weakly sensitive [45].

For reasons discussed in [57, 60], the correlations between different flow harmonics cannot be studied experimentally with the set of observables introduced in [36]. Based on [60], new flow observables obtained from multiparticle correlations, *Symmetric Cumulants* (SC), were introduced.

The SC observables are defined as:

$$\begin{aligned} \text{SC}(m, n) &\equiv \langle \langle \cos(m\varphi_1 + n\varphi_2 - m\varphi_3 - n\varphi_4) \rangle \rangle_c = \langle \langle \cos(m\varphi_1 + n\varphi_2 - m\varphi_3 - n\varphi_4) \rangle \rangle \\ &\quad - \langle \langle \cos[m(\varphi_1 - \varphi_2)] \rangle \rangle \langle \langle \cos[n(\varphi_1 - \varphi_2)] \rangle \rangle \\ &= \langle v_m^2 v_n^2 \rangle - \langle v_m^2 \rangle \langle v_n^2 \rangle, \end{aligned} \quad (3)$$

with the condition $m \neq n$ for two positive integers m and n (for details see Sec. IV C in [60]). In this article $\text{SC}(m, n)$ normalized by the product $\langle v_m^2 \rangle \langle v_n^2 \rangle$ [57, 61] is denoted by $\text{NSC}(m, n)$:

$$\text{NSC}(m, n) \equiv \frac{\text{SC}(m, n)}{\langle v_m^2 \rangle \langle v_n^2 \rangle}. \quad (4)$$

Normalized symmetric cumulants reflect only the strength of the correlation between v_m and v_n , while $\text{SC}(m, n)$ has contributions from both the correlations between the two different flow harmonics and the individual harmonics. In Eq. (4) the products in the denominator are obtained from two-particle correlations using a pseudorapidity gap of $|\Delta\eta| > 1.0$ which suppresses biases from few-particle nonflow correlations. For the two two-particle correlations which appear in the definition of $\text{SC}(m, n)$ in Eq. (3) the pseudorapidity gap is not needed, since nonflow is suppressed by construction in this observable. This was verified by HIJING model simulations in [57].

The ALICE measurements [57] have revealed that fluctuations of v_2 and v_3 are anti-correlated, while fluctuations of v_2 and v_4 are correlated for all centralities [57]. It was found that the details of the centrality dependence differ in the fluctuation-dominated (most central) and the geometry-dominated (mid-central) regimes [57]. The observed centrality dependence of $\text{SC}(4, 2)$ cannot be captured by models with constant η/s , indicating that the temperature dependence of η/s plays an important role. These results were also used to discriminate between different parameterizations of initial conditions. It was demonstrated that in the fluctuation-dominated regime (central collisions), MC-Glauber initial conditions with binary collision weights are favored over wounded nucleon weights [57]. The first theoretical studies of SC observables can be found in [58, 61–65].

3 Data Analysis

The data samples of Pb–Pb collisions at the centre-of-mass energy $\sqrt{s_{\text{NN}}} = 2.76$ TeV analyzed in this article were recorded by ALICE during the 2010 heavy-ion run of the LHC. Detailed descriptions of the ALICE detector can be found in [66–68]. The Time Projection Chamber (TPC) was used to reconstruct charged particle tracks and measure their momenta with full azimuthal coverage in the pseudorapidity range $|\eta| < 0.8$. Two scintillator arrays (V0A and V0C) which cover the pseudorapidity ranges $-3.7 < \eta < -1.7$ and $2.8 < \eta < 5.1$ were used for triggering and the determination of centrality [69]. The trigger conditions and the event selection criteria are identical to those described in [8, 69]. Approximately 10^7 minimum-bias Pb–Pb events with a reconstructed primary vertex within ± 10 cm from the nominal interaction point along the beam direction are selected. Only charged particles reconstructed in the TPC in $|\eta| < 0.8$ and $0.2 < p_T < 5$ GeV/c were included in the analysis. The charged track quality cuts

described in [8] were applied to minimize contamination from secondary charged particles and fake tracks. The track reconstruction efficiency and contamination were estimated from HIJING Monte Carlo simulations [70] combined with a GEANT3 [71] detector model and were found to be independent of the collision centrality. The reconstruction efficiency increases with transverse momenta from 70% to 80% for particles with $0.2 < p_T < 1$ GeV/c and remains constant at $(80 \pm 5)\%$ for $p_T > 1$ GeV/c. The estimated contamination by secondary charged particles from weak decays and photon conversions is less than 6% at $p_T = 0.2$ GeV/c and falls below 1% for $p_T > 1$ GeV/c. The p_T cut-off of 0.2 GeV/c reduces event-by-event biases due to small reconstruction efficiency at lower p_T , while the high p_T cut-off of 5 GeV/c reduces the effects of jets on the measured correlations. Reconstructed TPC tracks constrained to vertex are required to have at least 70 space points (out of a maximum of 159). Only tracks with a transverse distance of closest approach to the primary vertex less than 3 mm, both in the longitudinal and transverse directions, are accepted. This reduces the contamination from secondary tracks produced in the detector material, particles from weak decays, etc. Tracks with kinks (i.e. tracks that appear to change direction due to multiple scattering or K^\pm decays) were rejected.

4 Systematic Uncertainties

The systematic uncertainties are estimated by varying the event and track selection criteria. All systematic checks described here are performed independently. The $SC(m,n)$ values resulting from each variation are compared to ones from the default event and track selection described in the previous section, and differences are taken as the systematic uncertainty due to each individual source. The contributions from different sources were added in quadrature to obtain the total systematic uncertainty.

The event centrality was determined by the V0 detectors [72] with better than 2% resolution for the whole centrality range analyzed. The systematic uncertainty from the centrality determination was evaluated by using the TPC and Silicon Pixel Detector (SPD) [73] detectors instead of the V0 detectors. The systematic uncertainty on the symmetric cumulants which arises from the centrality uncertainty is about 3% both for $SC(5,2)$ and $SC(4,3)$, and 8% for $SC(5,3)$. As described in Sec. 3, the reconstructed vertex position along the beam axis (z -vertex) is required to be located within 10 cm of the nominal interaction to ensure uniform detector acceptance for tracks within $|\eta| < 0.8$. The systematic uncertainty from the z -vertex cut was estimated by reducing the z -vertex range to 8 cm and was found to be less than 3%.

The analyzed events were recorded with two settings of the magnet field polarity and the resulting data sets have almost equal numbers of events. Events with both magnet field polarities were used in the default analysis, and the systematic uncertainties were evaluated from the variation between each of the two magnetic field settings. The uncertainty due to the p_T dependence of the track reconstruction efficiency was also taken into account. Magnetic field polarity variation and reconstruction efficiency effects contribute less than 2% to the systematic uncertainty.

The systematic uncertainty due to the track reconstruction procedure was estimated from comparisons between results for the so-called standalone TPC tracks with the same parameters as described in Sec. 3, and tracks from a combination of the TPC and the Inner Tracking System (ITS) detectors with tighter selection criteria. To avoid non-uniform azimuthal acceptance due to dead zones in the SPD, and to get the best transverse momentum resolution, a hybrid track selection utilizing SPD hits and/or ITS refit tracks combined with TPC information was used. Then each track reconstruction strategy was evaluated by varying the threshold on parameters used to select the tracks at the reconstruction level. A systematic difference of up to 12% was observed in $SC(m,n)$ from the different track selections. In addition, we applied the like-sign technique to estimate nonflow contributions [8] to $SC(m,n)$. The difference between results obtained by selecting all charged particles and results obtained after either selecting only positively or only negatively charged particles was the largest contribution to the systematic uncertainty and is about 7% for $SC(4,3)$ and 20% for $SC(5,3)$.

Another large contribution to the systematic uncertainty originates from azimuthal non-uniformities in the reconstruction efficiency. In order to estimate its effects, we use the AMPT model (see Sec. 6) which has a uniform distribution in azimuthal angle. Detector inefficiencies were introduced to mimic the non-uniform azimuthal distribution in the data. For the observables $SC(5,2)$, $SC(5,3)$ and $SC(4,3)$ the variation due to non-uniform acceptance is about 9%, 17% and 11%, respectively. Overall, the systematic uncertainties are larger for $SC(5,3)$ and $SC(5,2)$ than for the lower harmonics of $SC(m,n)$. This is because v_n decreases with increasing n and becomes more sensitive to azimuthal modulation due to detector imperfections.

5 Results

The centrality dependence of the higher order harmonic correlations ($SC(4,3)$, $SC(5,2)$ and $SC(5,3)$) are presented in Fig. 1 and compared to the lower order harmonic correlations ($SC(3,2)$ and $SC(4,2)$) which were published in [57]. The correlation between v_3 and v_4 is negative, and similarly for v_3 and v_2 , while the other correlations are all positive, which reveals that v_2 and v_5 as well as v_3 and v_5 are correlated like v_2 and v_4 , while v_3 and v_4 are anti-correlated like v_3 and v_2 .

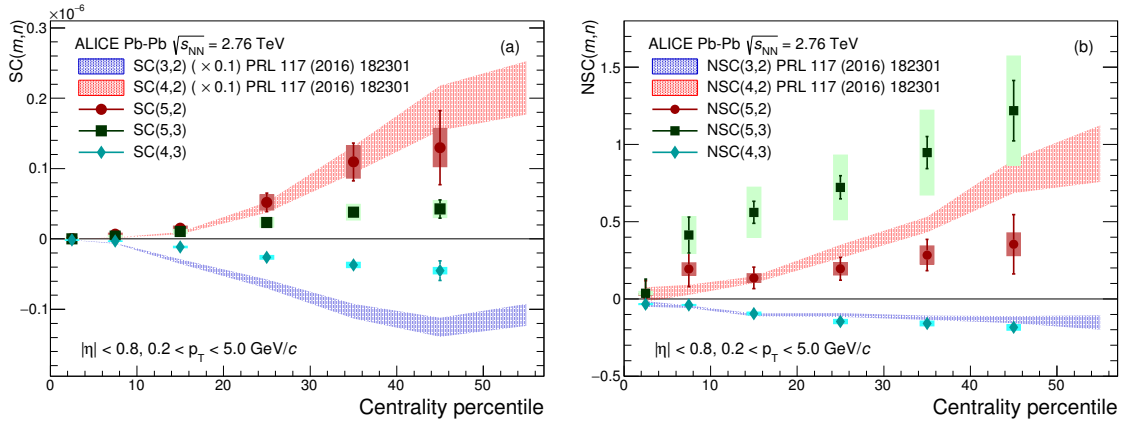


Fig. 1: The centrality dependence of $SC(m,n)$ (a) and $NSC(m,n)$ (b) with flow harmonics for $m = 3-5$ and $n = 2,3$ in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV. The lower order harmonic correlations ($SC(3,2)$, $SC(4,2)$, $NSC(3,2)$ and $NSC(4,2)$) are taken from [57] and shown as bands. The systematic and statistical errors are combined in quadrature for these lower order harmonic correlations. The $SC(4,2)$ and $SC(3,2)$ are downscaled by 0.1. Systematic uncertainties are represented with boxes for higher order harmonic correlations.

The higher order flow harmonic correlations are much smaller compared to the lower order harmonic correlations. In particular $SC(5,2)$ is 10 times smaller than $SC(4,2)$ and $SC(4,3)$ is about 20 times smaller than $SC(3,2)$.

Unlike $SC(m,n)$, the $NSC(m,n)$ results with the higher order flow harmonics show almost the same order of the correlation strength as the lower order flow harmonic correlations $NSC(3,2)$ or $NSC(4,2)$. This demonstrates the advantage of using the normalized SC observables in which the correlation strength between flow harmonics is not hindered by the differences in magnitudes of different flow harmonics. The $NSC(4,3)$ magnitude is comparable to $NSC(3,2)$ and one finds that a hierarchy, $NSC(5,3) > NSC(4,2) > NSC(5,2)$, holds for centrality range 20–50% within the errors as shown in Fig. 1(b). The $SC(5,2)$ magnitude is larger than $SC(5,3)$, but the normalized correlation between v_5 and v_3 is stronger than the normalized correlation between v_5 and v_2 . These results indicate that the lower order harmonic correlations are larger than higher order harmonic correlations, not only because of the correlation strength itself but also because of the strength of the individual flow harmonics.

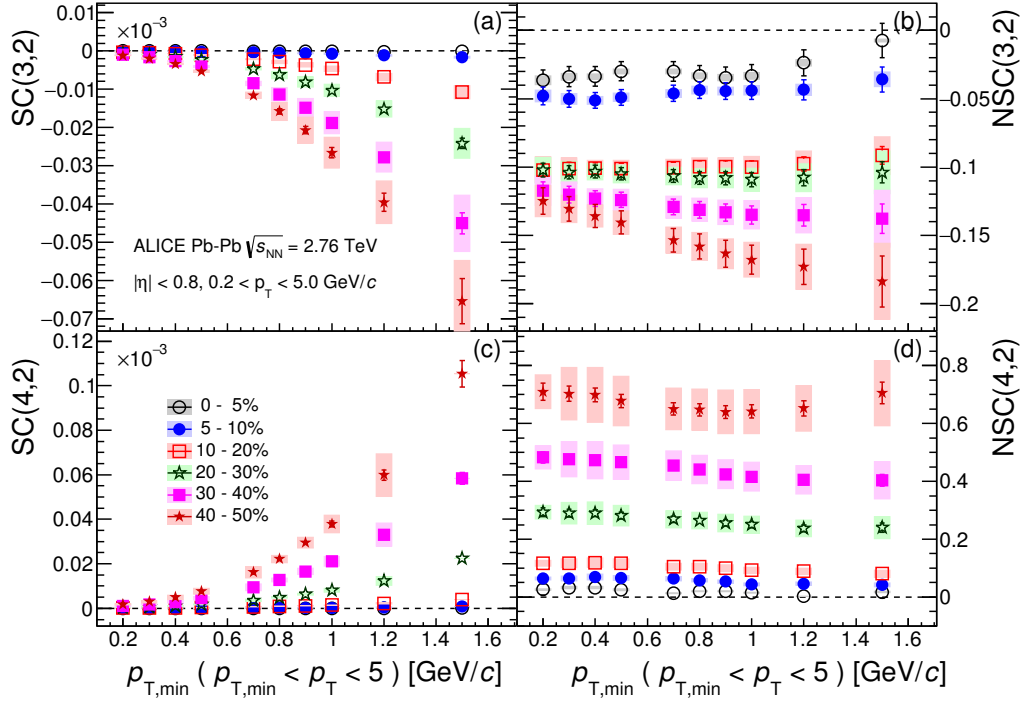


Fig. 2: SC(3,2) and SC(4,2) ((a) and (c)) as a function of minimum p_T cuts in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV are shown in the left panels. The NSC(3,2) and NSC(4,2) ((b) and (d)) are shown in the right panels. Systematic uncertainties are represented with boxes.

It can be seen in Fig. 1(a) that the lower order harmonic correlations as well as SC(5,2) increase non-linearly towards peripheral collisions. In the case of SC(5,3) and SC(4,3), the centrality dependence is weaker than for the other harmonic correlations. The NSC(5,3) observable shows the strongest normalized correlation among all harmonics while NSC(5,2) shows the weakest centrality dependence. Both NSC(3,2) and NSC(4,3) are getting more anti-correlated toward peripheral collisions and have the similar magnitude.

To study the p_T dependence of SC(m,n), we present the results as a function of the low p_T cut-off ($p_{T,min}$), instead of using independent p_T intervals; this decreases large statistical fluctuations in the results. Various minimum p_T cuts from 0.2 to 1.5 GeV/c are applied. The p_T dependent results for SC(3,2) and SC(4,2) as a function of minimum p_T cuts are shown in Figs. 2a and 2c. The strength of SC(m,n) becomes larger as $p_{T,min}$ increases. The centrality dependence is stronger with higher $p_{T,min}$ cuts, with SC(m,n) getting much larger as centrality or $p_{T,min}$ increases. The NSC(3,2) and NSC(4,2) observables with different $p_{T,min}$ are shown in Figs. 2b and 2d. The strong $p_{T,min}$ dependence observed in SC(m,n) is not seen in NSC(m,n). This indicates that the p_T dependence of SC(m,n) is dominated by the p_T dependence of the individual flow harmonics $\langle v_n \rangle$. The $p_{T,min}$ dependence of NSC(3,2) is not clearly seen and it is consistent with no $p_{T,min}$ dependence within the statistical and systematic errors for the centrality range 0–30%, while showing a moderate increase of anti-correlation with increasing $p_{T,min}$ for the 30–50% centrality. The NSC(4,2) observable shows a moderate decreasing trend as $p_{T,min}$ increases. These observations are strikingly different from the p_T dependence of the individual flow harmonics, where the relative flow fluctuations $\sigma_{v_2}/\langle v_2 \rangle$ [74] are independent of transverse momentum up to $p_T \sim 8$ GeV/c (see Fig. 3 in Ref. [75]).

As discussed in Sec. 2, the NSC(m,n) observables are normalized by the product $\langle v_m^2 \rangle \langle v_n^2 \rangle$. These products are obtained from two-particle correlations using a pseudorapidity gap of $|\Delta\eta| > 1.0$. In this

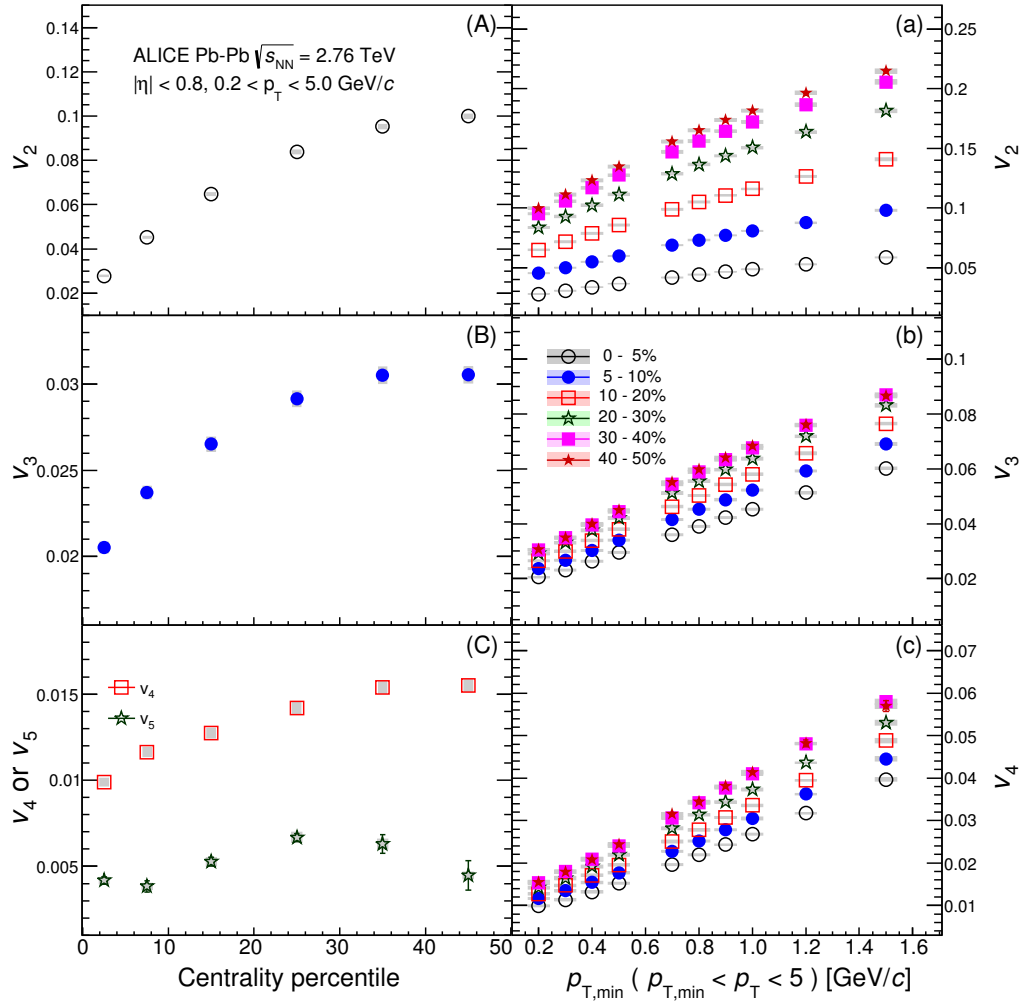


Fig. 3: The individual flow harmonics v_n for $n = 2-5$ in Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV are shown in the left panels ((A), (B) and (C)). v_4 and v_5 are shown in the same panel (C). The $p_{T,min}$ dependence of v_n for $n = 2-4$ is shown in the right panels ((a), (b) and (c)).

paper we shortly denote the p_T integrated $v_n\{2, |\Delta\eta| > 1\}$ as v_n . The individual flow harmonics v_n used in calculations of the NSC observables are shown in Fig. 3. The centrality dependence of v_n for $n = 2-5$ is shown in the left panels ((A), (B) and (C)) of Fig. 3. The v_n values ($n < 5$) are equivalent to those in [11]. The 5th order flow harmonic v_5 is shown in panel (C). The $p_{T,\min}$ dependence of v_n for $n = 2-4$ is shown in the right panels ((a), (b) and (c)) of Fig. 3 in all centrality ranges relevant to the measured NSC(m,n) observables.

6 Model Comparisons

We have performed a systematic comparison of the centrality and transverse momentum dependence of the SC(m,n) and NSC(m,n) to the event-by-event EKRT+viscous hydrodynamics [45], VISH2+1 [76, 77], and the AMPT [63, 78, 79] models. Comparisons for v_n coefficients with the model calculations are presented in Appendix A.

In the event-by-event EKRT+viscous hydrodynamic calculations [45], the initial energy density profiles are calculated using a next-to-leading order perturbative-QCD+saturation model [80, 81]. The subsequent space-time evolution is described by relativistic dissipative fluid dynamics with different parameterizations for the temperature dependence of the shear viscosity to entropy density ratio $\eta/s(T)$. This model gives a good description of the charged hadron multiplicity and the low p_T region of the charged hadron spectra at RHIC and the LHC (see Figs. 11-13 in [45]). Each of the $\eta/s(T)$ parameterizations is adjusted to reproduce the measured v_n from central to mid-peripheral collisions (see Fig. 15 in [45] and Appendix A).

The VISH2+1 [76, 77] event-by-event calculations for relativistic heavy-ion collisions are based on (2+1)-dimensional viscous hydrodynamics which describes the QGP phase and the highly dissipative and off-equilibrium late hadronic stages with fluid dynamics. By tuning transport coefficients and decoupling temperature for a given scenario of initial conditions, it can describe the p_T spectra and different flow harmonics at RHIC and the LHC [20, 76, 82, 83] energies. Three different types of initial conditions [58] (MC-Glauber, MC-KLN and AMPT) along with different constant η/s values have been used for our data to model comparisons. Traditionally, the Glauber model constructs the initial entropy density from the wounded nucleon and binary collision density profiles [84]. The KLN model assumes that the initial energy density is proportional to that of the initial gluons calculated from the corresponding k_T factorization formula [85]. In Monte Carlo versions MC-Glauber and MC-KLN [86–88] of these models an additional initial state fluctuations are introduced through position fluctuations of individual nucleons inside the colliding nuclei. For the AMPT initial conditions [83, 89, 90], the fluctuating energy density profiles are constructed from the energy distribution of individual partons, which fluctuate in both momentum and coordinate space. Compared with the MC-Glauber and MC-KLN initial conditions, the additional Gaussian smearing in the AMPT initial conditions gives rise to non-vanishing initial local flow velocities [89].

Even though thermalization could be achieved shortly in collisions of very large nuclei and/or at extremely high energy [91], the dense matter created in heavy-ion collisions may not reach full thermal or chemical equilibrium due to its finite size and short lifetime. To address such non-equilibrium many-body dynamics, the AMPT model [78, 92, 93] has been developed, which includes both initial partonic and final hadronic interactions and the transition between these two phases of matter. The initial conditions in the AMPT are given by the spatial and momentum distributions of minijets and soft strings from the HIJING model [70, 94]. For the data comparisons three different configurations of the AMPT model has been used: the default one and string melting with and without hadronic rescattering. The input parameters used in all configurations are: $\alpha_s = 0.33$, a partonic cross-section of 1.5 mb, while the Lund string fragmentation parameters were set to $\alpha = 0.5$ and $b = 0.9 \text{ GeV}^{-2}$. In the default configuration, partons are recombined with their parent strings when they stop interacting. The resulting strings are

later converted into hadrons using the Lund string fragmentation model [95, 96]. In the string melting configuration, the initial strings are melted into partons whose interactions are described by the ZPC parton cascade model [97]. These partons are then combined into the final state hadrons via a quark coalescence model. In both configurations, the dynamics of the subsequent hadronic matter is described by a hadronic cascade based on A Relativistic Transport (ART) model [98] which includes resonance decays. The string melting configuration of the AMPT without hadronic rescattering was used to study the influence of the hadronic phase on the development of the anisotropic flow. Even though the string melting version of AMPT [78, 99] reasonably well reproduces particle yields, p_T spectra, and v_2 of low p_T pions and kaons in central and mid-central Au–Au collisions at $\sqrt{s_{NN}} = 200$ GeV and Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV [79], it was seen clearly in a recent study [100] that it fails to quantitatively reproduce the flow harmonics of identified hadrons (v_2 , v_3 , v_4 and v_5) at $\sqrt{s_{NN}} = 2.76$ TeV. It turns out that the radial flow in AMPT is 25% lower than that measured at the LHC, which is responsible for this quantitative disagreement [100]. The details of the AMPT configurations used in this article and the comparisons of p_T -differential v_n for pions, kaons and protons to the data can be found in [100].

6.1 Centrality Dependence of $SC(m,n)$ and $NSC(m,n)$

Comparison to event-by-event EKRT+viscous hydrodynamic predictions with various parameterizations of the temperature dependence of $\eta/s(T)$ was shown in Fig. 2 of Ref. [57]. It was demonstrated that $NSC(3,2)$ is sensitive mainly to the initial conditions, while $NSC(4,2)$ is sensitive to both the initial conditions and the system properties, which is consistent with the predictions from [36]. The model calculations for $NSC(4,2)$ observable show that it has better sensitivity for different $\eta/s(T)$ parameterizations but they cannot describe neither the centrality dependence nor the absolute values. The discrepancy between data and theoretical predictions indicates that the current understanding of initial conditions in models of heavy-ion collisions needs to be revisited to further constrain $\eta/s(T)$. The measurement of $SC(m,n)$ and $NSC(m,n)$ can provide new constraints for the detailed modeling of fluctuating initial conditions.

The calculations for the two sets of parameters which describe the lower order harmonic correlations best are compared to the data in Fig. 4. As can be seen in Fig. 1 from Ref. [45], for the “param1” parameterization the phase transition from the hadronic to the QGP phase occurs at the lowest temperature, around 150 MeV. This parameterization is also characterized by a moderate slope in $\eta/s(T)$ which decreases (increases) in the hadronic (QGP) phase. The model calculations in which the temperature of the phase transition is larger than for “param1” are ruled out by the previous measurements [57]. As shown in Fig. 4, the correlations between v_5 and v_2 are well described for all available centralities and similarly for the correlations between v_5 and v_3 within the errors. In the case of v_4 and v_3 , the same models underestimate the anti-correlation in the data significantly in mid-central collisions and it fails similarly for the anti-correlation between v_3 and v_2 .

The comparison to the VISH2+1 calculation [58] is shown in Fig. 5. All calculations with large η/s regardless of the initial conditions ($\eta/s = 0.2$ for MC-KLN and MC-Glauber initial conditions and $\eta/s = 0.16$ for AMPT initial conditions) fail to capture the centrality dependence of the $SC(m,n)$ observables of all orders, shown in the left panels in Fig. 5. Among the calculations with small η/s ($\eta/s = 0.08$), the one with the AMPT initial conditions describes the data better than the ones with other initial conditions for all $SC(m,n)$ observables measured, but it cannot describe the data quantitatively for most of the centrality ranges.

However, $NSC(4,2)$ is sensitive both to the initial conditions and the η/s parameterizations used in the models. Even though $NSC(4,2)$ favors both AMPT initial conditions with $\eta/s = 0.08$ and MC-Glauber initial conditions with $\eta/s = 0.20$, $SC(4,2)$ can only be described by models with smaller η/s . Hence the calculation with large $\eta/s = 0.20$ is ruled out. We conclude that η/s should be small and that AMPT initial conditions are favored by the data. The $NSC(5,2)$ and $NSC(5,3)$ observables are quite sensitive

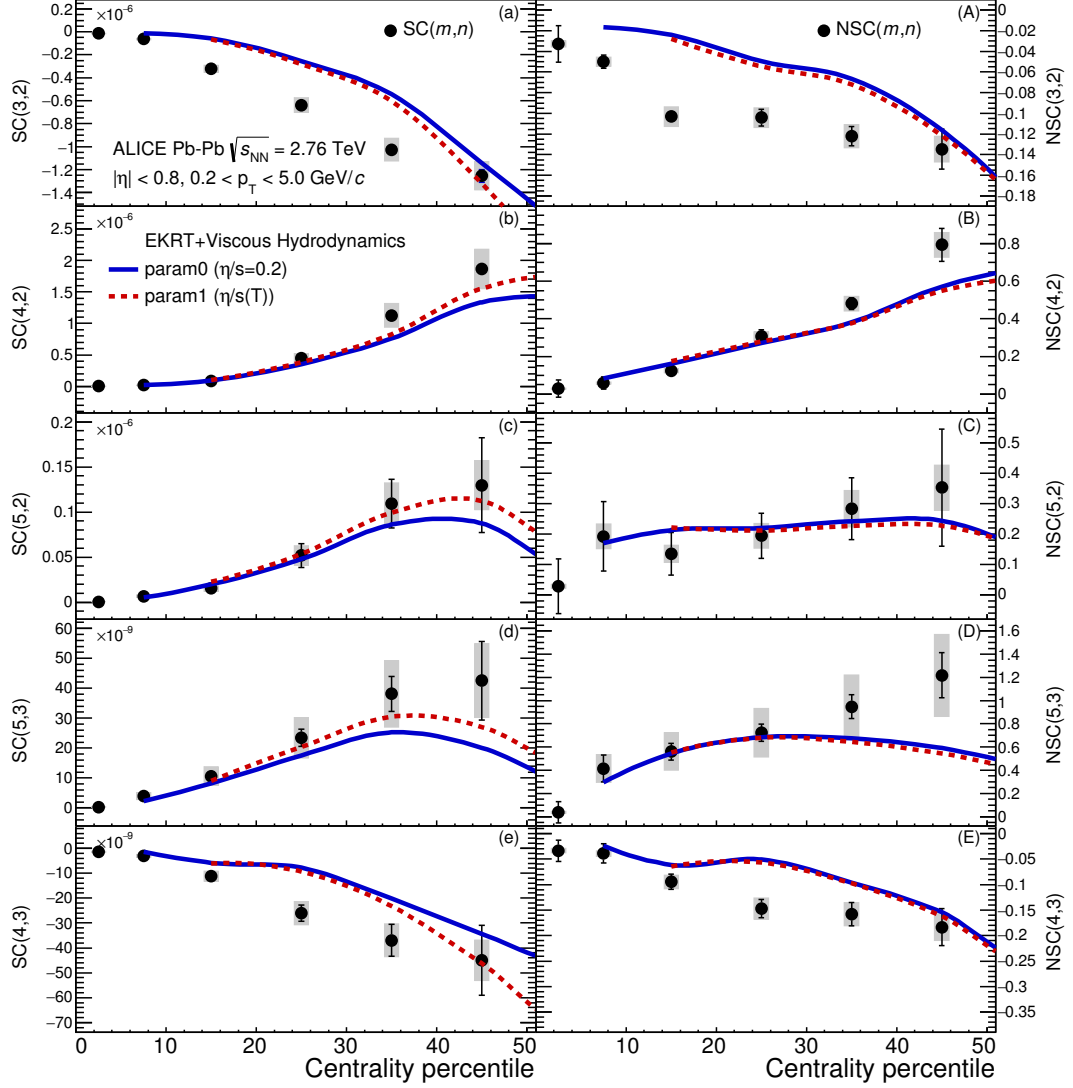


Fig. 4: The centrality dependence of $SC(m,n)$ and $NSC(m,n)$ in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV. Results are compared to the event-by-event EKRT+viscous hydrodynamic calculations [45]. The lines are hydrodynamic predictions with two different $\eta/s(T)$ parameterizations. Left (right) panels show $SC(m,n)$ ($NSC(m,n)$).

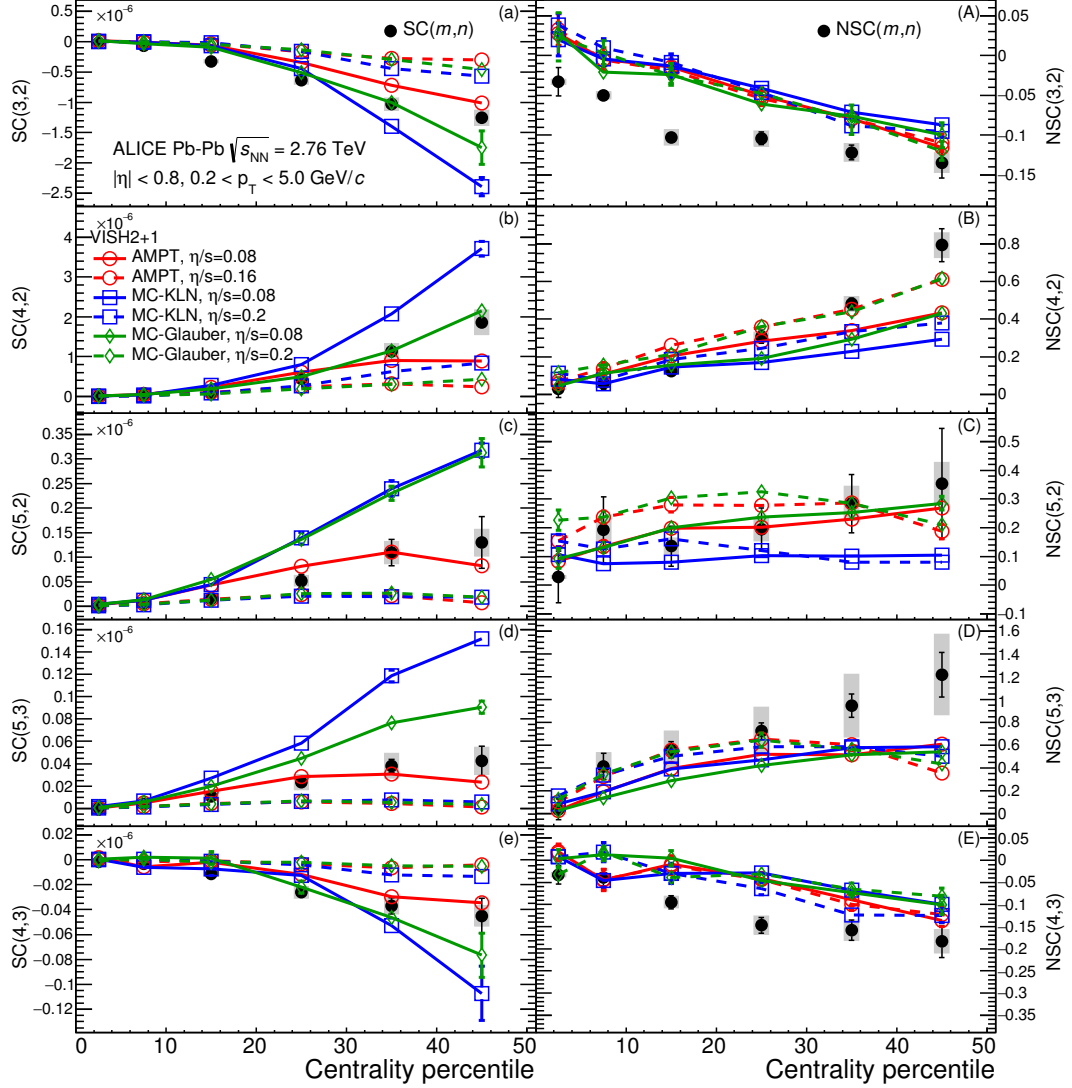


Fig. 5: The centrality dependence of $SC(m,n)$ and $NSC(m,n)$ in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV. Results are compared to various VISH2+1 calculations [58]. Three initial conditions from AMPT, MC-KLN and MC-Glauber are drawn as different colors and markers. The η/s parameters are shown as different line styles, the small shear viscosity ($\eta/s = 0.08$) are shown as solid lines, and large shear viscosities ($\eta/s = 0.2$ for MC-KLN and MC-Glauber and 0.16 for AMPT) are drawn as dashed lines. Left (right) panels show $SC(m,n)$ ($NSC(m,n)$).

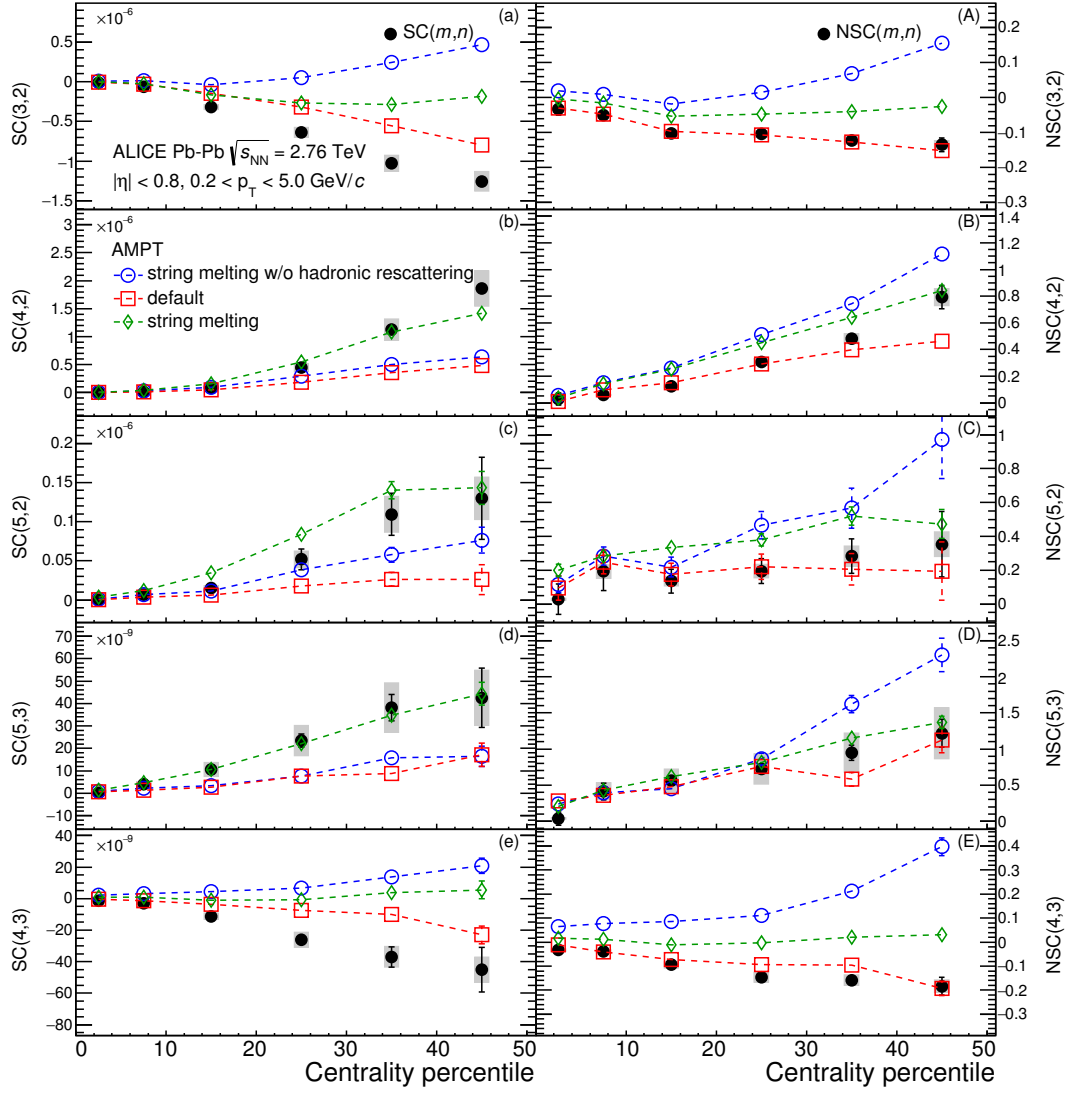


Fig. 6: The centrality dependence of $SC(m,n)$ and $NSC(m,n)$ in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV. Results are compared to various AMPT models. Left (right) panels show $SC(m,n)$ ($NSC(m,n)$).

to both the initial conditions and the η/s parameterizations. The $SC(4,3)$ results are clearly favored by smaller η/s values but $NSC(4,3)$ cannot be described by these models quantitatively.

The $SC(m,n)$ and $NSC(m,n)$ observables calculated from AMPT simulations are compared with data in Fig. 6. For $SC(3,2)$, the calculation with the default AMPT settings is closest to the data, but none of the AMPT configurations can describe the data fully. The third version based on the string melting configuration without the hadronic rescattering phase is also shown. The hadronic rescattering stage makes both $SC(3,2)$ and $NSC(3,2)$ stronger in the string melting AMPT model but not enough to describe the data. Further investigations proved why the default AMPT model can describe $NSC(3,2)$ fairly well but underestimates $SC(3,2)$. By taking the differences in the individual flow harmonics (v_2 and v_3) between the model and data into account, it was possible to recover the difference in $SC(3,2)$ between the data and the model. The discrepancy in $SC(3,2)$ can be explained by the overestimated individual v_n values as reported in [100] in all centrality ranges.

In the case of $SC(4,2)$, the string melting configuration of the AMPT model can describe the data fairly

well while the default configuration underestimates it. The NSC(4,2) observable is slightly overestimated by the string melting setting which can describe SC(4,2) but the default AMPT configuration can describe the data better. The influence of the hadronic rescattering phase on NSC(4,2) is opposite to other observables (SC(3,2), NSC(3,2) and SC(4,2)). The hadronic rescattering makes NSC(4,2) slightly smaller. It should be noted that the agreement with SC(m,n) should not be overemphasized since there are discrepancies in the individual v_n between the AMPT models and the data as was demonstrated for SC(3,2). Hence the simultaneous description of SC(m,n) and NSC(m,n) should give better constraints on the parameters in AMPT models. The string melting AMPT model describes SC(5,3) and NSC(5,3) well. However, the same setting overestimates SC(5,2) and NSC(5,2). The default AMPT model can describe NSC(5,3) and NSC(5,2) fairly well as in the case of NSC(3,2) and NSC(4,2). In the case of SC(4,3), neither of the settings can describe the data but the default AMPT model comes the closest to the data. The NSC(4,3) observable is well described by the default AMPT model but cannot be reproduced by the string melting AMPT model. In summary, the default AMPT model describes well the normalized symmetric cumulants (NSC(m,n)) from lower to higher order harmonic correlations while the string melting AMPT model overestimates NSC(3,2) and NSC(5,2) and predicts a very weak correlation both for NSC(3,2) and NSC(4,3).

As discussed in Sec. 5, a hierarchy NSC(5,3) > NSC(4,2) > NSC(5,2) holds for centrality ranges > 20% within the errors. Except for the 0–10% centrality range, we found that the same hierarchy also holds in the hydrodynamic calculations and the AMPT models explored in this article. While NSC(5,2) is smaller than NSC(5,3), SC(5,2) is larger than SC(5,3). The observed inverse hierarchy, SC(5,2) > SC(5,3), can be explained by different magnitudes of the individual flow harmonics ($v_2 > v_3$). This can be attributed to the fact that flow fluctuations are stronger for v_3 than v_2 [14]. This was claimed in Ref. [58] and also seen in Ref. [101] based on the AMPT model calculations. NSC(m,n) correlators increase with larger η/s in hydrodynamic calculations in the 0–30% centrality range in the same way as the event plane correlations [102, 103]. In semi-peripheral collisions (> 40%), the opposite trend is observed.

We list here the important findings from the model comparisons to the centrality dependence of SC(m,n) and NSC(m,n):

- (i) The NSC(3,2) observable is sensitive mainly to the initial conditions, while the other observables are sensitive to both the initial conditions and the temperature dependence of η/s .
- (ii) The correlation strength between v_3 and v_2 and between v_4 and v_3 (SC(3,2), SC(4,2), NSC(3,2) and NSC(4,3)) is significantly underestimated in hydrodynamic model calculations in mid-central collisions.
- (iii) All the VISH2+1 model calculations with large η/s fail to capture the centrality dependence of the correlations regardless of the initial conditions.
- (iv) Among the VISH2+1 model calculations with small η/s ($\eta/s = 0.08$), the one with the AMPT initial conditions captures the data qualitatively, but not quantitatively for most of the centrality ranges.
- (v) The default AMPT model can describe the normalized symmetric cumulants (NSC(m,n)) quantitatively for most centralities while the string melting AMPT model fails to describe them.
- (vi) A hierarchy NSC(5,3) > NSC(4,2) > NSC(5,2) holds for centrality ranges > 20% within the errors. This hierarchy is reproduced well both by hydrodynamic and AMPT model calculations.

The agreement of various model calculations with the data is quantified using the log-likelihood or least squares- χ^2 method [104] by calculating the χ^2 :

$$\chi^2 = \sum_{i=1}^N \frac{(y_i - f_i)^2}{\sigma_i^2}, \quad (5)$$

where y_i (f_i) is a measurement (model) value in a centrality bin i . The systematical and statistical errors are combined in quadrature $\sigma_i = \sqrt{\sigma_{i,stat}^2 + \sigma_{i,syst}^2}$. The total number of data samples N in Eq. 5 is 4, which corresponds to the number of bins in the centrality range 10–50% used in χ^2 calculations. The χ^2 for model calculations which are best in describing the SC observables for each of the three different types of models is shown in Fig. 7.

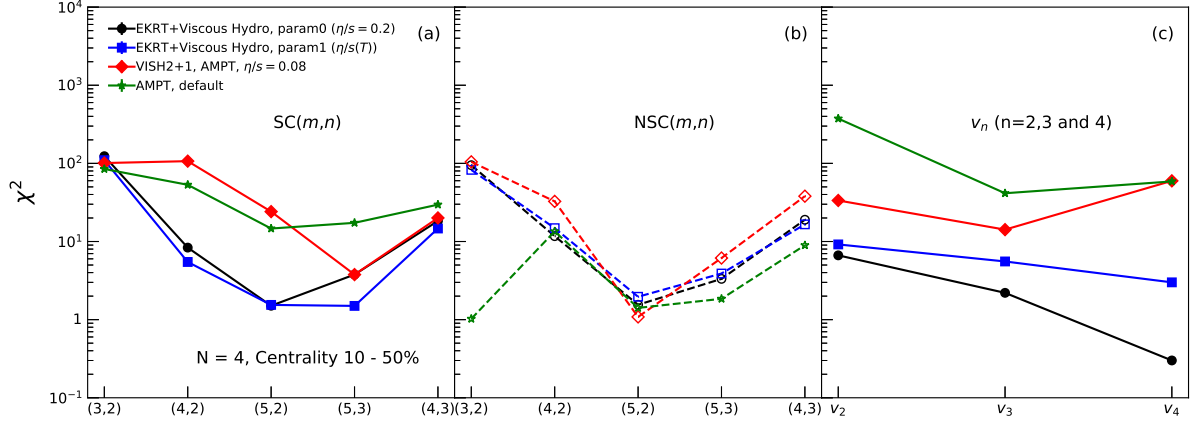


Fig. 7: The χ^2 values calculated by Eq. 5 are shown for $SC(m,n)$ (a), $NSC(m,n)$ (b) and individual harmonics v_n (c). Results are for model calculations which are best in describing the SC observables for each of the three different types of models.

The results for $SC(m,n)$ and $NSC(m,n)$ are presented in Fig. 7a and 7b, respectively. The χ^2 values for the individual flow harmonics v_n ($n = 2, 3$ and 4) are shown in Fig. 7c. We found that in case of the calculations from VISH2+1 with AMPT initial conditions ($\eta/s = 0.08$) and default AMPT models, the χ^2 values for $SC(m,n)$ are larger than those for $NSC(m,n)$. This reflects the fact that the individual flow harmonics v_n are not well described by those models compared to event-by-event EKRT+viscous hydrodynamics. This is quantified in Fig. 7c, where the χ^2 values for v_n are much larger both for VISH2+1 and default AMPT calculations than event-by-event EKRT+viscous hydrodynamics. The default AMPT model gives best χ^2 values for $NSC(m,n)$, especially for $NSC(3,2)$. However the χ^2 values of this model are largest for v_n among the models especially for v_2 .

The χ^2 values for v_2 and v_3 are significantly smaller than those for $SC(3,2)$ and $NSC(3,2)$ for all the hydrodynamic calculations. The χ^2 values for $SC(4,2)$ and $NSC(4,2)$ from event-by-event EKRT+viscous hydrodynamics are comparable to that for v_2 but larger than for v_4 . This illustrates that the SC observables together with the individual flow harmonics provide better constraints for the model parameters than each of them individually.

Once the calculations with an average over the temperature evolution of the expanding medium (constant η/s parameterization ($\eta/s = 0.20$, “param0”)) is compared to the temperature dependent η/s parameterization (“param1”, where η/s is the lowest ($\eta/s = 0.12$) at the temperature around 150 MeV), in case of v_n , the χ^2 values with “param0” are smaller than those with “param1” but the opposite trend is observed for $SC(m,n)$. Furthermore we observe a clear separation of two parameterizations for $SC(5,3)$ which is even larger than $SC(4,2)$. These observations support that with the SC observables one can constraint the temperature dependent $\eta/s(T)$, taking one step further from estimating the simple average quantity η/s .

Even though the calculations from event-by-event EKRT+viscous hydrodynamics give best χ^2 values both for $SC(m,n)$ and $NSC(m,n)$, the χ^2 values are large especially for the observables which include v_3 . Even with the best model calculations from theoretical models, the χ^2 value varies a lot depending on the model parameters and/or different order SC observables, which implies that the different order harmonic correlations have different sensitivity to the initial conditions and the system properties.

6.2 Transverse Momentum Dependence of NSC(3,2) and NSC(4,2)

The NSC(3,2) and NSC(4,2) observables as a function of $p_{T,\min}$ are compared to the AMPT simulations in Fig. 8 and Fig. 9, respectively. The observed p_T dependence for NSC(3,2) in mid-central collisions is also seen in AMPT simulations for higher $p_{T,\min}$. The default configuration of the AMPT is well reproducing NSC(3,2), while the other AMPT configurations predict a very strong p_T dependence above 1 GeV/c and cannot describe the magnitudes of both NSC(3,2) and NSC(4,2) simultaneously. In the case of NSC(3,2), the default AMPT model describes the magnitude and p_T dependence well in all collision centralities except for 40–50% where the model underestimates the data and shows a stronger p_T dependence than the data. As for NSC(4,2), the default AMPT configuration which describes NSC(3,2) can also reproduce the data well except for the 10–20% and 40–50% centralities. Comparison of the string melting AMPT configuration with and without hadronic rescattering suggests that a very strong p_T dependence as well as the correlation strength are weakened by the hadronic rescattering. Consequently, the observed weak p_T dependence may be due to hadronic rescattering. To clarify, the relative contributions to the final state particle distributions from partonic and hadronic stages need further study.

The event-by-event EKRT+viscous hydrodynamic calculations are also compared to the data in Fig. 8 and Fig. 9. In the case of NSC(3,2), the hydrodynamic calculations underestimate the magnitude of the data as discussed in Sec. 6.1 and show very weak p_T dependence for all centralities. The p_T dependence of NSC(3,2) is well captured by the model calculations in all collision centralities except for 40–50% where the data shows stronger p_T dependence than the models. The difference between the model calculations with the two different parameterizations of $\eta/s(T)$ is very small. As for NSC(4,2), the model calculations overestimate the magnitude of the data in the 5–20% centralities and underestimate it in the centrality range 30–50%. However, the p_T dependence is well described by the model calculations in all centrality ranges, while the difference of the model results for the two parameterizations in most centralities is rather small.

The observed moderate p_T dependence in mid-central collisions both for NSC(3,2) and NSC(4,2) might be an indication of possible viscous corrections to the equilibrium distribution at hadronic freeze-out, as predicted in [36]. The comparisons to hydrodynamic models can further help to understand the viscous corrections to the momentum distributions at hadronic freeze-out [45, 52, 54–56].

7 Summary

In this article, we report the centrality dependence of correlations between the higher order harmonics (v_4 , v_5) and the lower order harmonics (v_2 , v_3) as well as the transverse momentum dependence of the correlations between v_3 and v_2 and between v_4 and v_2 . The results are presented in terms of the Symmetric Cumulants $SC(m,n)$. It was demonstrated earlier in [57] that $SC(m,n)$ is insensitive to nonflow effects and independent of symmetry plane correlations.

We have found that fluctuations of SC(3,2) and SC(4,3) are anti-correlated in all centralities while fluctuations of SC(4,2), SC(5,2) and SC(5,3) are correlated for all centralities. These measurements were compared to various hydrodynamic model calculations with different initial conditions as well as different parameterizations of the temperature dependence of η/s . It is found that the different order harmonic correlations have different sensitivities to the initial conditions and the system properties. Therefore they have discriminating power in separating the effects of η/s from the initial conditions on the final state

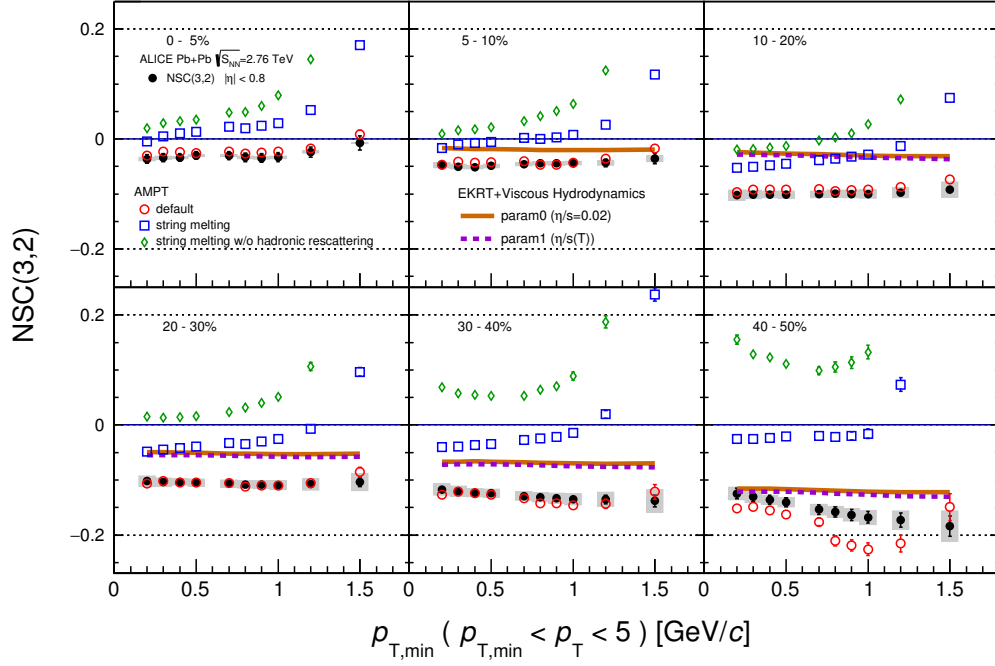


Fig. 8: NSC(3,2) as a function of the minimum p_T cut in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV is compared to various AMPT configurations and event-by-event EKRT+viscous hydrodynamic calculations [45].

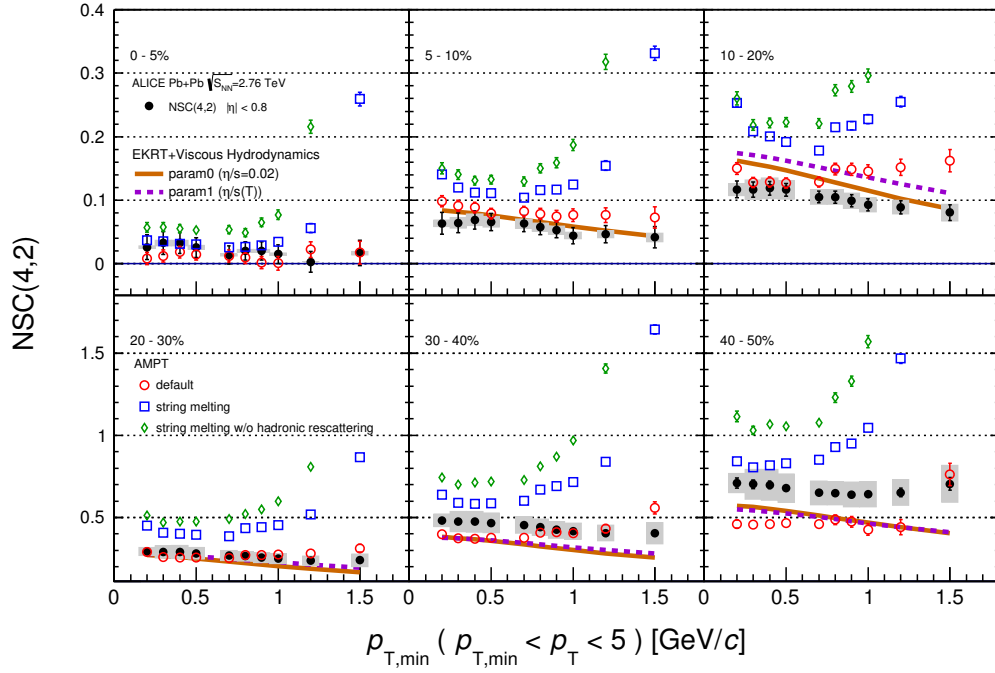


Fig. 9: NSC(4,2) as a function of the minimum p_T cut in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV is compared to various AMPT configurations and event-by-event EKRT+viscous hydrodynamic calculations [45].

particle anisotropies. The comparisons to VISH2+1 calculations show that all the models with large η/s , regardless of the initial conditions, fail to capture the centrality dependence of higher order correlations. Based on the tested model parameters, the data favors small η/s and the AMPT initial conditions.

A quite clear separation of the correlation strength for different initial conditions is observed for these higher order harmonic correlations compared to the lower order. The default configuration of the AMPT model describes well the normalized symmetric cumulants ($NSC(m,n)$) for most centralities and for most combinations of harmonics which were considered. Finally, we have found that v_3 and v_2 as well as v_4 and v_2 correlations have moderate p_T dependence in mid-central collisions. This might be an indication of possible viscous corrections to the equilibrium distribution at hadronic freeze-out. Together with the measurements of individual harmonics the new results for $SC(m,n)$ and $NSC(m,n)$ can be used to further optimize model parameters and put better constraints on the initial conditions and the transport properties of nuclear matter in ultra-relativistic heavy-ion collisions.

Acknowledgements

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A Model Comparisons of the Individual Flow Harmonics v_n

As discussed in Sec. 2, $NSC(m,n)$ is expected to be insensitive to the magnitudes of v_m and v_n but $SC(m,n)$ has contributions from both the correlations between the two different flow harmonics and the individual harmonics v_n . Therefore, it is important to check how well the theoretical models used in Sec. 6 describe the measured v_n data shown in Sec. 5. The v_n is measured for charged particles in the pseudorapidity range $|\eta| < 0.8$ and the transverse momentum range $0.2 < p_T < 5.0$ GeV/c as a function of collision centrality [11].

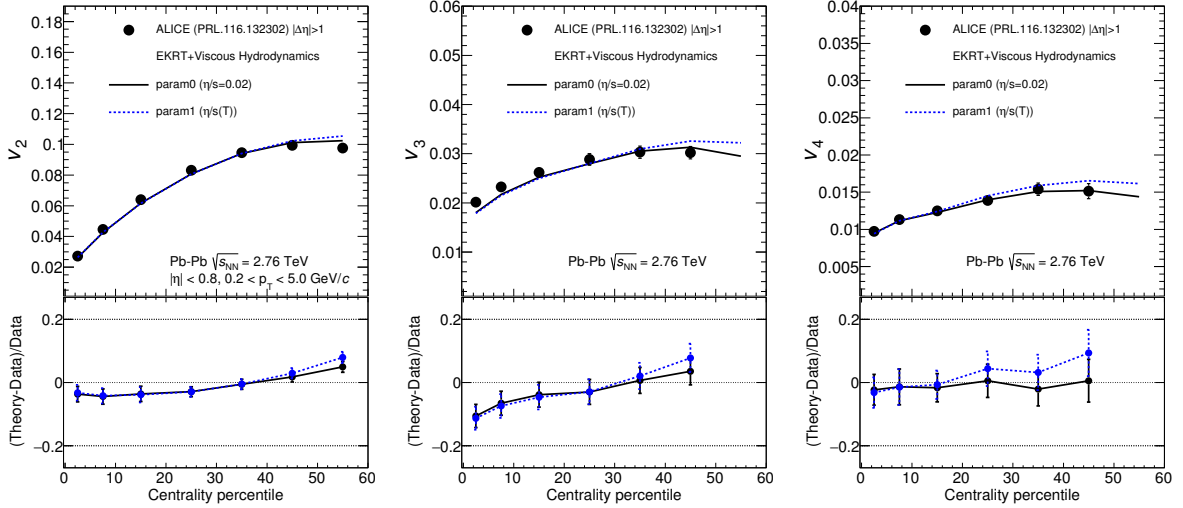


Fig. A.1: The individual flow harmonics v_n for $n = 2-4$ in Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV [11]. Results are compared to the event-by-event EKRT+viscous hydrodynamic calculations [45] for two different $\eta/s(T)$ parameterizations, labeled in the same way as in [45].

The measured v_n for $n = 2-4$ in Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV are compared to the event-by-event EKRT+viscous hydrodynamic calculations [45] in Fig. A.1. In these calculations the initial conditions and η/s parameterizations are chosen to reproduce the LHC v_n data. The calculations capture the centrality dependence of v_n in the central and mid-central collisions within 5% for v_2 and 10% for v_3 and v_4 .

The VISH2+1 calculations with various initial conditions and η/s parameters are compared to the v_n data in Fig. A.2. Neither MC-Glauber nor MC-KLN initial conditions can simultaneously describe v_2 , v_3 and v_4 . In particular, for MC-Glauber initial conditions, VISH2+1 with $\eta/s = 0.08$ can describe well v_2 from central to mid-central collisions, but overestimates v_3 and v_4 for the same centrality ranges. For MC-KLN initial conditions, VISH2+1 with $\eta/s = 0.20$ reproduces v_2 but underestimates v_3 and v_4 for the presented centrality regions. The calculations with AMPT initial conditions improves the simultaneous descriptions of v_n ($n = 2, 3$ and 4). The overall difference to the data is quite large if all the model settings are considered, about 30% for v_n ($n = 2$ and 3) and 50% for v_4 . The calculations with AMPT initial conditions reproduce the observed centrality dependence with an accuracy of 10–20%.

The AMPT calculations with various configurations are compared to the v_n data in Fig. A.3. The string melting version of AMPT [78, 99] reasonably reproduces v_n as shown in Fig. A.3 within 20% for v_2 and 10% for v_3 and v_4 . The version based on the string melting configuration without the hadronic rescattering phase underestimates the data compared to the calculations with the string melting version of AMPT, which demonstrates that a large fraction of the flow is developed during the late hadronic rescattering stage in the string melting version of AMPT. The default version of AMPT underestimates v_n for $n =$

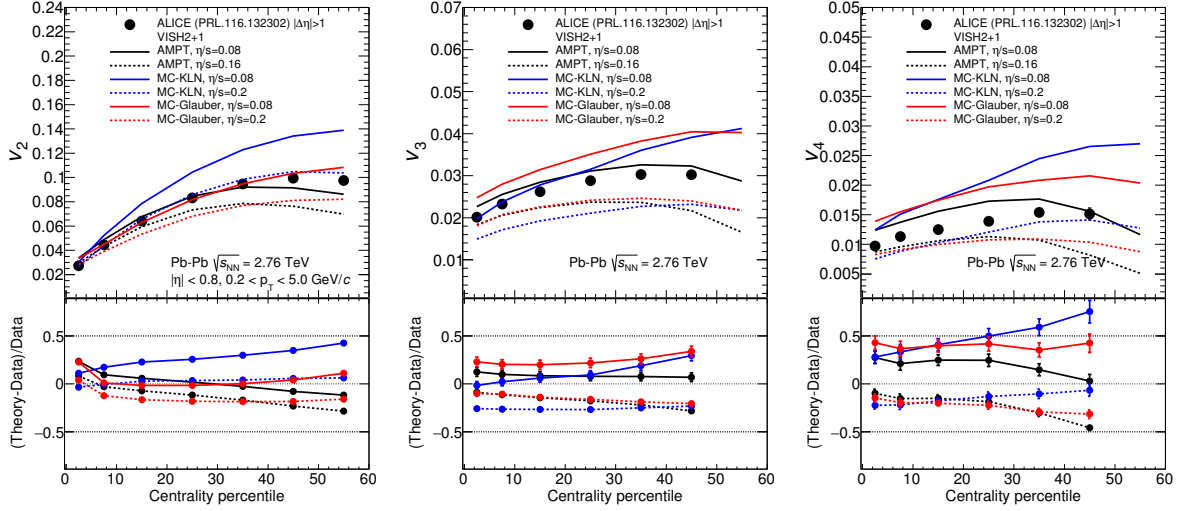


Fig. A.2: The individual flow harmonics v_n for $n = 2-4$ in Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV [11]. Results are compared to various VISH2+1 calculations [58]. Three initial conditions from AMPT, MC-KLN and MC-Glauber are shown in different colors. The results for different η/s values are shown as different line styles, the small shear viscosity ($\eta/s = 0.08$) are shown as solid lines, and large shear viscosities ($\eta/s = 0.2$ for MC-KLN and MC-Glauber and, 0.16 for AMPT) are drawn as dashed lines.

2–4 by $\approx 20\%$. It should be noted that the default AMPT model can describe the normalized symmetric cumulants ($NSC(m,n)$) quantitatively for most centralities while the string melting AMPT model fails to describe them.

Finally, few selected calculations from three theoretical models which describe the v_n data best are shown in Fig. A.4. The calculations from event-by-event EKRT+viscous hydrodynamics, VISH2+1 with AMPT initial conditions ($\eta/s = 0.08$) and the string melting version of AMPT give the best description of the individual flow harmonics v_n ($n = 2, 3$ and 4) with an accuracy of 5–20%. The centrality dependence differs in the three models as well as in the different order flow harmonics. Together with $SC(m,n)$ and $NSC(m,n)$, the simultaneous description of individual flow harmonics v_n at all orders is necessary to further optimize model parameters and put better constraints on the initial conditions and the transport properties of nuclear matter in ultra-relativistic heavy-ion collisions.

B The ALICE Collaboration

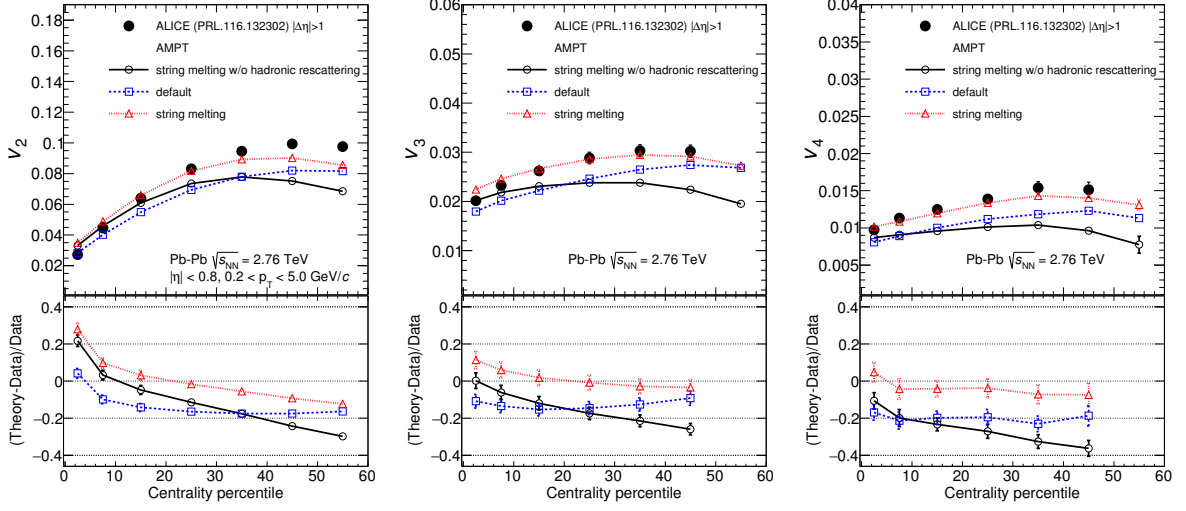


Fig. A.3: The individual flow harmonics v_n ($n = 2, 3$ and 4) in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV [11] are compared to various AMPT models.

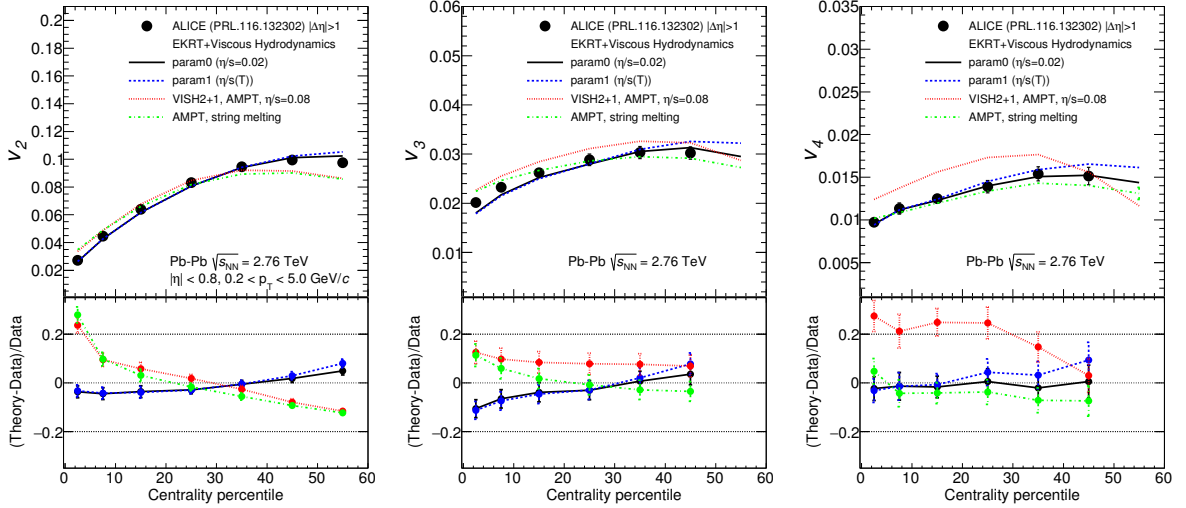


Fig. A.4: The individual flow harmonics v_n for $n = 2-4$ in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV [11]. The results are compared with selected calculations from three different types of models which are best in describing v_n coefficients.