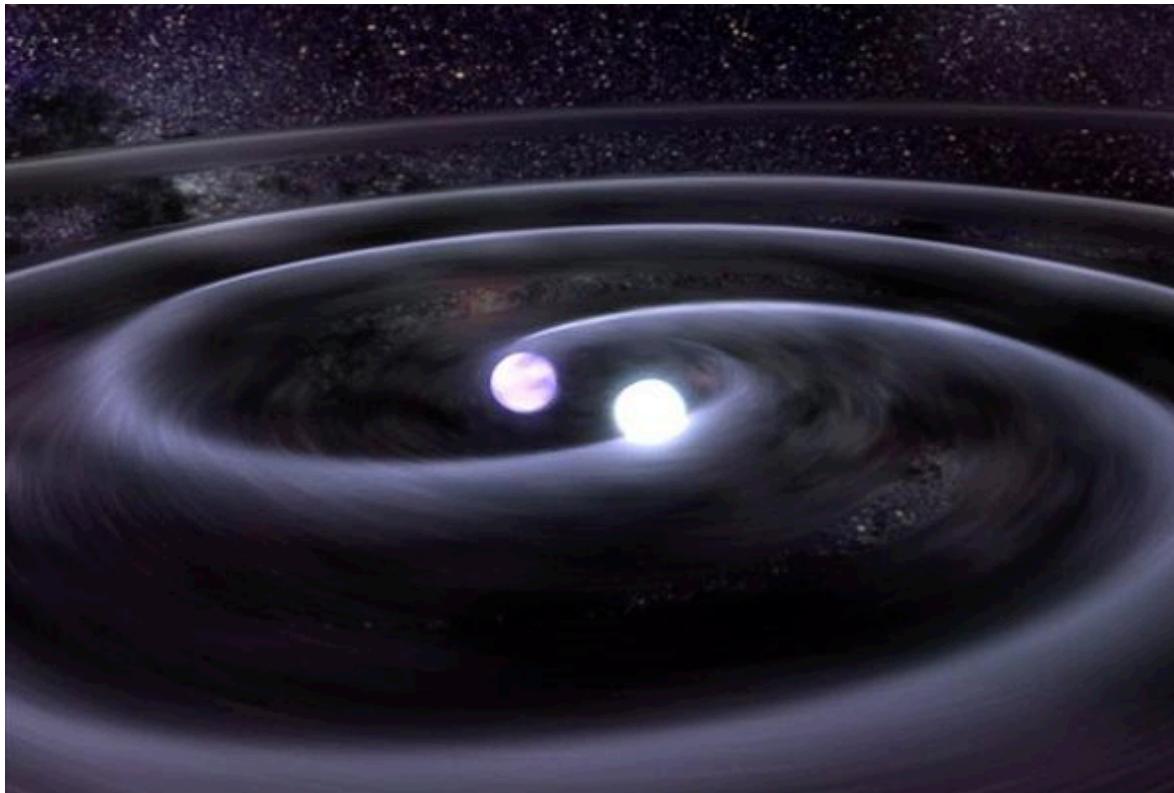


Merging Compact Binaries



Dong Lai
Cornell University

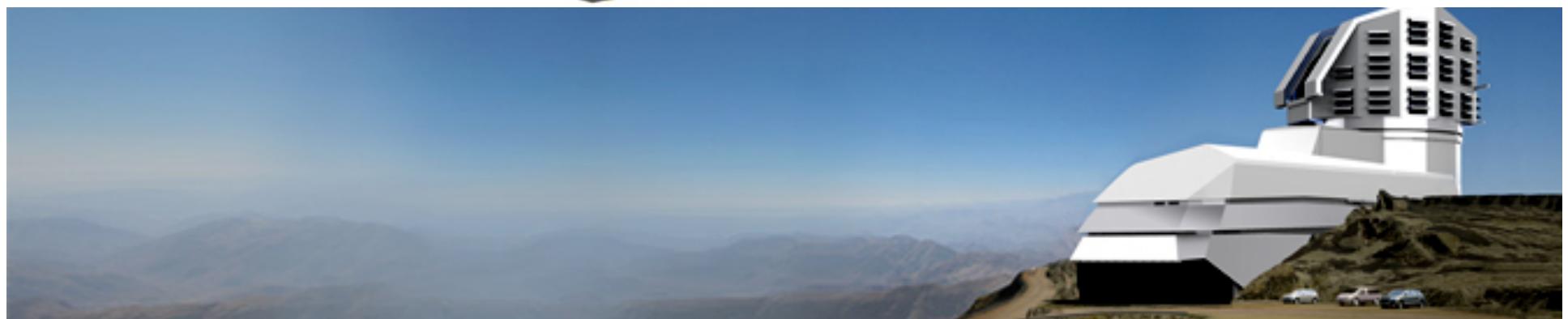
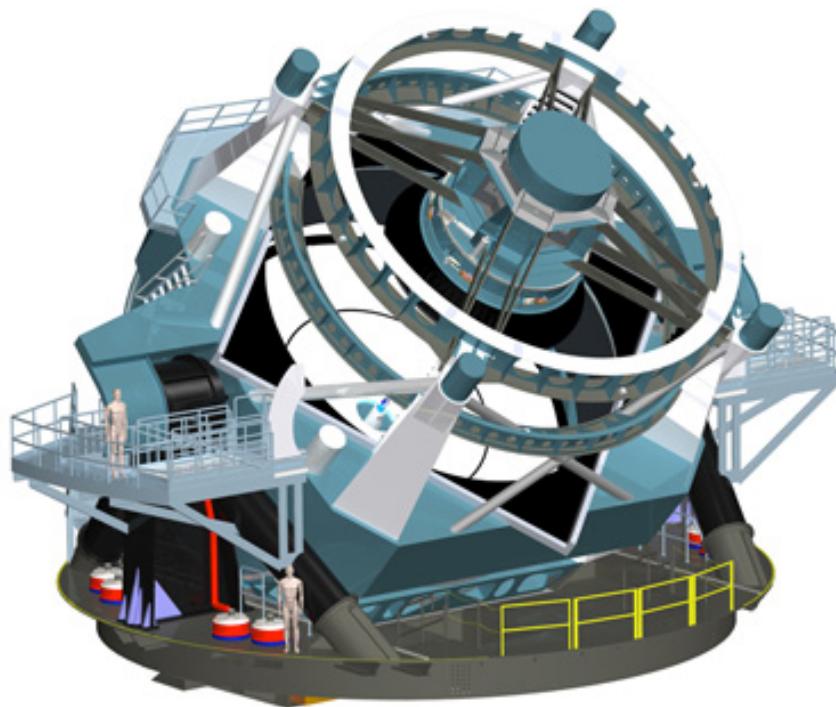
Astronomical Institute, Prague, 7/2/2013

Merging Compact Binaries

- 1. Neutron Star/Black Hole Binaries**
- 2. White Dwarf Binaries**

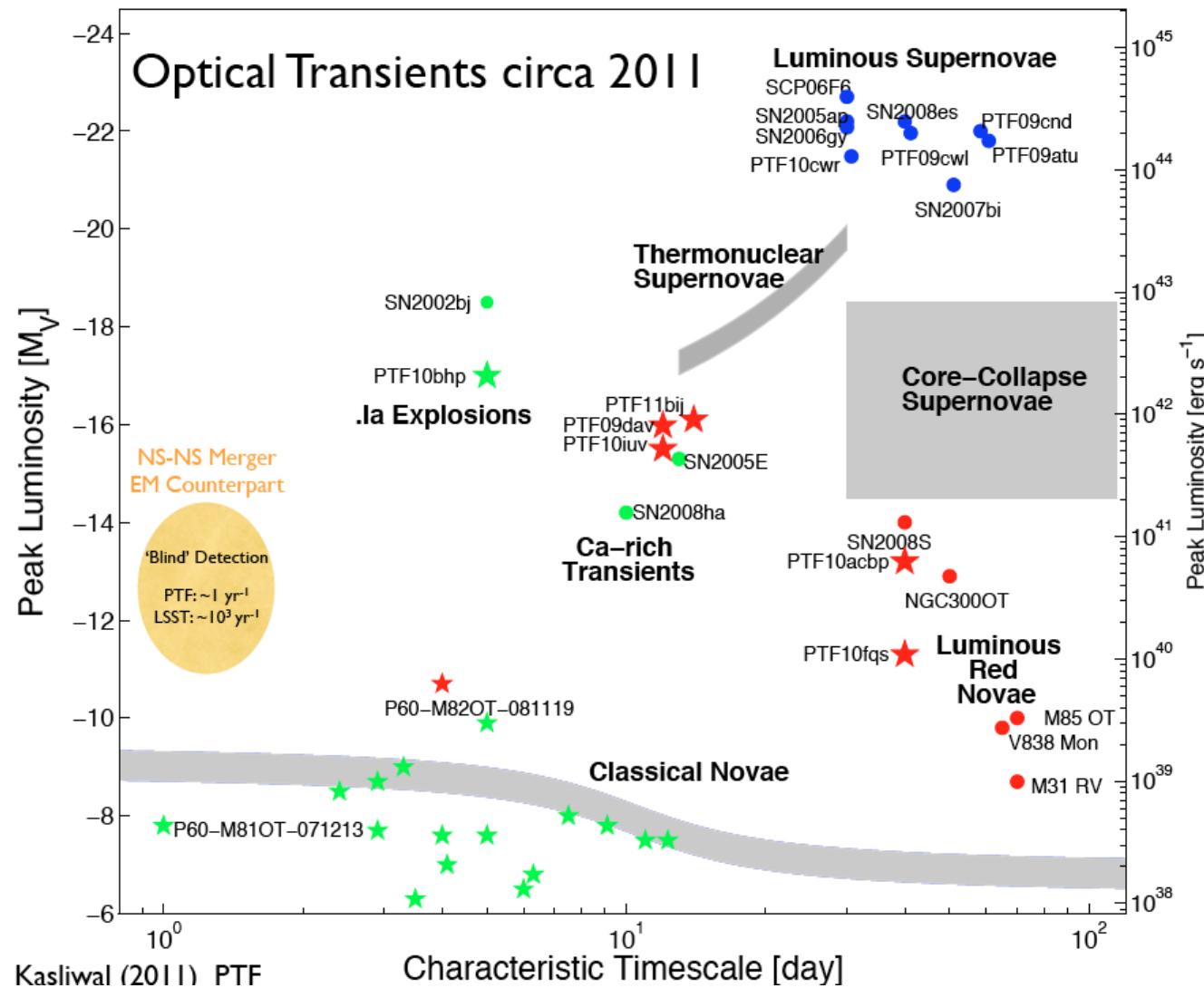
Transient & Variable Universe

Wide-field, fast imaging telescopes in optical: **PTF, Pan-Starrs, LSST**

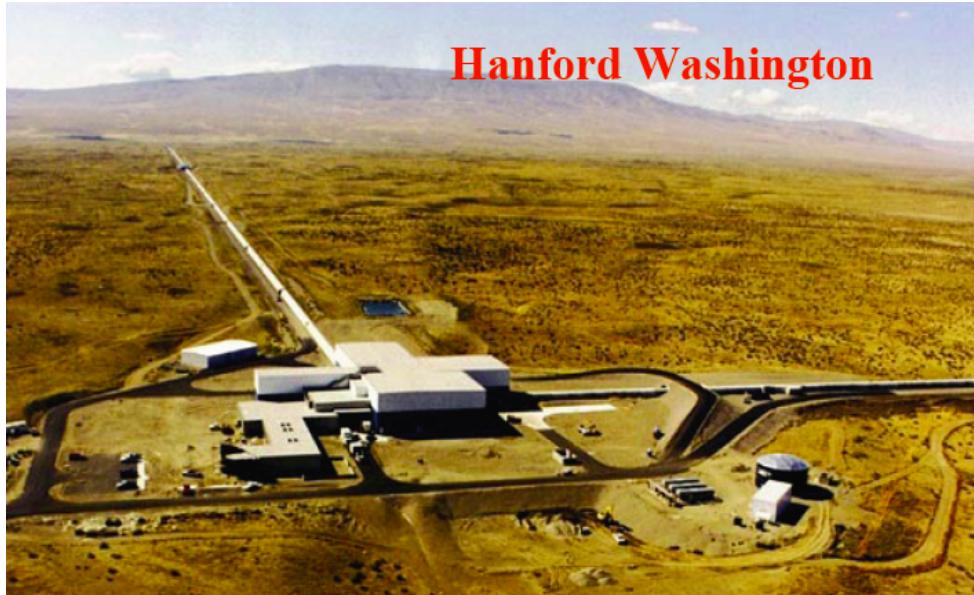


Transient & Variable Universe

Wide-field, fast imaging telescopes in optical: **PTF, Pan-Starrs, LSST**



Gravitational Wave Astronomy



Hanford Washington

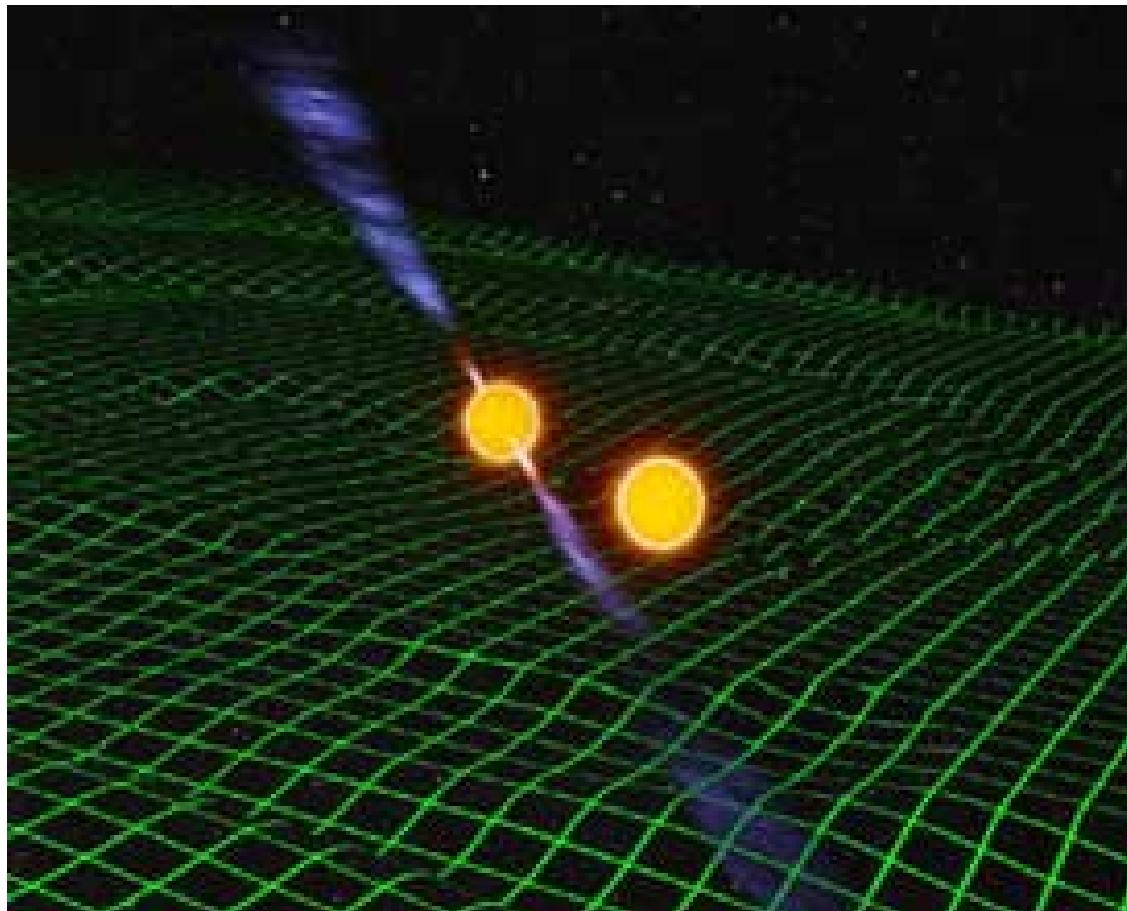


Livingston,
Louisiana

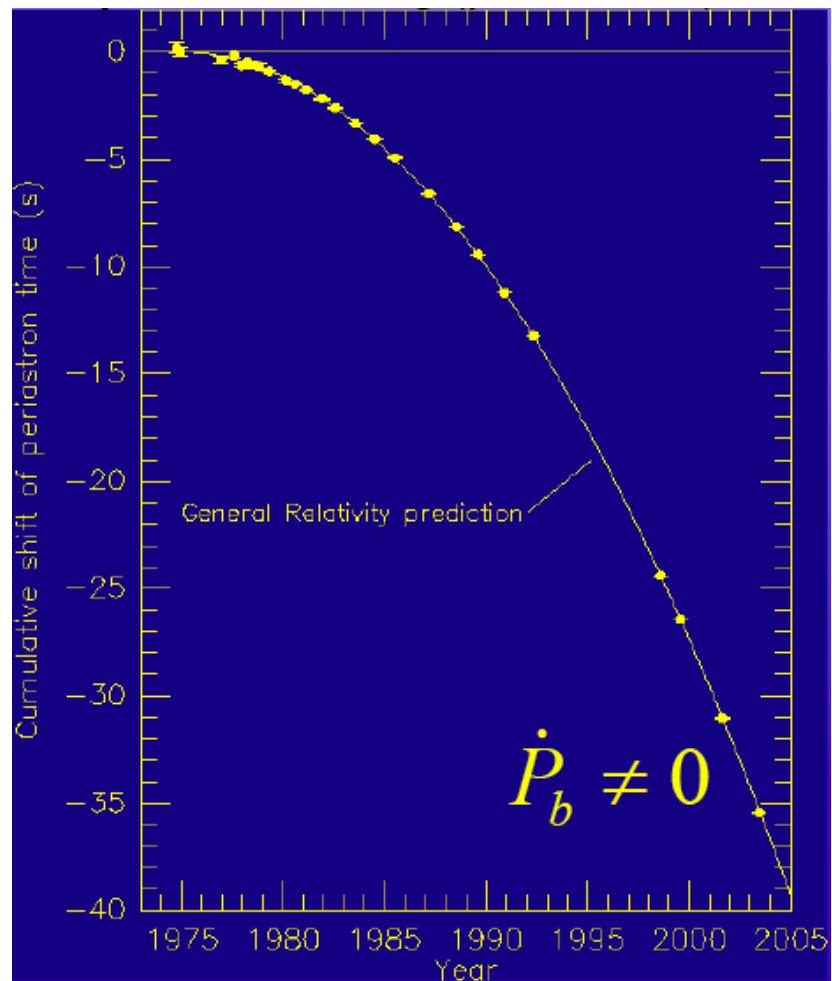
LIGO
VIRGO



Merging NS and BH Binaries



NS/NS Binaries: Binary Pulsars



$$\dot{N}_{\text{merge}} = 10^{-5} - 3 \times 10^{-4} \text{ yr}^{-1} \text{ per galaxy}$$

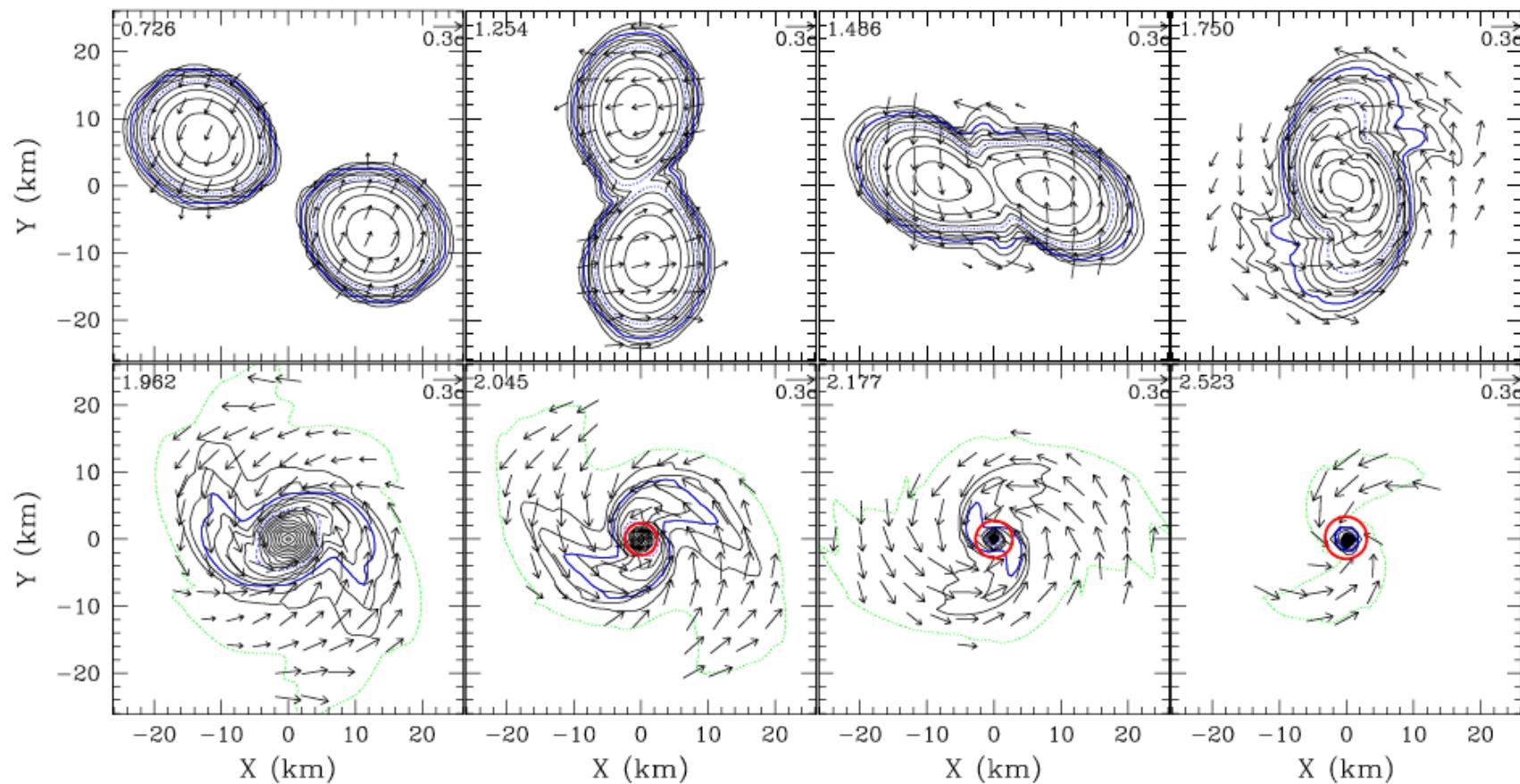


Nobel Prize 1993

Taylor & Weisberg 2005

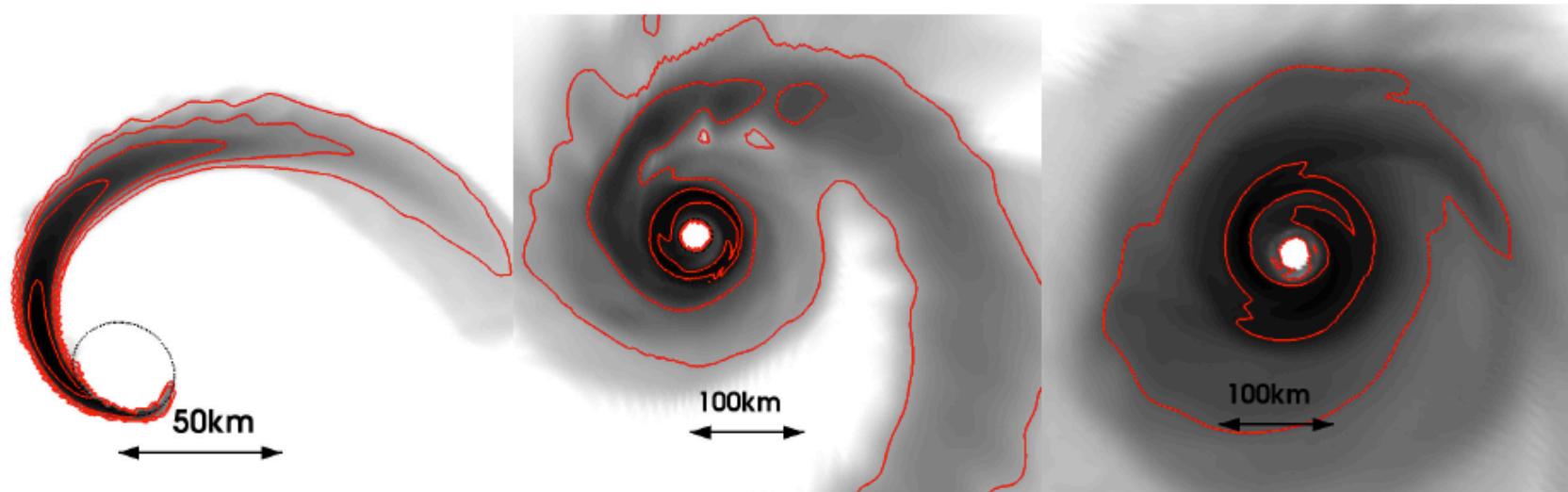
(Based on 3 systems in Galaxy that will merge within Hubble time;
No observed NS/BH and BH/BH yet !)

NS-NS Merger



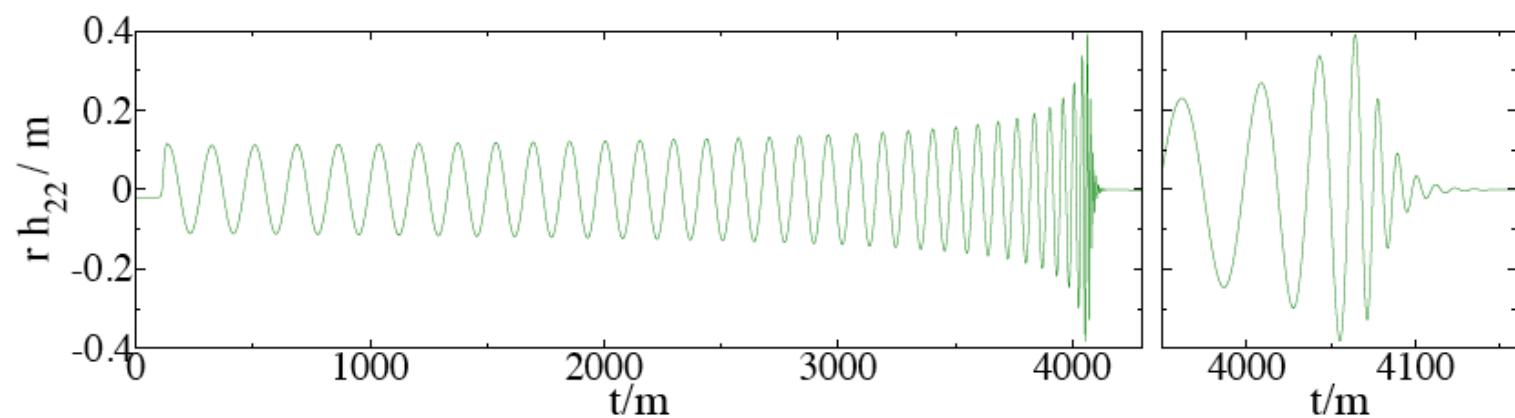
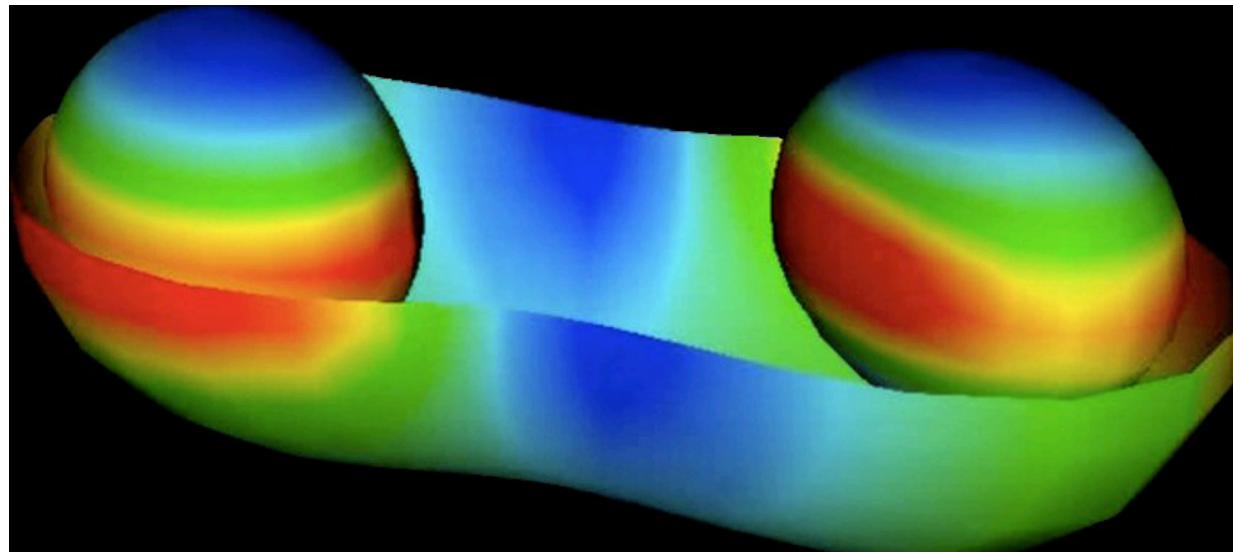
Shibata et al. 2006

BH-NS Merger



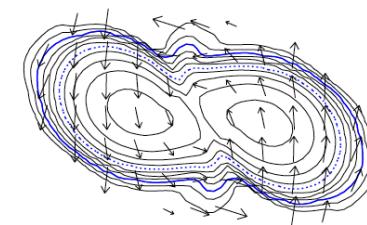
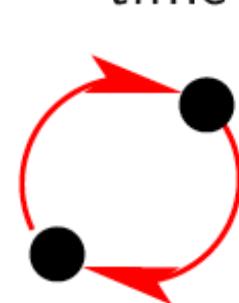
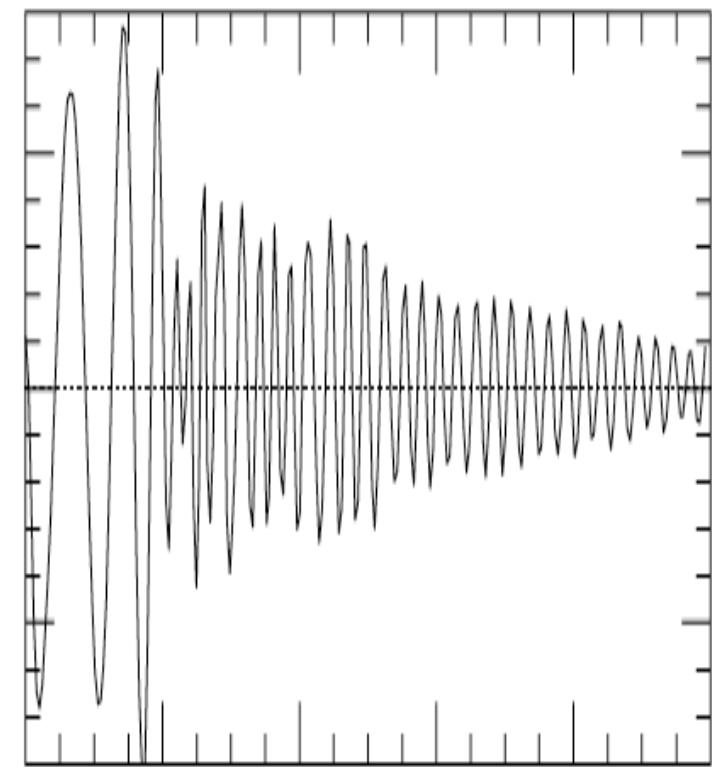
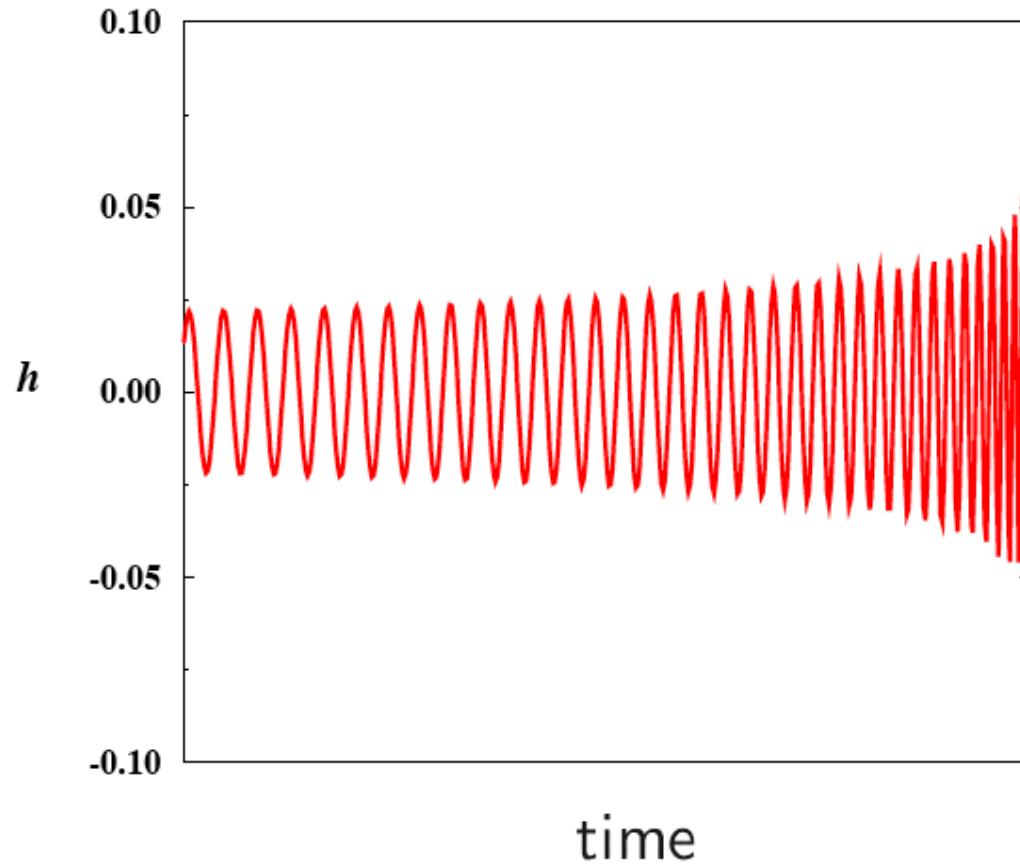
F. Foucart et al (Cornell) 2011,13

BH-BH Merger



Cornell-Caltech collaboration

The last few minutes: Gravitational Waveform



Gravitational Waves

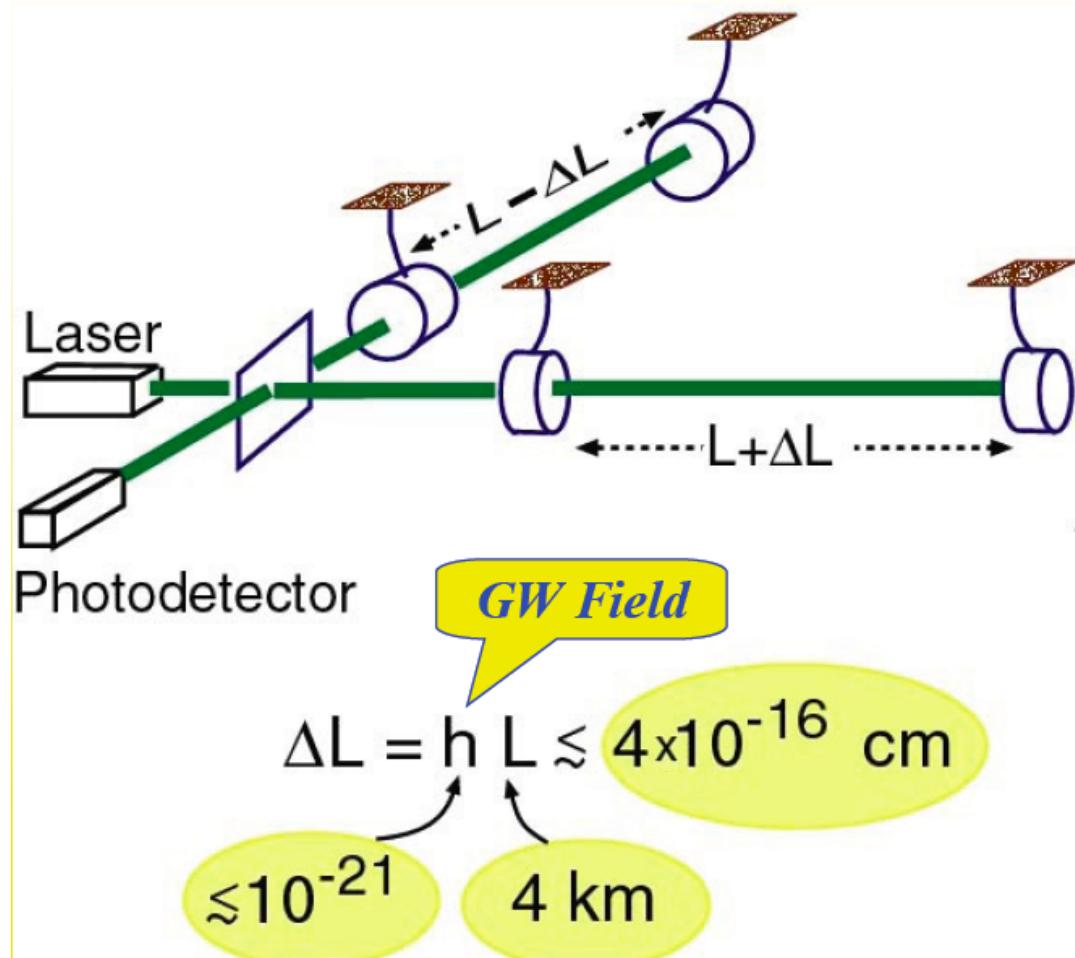
- Warpage of Spacetime
- Generated by time-dependent quadrupoles

$$h \sim \frac{G}{c^4} \frac{\ddot{Q}}{D}$$

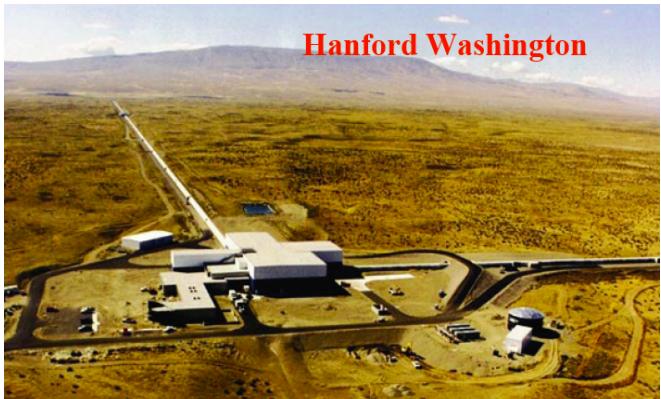
- Detector response to passage of GWs:



Gravitational Wave Interferometer



Kip Thorne



Hanford Washington

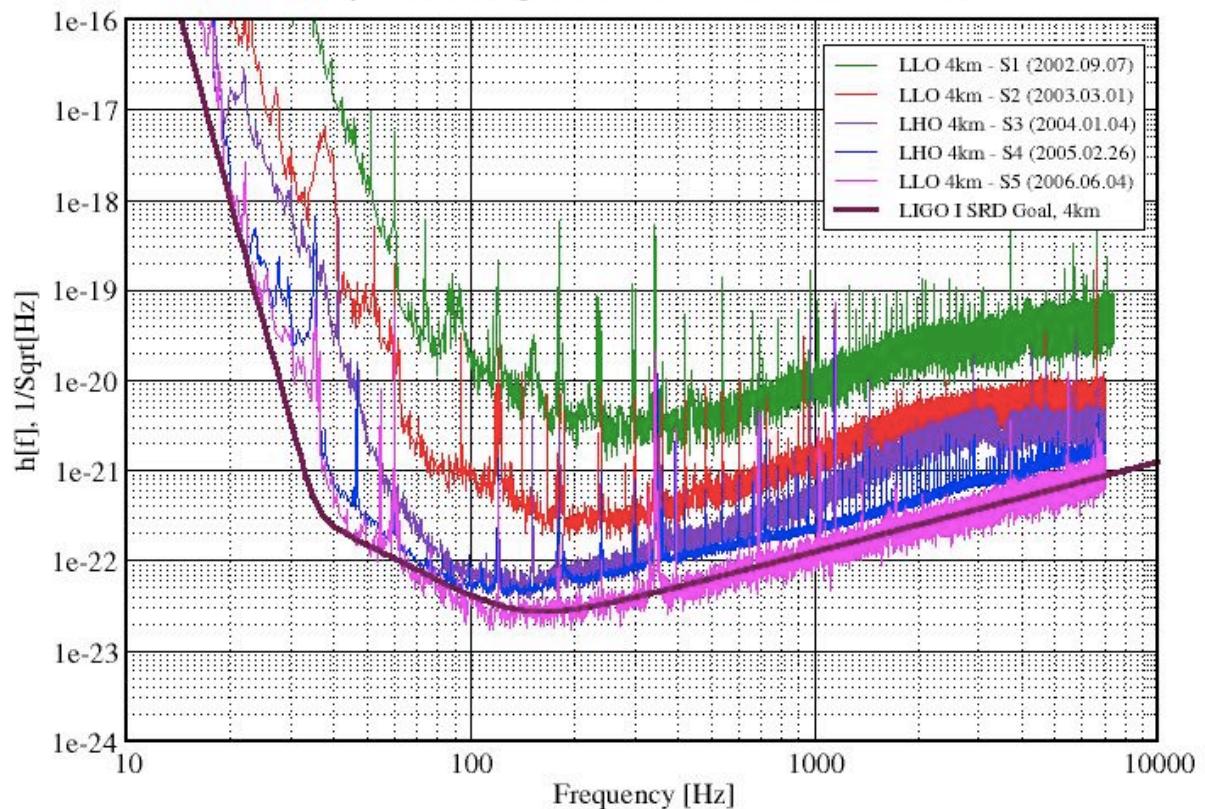


Livingston,
Louisiana



Best Strain Sensitivities for the LIGO Interferometers

Comparisons among S1 - S5 Runs LIGO-G060009-02-Z



iLIGO: reached $h \sim 10^{-21}$ (2006)

eLIGO: $h \sim 1/2$ smaller (taking/analysing data)

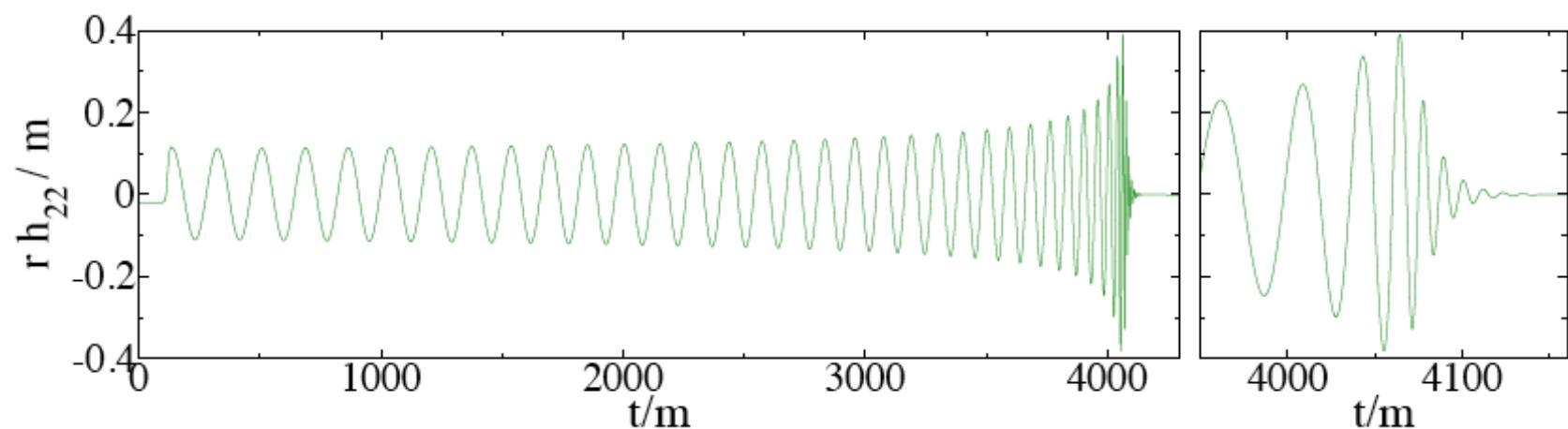
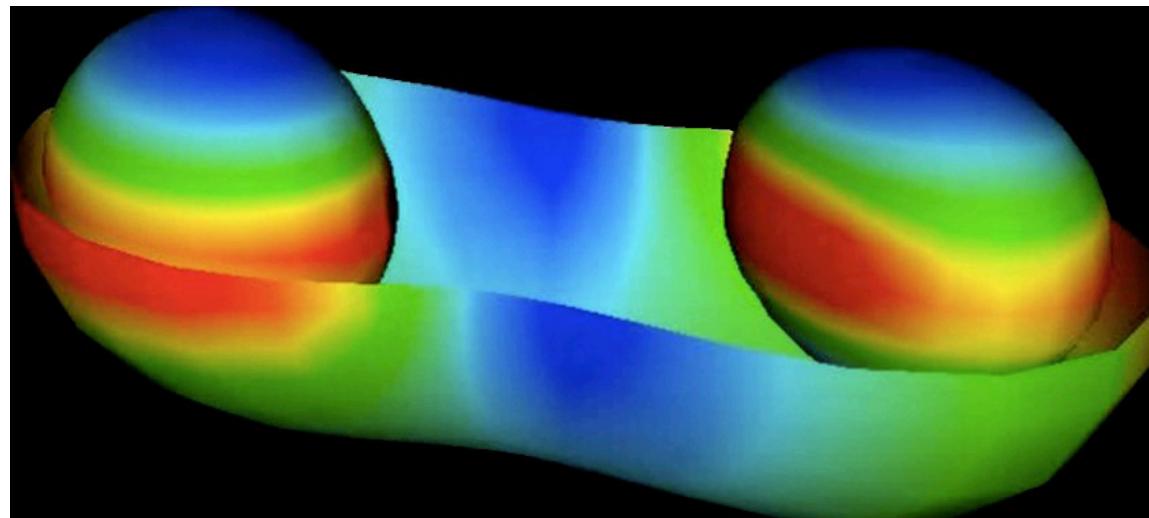
aLIGO: $h \sim 1/10$ smaller (2018?)

Compact Binary Inspiral Rates, yr^{-1}

	<i>FROM</i>	<i>Initial LIGO</i>	<i>Enhanced</i>	<i>Advanced</i>
<i>NS/NS</i>	Observed binary pulsars - Kalogera et al	.007 - .04 - .13	.06 - .3 - 1	20 - 1200 - 4000
<i>NS/BH</i>	Bethe/Brown/ Lee	.14 - .8 - 3	1 - 6 - 24	400 - 2400 - 10,000
<i>NS/NS</i> <i>or</i> <i>NS/BH</i>	Short γ burst afterglows: Nakar et al	0.001 - 0.3 \sim 0.1 γ -GW coincidences	0.01 - 3 \sim 0.8 γ -GW coincidences	2 - 30 \sim 300 γ -GW coincidences
<i>BH/BH</i>	Short γ burst afterglows: Nakar et al	0.01 - 3 \sim 0.3 γ -GW coincidences	0.1 - 30 \sim 2.4 γ -GW coincidences	20 - 1000 \sim 1000 γ -GW coincidences
				27

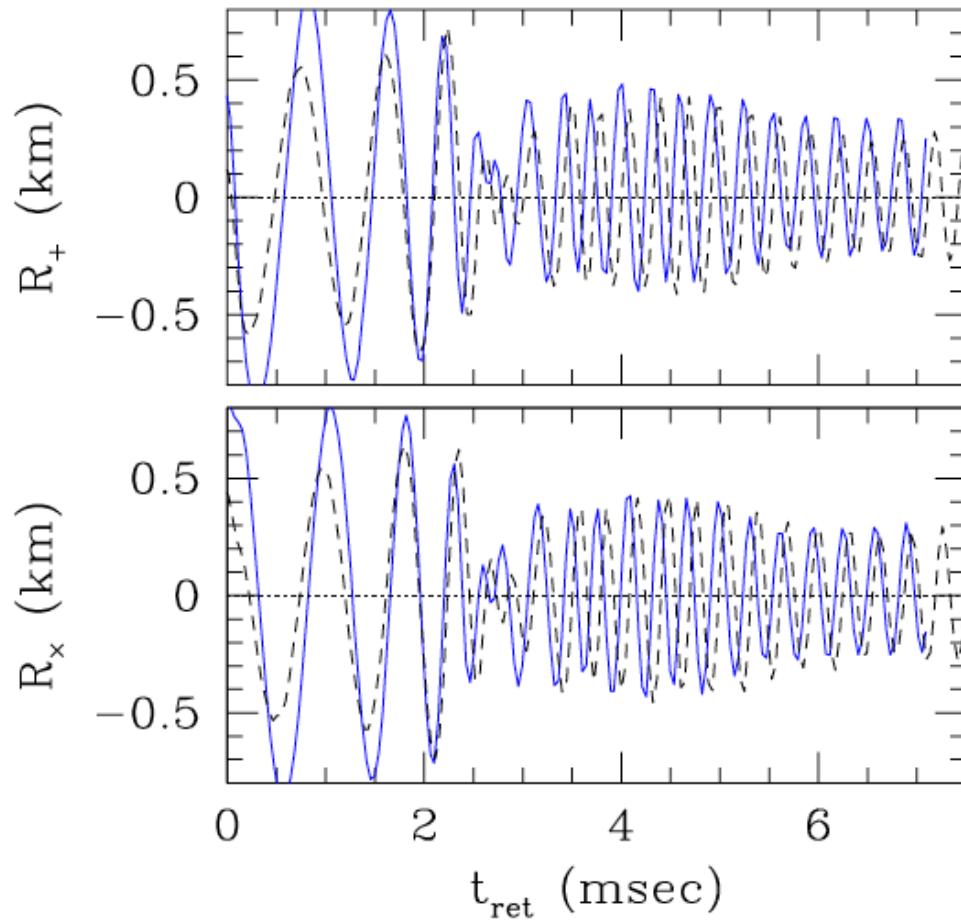
Kip Thorne

Gravitational waves probe nonlinear gravity

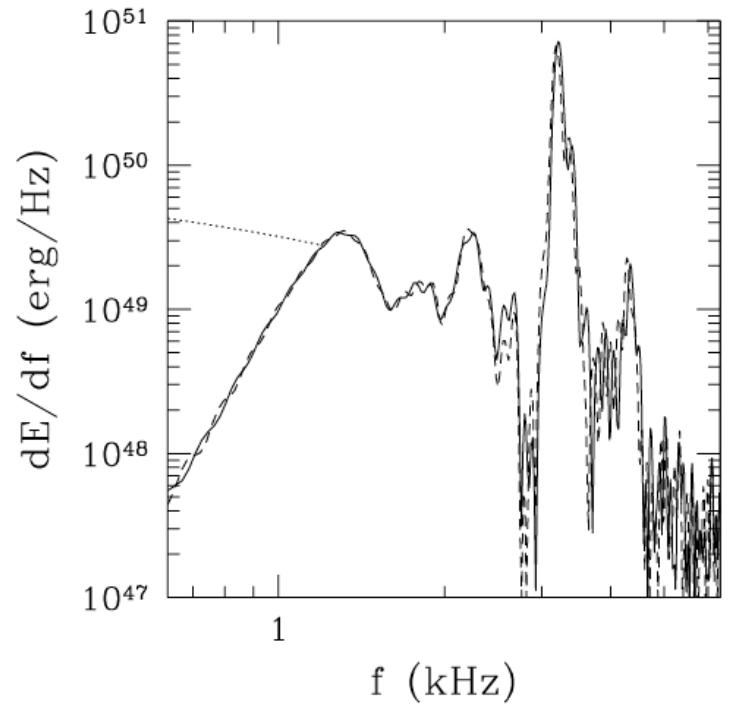


Cornell-Caltech collaboration

Gravitational waves probe NS EOS



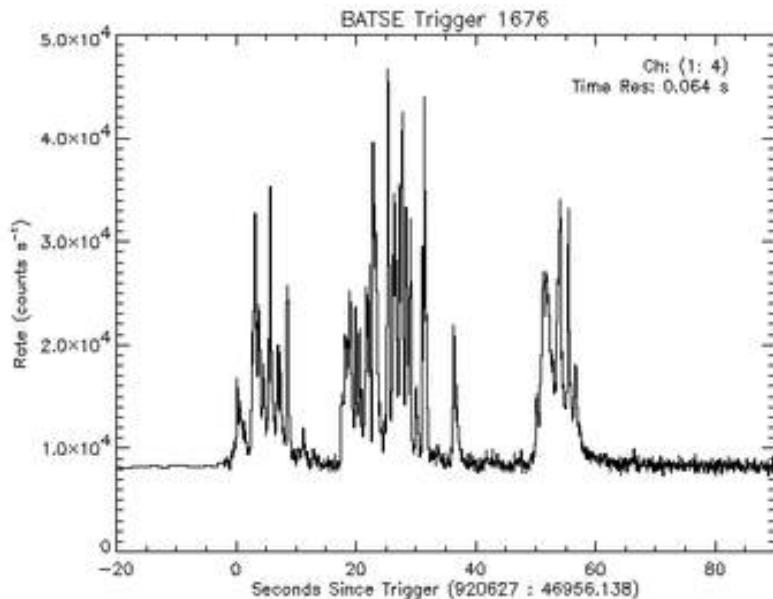
Masses well measured from inspiral waveform
Final cut-off frequency $\sim (GM/R^3)^{1/2}$



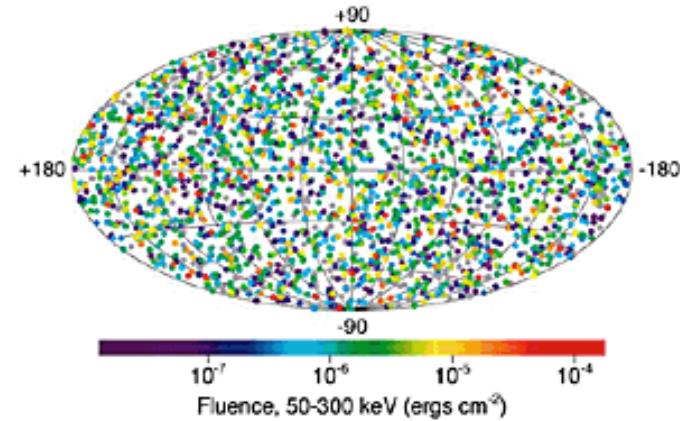
Cutler et al. '92;
DL & Wiseman '96;
Shibata et al.' 06;
...
Bauswein, Janka...' 12

NS/NS and NS/BH Mergers: Electromagnetic Counterparts

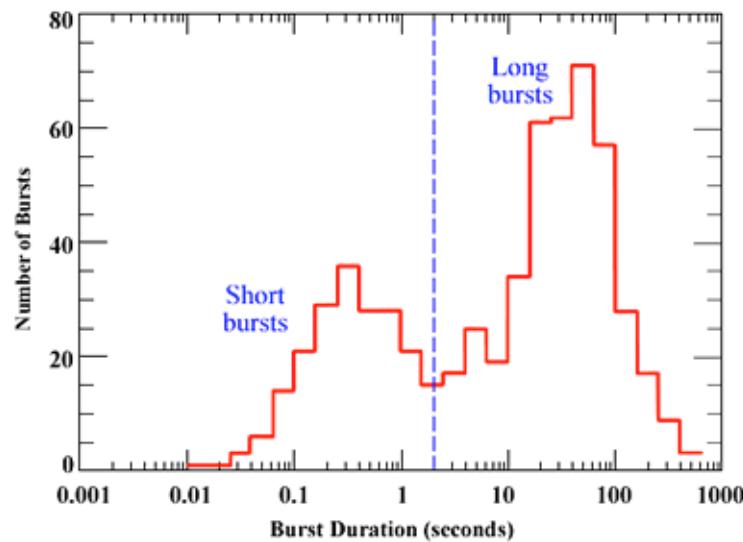
Gamma-Ray Bursts



2704 BATSE Gamma-Ray Bursts



Gamma-ray bursts come from all directions.



--Bursts of 0.1-10 MeV gamma-rays

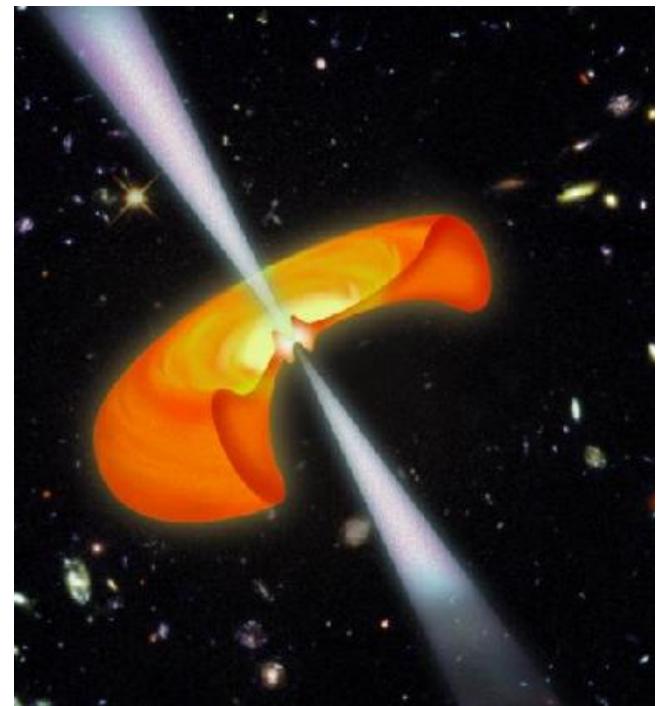
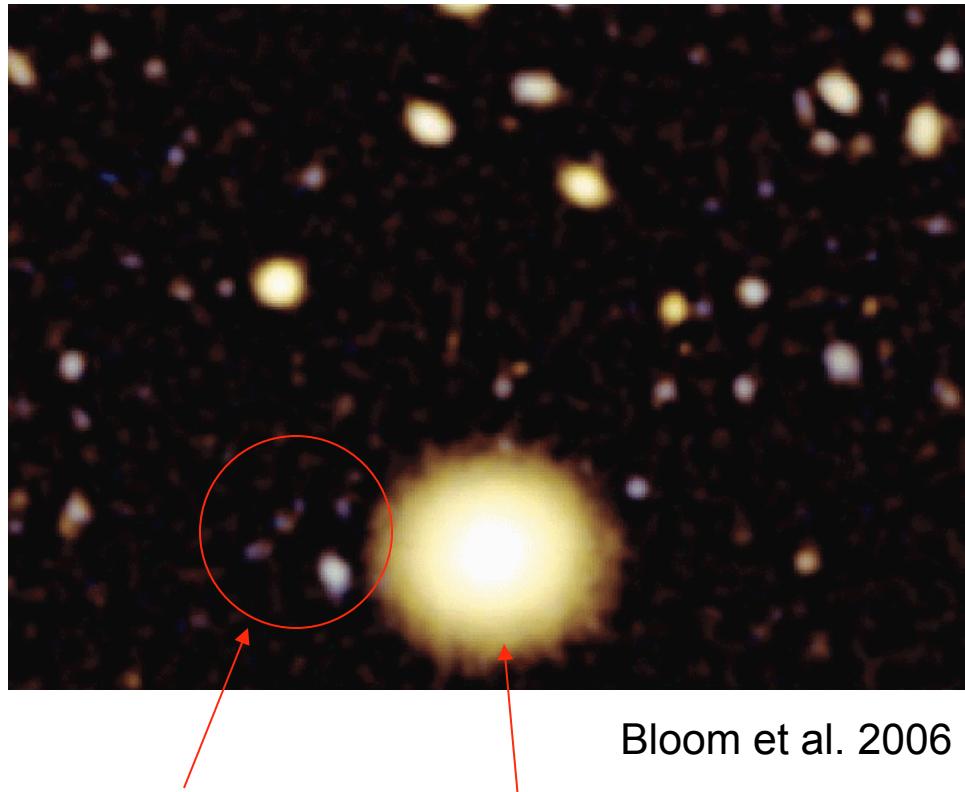
--From all directions, $z \sim 0.1-10$

--Very energetic $\sim 10^{48-55}$ erg

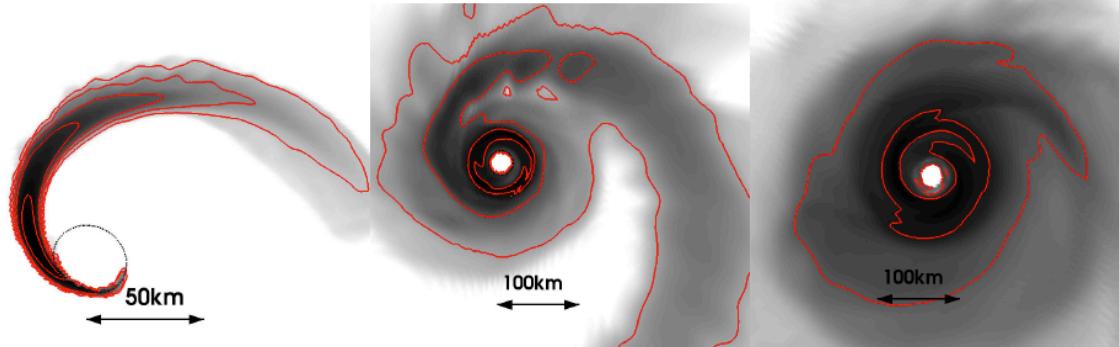
--Rare: GRB rate $\sim 10^{-6}/\text{yr/galaxy}$

-- "Long" ($\sim 30\text{s}$) and "short" ($\sim 0.3\text{s}$)

Merging NS/BH (or NS/NS?): Central Engine of Short GRBs



Merging NS/BH and NS/NS: Optical/IR Transients (?)



Foucart et al. (Cornell-Caltech)

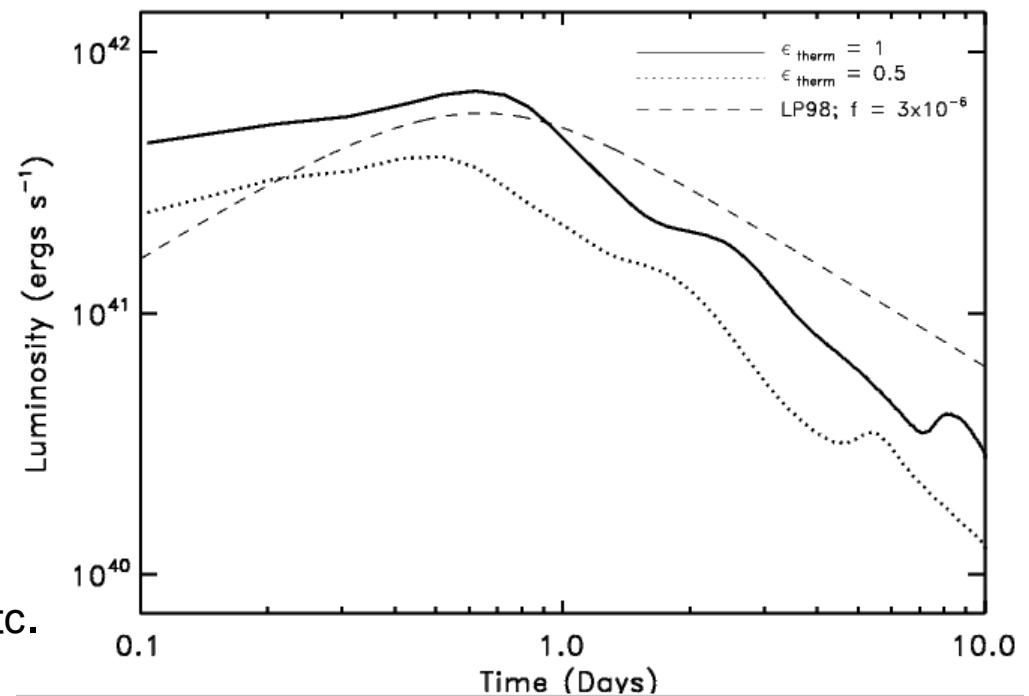
NS tidal ejecta $10^{-3} - 10^{-2} M_{\odot}$ (?)

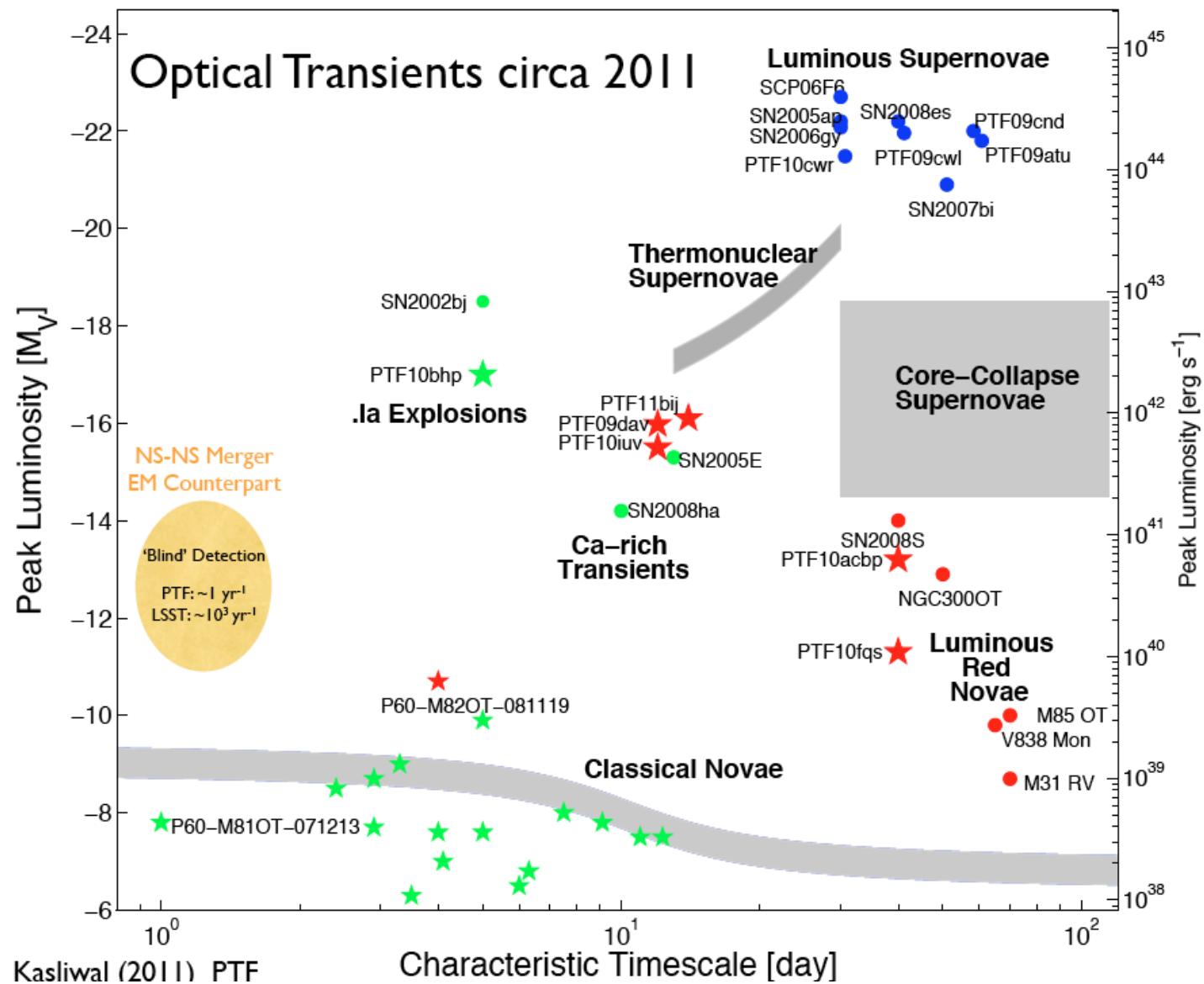
Ejecta evolution:

Initially mostly hot neutrons,
Decompression (cooling),
Nuclear reactions==> heating

$L \sim 3 \times 10^{41} \text{ erg s}^{-1}$ at $t \sim 1 \text{ day}$
 $T \sim 10^4 \text{ K}$ (optical)

Matzger, Quataert, etc.





Pre-Merger Phase:

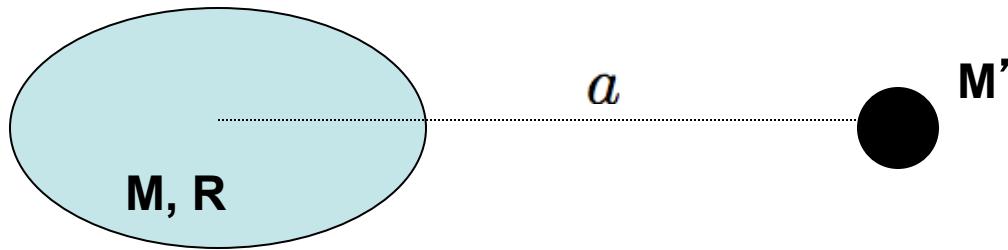
Anything interesting?

Pre-Merger Phase: Non-magnetic Neutron Stars

Tides

- Equilibrium tides**
- Dynamical tides**

Equilibrium Tide



$$V = -\frac{MM'}{a} - \mathcal{O}\left(k_2 \frac{M'^2 R^5}{a^6}\right) \quad k_2 = \text{Love number}$$

$$dN_{\text{GW}} = dN_{\text{GW}}^{(0)} \left[1 - \mathcal{O}\left(k_2 \frac{M'R^5}{Ma^5}\right) \right] \quad (\text{Missing GW cycles})$$

==> Important only at small separation (just prior to merger)

(Bildsten & Cutler 1992; Kochenek 92; DL, Rasio & Shapiro , etc)

Numerical GR Quasi-equilibrium NS binary sequence

(Baumgarte, Shapiro, Teukolsky, Shibata, Meudon group, etc. 1990s--200x)

Recent (semi-analytic) GR calculation of tidal effect

(Hinderer, Flanagan, Poisson, Damour, Penner, Andersson, Jones, etc., 2008+)

Dynamical Tides: Excitations of Internal Waves/Modes

NS has low-frequency oscillation modes:
g-modes (~ 100 Hz) (depends on symmetry energy)
inertial modes (incl. r-modes),...

Resonance: $\omega_\alpha = m\Omega_{\text{orb}}, \quad m = 2, 3, \dots$

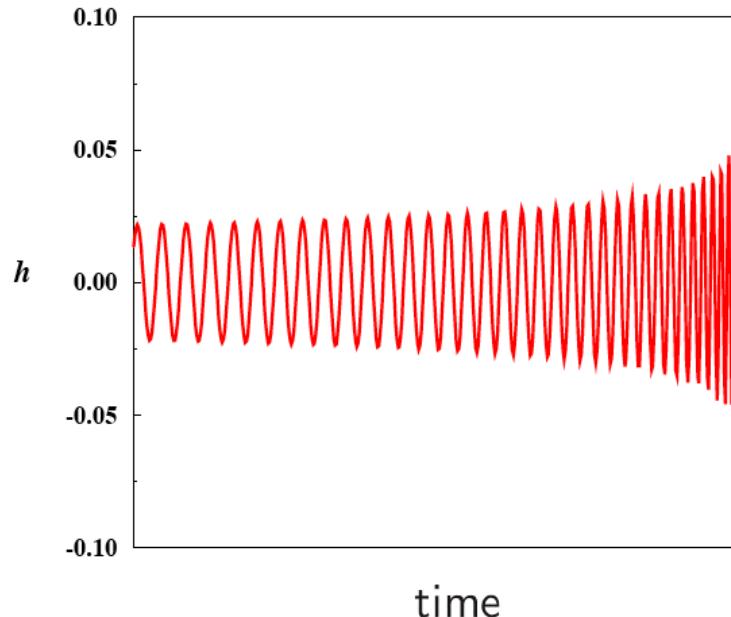
Dynamical Tides: Excitations of Internal Waves/Modes

==> Probe NS EOS using Inspiral Waveform

Rosonant tidal excitations of NS modes during inspiral

==> transfer orbital energy to NS

==> Missing GW cycles



Resonant Excitations of NS Oscillations During Inspiral

Non-rotating NS:

G-mode (Reisenegger & Goldreich 94; Shibata 94; DL 94)

Rotating NS:

G-mode, F-mode, R-mode (Ho & DL 99)

Inertial modes (DL & Wu 06)

R-mode (excited by gravitomagnetic force; Racine & Flanagan 06)

Results:

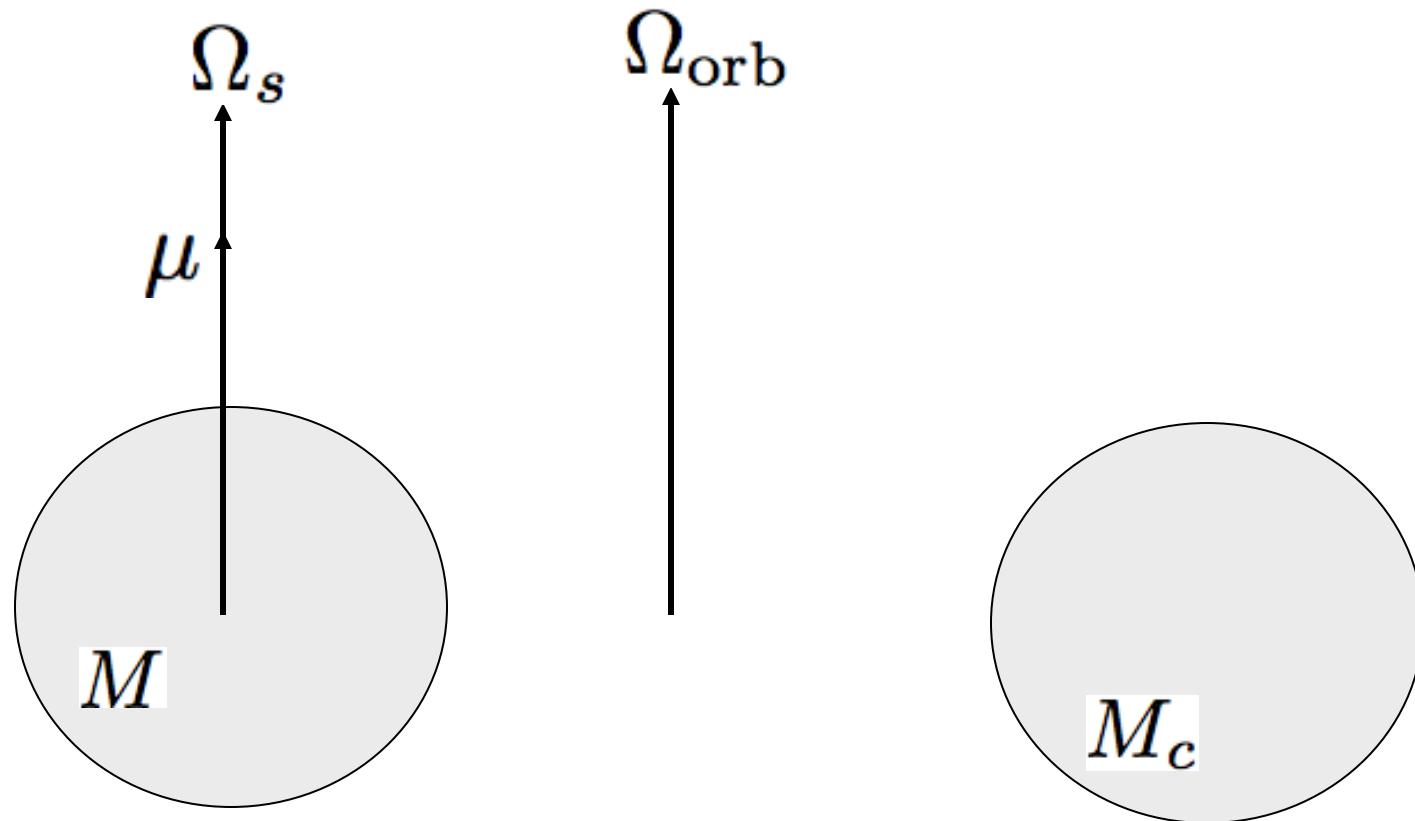
- For $R=10$ km NS, the number of missing cycles < 0.1 , barely measurable (unless NS is rapidly rotating)
- Number of missing cycles $\Delta N \propto R^4$ (g mode) or $R^{3.5}$ (r mode)
Important for larger NS
- **G-modes:** No law that requires ΔN should be < 1 ! (Newton & DL 2013)
- **Crustal modes:** Could shatter crust, pre-cursor of short GRB (Tsang et al. 12)

Pre-Merger Phase: Magnetic NSs

Cf. Double Pulsars: PSR J0737-3039

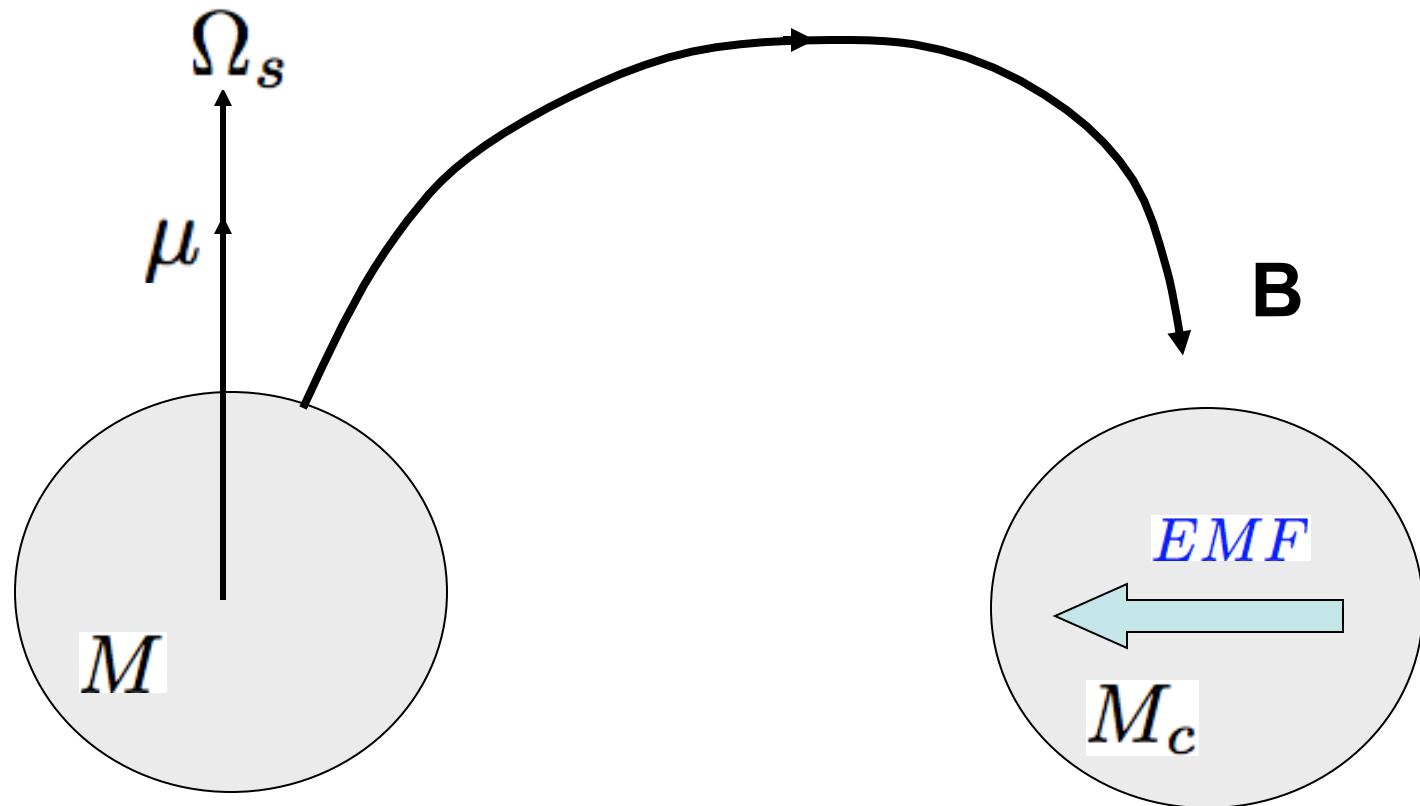
pulsar A: $\sim 10^{10}$ G

pulsar B: \sim a few $\times 10^{12}$ G



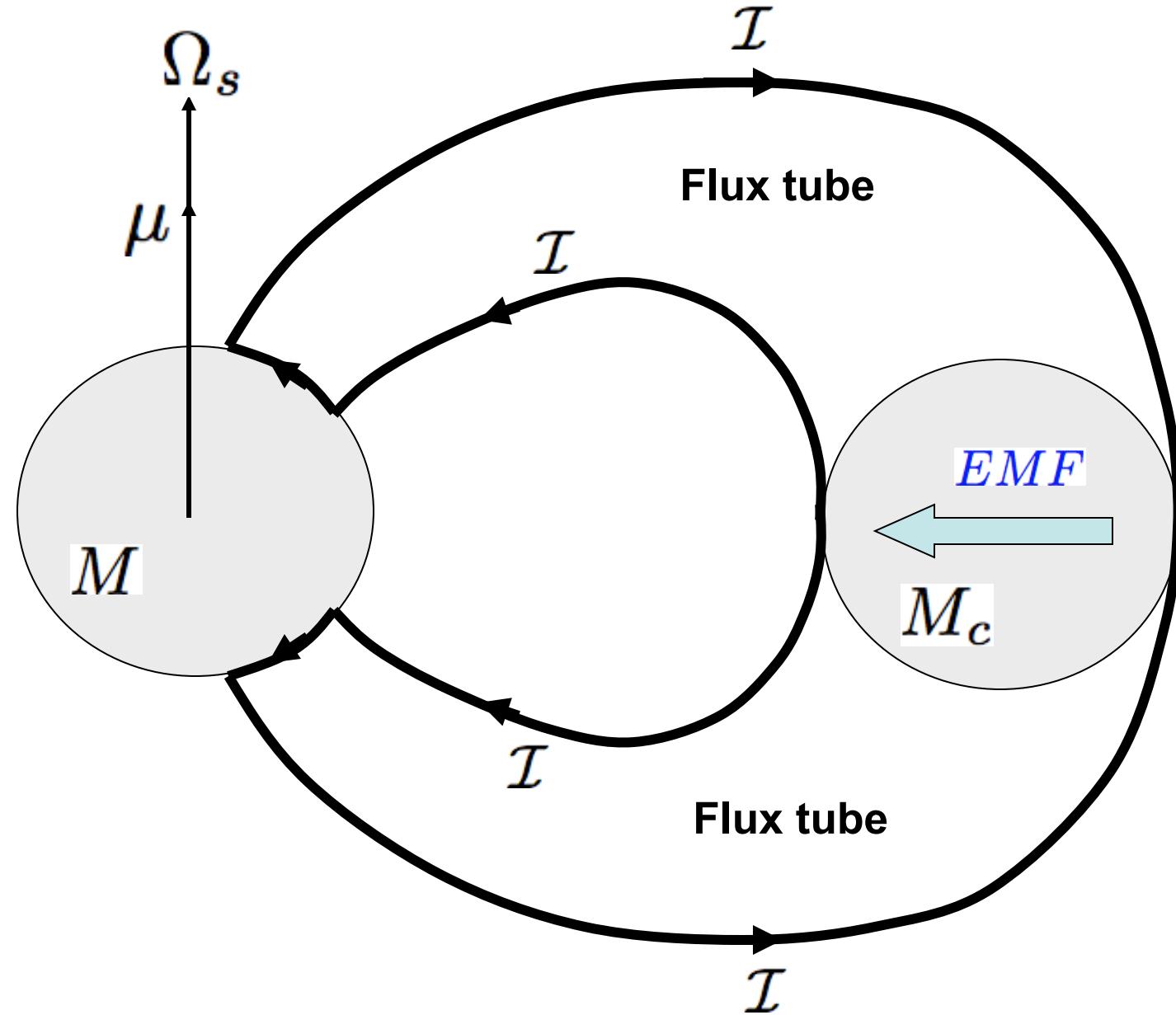
Consider a binary with

- magnetic NS ($>10^{12}$ G) + non-magnetic NS
- embedded in a tenuous plasma (magnetosphere)

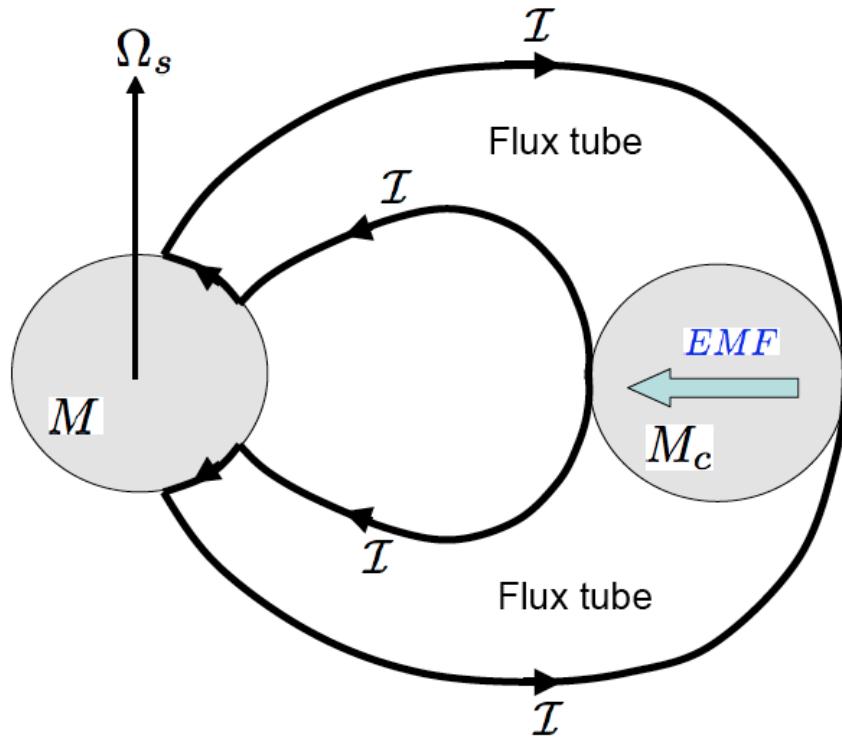


$$\text{EMF} : \Phi = 2R_c \left| \frac{\mathbf{v}}{c} \times \mathbf{B} \right|$$

e.g. $\Phi \sim 10^{13}$ Volt at $f_{\text{orb}} = 20$ Hz



DC Circuit Powered by Orbital Motion



$$\text{EMF} : \Phi = \frac{2\mu R_c}{ca^2}(\Omega_{\text{orb}} - \Omega_s)$$

$$\text{Current} : \mathcal{I} = \frac{\Phi}{\mathcal{R}}$$

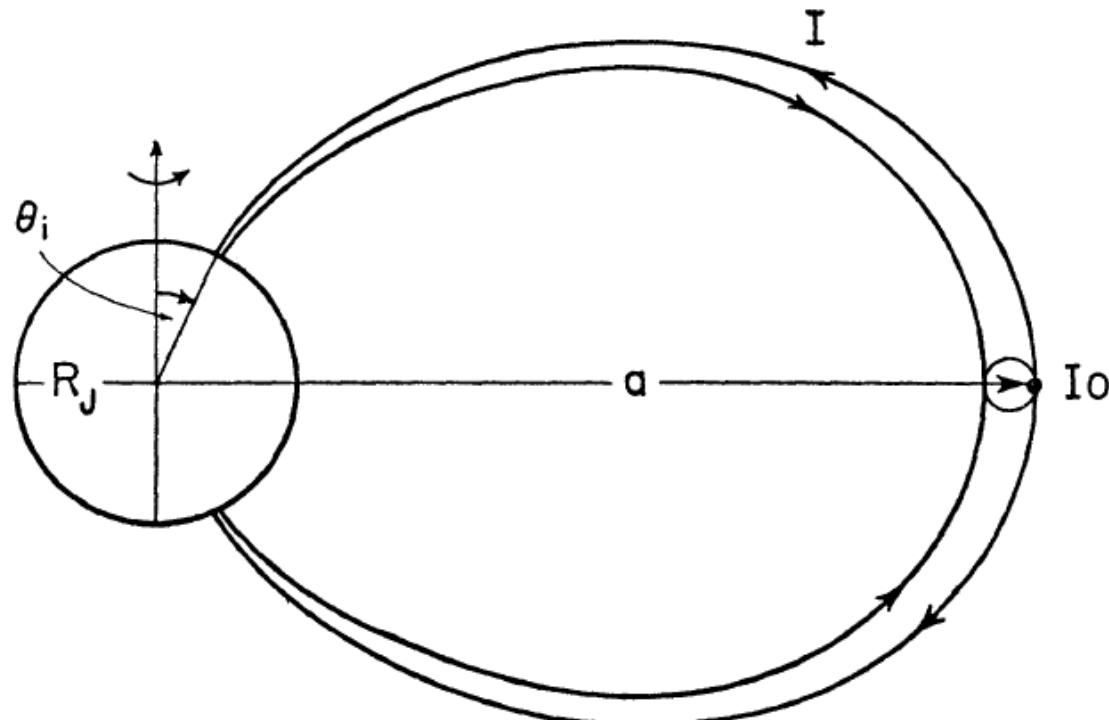
$$\text{Dissipation} : \dot{E}_{\text{diss}} = \frac{\Phi^2}{\mathcal{R}}$$

THE ASTROPHYSICAL JOURNAL, Vol. 156, April 1969
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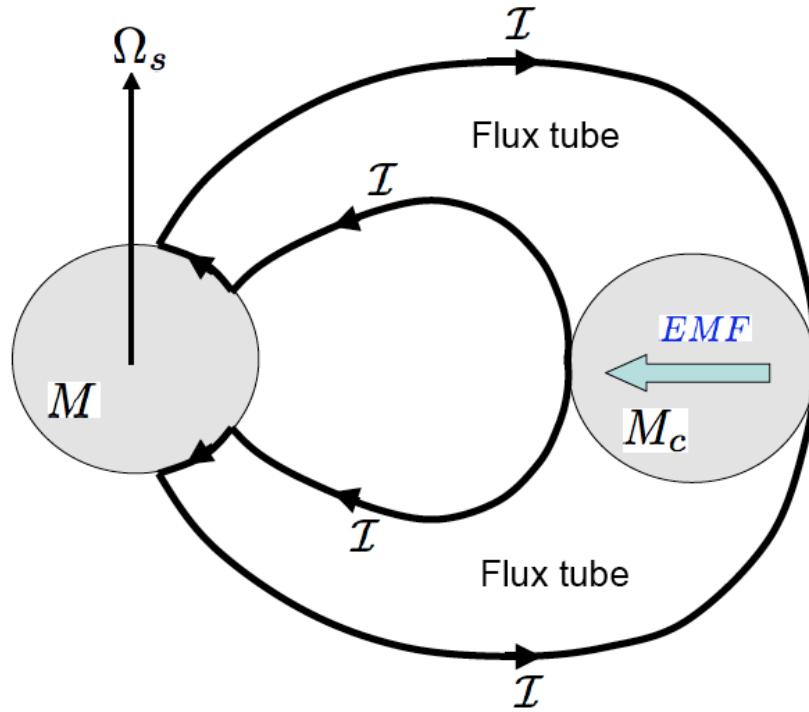
IO, A JOVIAN UNIPOLAR INDUCTOR

PETER GOLDREICH*
California Institute of Technology

AND
DONALD LYNDEN-BELL
Royal Greenwich Observatory



DC Circuit Powered by Orbital Motion



Applications to:

WD-WD Binaries

(K.Wu et al. 02,09; Dall'Osso, Israel, Stella 06,07)

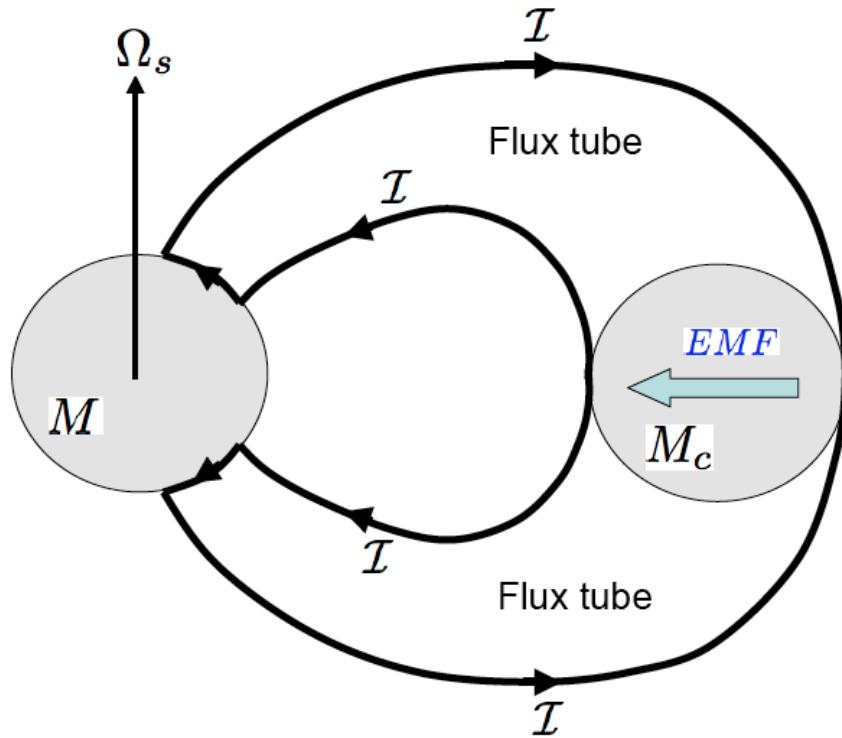
NS-NS, NS-BH Binaries

(Hansen & Lyutikov 01; McWilliams & Levin 11; Piro 1

Exoplanetary systems (Laine & Lin 12,...)

Caution: Some of these were wrong

DC Circuit Powered by Orbital Motion



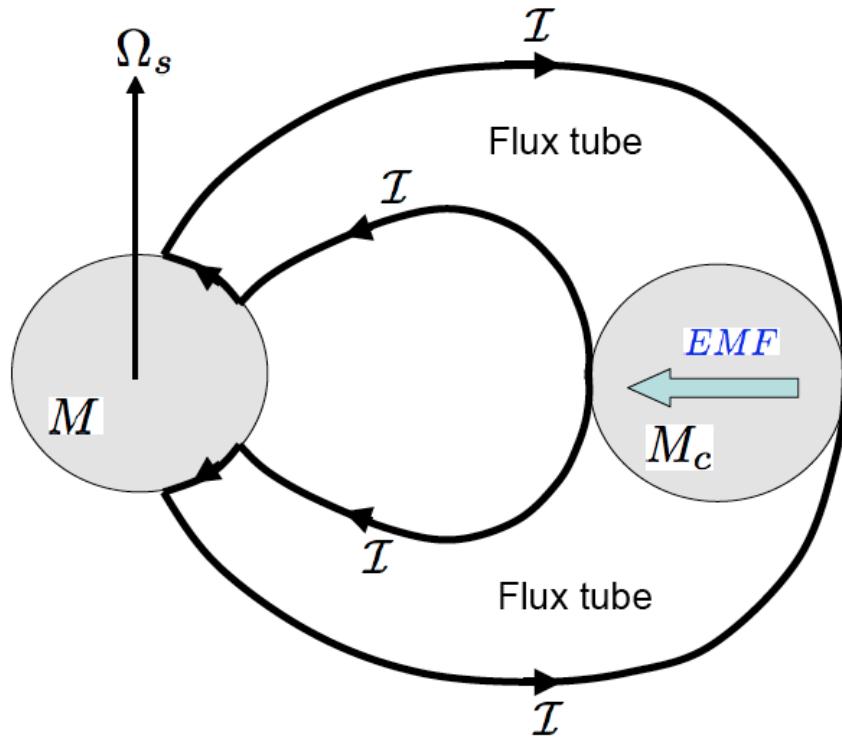
$$\text{EMF} : \Phi = \frac{2\mu R_c}{ca^2}(\Omega_{\text{orb}} - \Omega_s)$$

$$\text{Current} : \mathcal{I} = \frac{\Phi}{\mathcal{R}}$$

$$\text{Dissipation} : \dot{E}_{\text{diss}} = \frac{\Phi^2}{\mathcal{R}}$$

Results depend on the resistance: \mathcal{R}

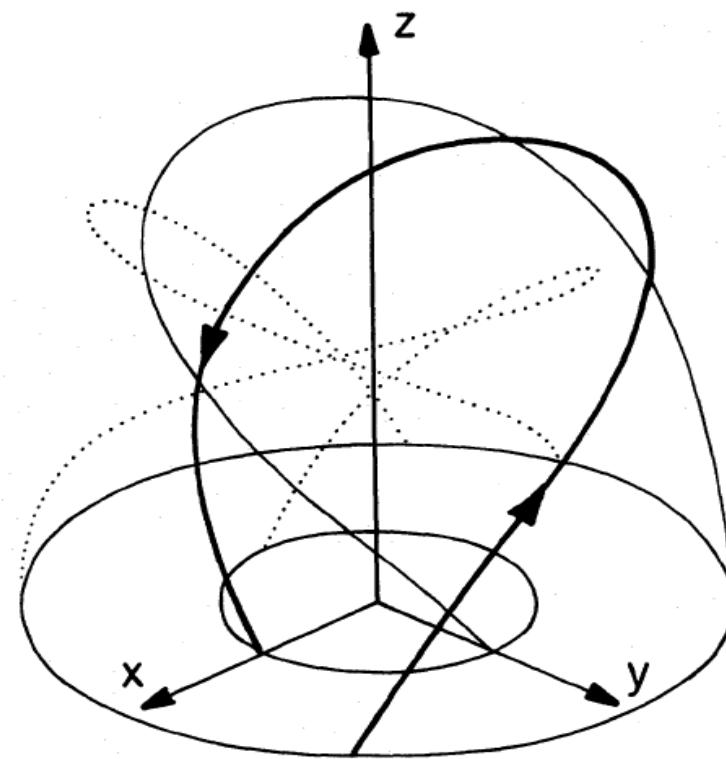
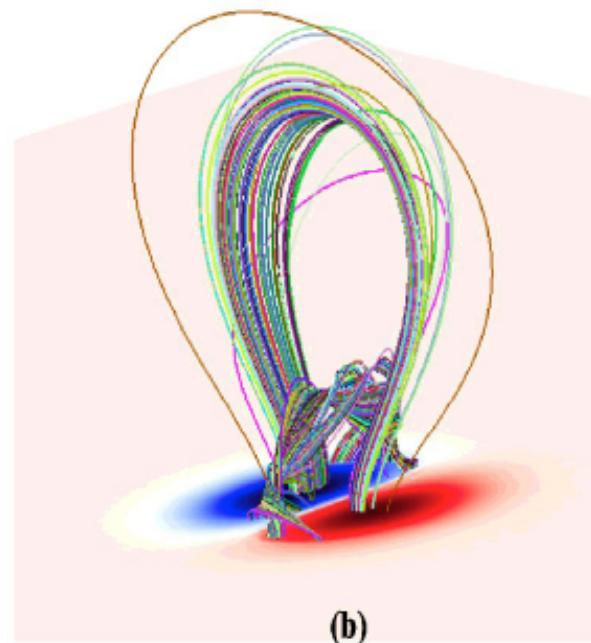
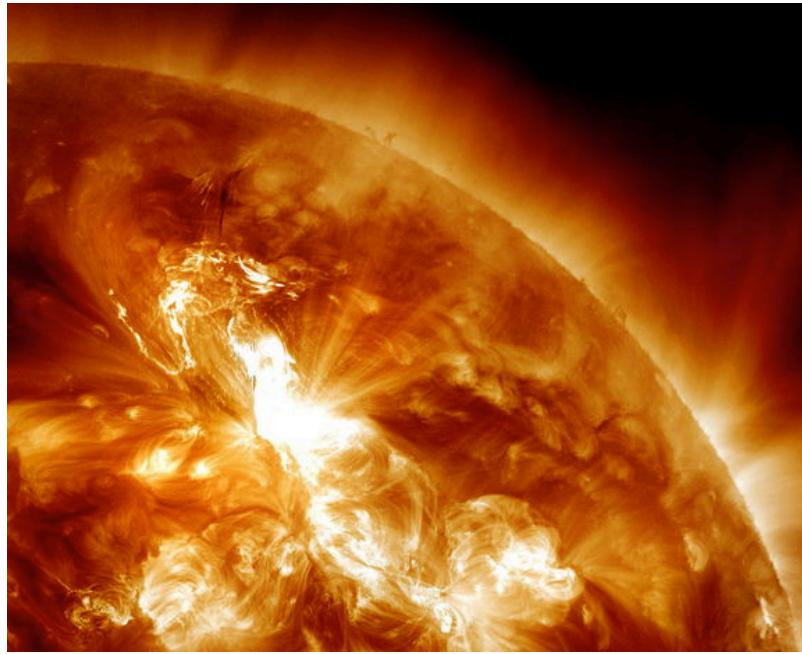
DC Circuit Powered by Orbital Motion



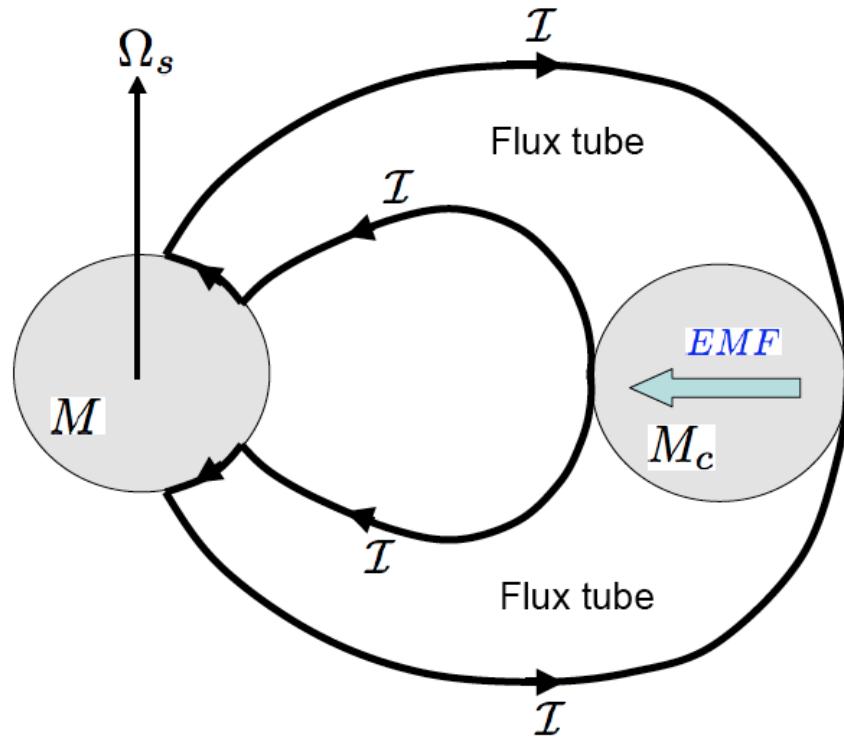
Problems with small \mathcal{R}
(\rightarrow large \mathcal{I}):

Flux tube is twisted

$$\frac{|B_\phi|}{|B_z|} \sim \frac{16 v/c^2}{\mathcal{R}}, \quad v = (\Omega_{\text{orb}} - \Omega_s)a$$



DC Circuit Powered by Orbital Motion



$$\frac{|B_\phi|}{|B_z|} \sim \frac{16v/c^2}{\mathcal{R}}, \quad v = (\Omega_{\text{orb}} - \Omega_s)a$$

Circuit will break when $|B_\phi|/|B_z| \gtrsim 1$

Energy Dissipation in the Magnetosphere of Pre-merging NS Binary

DL 2012

$$\dot{E}_{\max} \simeq 7 \times 10^{44} \left(\frac{B_{\text{NS}}}{10^{13} \text{ G}} \right)^2 \left(\frac{a}{30 \text{ km}} \right)^{-13/2} \text{ erg s}^{-1}$$

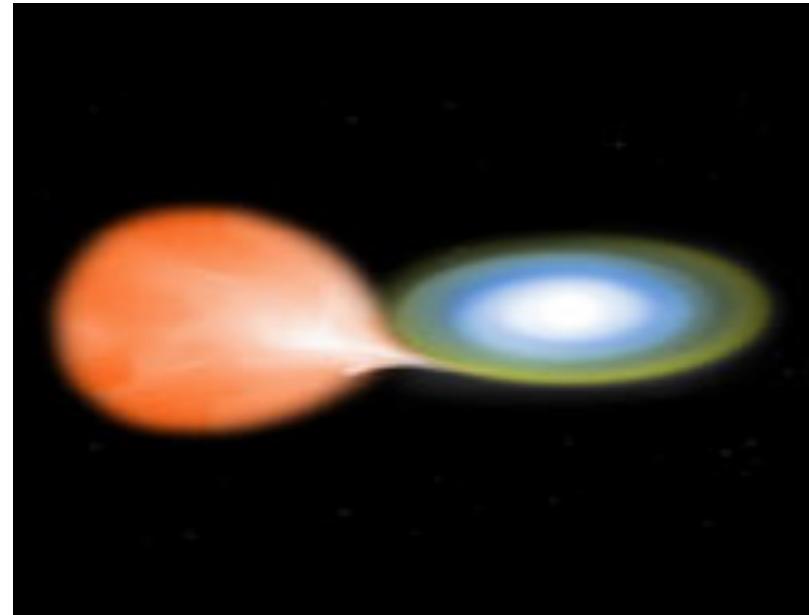
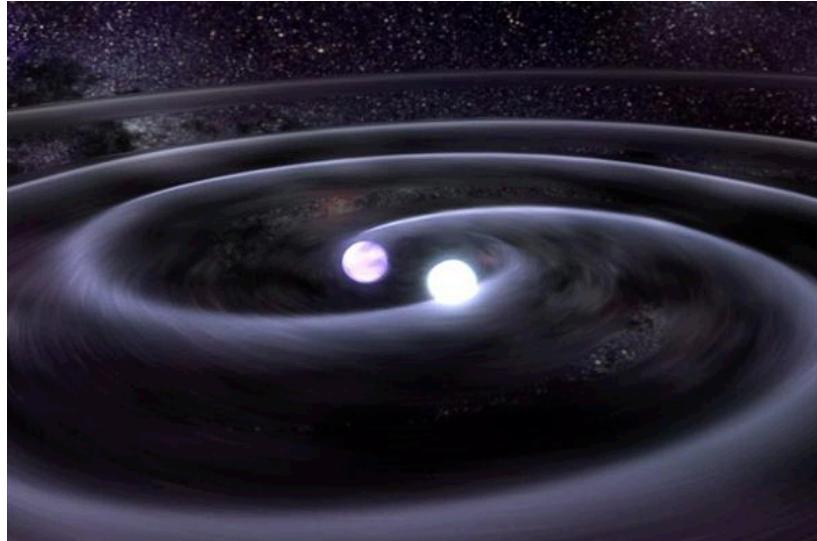
Actual dissipation rate:

$$\dot{E} \sim 2 \times 10^{44} \left(\frac{B_{\text{NS}}}{10^{13} \text{ G}} \right)^2 \left(\frac{a}{30 \text{ km}} \right)^{-7} \text{ erg s}^{-1}$$

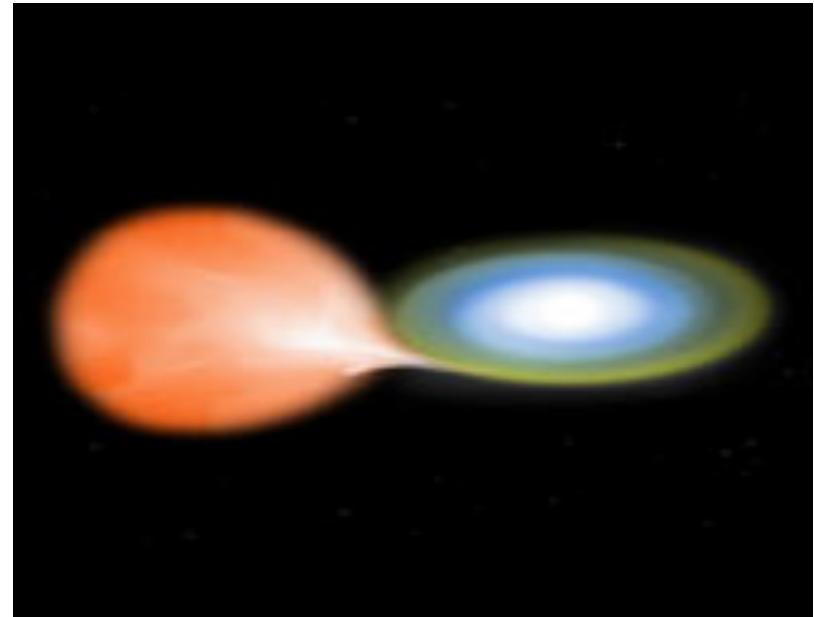
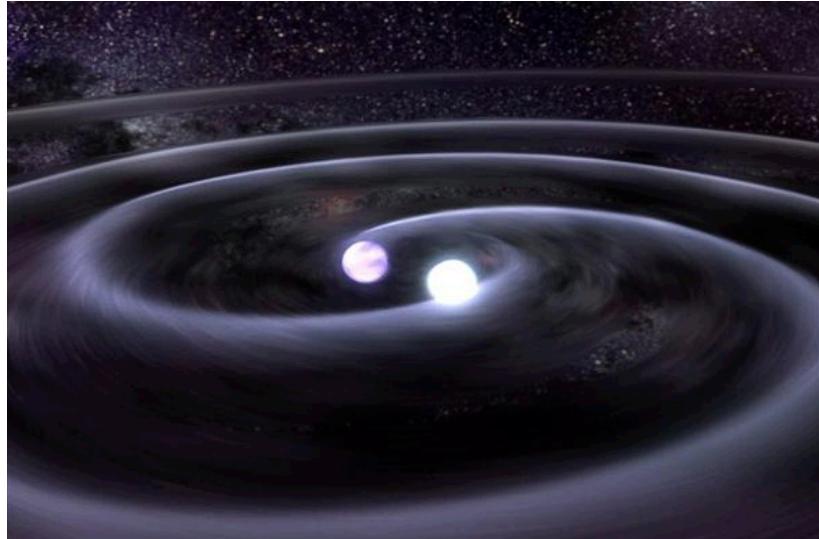
- This Edot will not affect orbital decay rate (GW signal)
- Radio emission prior to binary merger (?) cf. Vietri 96; Hansen & Lyutikov 01

cf. isolated pulsars: $\dot{E} \simeq 10^{33} \left(\frac{B_{\text{NS}}}{10^{13} \text{ G}} \right)^2 \left(\frac{P}{1 \text{ s}} \right)^{-4} \text{ erg s}^{-1}$

Compact White Dwarf Binaries (mins - hour)



Compact White Dwarf Binaries (mins - hour)



-- Dominant sources of gravitational waves (10^{-4} -0.1 Hz)

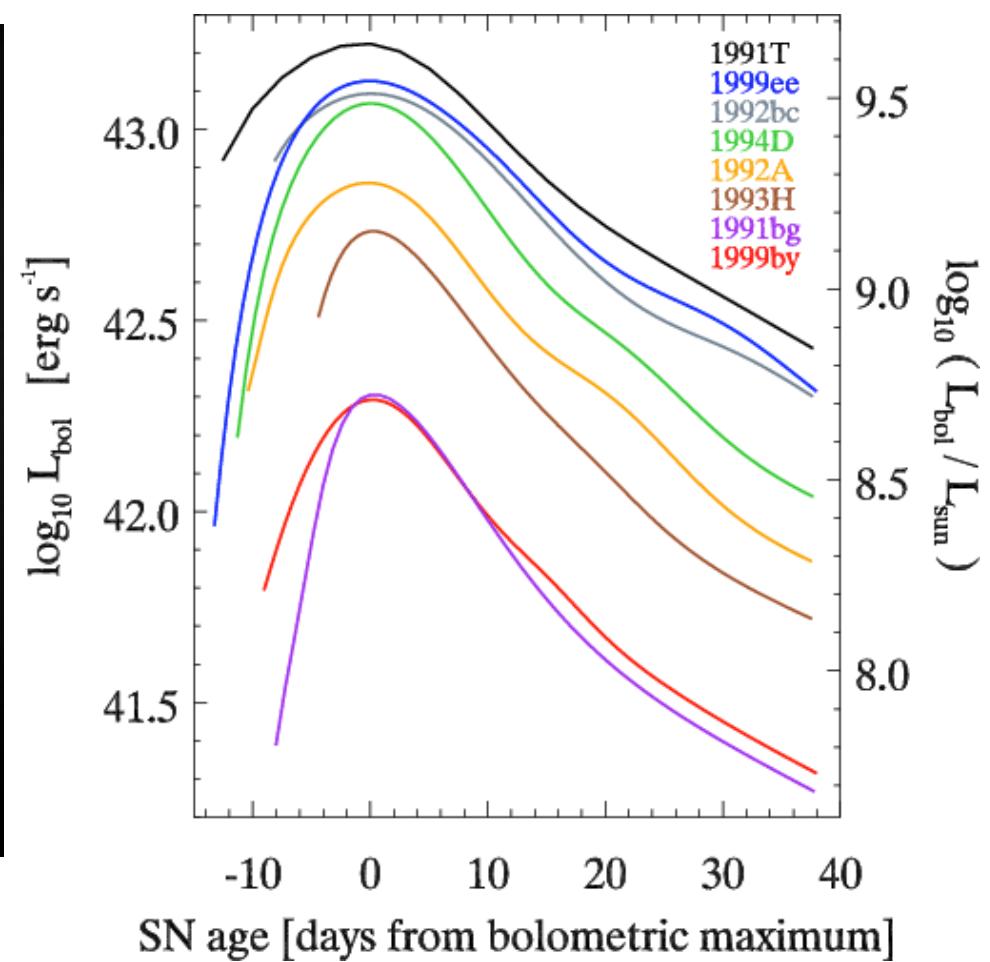
Space interferometer (eLISA-NGO??)

-- Lead to various outcomes:

R CrB stars, AM CVn binaries, transients

If total mass $\sim 1.4M_{\text{sun}}$: AIC => NS or SN Ia

Type Ia Supernovae



Type Ia Supernovae

Thermonuclear explosion of CO white dwarfs of $\sim 1.4 M_{\text{sun}}$

Progenitors ??

WD + non-deg. star: “Single-degenerate” Scenario

WD + WD merger: “Double-degenerate” Scenario

WD + WD collision ?

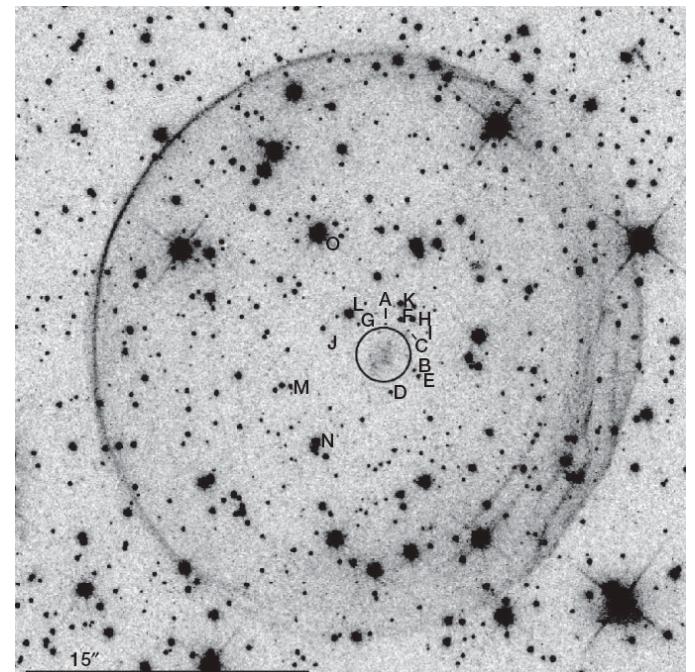
Various arguments for/against each scenario:

Rates, super-soft sources, delay time...

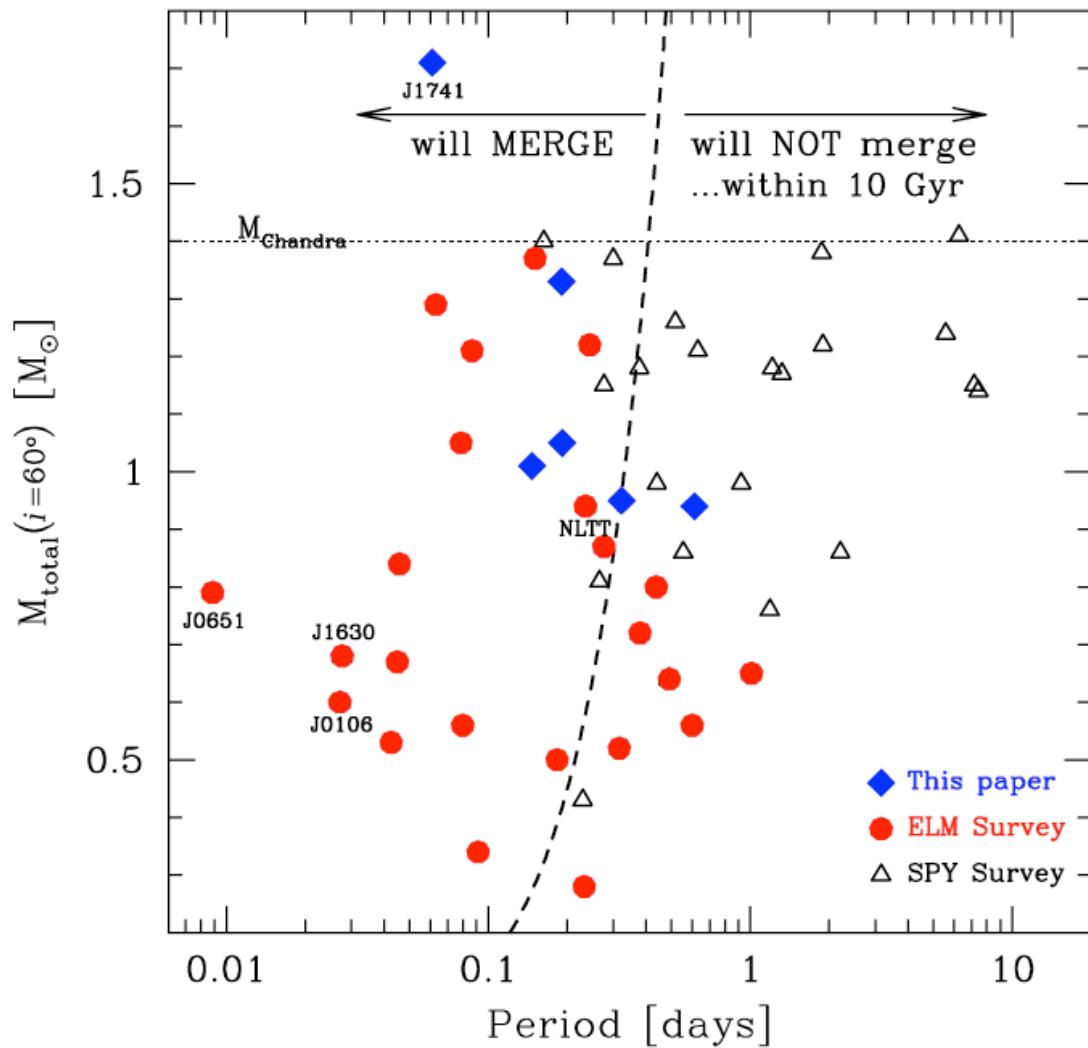
Recent observations in favor of DD:

e.g., Absence of ex-companion stars
in SN Ia remnant SNR 0509-67.5
==> rule out V=26.9

Schaefer & Pagnotta 2012
(cf Di Stefano & Kilic 2012)

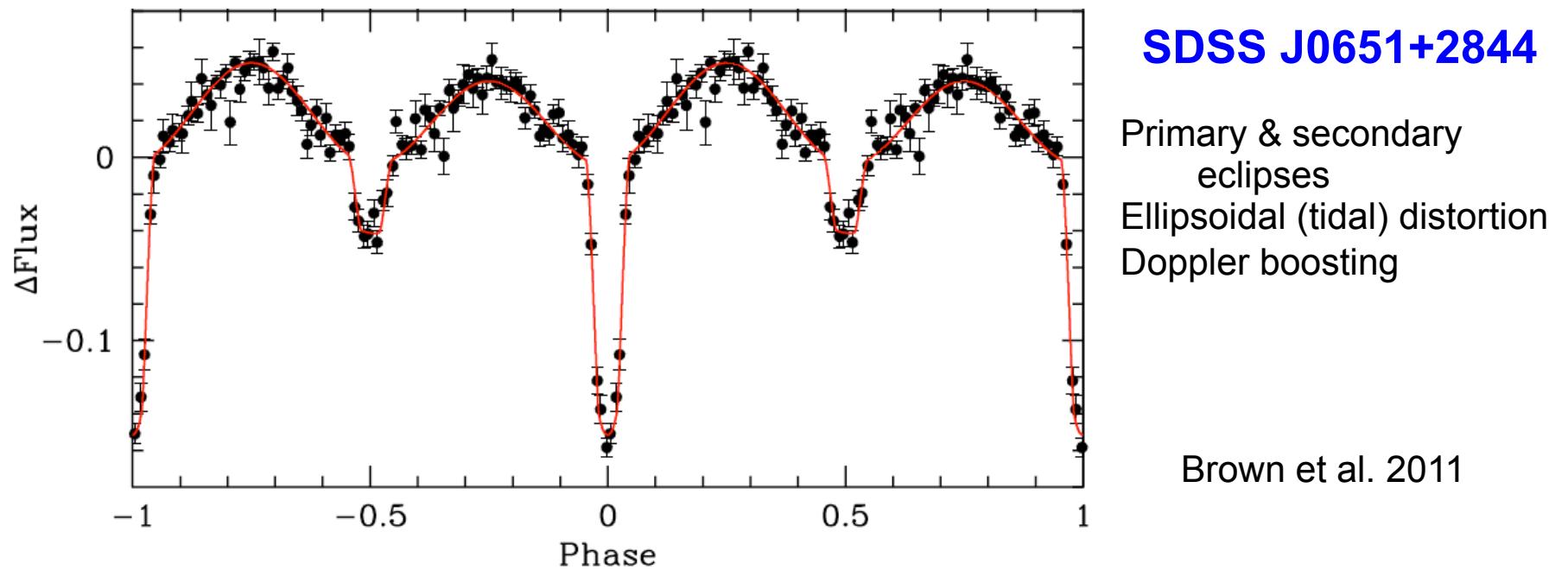


Radial Velocity Surveys of Compact WD Binaries



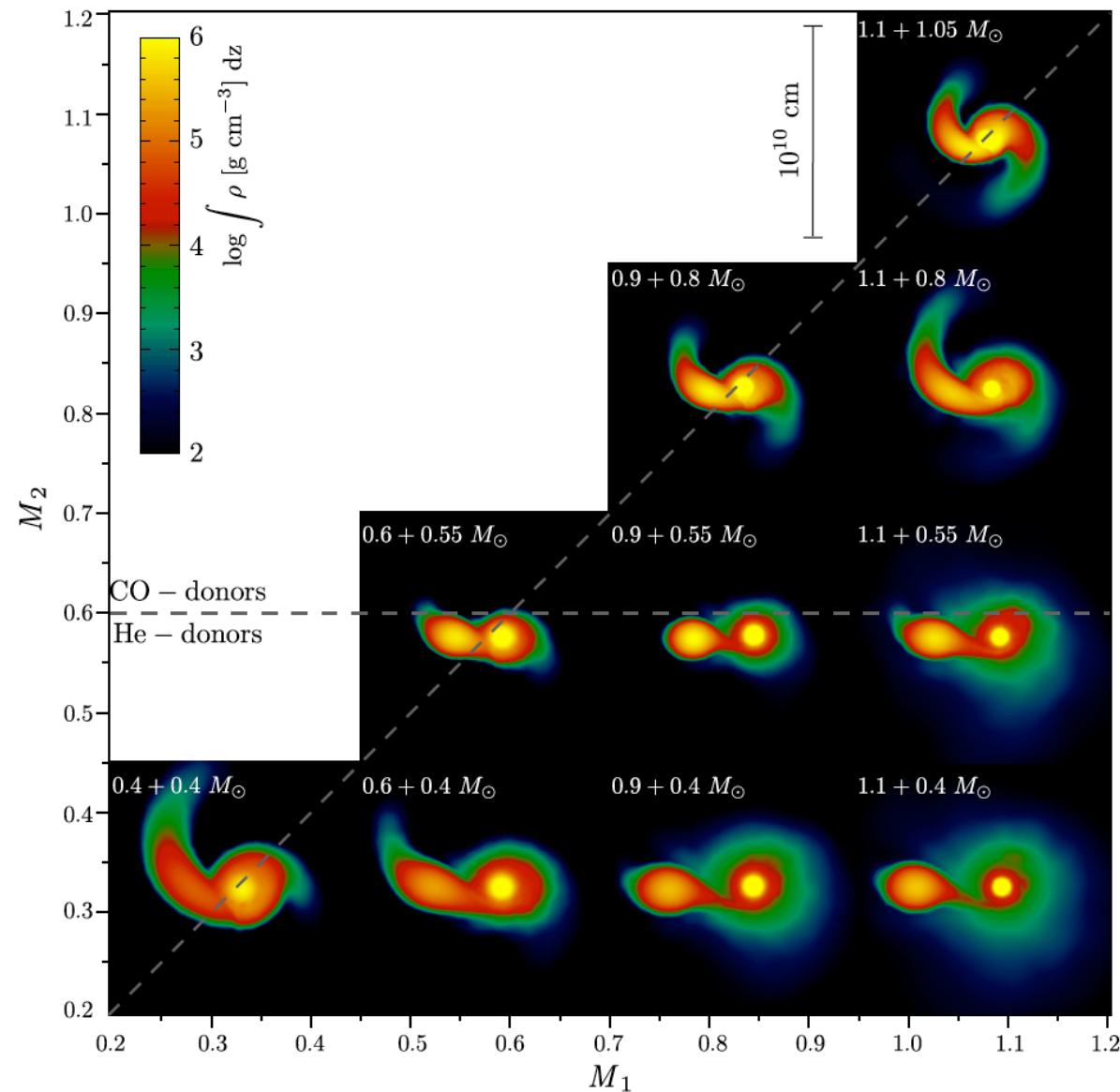
Brown et al. 2012

12 min orbital period double WD eclipsing binary

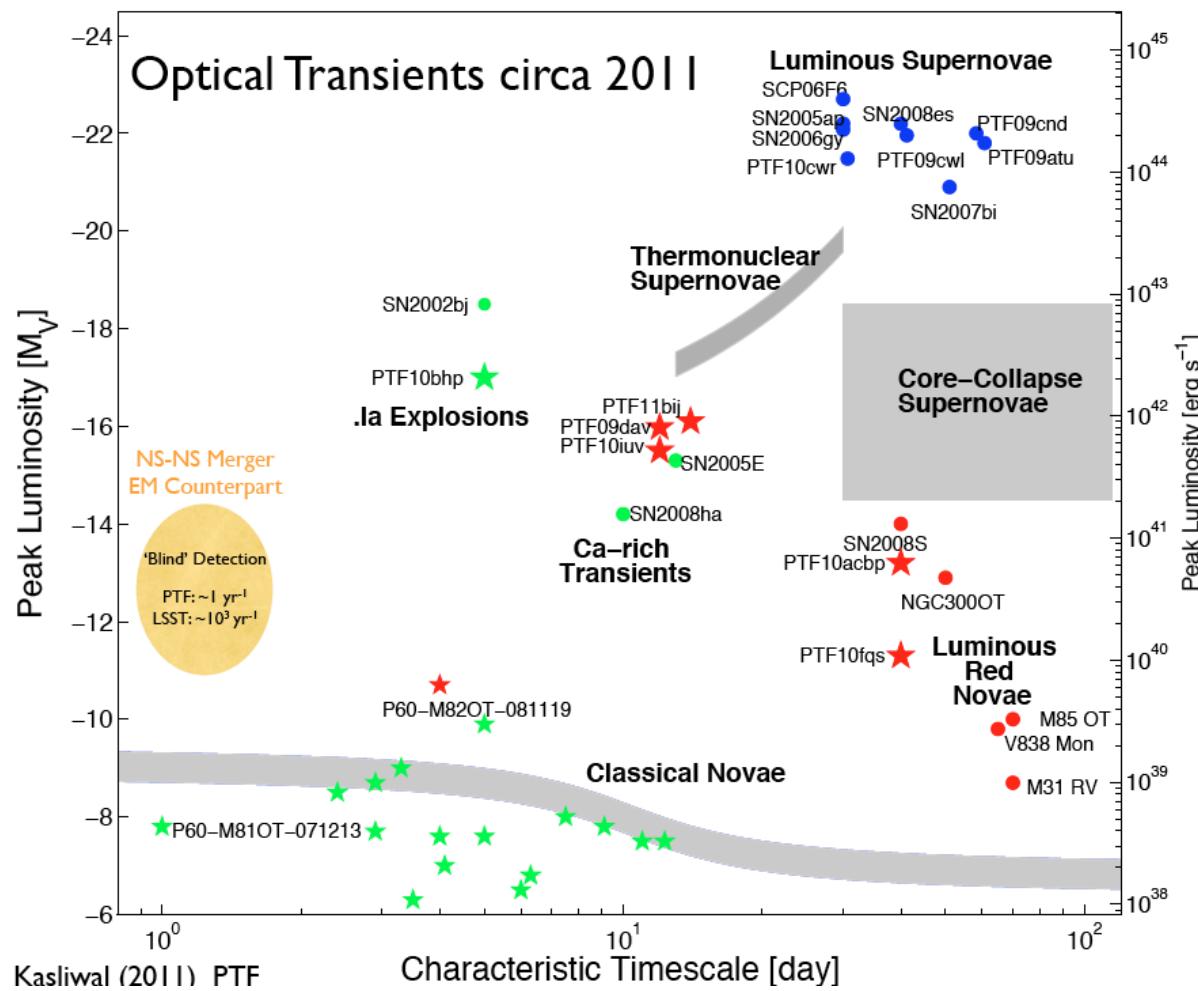


- will merge in 0.9 Myr
- large GW strain ==> (LISA)
- orbital decay measurable from eclipse timing (Hermes et al 2012)

WD Binary Merger



Dan, Rosswog, et al. 2012



WD binary merger: Outcome depends on WD masses, composition, and pre-merger conditions

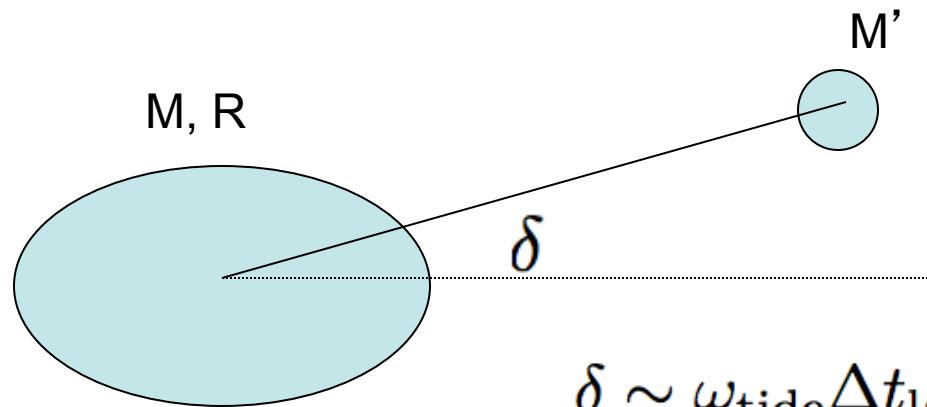
Dynamical Tides in Compact WD Binaries

Jim Fuller & DL

Issues:

- Spin-orbit synchronization?
- Tidal dissipation and heating?
- Effect on orbital decay rate? (e.g. eLISA-NGO)

Equilibrium Tide



$$\delta \sim \omega_{\text{tide}} \Delta t_{\text{lag}} \sim 1/Q$$

$$\omega_{\text{tide}} = 2(\Omega_{\text{orb}} - \Omega_s)$$

$$\text{Torque} \sim G \left(\frac{M'}{a^3} \right)^2 R^5 \delta$$

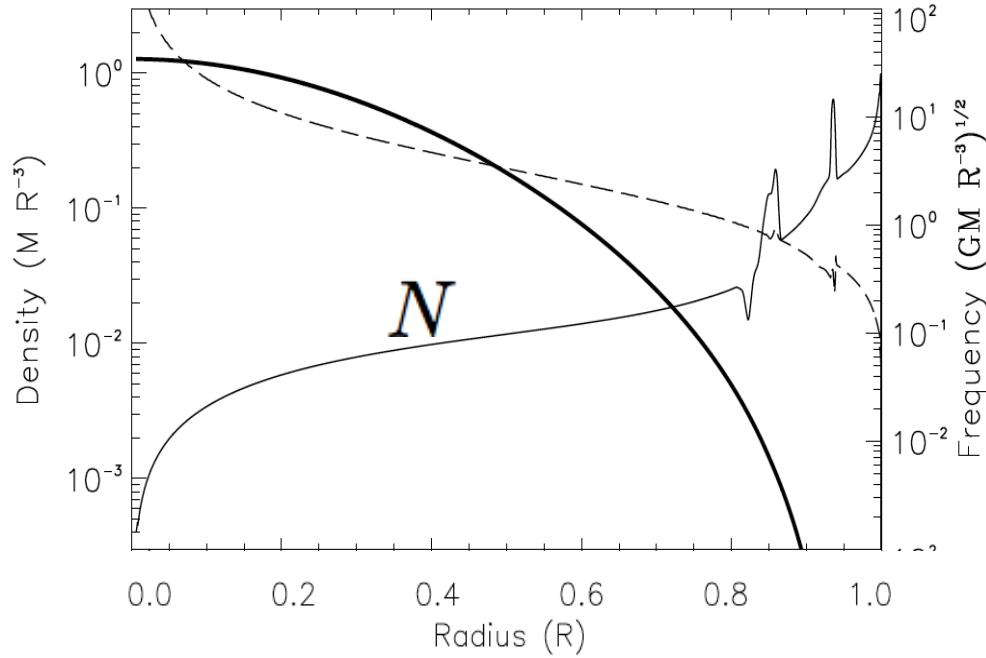
$$\dot{E}_{\text{tide}} = \text{Torque} \cdot \Omega$$

Problems:

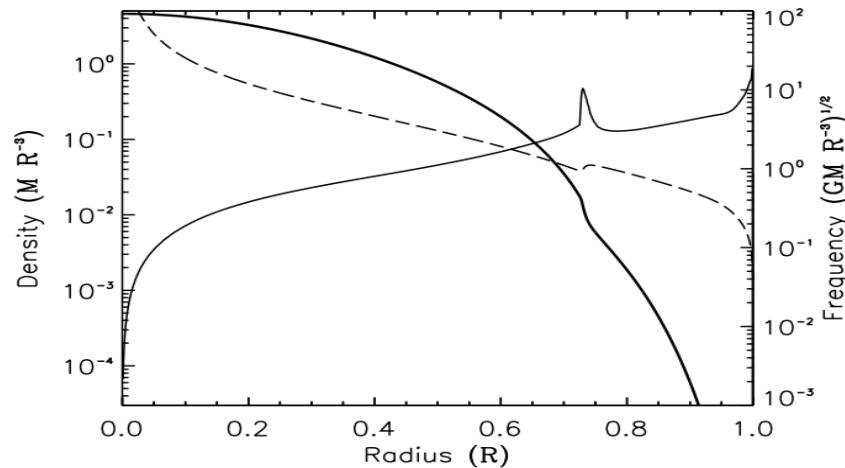
- Parameterized theory
- The physics of tidal dissipation is more complex:

Excitation/damping of internal waves/modes (Dynamical Tides)

Wave Propagation inside White Dwarf



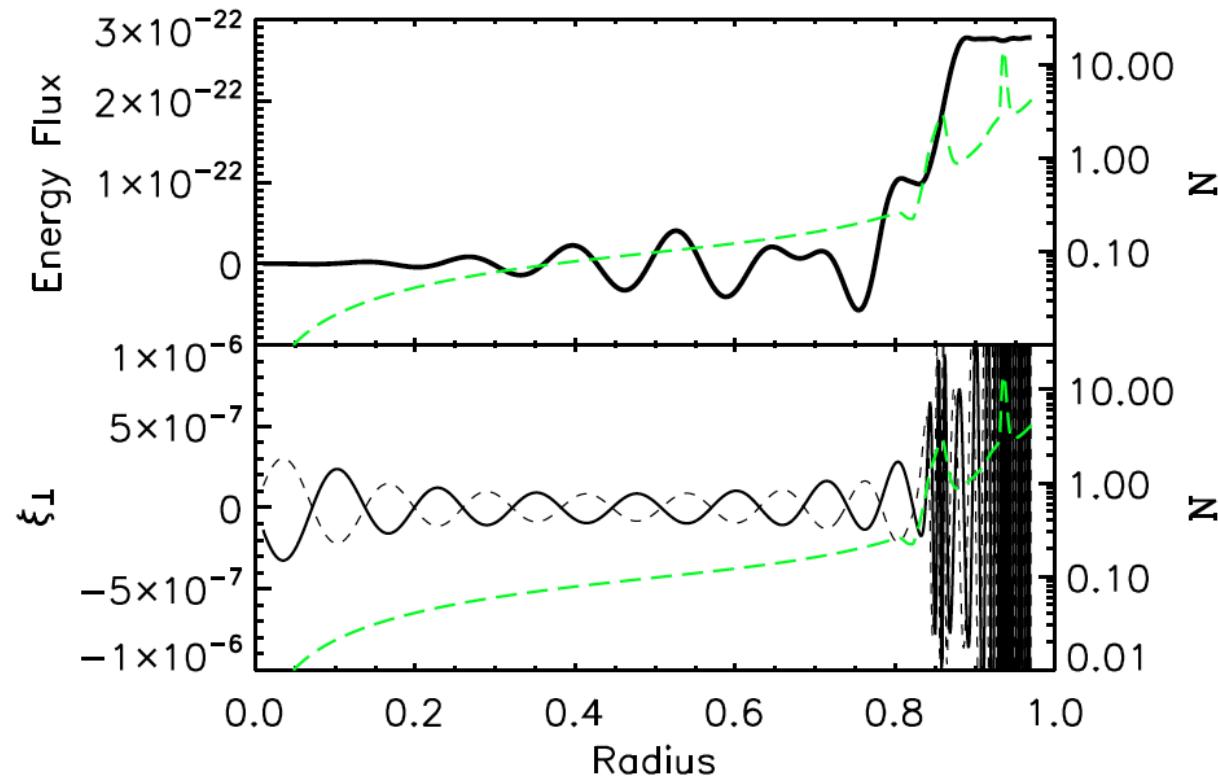
CO WD
 $0.6M_{\odot}, 8720\text{ K}$



He-core WD
 $0.3M_{\odot}, 12000\text{ K}$

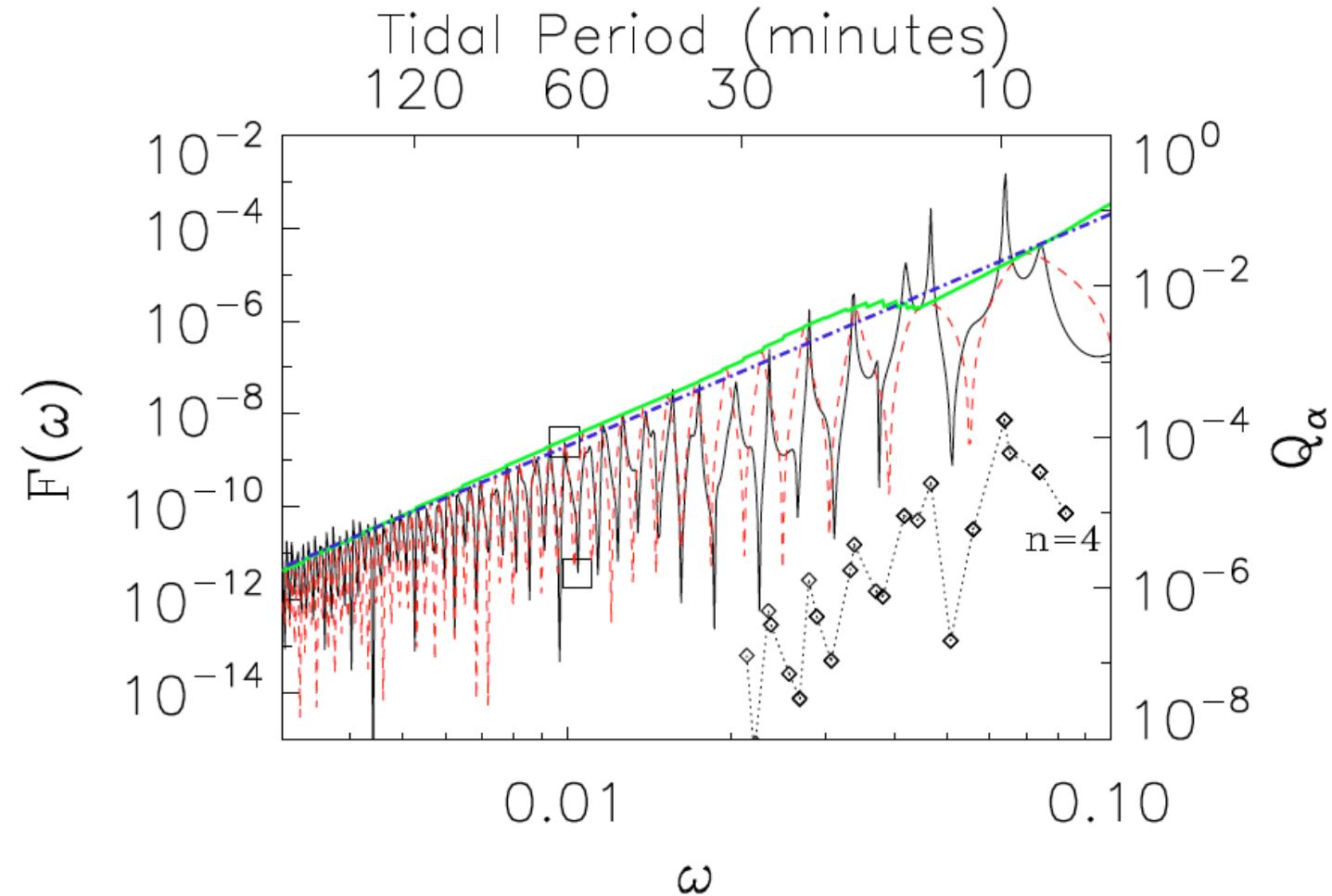
“Continuous” Excitation of Gravity Waves

Waves are excited in the interior/envelope, propagate outwards and dissipate near surface

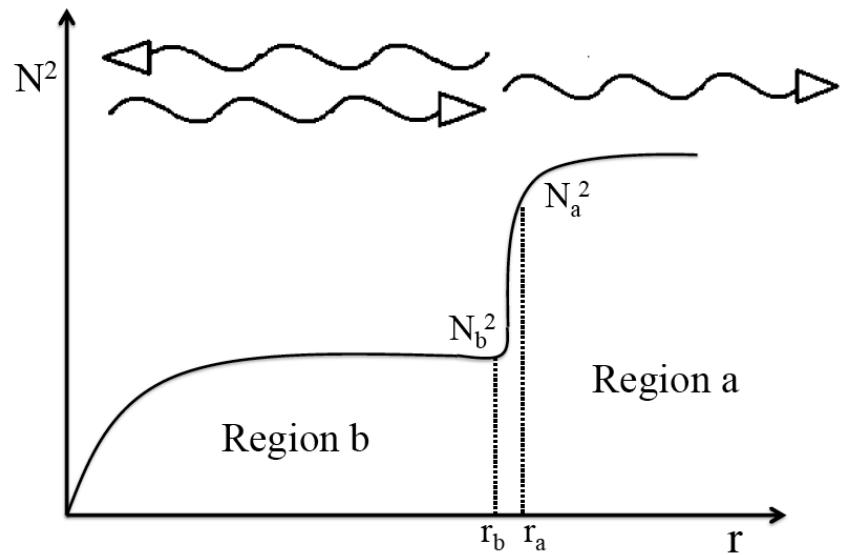
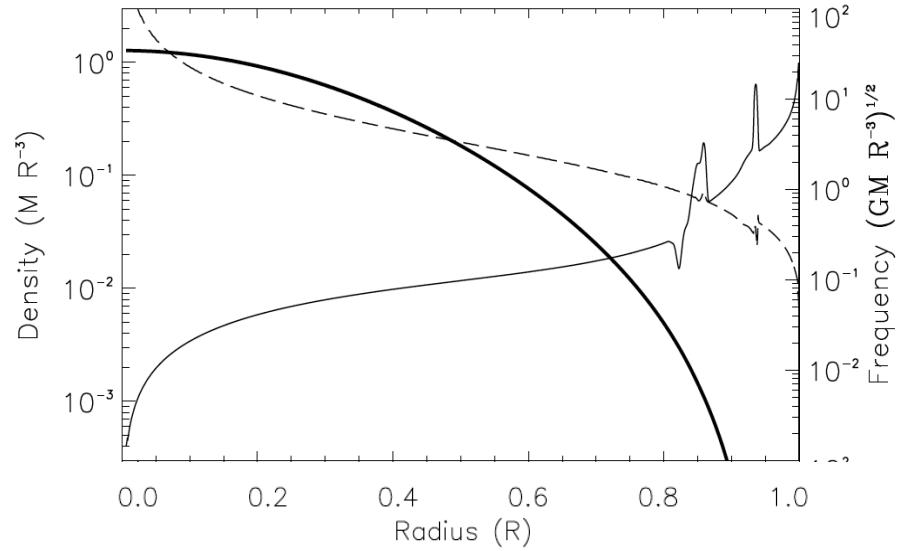


$$M = 0.6M_{\odot}, \quad \omega = 0.01$$

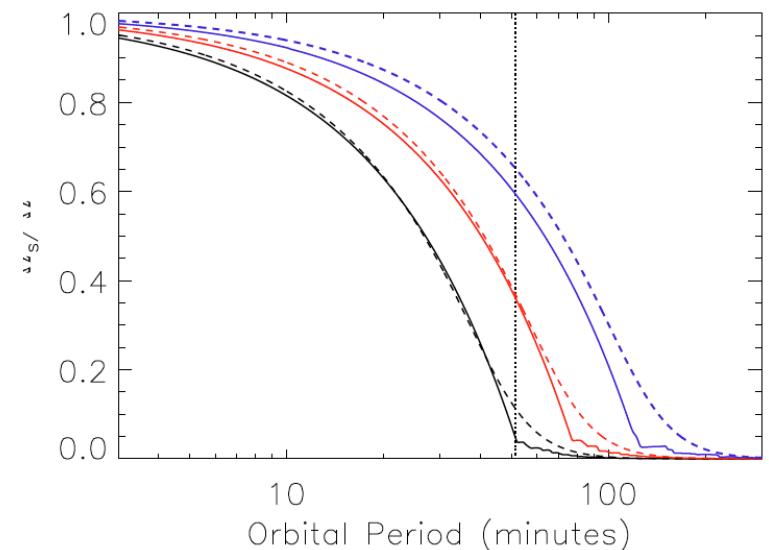
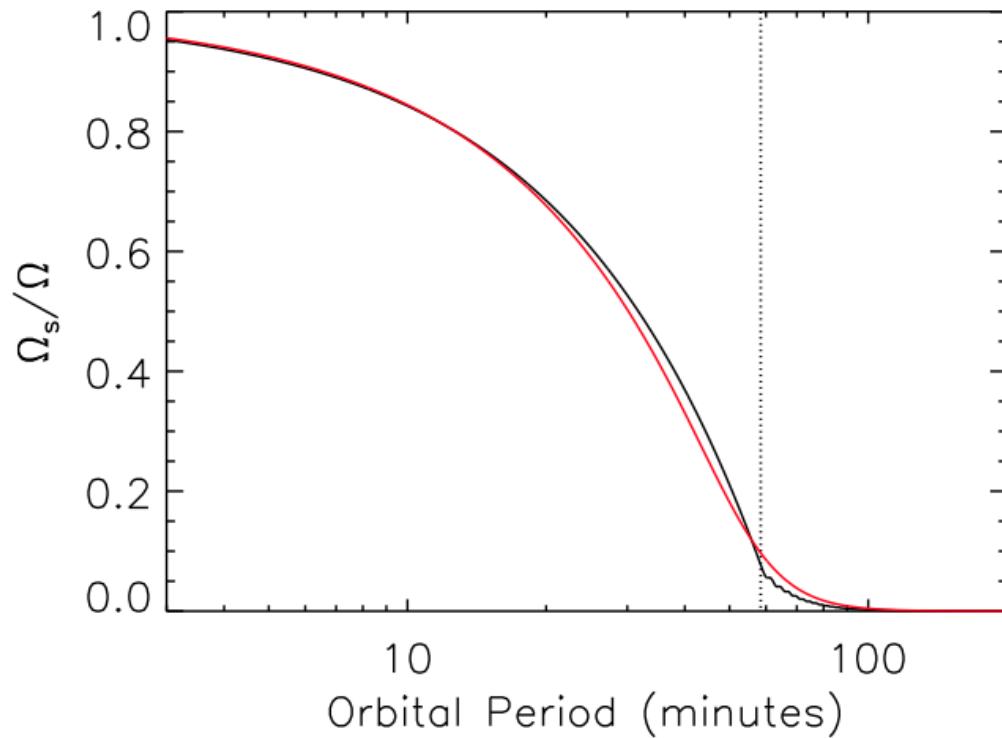
$$\text{Torque} = G \left(\frac{M'}{a^3} \right)^2 R^5 F(\omega)$$



Why is $F(\omega)$ not smooth ?



Spin-Orbit Synchronization

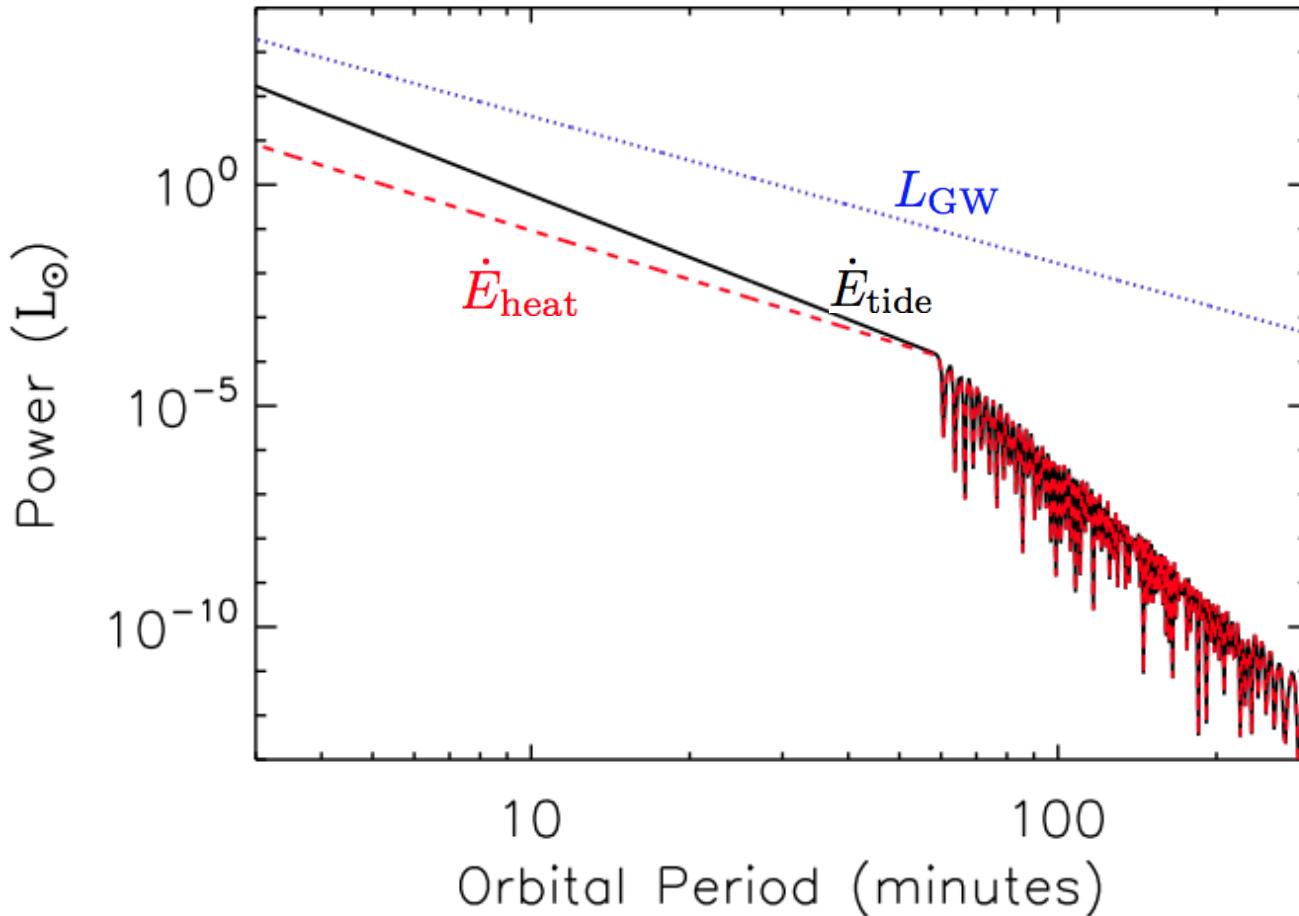


Critical orbital Ω_c : $\dot{\Omega}_s = \frac{\text{Torque}}{I} \simeq \dot{\Omega}_{\text{orb}} = \frac{3\Omega_{\text{orb}}}{2t_{\text{GW}}}$

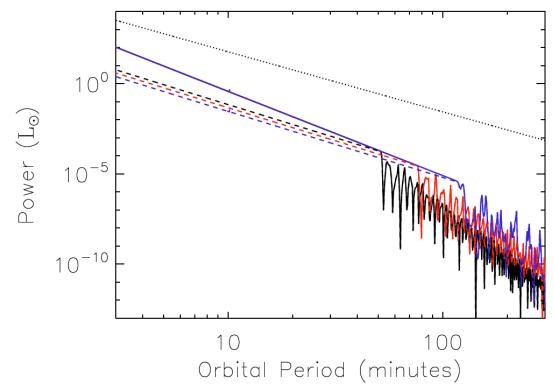
For $\Omega_{\text{orb}} > \Omega_c$: $\dot{\Omega}_s > \dot{\Omega}_{\text{orb}}$

$$\dot{\Omega}_s - \dot{\Omega}_{\text{orb}} \ll \dot{\Omega}_{\text{orb}} \implies \dot{E}_{\text{tide}} = \Omega_{\text{orb}} T \simeq \frac{3I\Omega_{\text{orb}}^2}{2t_{\text{GW}}}$$

Tidal Heating Rate



$$\dot{E}_{\text{heat}} = \dot{E}_{\text{tide}} \left(1 - \frac{\Omega_s}{\Omega_{\text{orb}}} \right)$$



Consequences of Tidal Heating

Depend on where the heat is deposited ...

If deposited in shallow layer:

thermal time short

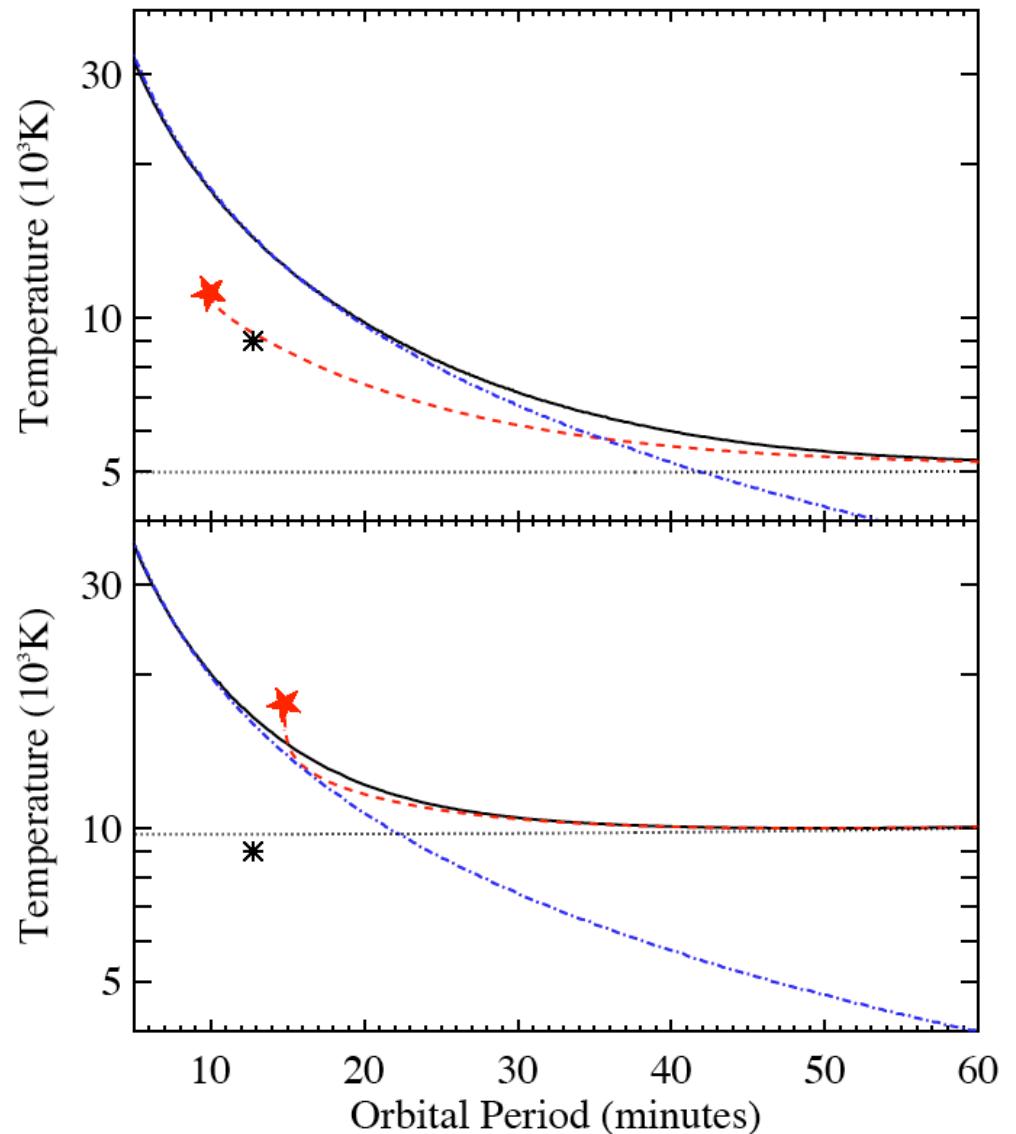
\Rightarrow change T_{eff}

If deposited in deeper layer:

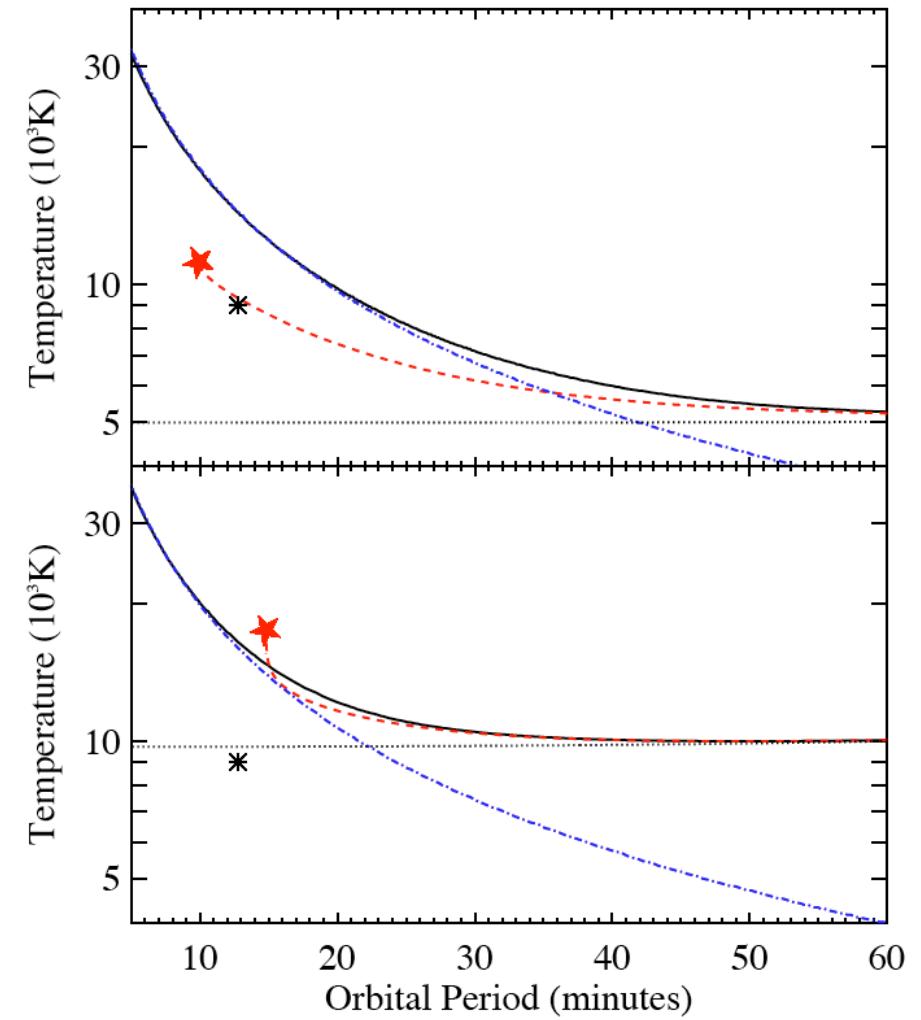
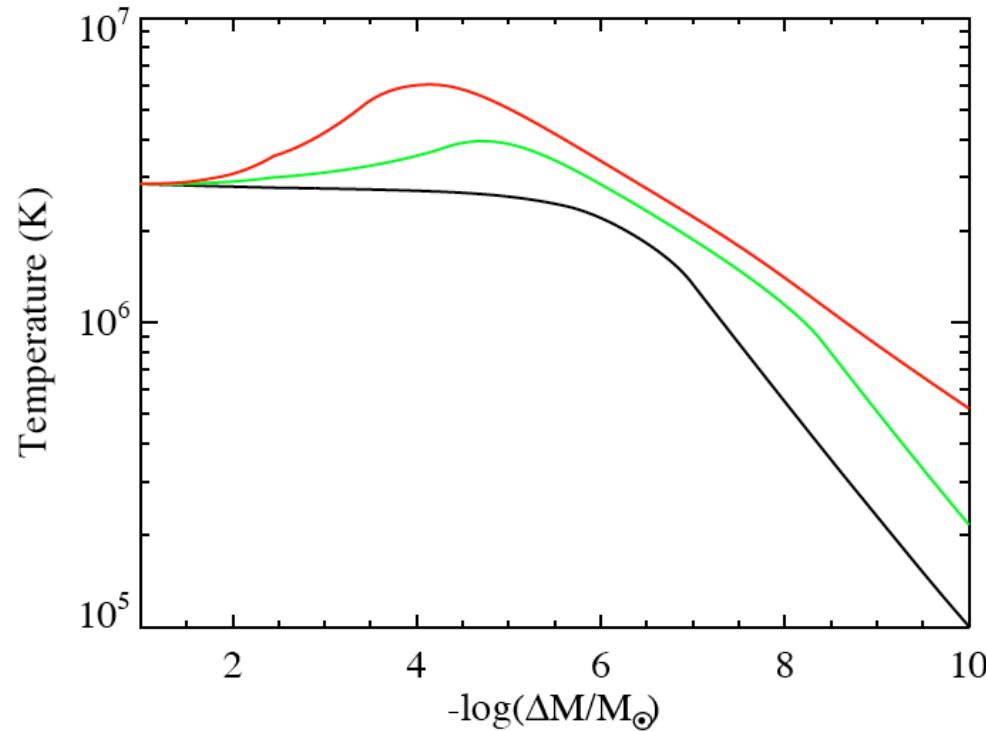
thermal time longer than orbital

\Rightarrow nuclear flash

“Tidal Nova” !



Tidal Nova



Summary

- **Merging NS and BH Binaries:**

- will likely be detected this decade by LIGO/VIRGO
- Most severe test of gravity; also probe NS EOS
- EM counterparts: GRBs, Optical/IR detectable by LSST (?),
Precursors (??)

- **Merging WD Binaries:**

- Being detected in recent/ongoing surveys
- Produce various outcomes: e.g., SN Ia
- Transient sources (PTF, LSST)
- Pre-merger: Tidal Nova
- Low-frequency GW sources (LISA??)

Neutron Stars:

Different Observational Manifestations

- Isolated NSs
- Accreting NSs
- Merging NSs

Isolated Neutron Stars

Radio pulsars: $P, \dot{P} \Rightarrow$

Most pulsars : $B \sim 10^{12-13}$ G

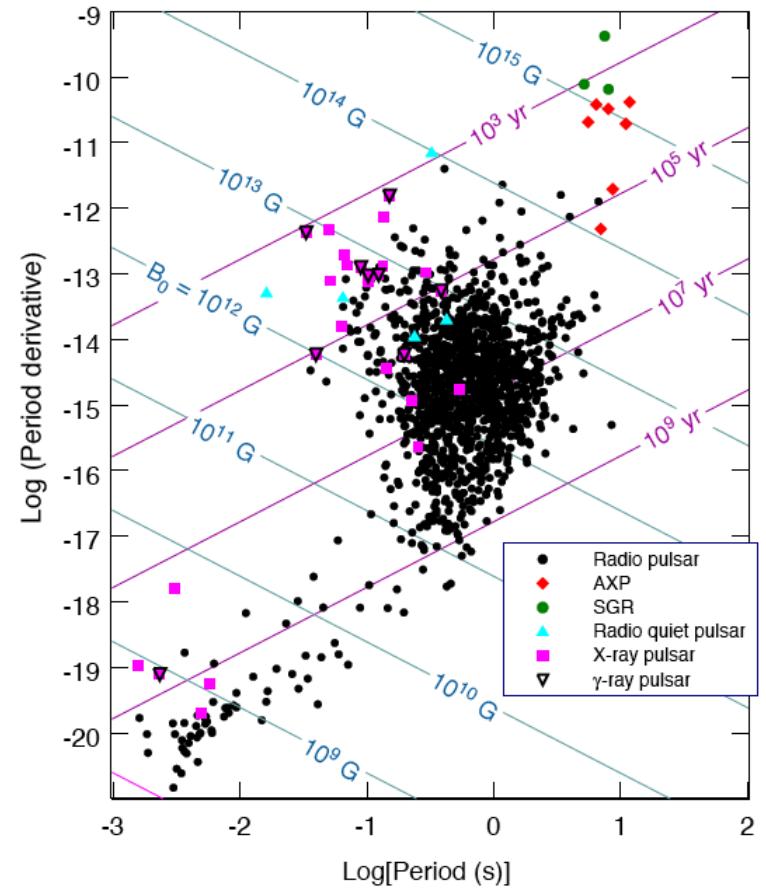
Millisecond pulsars : $B \sim 10^{8-9}$ G

High – B radio pulsars : $B \sim 10^{14}$ G

Radiation at all wavelengths:
radio, IR, optical, X-rays, Gamma-rays

New Odd Behaviors:

- RRATs (rotating radio transients)
radio bursts (2-30 ms), quiescence (min-hrs);
period \sim sec
- Intermittent Pulsars (“Sometimes a pulsar”)
e.g. PSR B1931+24: “on” for \sim a week,
“off” for \sim a month



Magnetars

Neutron stars powered by superstrong magnetic fields ($B > 10^{14}$ G)

Soft Gamma-Ray Repeaters (SGRs) (7+4 systems)

Anomalous X-ray Pulsars (AXPs) (9+3 systems)

Even in quiescence, $L \sim 10^{34-36}$ erg s $^{-1}$ $\gg I\Omega\dot{\Omega}$

T ~ 0.5 keV, but significant emission up to ~ 100 keV (e.g. Kuiper et al.) \Rightarrow Magnetar corona

AXP/SGR bursts/flares (e.g. Kaspi, Gavriil, Kouveliotou, Woods, etc)

Giant flares in 3 SGRs

12/04 flare of SGR1806-20 has $E > 10^{46}$ erg

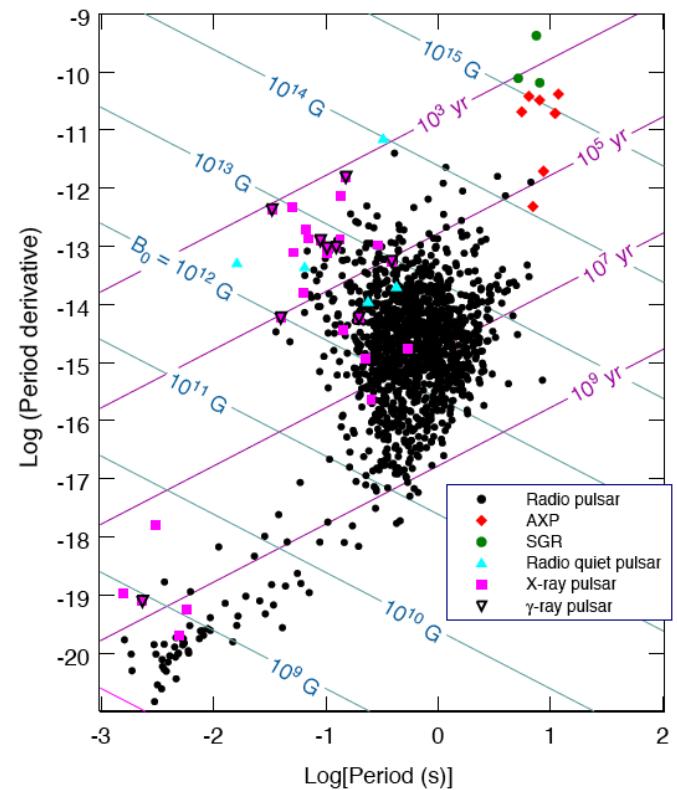
QPOs during giant flares (e.g. Israel, Strohmayer, Watts, etc)

Magnetars do not show persistent radio emission

Connection with high-B radio pulsars?

Note:

- Transient magnetars: Radio emission triggered by X-ray outbursts
XTE J1810-197, 1E 1547.0-5408 (Camilo et al. 2007)
- PSR J1622-4950 has $B \sim 3 \cdot 10^{14}$ G, but $L_x \sim L_{sd}/4$ (Levin et al 2010)

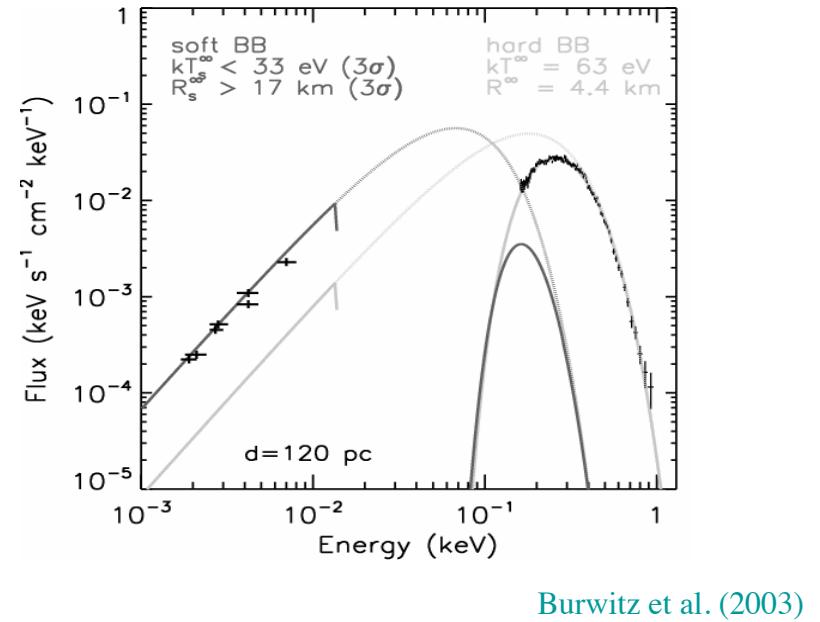


Thermally Emitting Isolated NSs

“Perfect” X-ray blackbody:
RX J1856.5-3754

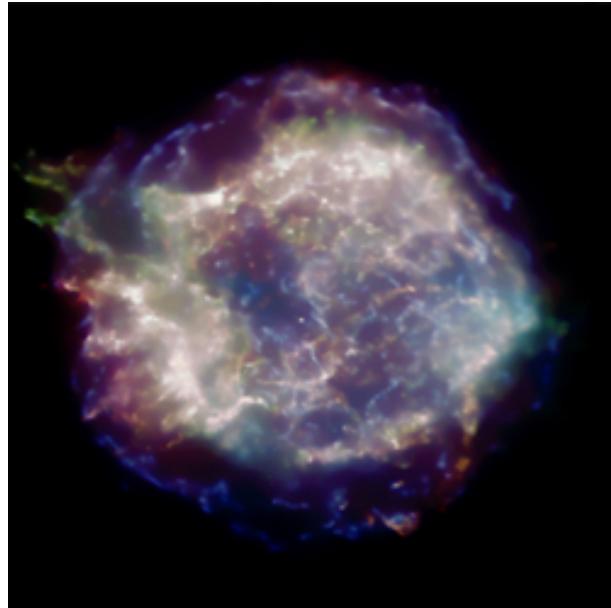
Spectral lines detected:
(e.g., van Kerkwijk & Kaplan 06; Haberl 06)

- RXJ1308+2127 (0.2-0.3 keV)
- RXJ1605+3249 (~0.45 keV)
- RXJ0720-3125 (~0.3 keV)
- RXJ0420-5022 (~0.3 keV)?
- RXJ0806-4123 (~0.5 keV)?
- RBS 1774 (~0.7 keV)?



$\implies B \sim 10^{13-14}$ G? magnetar descendant & off-beam radio pulsar?

Central Compact Objects (CCOs) in SNRs

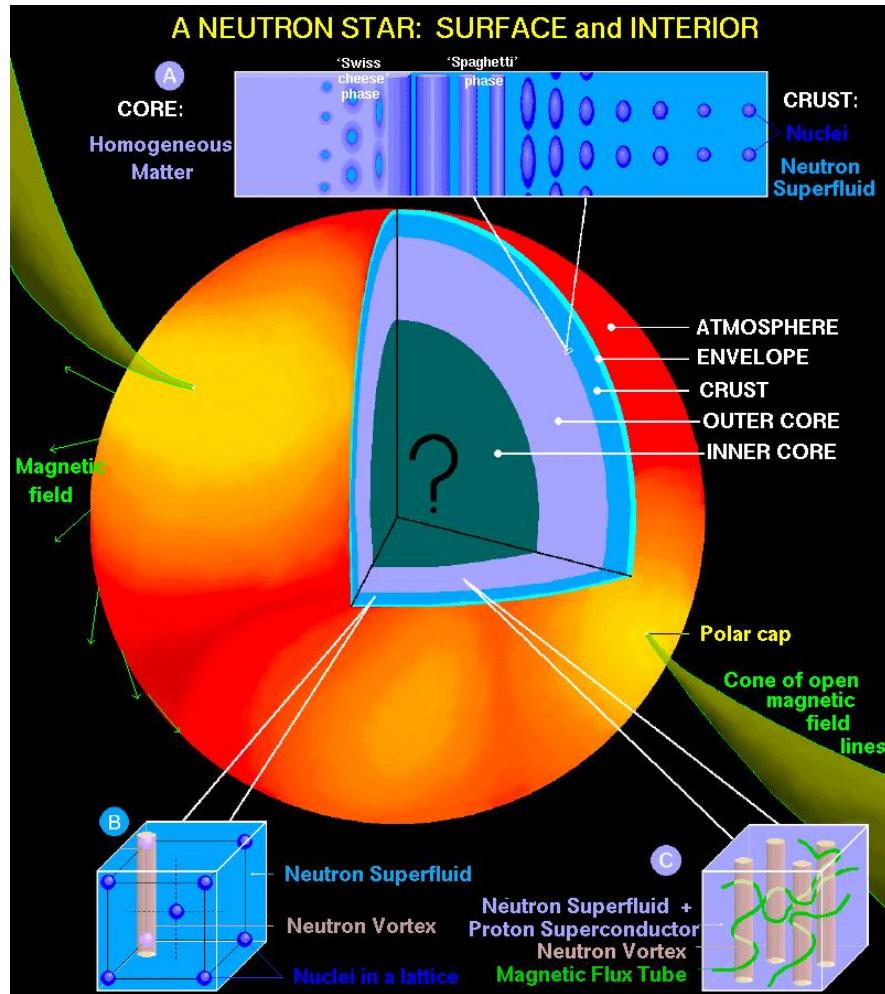


CCO	SNR	Age (kyr)	d (kpc)	P (s)	f_p^a (%)	B_s (10^{11} G)
RX J0822.0 – 4300	Puppis A	3.7	2.2	0.112	11	<9.8
CXOU J085201.4 – 461753	G266.1 – 1.2	1	1	...	<7	...
1E 1207.4 – 5209	PKS 1209 – 51/52	7	2.2	0.424	9	<3.3
CXOU J160103.1 – 513353	G330.2 + 1.0	$\gtrsim 3$	5	...	<40	...
1WGA J1713.4 – 3949	G347.3 – 0.5	1.6	1.3	...	<7	...
CXOU J185238.6 + 004020	Kes 79	7	7	0.105	64	0.31
CXOU J232327.9 + 584842	Cas A	0.33	3.4	...	<12	...

Halpern & Gotthelf 2010

Small surface dipole field ... (are they “anti-magnetars”?)

Neutron Star Structure



D. Page

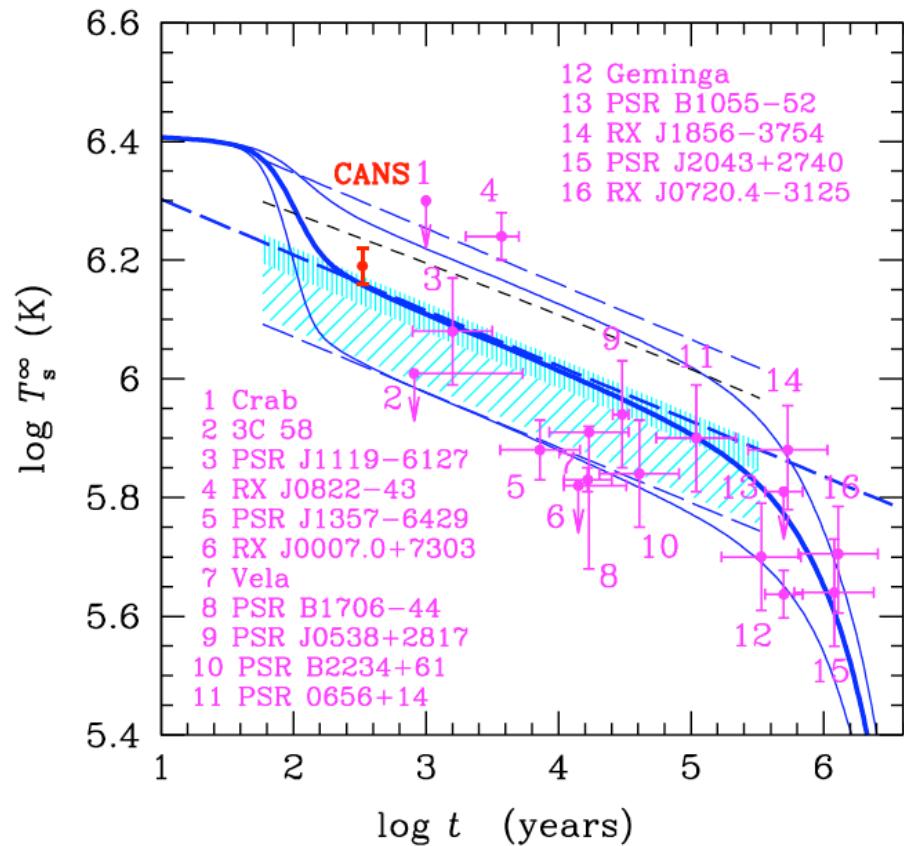
Highlight #1: Constrain NS Interior physics by Cooling

Surface emission has been
Detected in ~20 NSs

Probe the interior of NS
(EOS, exotic processes)

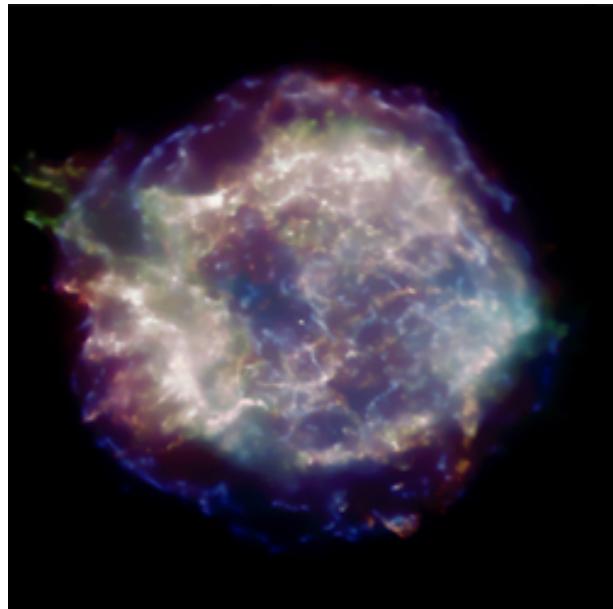
Difficulties:

Many parameters and theoretical
Models/processes...

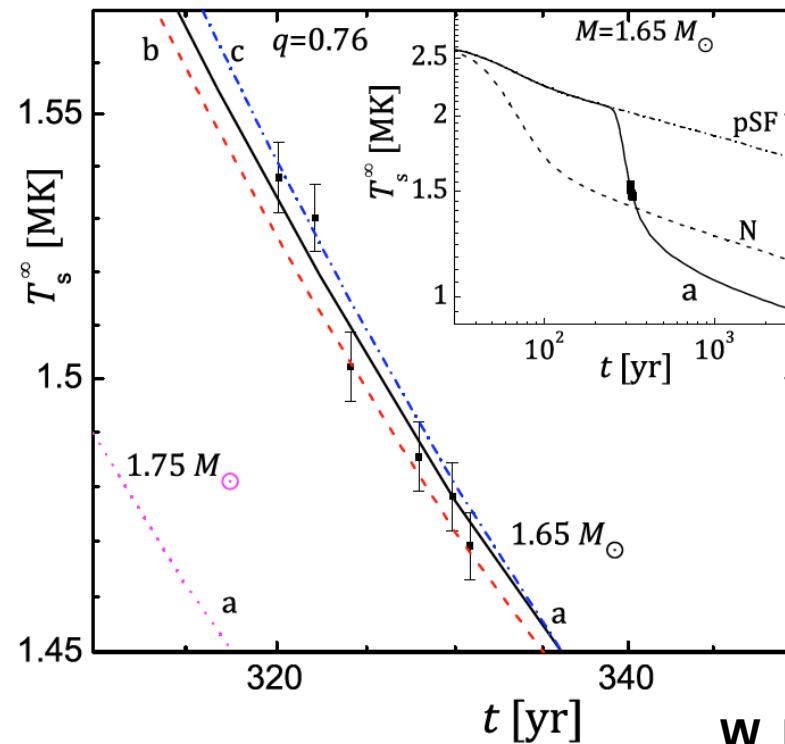


Yakovlev et al. 2011

NS in Cas A SNR: Evidence of Superfluidity



Cas A SNR: age 330 yrs
CCO (NS) first discovered in 1999
Many observations since then...



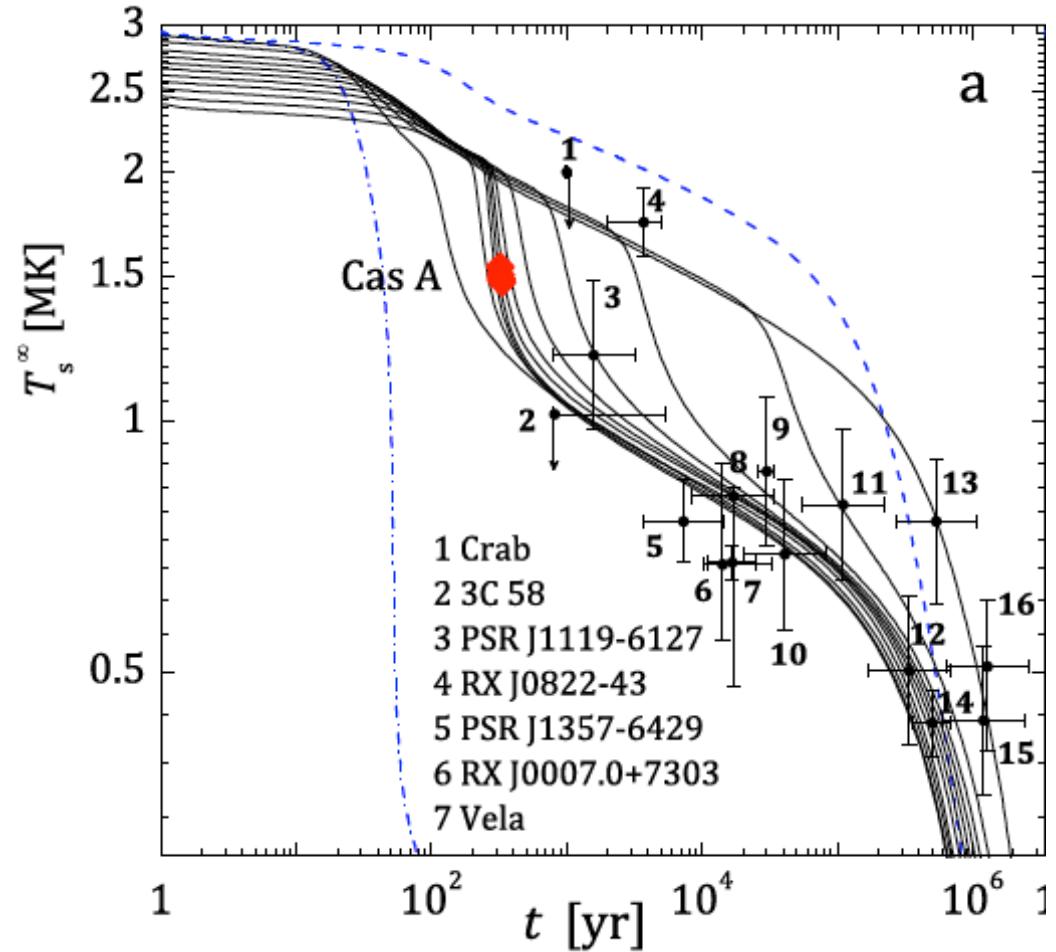
W. Ho & Heinke 2010

Shternin et al. 2011

NS in Cas A SNR: Evidence for Superfluidity

Decline in T is quite large for $t \sim 300$ yrs.

Solution: Internal T drops below T_{crit} at 300 yrs ==> neutrons become superfluid at 300 yrs, leading to sudden Cooling.



Strongest evidence of neutron superfluidity in NS core.

(Shternin et al. 2011; Page et al 2011)

Highlight #2: Probing QED Processes in Superstrong B Fields

- One photon pair production:

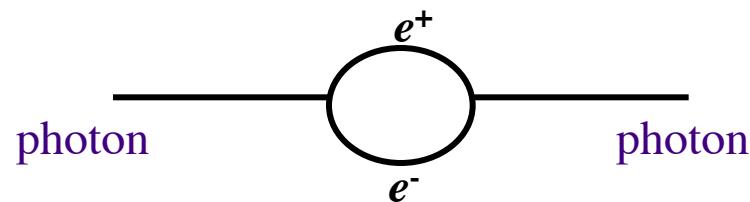
$$\gamma \longrightarrow e^+ + e^-$$

- Photon splitting:

$$\gamma \longrightarrow \gamma + \gamma$$

- **Vacuum birefringence:**

(photon propagation affected by B field)

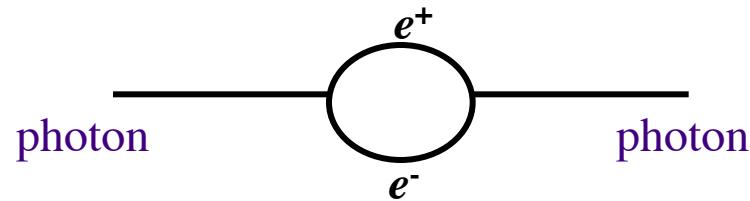


Context: Modeling Radiation from Magnetic NS Atmospheres

NS Atmospheres:

- Outermost ~cm of the star
- Density $0.1\text{-}10^3 \text{ g/cm}^3$: nonideal, partially ionized, magnetic plasma
- **Effect of QED: Vacuum polarization**

Vacuum Polarization in Strong B



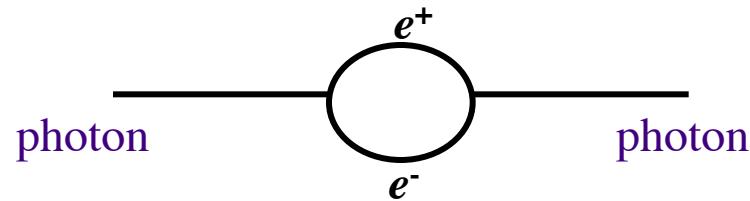
Heisenberg & Euler,
Weisskopf, Schwinger,
Adler...

Important when B is of order or larger than

$$B_Q = 4.4 \times 10^{13} \text{ G}$$

at which $\hbar\omega_{ce} = \hbar \frac{eB}{m_e c} = m_e c^2$

Vacuum Polarization in Strong B



Heisenberg & Euler,
Weisskopf, Schwinger,
Adler...

Dielectric tensor: $\boldsymbol{\epsilon} = \mathbf{I} + \Delta\boldsymbol{\epsilon}_{\text{vac}}$

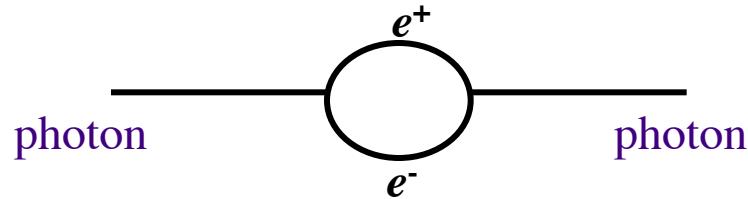
$$|\Delta\epsilon_{\text{vac}}| \sim 10^{-4}(B/B_Q)^2, \text{ with } B_Q = 4.4 \times 10^{13} \text{ G}$$

Two photon modes:

Ordinary mode (\parallel)

Extraordinary mode (\perp)

Vacuum Polarization in Strong B



Heisenberg & Euler,
Weisskopf, Schwinger,
Adler...

Dielectric tensor: $\boldsymbol{\epsilon} = \mathbf{I} + \Delta\boldsymbol{\epsilon}_{\text{vac}}$

$$|\Delta\epsilon_{\text{vac}}| \sim 10^{-4}(B/B_Q)^2, \text{ with } B_Q = 4.4 \times 10^{13} \text{ G}$$

Two photon modes:

Ordinary mode (//)

Extraordinary mode (\perp)

On the other hand...

Magnetic Plasma by itself (without QED) is birefringent:

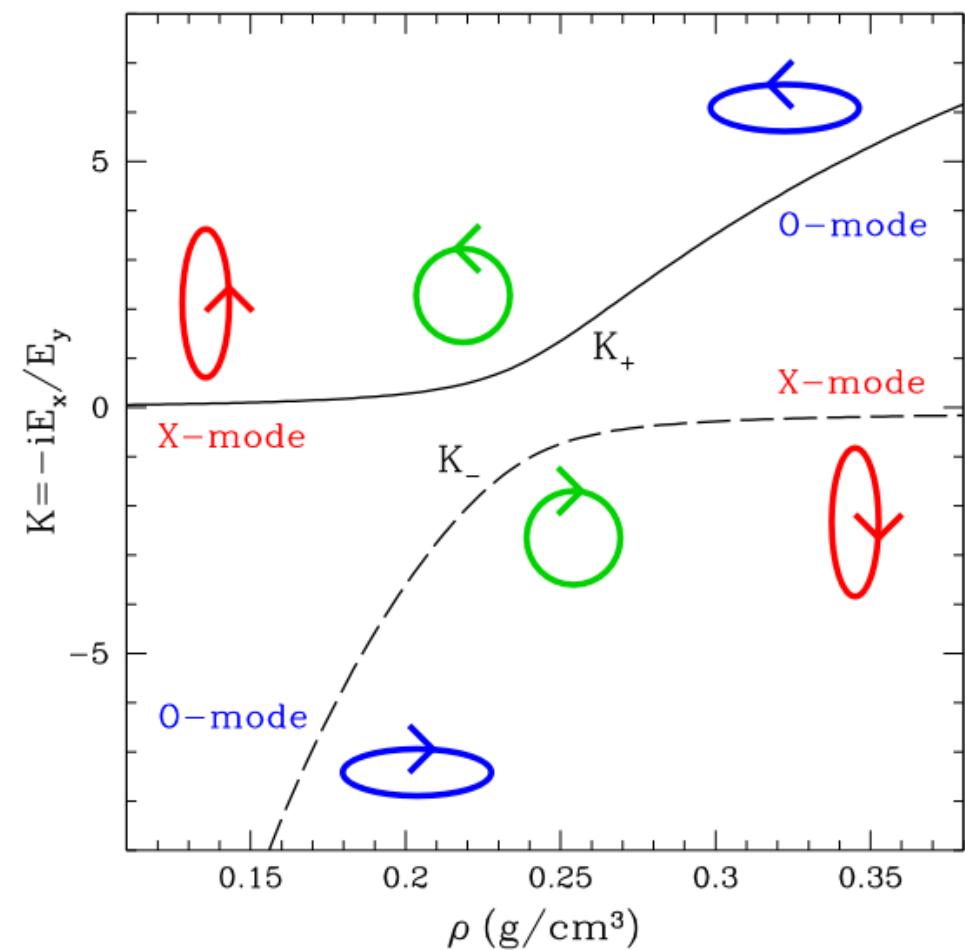
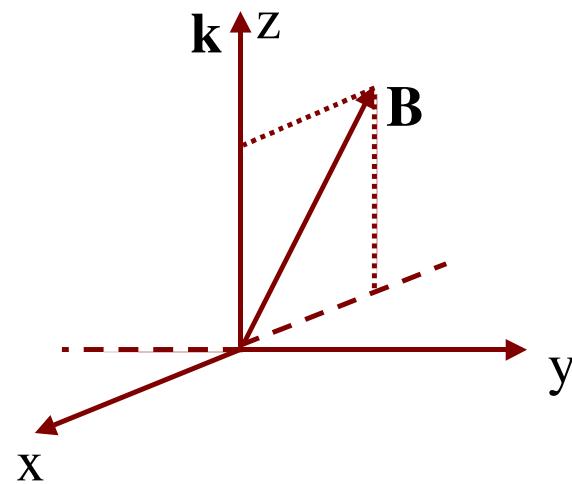
Ordinary mode



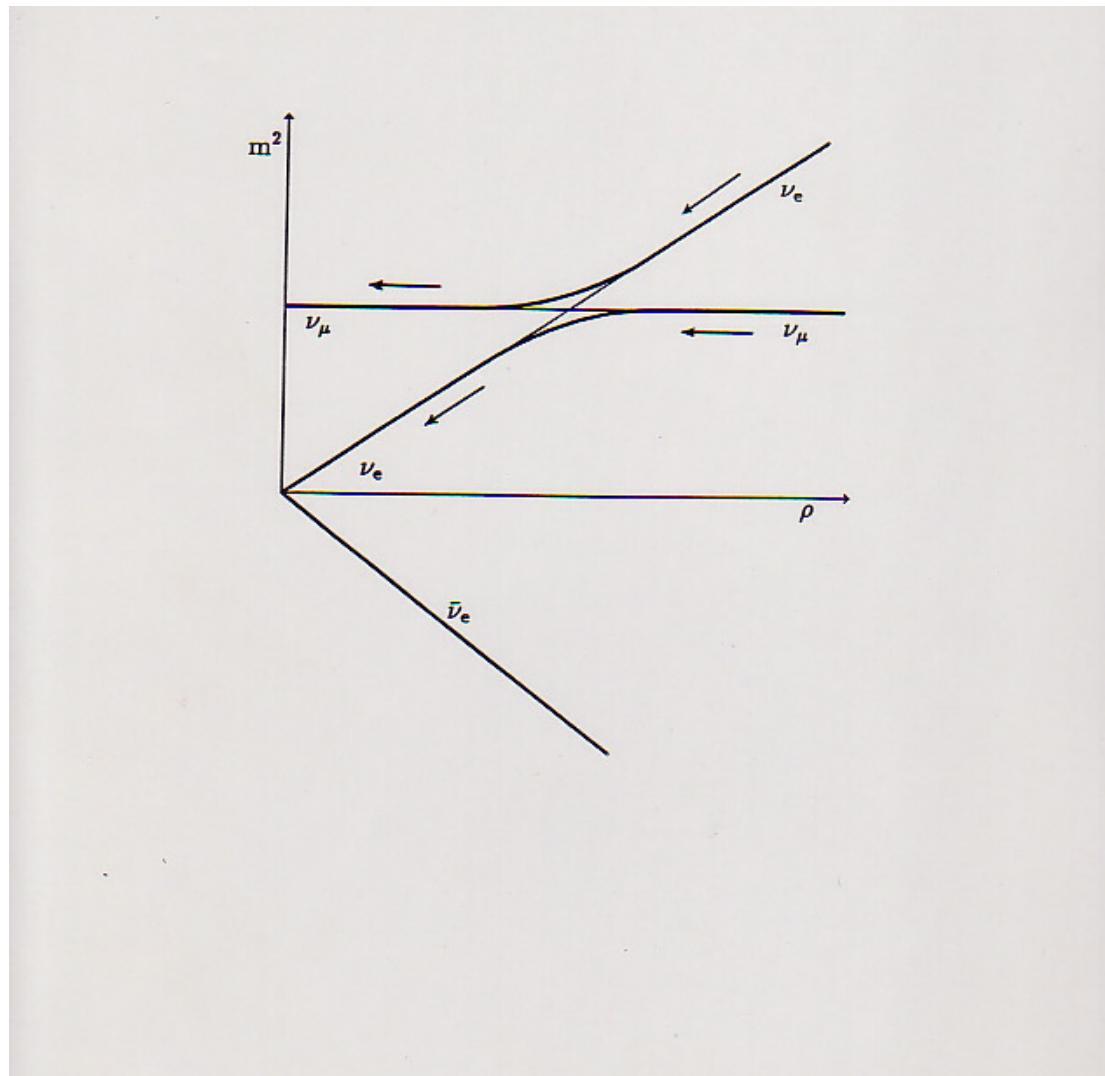
Extraordinary mode



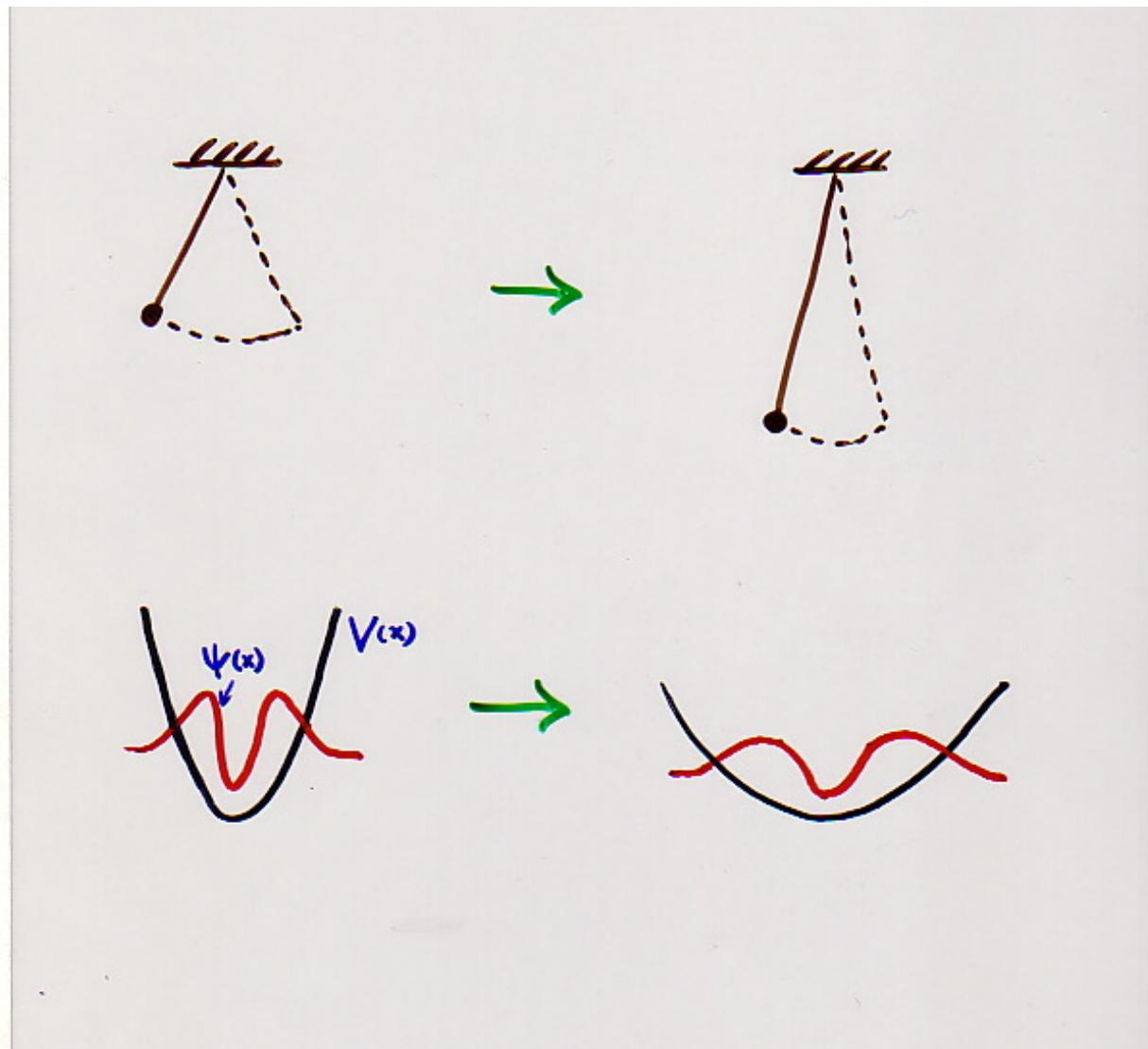
“Plasma+Vacuum” ==> Vacuum resonance



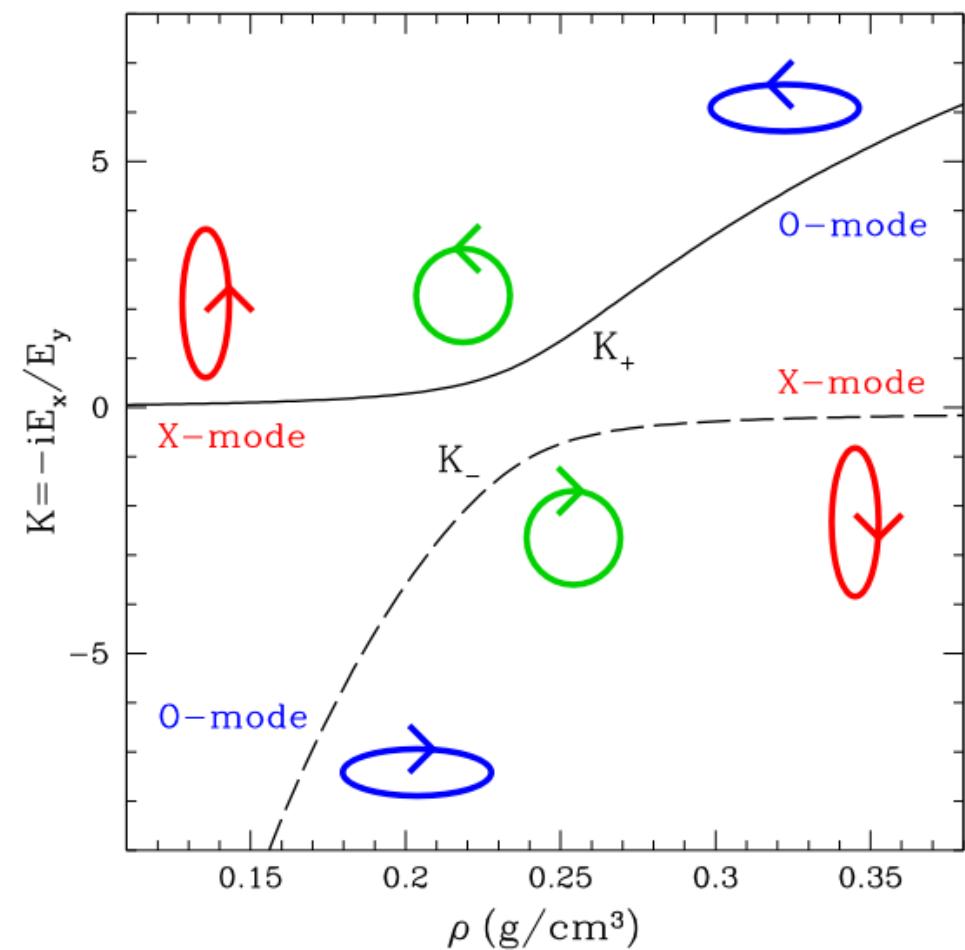
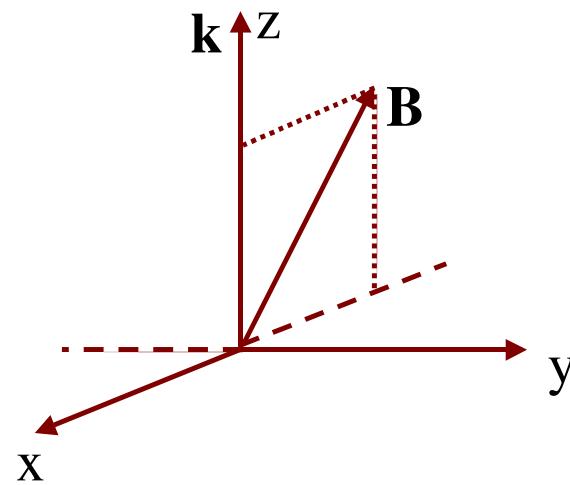
Mikheyev-Smirnov-Wolfenstein (MSW) Neutrino Oscillation



Adiabatic Evolution of a Quantum State



“Plasma+Vacuum” ==> Vacuum resonance

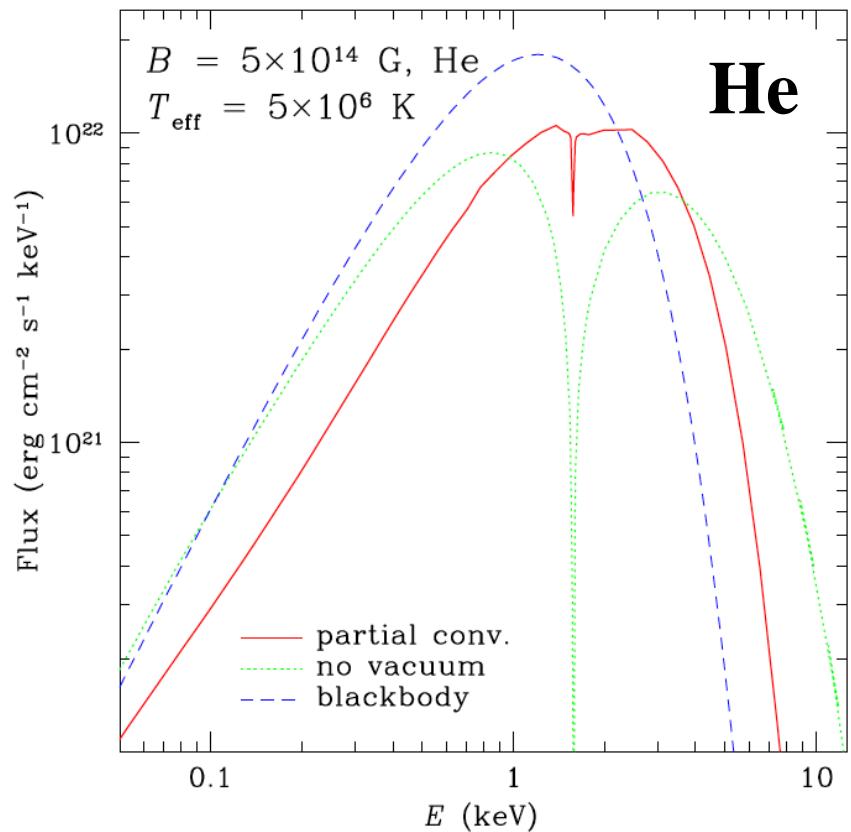
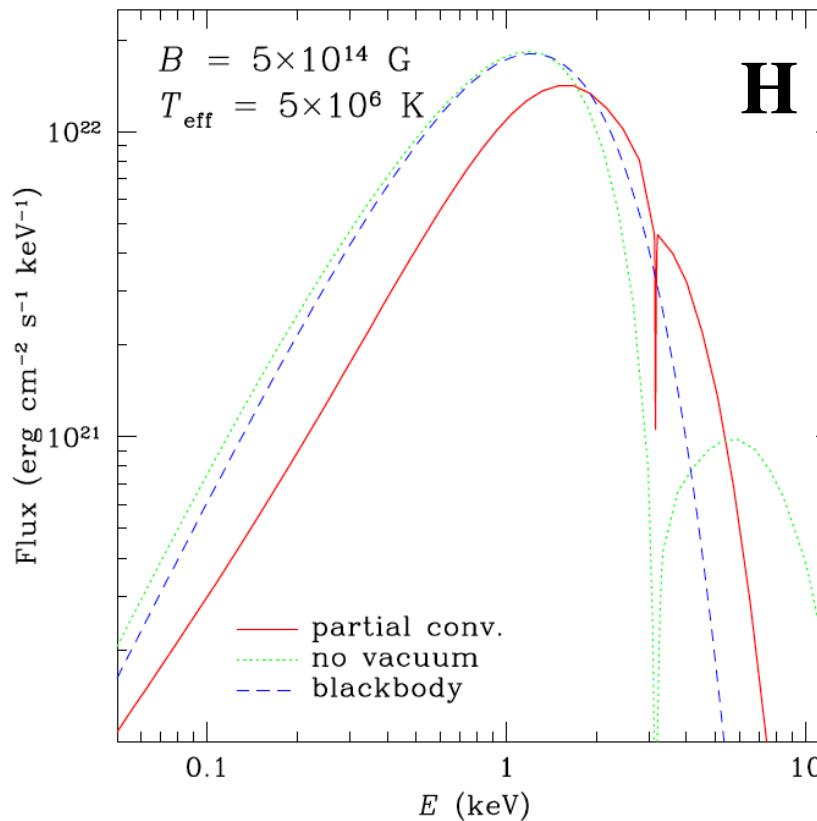


Why do we care?

The two photon modes have very different opacities

- => Mode conversion can affect radiative transfer significantly
- => Spectrum and polarization signal from the NS

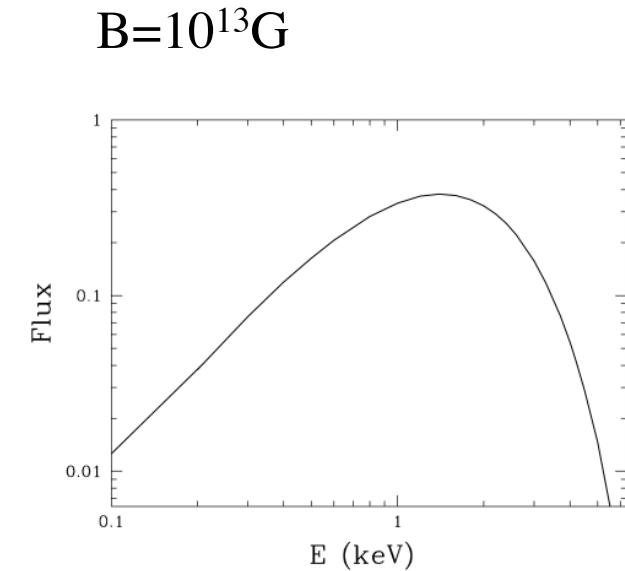
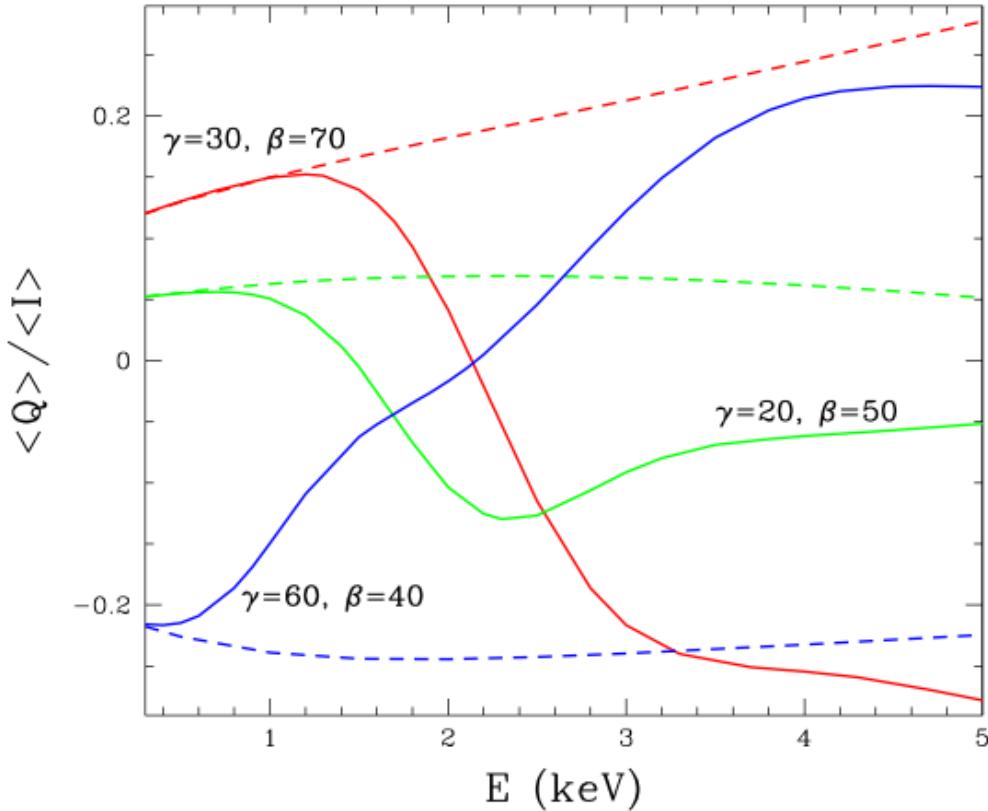
For $B > 10^{14}$ G, vacuum polarization strongly affects spectrum



Matt Van Adelsberg & DL 2006

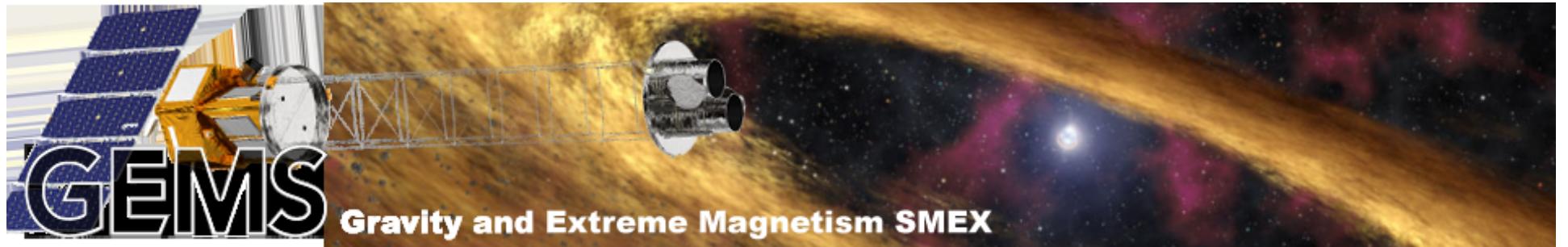
==> Magnetars do not show absorption features in thermal emission
QED at work!

Even for modest B's, vacuum resonance produces unique polarization signals

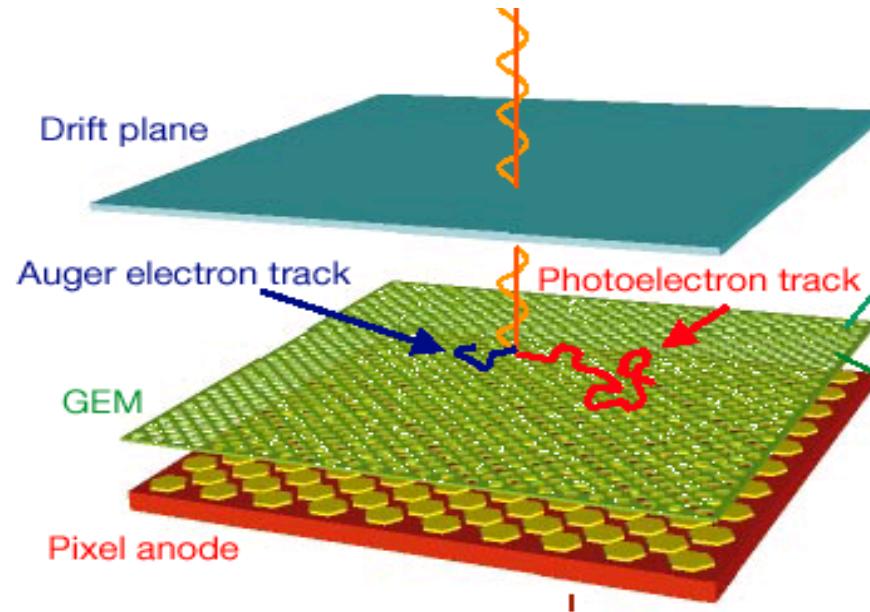


“boring” spectrum & lightcurve,
but interesting/nontrivial polarization spectrum!

==> X-ray polarimeters



X-ray polarization
Measurement:



GEMS: J. Swank, GSFC (PI)

Technical slides: QED Effect in NS Atmosphere

Dielectric tensor of magnetized plasma including vacuum polarization

$$\boldsymbol{\mathcal{E}} = \mathbf{I} + \Delta\boldsymbol{\mathcal{E}}^{(\text{plasma})} + \Delta\boldsymbol{\mathcal{E}}^{(\text{vac})}$$

where $\Delta\boldsymbol{\mathcal{E}}^{(\text{vac})} \sim 10^{-4} (B/B_Q)^2 f(B)$, with $B_Q = 4.4 \times 10^{13} \text{ G}$, $f(B) \sim 1$

cf. Gnedin, Pavlov & Shibano 1978;
Meszaros & Ventura 1978, etc

Vacuum resonance:

$$\Delta\boldsymbol{\mathcal{E}}^{(\text{plasma})} + \Delta\boldsymbol{\mathcal{E}}^{(\text{vac})} \sim 0$$

depends on $-(\omega_p/\omega)^2 \propto \rho/E^2$

$$\rightarrow \rho_{\text{vac}} = 1.0 B_{14}^2 f(B)^{-1} (E/1 \text{ keV})^2 \text{ g cm}^{-3}$$

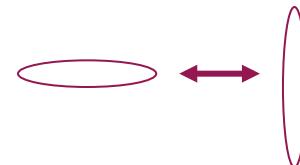
At resonance, X-mode and O-mode are “similar”

Adiabatic Condition:

$$|n_1 - n_2| \gtrsim (\dots) |d\rho/dr|$$

$$\rightarrow E \gtrsim E_{\text{ad}} = 2.5 (\tan \theta_B)^{2/3} \left(\frac{1 \text{ cm}}{H} \right)^{1/3} \text{ keV}$$

Photons with $E > 2 \text{ keV}$, mode conversion



Photons with $E < 2 \text{ keV}$, no mode conversion

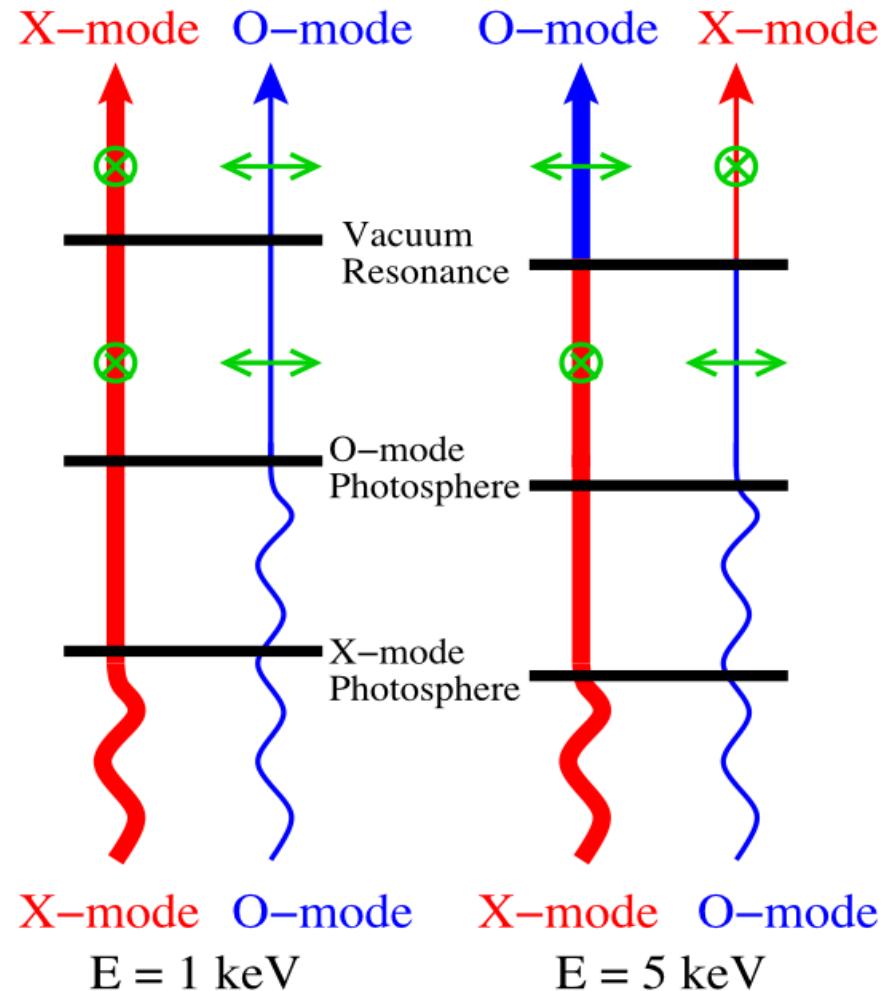
In general, nonadiabatic “jump” probability

$$P_{\text{jump}} = \exp [-(\pi/2) (E/E_{\text{ad}})^3]$$

(Landau-Zener formula)

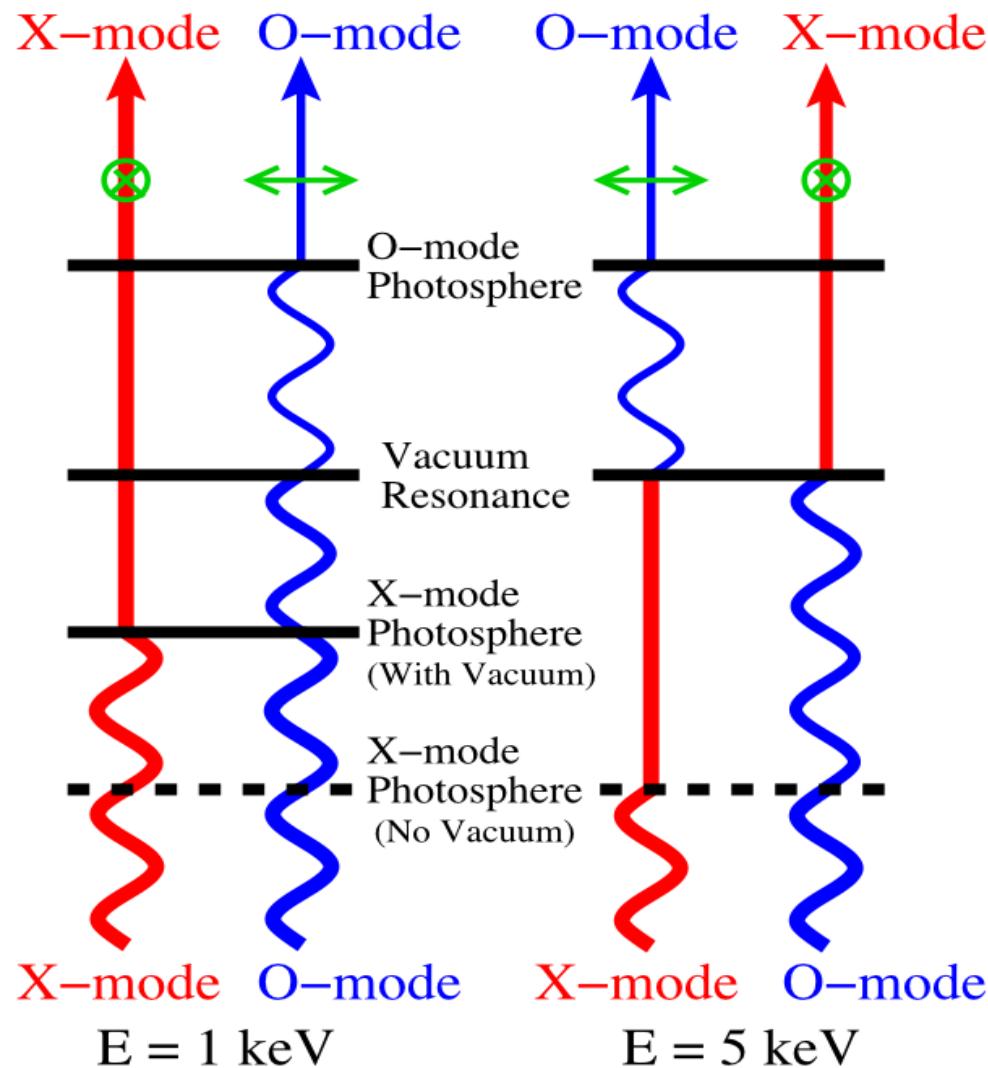
For $B < 7 \times 10^{13} T_6^{-1/8} E_1^{-1/4}$ G:

Vacuum resonance lies outside both photospheres



For $B > 7 \times 10^{13} T_6^{-1/8} E_1^{-1/4}$ G:

Vacuum resonance lies between the two photospheres



Highlight #3: Matter in Strong Magnetic Fields

Critical Field:

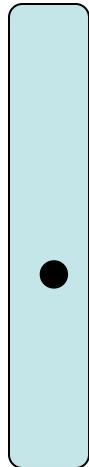
$$\hbar\omega_{ce} = \hbar \frac{eB}{m_e c} = \frac{e^2}{a_0} \implies B = B_0 = 2.35 \times 10^9 \text{ G}$$

Strong field: $B \gg B_0$

Property of matter is very different from zero-field

Atoms and Molecules

Strong B field significantly increases the binding energy of atoms



For $b = \frac{B}{B_0} \gg 1$, $B_0 = 2.35 \times 10^9$ G

$$|E| \propto (\ln b)^2$$

E.g. $|E| = 160$ eV at 10^{12} G

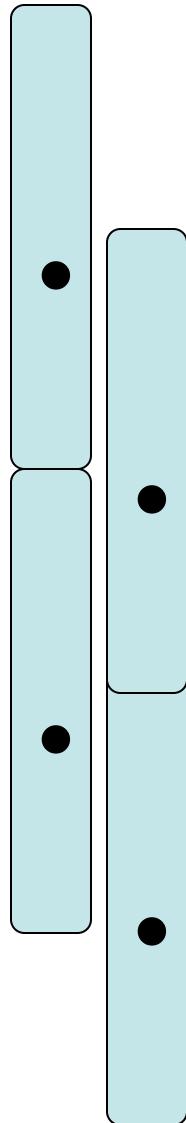
$$|E| = 540$$
 eV at 10^{14} G

Atoms combine to form molecular chains:

E.g. H_2, H_3, H_4, \dots

Condensed Matter

Chain-chain interactions lead to formation of 3D condensed matter



Binding energy per cell $|E| \propto Z^{9/5} B^{2/5}$

Zero-pressure density
 $\simeq 10^3 A Z^{3/5} B_{12}^{6/5} \text{ g cm}^{-3}$

Cohesive energy of condensed matter:

- Strong B field increases the binding energy of atoms and condensed matter

$$\text{For } b = \frac{B}{B_0} \gg 1, \quad B_0 = 2.35 \times 10^9 \text{ G}$$

Energy of atom: $\sim (\ln b)^2$

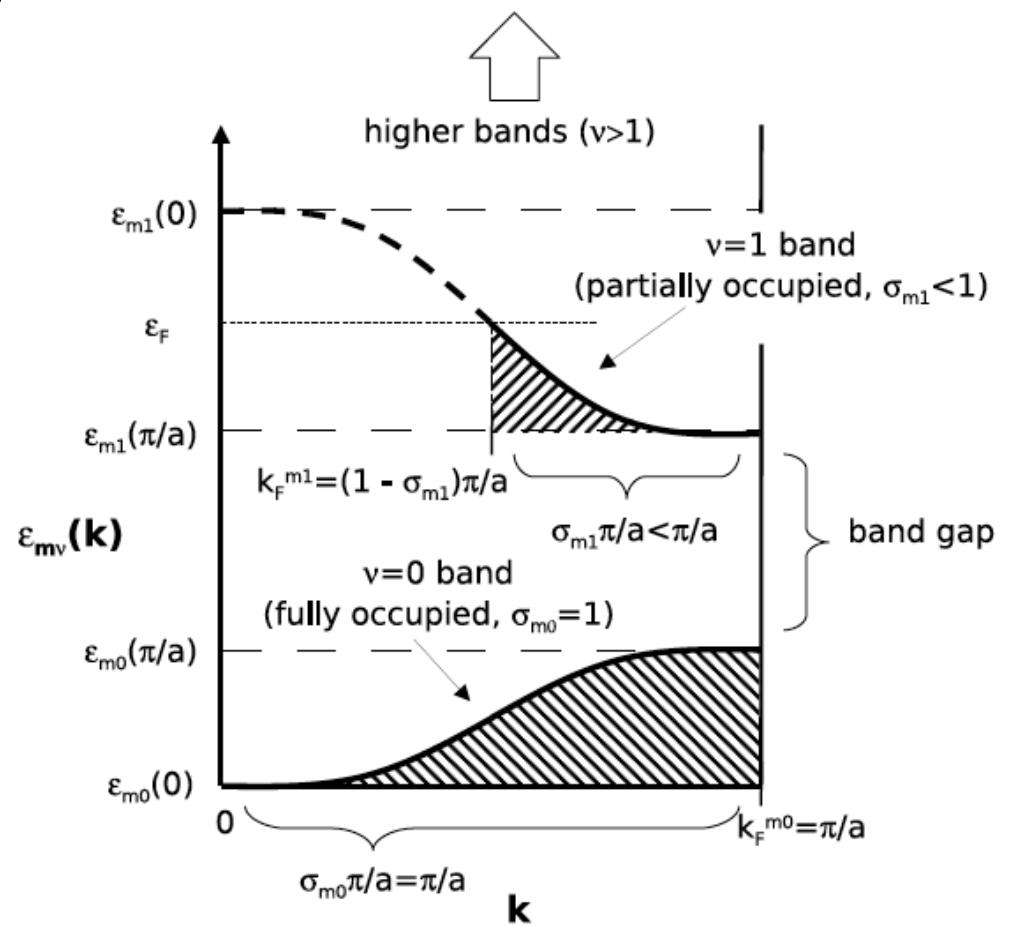
Energy of zero-pressure solid: $\sim b^{0.4}$

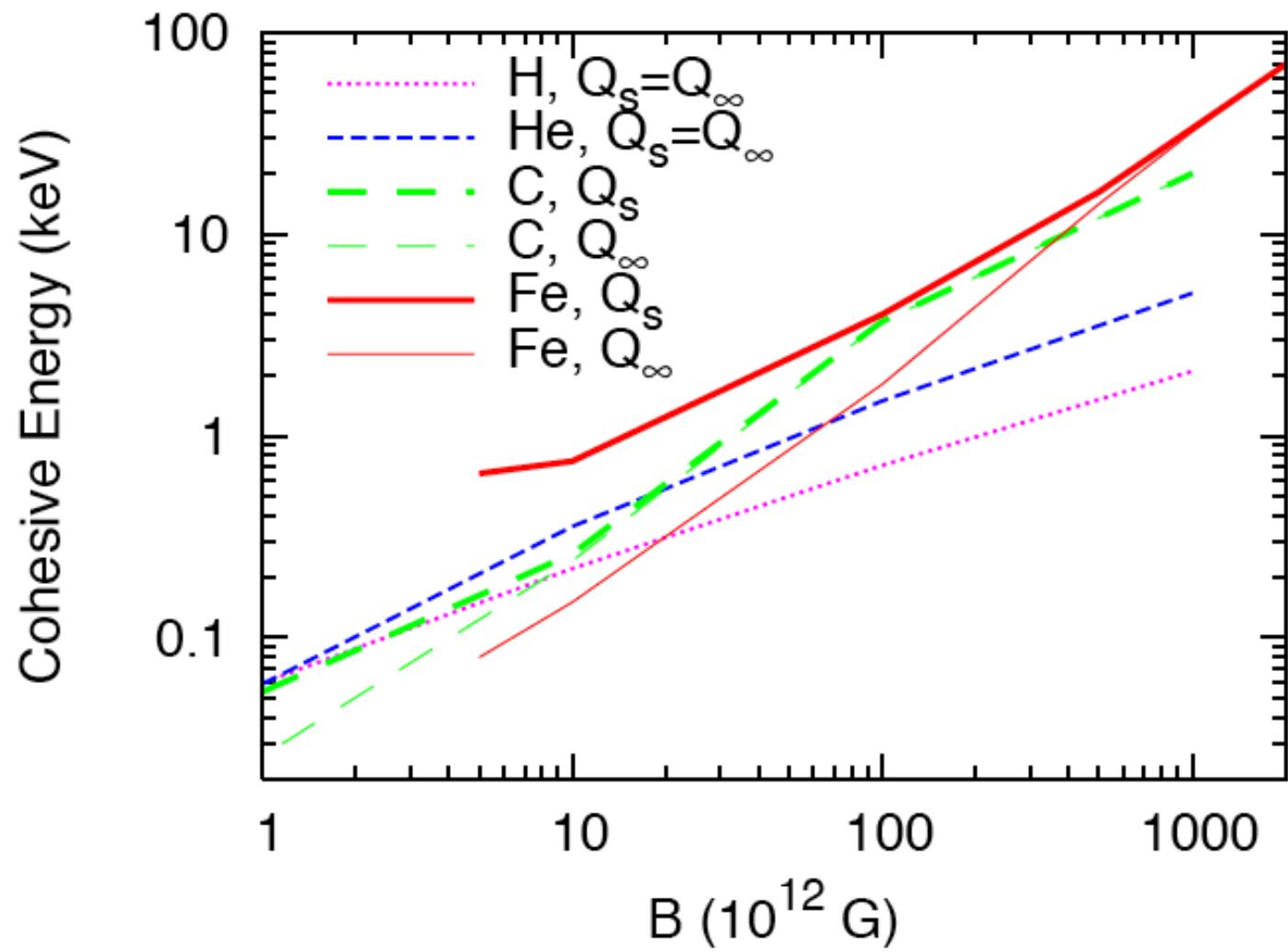
==> Expect condensed solid to have large cohesive energy

- Quantitative Calculations are needed ...

New calculations (Zach Medin & DL 2007)

- Density functional theory
- Accurate exchange-correlation energy
- Accurate treatment of band structure
- Extend to $\sim 10^{15}$ G

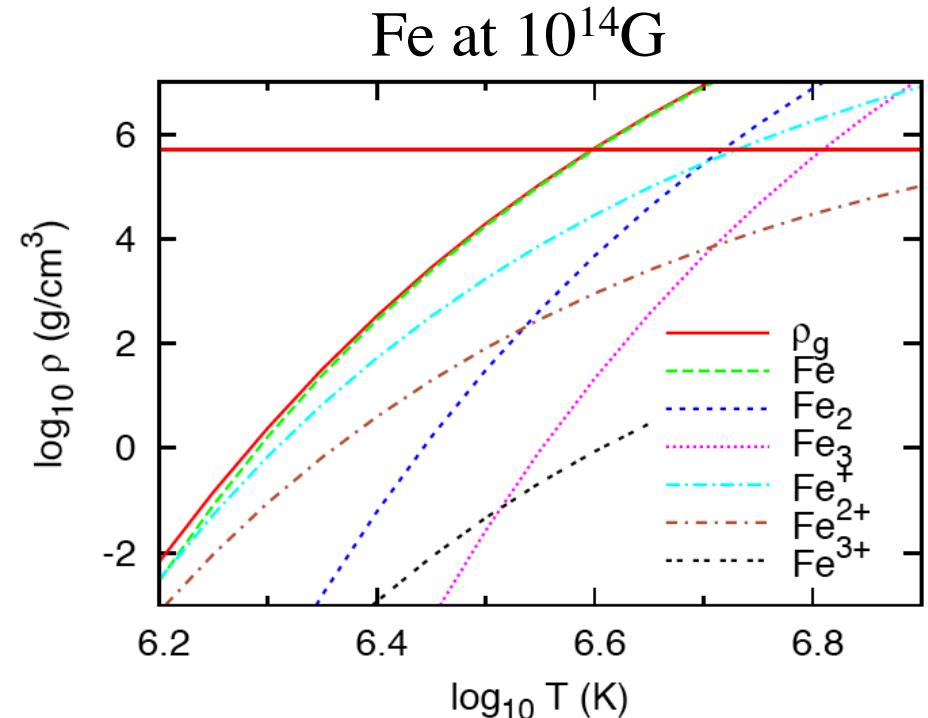
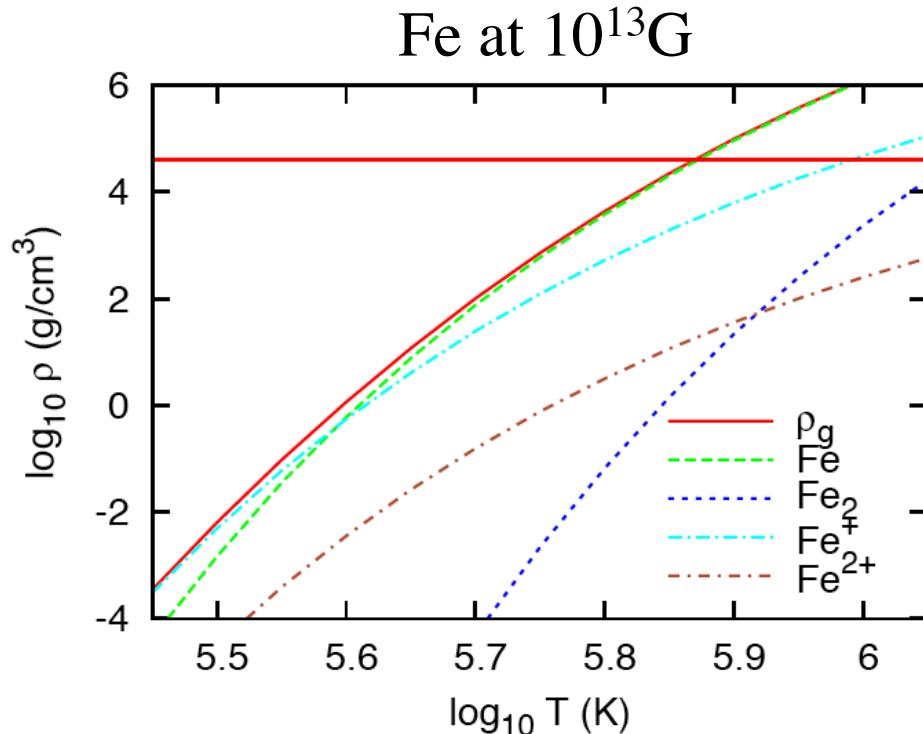




Implications...

Surface condensation of isolated NSs

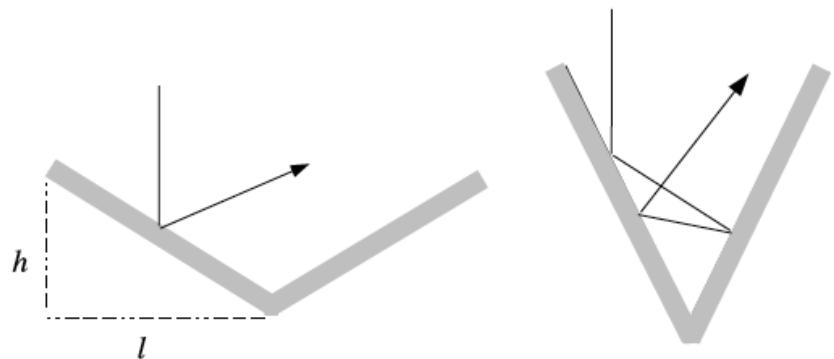
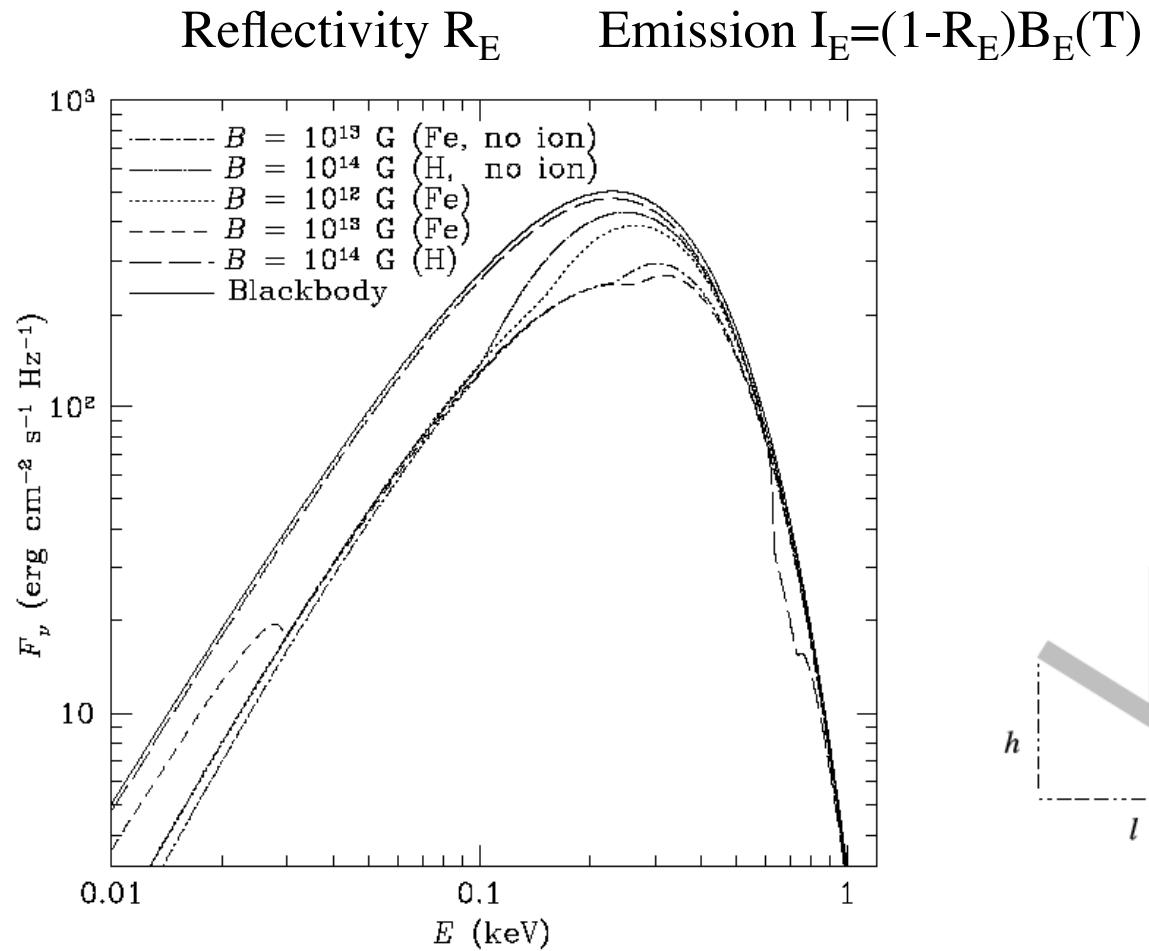
Saturated Vapor of Condensed NS Surface:



Zach Medin & DL 2007

For a given B , below $T_{\text{crit}}(B)$,
NS surface is in condensed form (with little vapor above)

Emission from condensed NS surface resembles a featureless blackbody

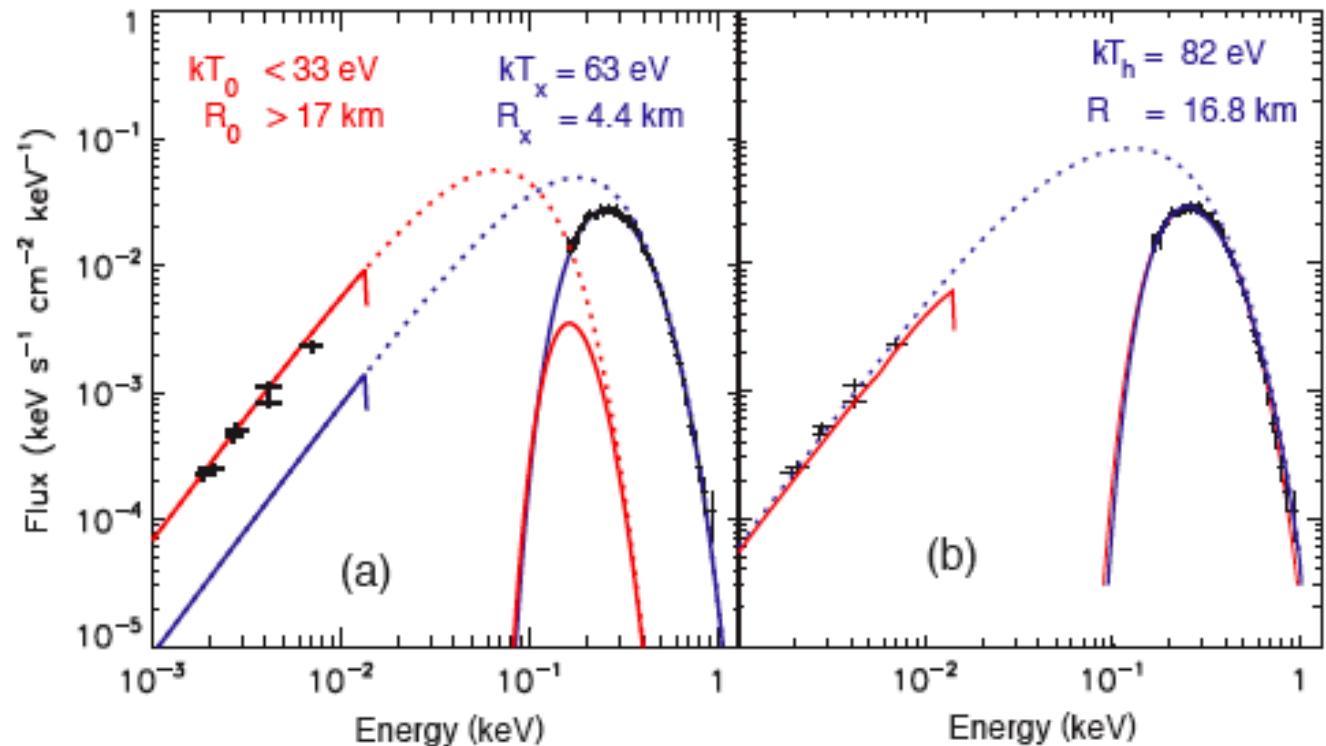


Thermally Emitting Isolated NSs

“Perfect” X-ray blackbody:

RX J1856.5-3754

($T \sim 60$ eV)



Burwitz et al. 03, Trumper et al 04

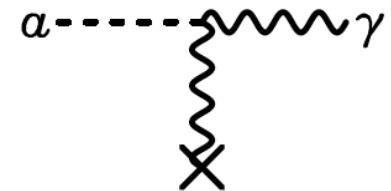
May be explained by emission from condensed surface

Highlight #4: Probing Axions with Magnetic NSs

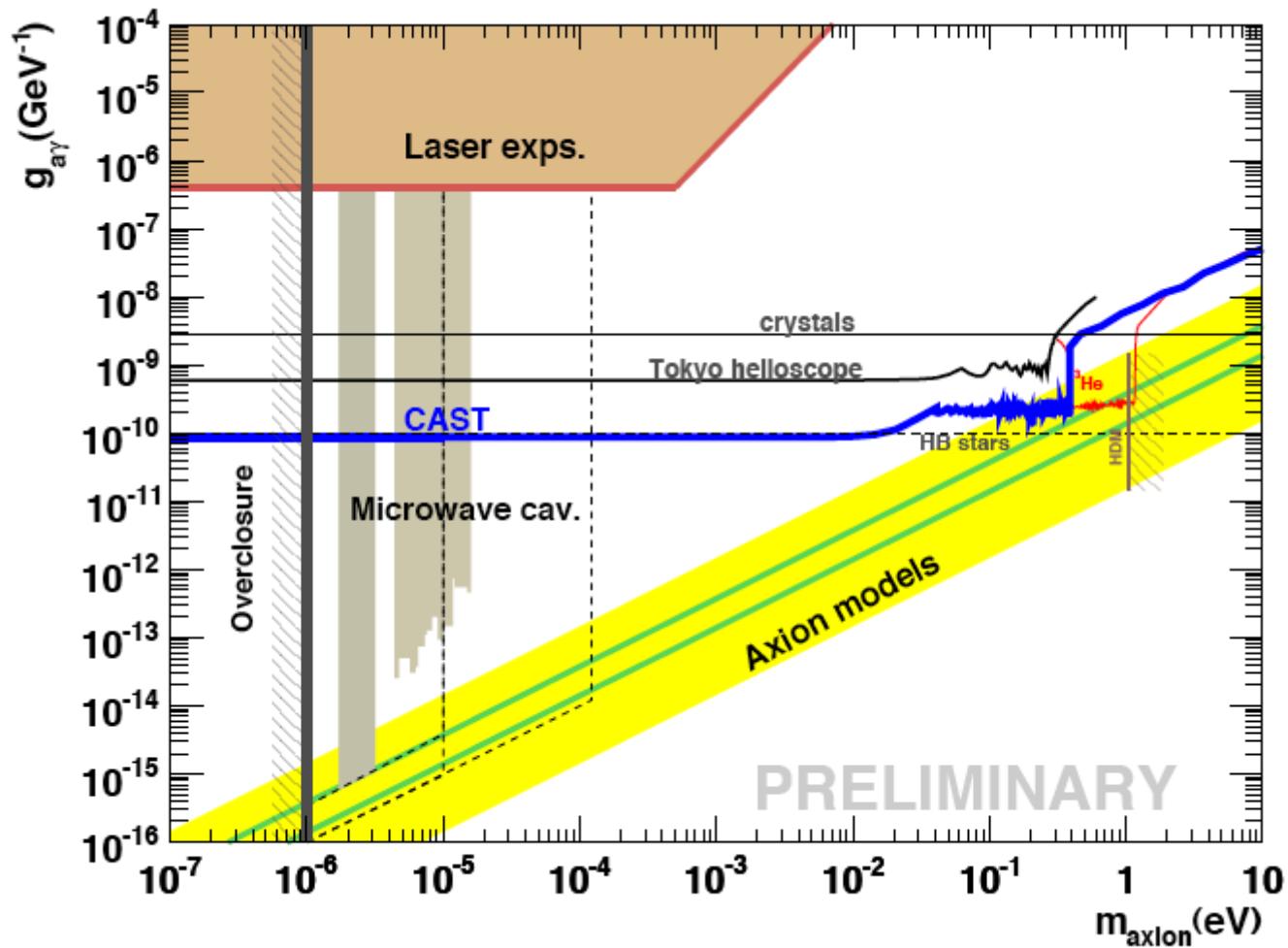
Axions: pseudoscalar particles, arise in the Peccei-Quinn solution of the strong CP problem; could be dark matter candidates

Can be produced or detected through the **Primakoff process:**

$$\mathcal{L}_{a\gamma} = -\frac{1}{4} g_{a\gamma} F_{\mu\nu} \tilde{F}^{\mu\nu} a = g_{a\gamma} \mathbf{E} \cdot \mathbf{B} a$$



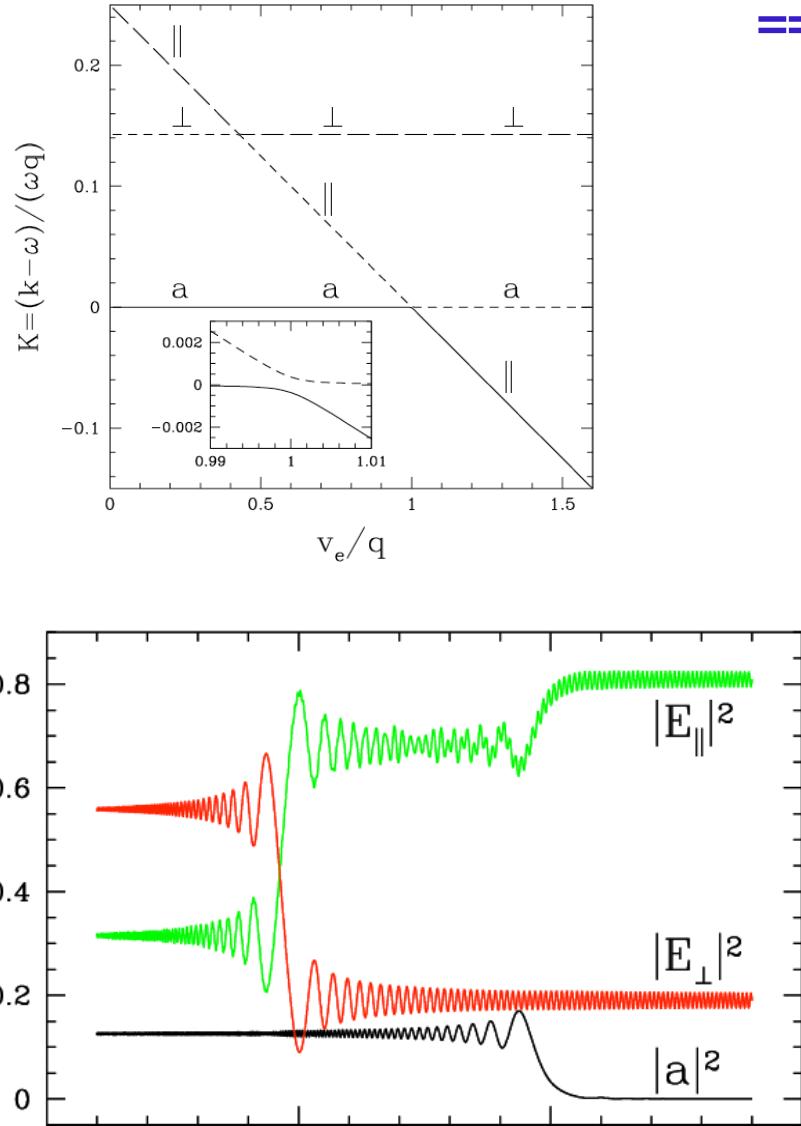
Current constraints on axion mass and coupling parameter



arXiv:0810.1874 (CAST collaboration)

Photon-Axion Conversion in Magnetic Neutron Stars

In the magnetized plasma of NSs, photons (\parallel -polarization component) can convert (resonantly a la MSW) into axions



\Longrightarrow modify radiation spectra
and polarization signals

DL & Heyl 2007; Perna et al. 2011

