

NS Seismology and Merging Binaries

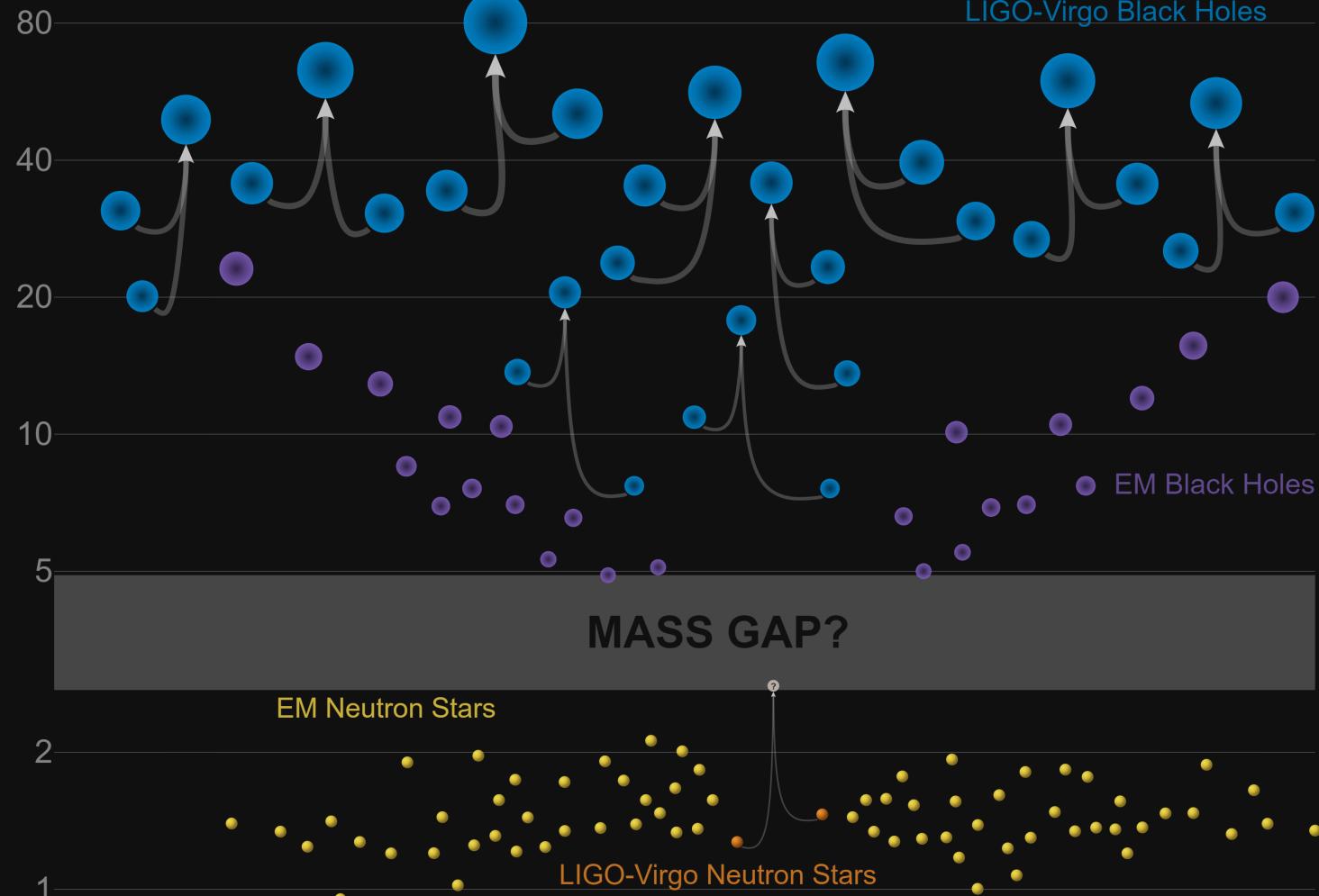
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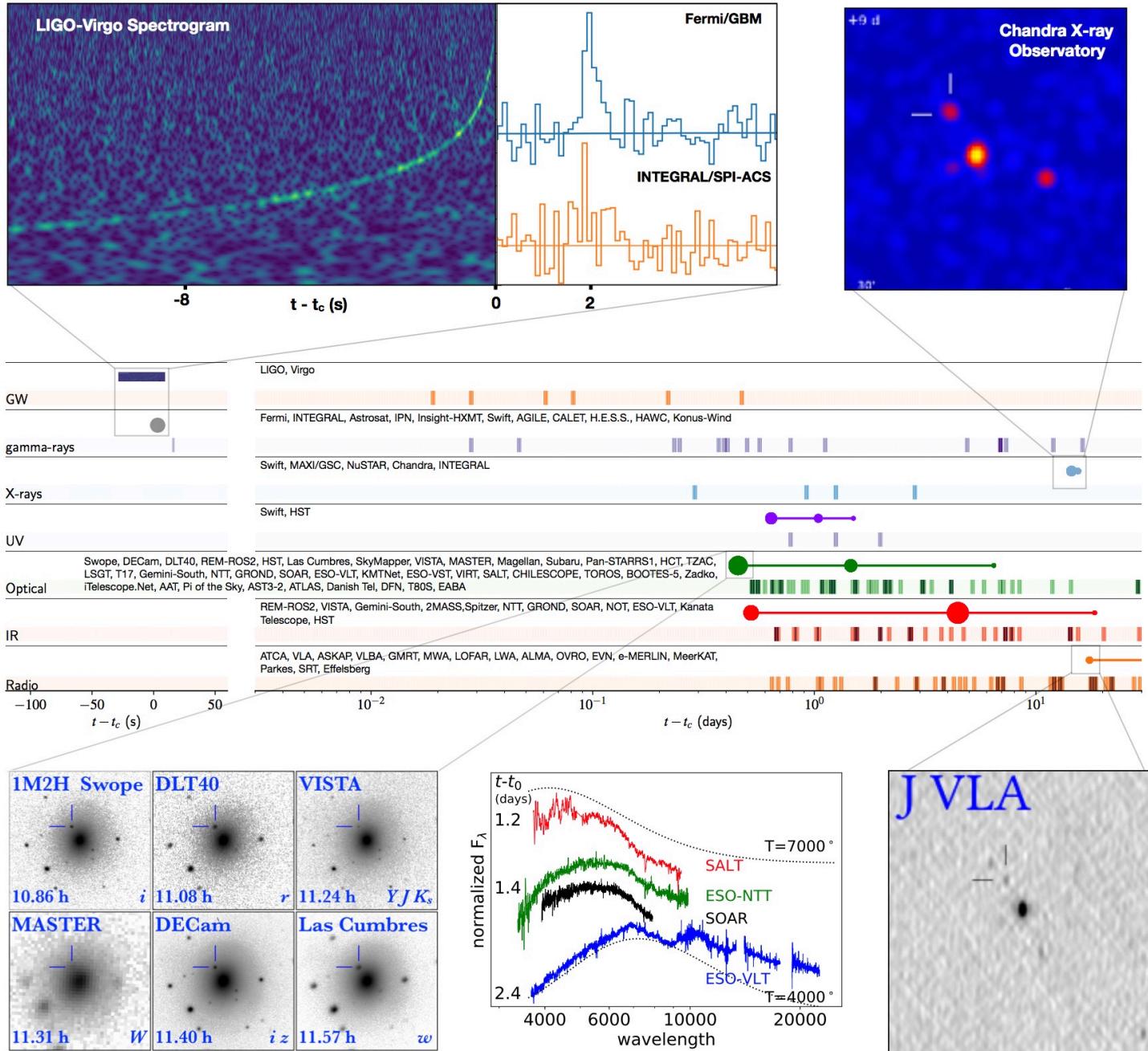
Masses in the Stellar Graveyard

in Solar Masses



Updated 2018-12-01
LIGO-Virgo | Frank Elavsky | Northwestern

GW170817 / AT2017gfo



Probe Nuclear EOS from NS Binary Mergers

Fate of the remnant of NS/NS

maximum TOV mass

Final merger GW of NS/NS or NS/BH

characteristic cut-off frequency $\sim(GM/R^3)^{1/2}$

Tidal effects on GW during Inspiral

Tidal Seismology in Merging NS Binaries

→ Probe NS EOS

- Quasi-Equilibrium tides
 - f-mode, p-modes
- Resonant (Dynamical) tides
 - g-modes, r-modes, mixed modes

Nonradial Oscillation Modes of NS

Acoustic modes (p-modes):

Stationary sound waves, depends on sound speed $c_s = (\partial P / \partial \rho)^{1/2}$
f-mode: global acoustic mode (no radial node), $\sigma \sim \sqrt{G \bar{\rho}} \sim 2$ kHz

G-modes (gravity modes = buoyancy modes):

Arise from stable stratification:

Density jumps in crust: interface modes

Core composition gradient: $x=n_p/n$ increases with density

$$N = g \left[\left(\frac{\partial \rho}{\partial x} \right)_P \left(\frac{dx}{dP} \right) \right]^{1/2} \quad \sim 100 \text{ Hz or less}$$

depends on symmetry energy

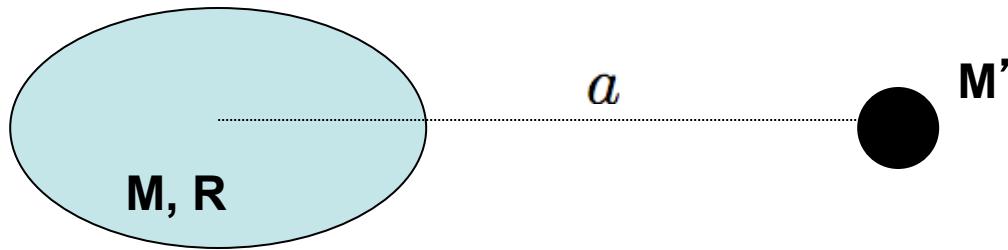
depends on superfluidity...

R-modes (inertial modes, Rossby modes):

arise from Coriolis force

$$\sigma \sim \Omega_s$$

“Quasi-Equilibrium” Tide (F-mode Distortion)



Tidally induced quadrupole $\sim k_2 \frac{M'R^5}{a^3}$ k_2 = Love number

$$V = -\frac{MM'}{a} - \mathcal{O}\left(k_2 \frac{M'^2 R^5}{a^6}\right)$$

$$\rightarrow d\Phi = d\Phi^{(0)} \left[1 - \mathcal{O}\left(\frac{k_2 R^5}{a^5}\right) \right] \quad (\text{GW Phase Shift})$$

→ Important only at small separation (just prior to merger)

(Kochanek; Bildsten & Cutler; Lai, Rasio & Shapiro, etc 1990s)

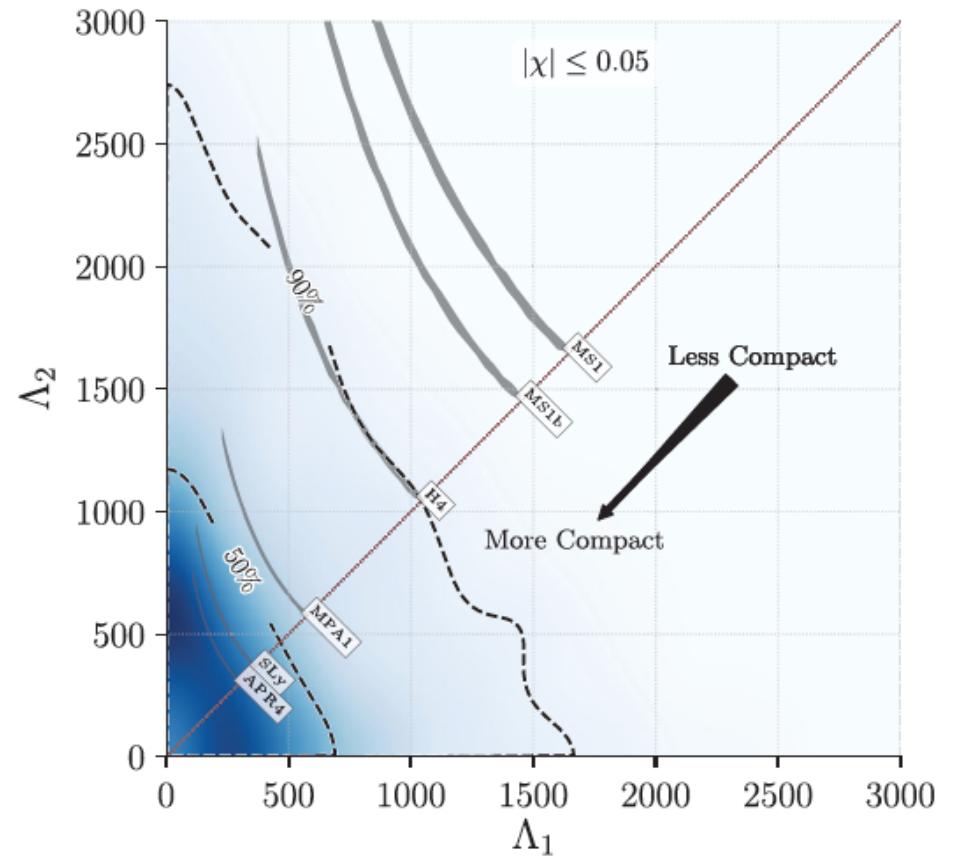
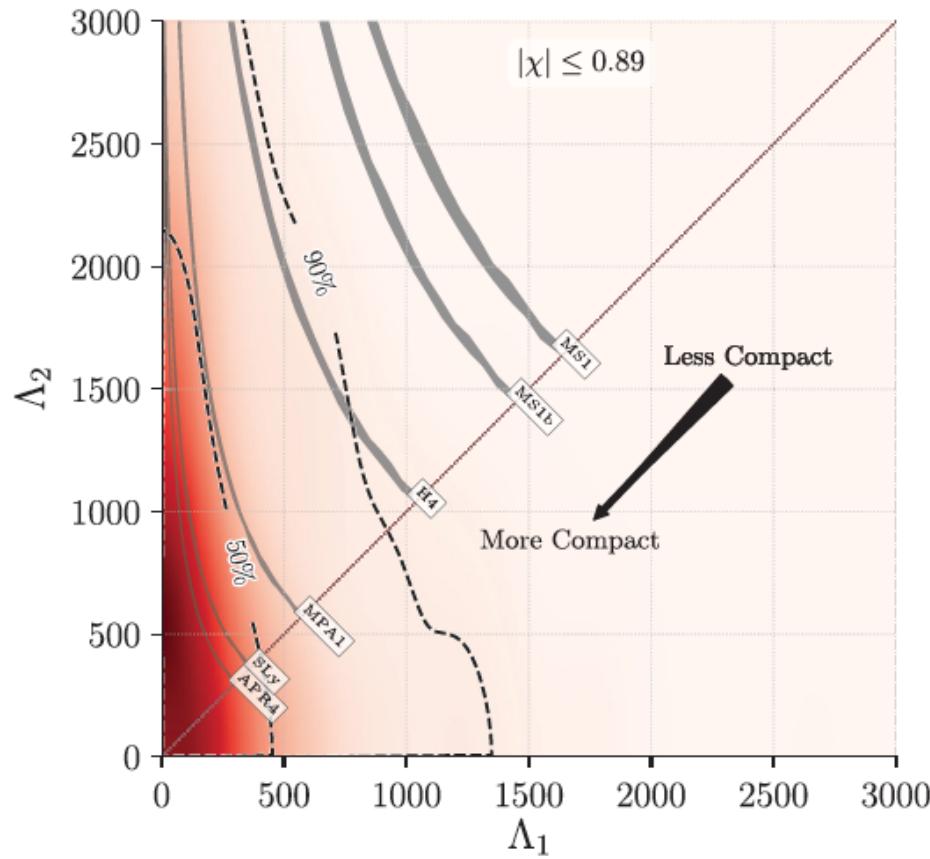
Numerical GR Quasi-equilibrium NS binary sequence

(Baumgarte, Shapiro, Teukolsky, Shibata, Meudon group, etc. 1990s--200x)

Recent (semi-analytic) GR calculation of tidal effect

(Hinderer, Flanagan, Poisson, Damour, Penner, Andersson, Jones, etc., 2008+)

GW170817:



$$\Lambda_2^{(i)} = \frac{2}{3} k_2^{(i)} \left[\left(\frac{c^2}{G} \right) \left(\frac{R_i}{M_i} \right) \right]^5$$

Quasi-Equilibrium Tide = F-mode ($l = 2$)

$$\ddot{\xi} + \sigma^2 \xi \sim e^{i2\Omega t}$$
$$\rightarrow \xi \sim \frac{e^{i2\Omega t}}{\sigma^2 - (2\Omega)^2}$$

The usual tidal distortion calculation assumes $\sigma_f \gg 2\Omega$

Dynamical correction to “equilibrium tide”:

$$d\Phi = d\Phi^{(0)} \left[1 - \mathcal{O} \left(\frac{k_2 R^5}{a^5} \right) \frac{1}{1 - 4\Omega^2/\sigma_f^2} \right]$$

Resonant Tides: Excitations of Internal Modes

NS has low-frequency oscillation modes:

g-modes (~ 100 Hz)

r-modes

$$\ddot{\xi} + \sigma_\alpha^2 \xi \sim e^{im\Omega t}$$

Resonance: $\sigma_\alpha = m\Omega_{\text{orb}}, \quad m = 2, 3, \dots$

Rosonant tidal excitations of NS modes during inspiral

→ transfer orbital energy to NS

→ GW phase shift

Resonant Excitations of NS Oscillations During Inspiral

$$\Delta\Phi \propto \frac{(Q_{\alpha,lm})^2}{\epsilon_\alpha \sigma_\alpha}, \quad Q_{\alpha,lm} = \int d^3x (\delta\rho_\alpha)^\star r^l Y_{lm}$$

Non-rotating NS:

G-mode (Reisenegger & Goldreich 94; DL 94; Shibata 94)

Rotating NS:

G-mode, F-mode, R-mode (Ho & DL 99)

Inertial modes (DL & Wu 06)

R-mode (excited by gravitomagnetic force; Racine & Flanagan 06)

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Recent works:

- Superfluid NSs: Yu & Weinberg 2017
- Gravity-inertial modes, scalings (parameterized EOS): Xu & DL 2017



Wenrui Xu
(Cornell → Princeton)

G-modes

crustal density continuities

stable composition stratification of core: symmetry energy, superfluidity

Parameterize the uncertainties:

$$P \propto \rho^\gamma, \quad \Gamma = (\partial \ln P / \partial \ln \rho)_{\text{ad}}$$

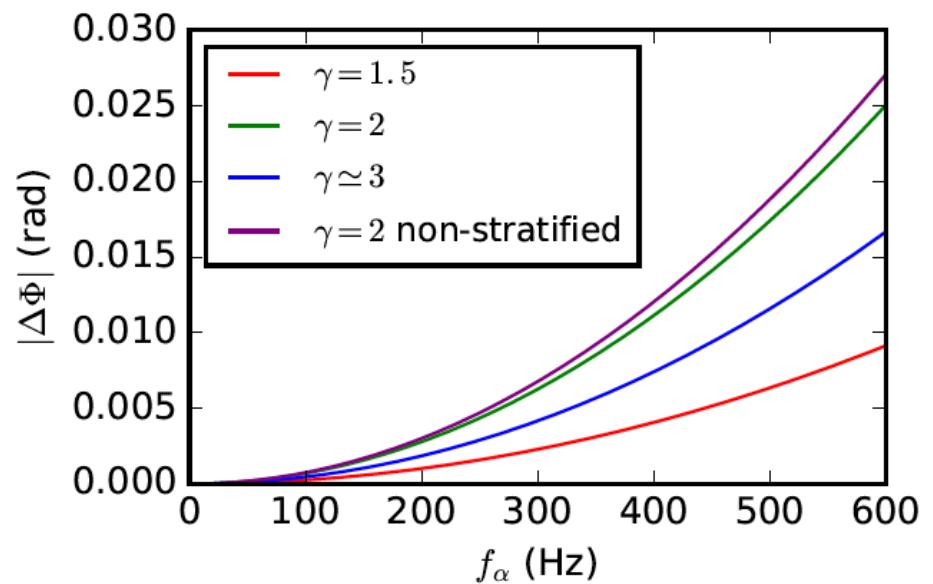
$$\omega_\alpha \propto (\Gamma - \gamma)^{1/2} M^{1/2} R^{-3/2}$$

$$Q_{\alpha,2m} \propto \Gamma - \gamma,$$

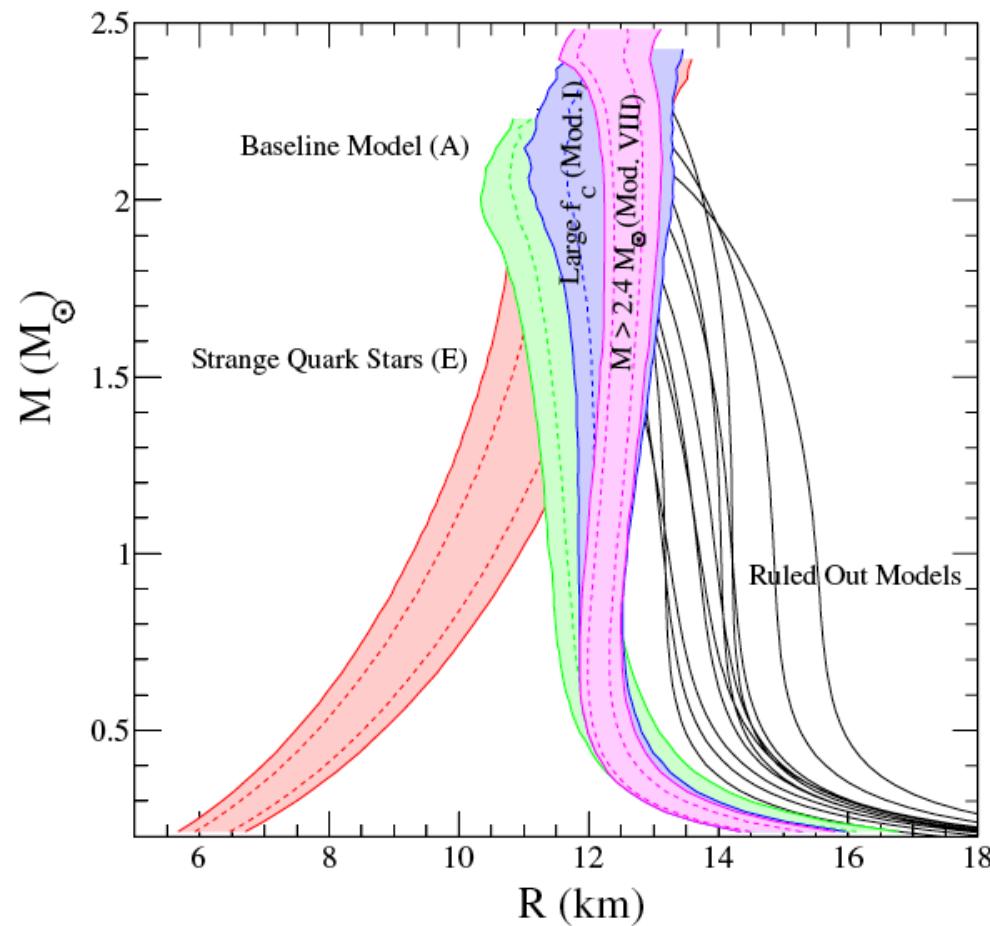
$$\begin{aligned} \Delta\Phi = & -0.060 \left(\frac{R_{10}^8}{M_{1.4}^6} \right) \frac{2}{q(1+q)} \left(\frac{f_\alpha}{100 \text{ Hz}} \right)^2 \\ & \times \left(\frac{\bar{f}_\alpha}{100 \text{ Hz}} \right)^{-4} \left(\frac{\bar{Q}_{\alpha,22}}{10^{-3}} \right)^2, \end{aligned}$$

Xu & DL 17

Note: $M = 1.2M_\odot$, $R = 13$ km NS
 → x 21 Important !



NS Mass-Radius



Measured NS mass:
1.17–2.0 Solar mass

Steiner et al 2013

G-mode excitation during binary inspiral

$$\Delta\Phi = -0.060 \left(\frac{R_{10}^8}{M_{1.4}^6} \right) \frac{2}{q(1+q)} \left(\frac{f_\alpha}{100 \text{ Hz}} \right)^2 \\ \times \left(\frac{\bar{f}_\alpha}{100 \text{ Hz}} \right)^{-4} \left(\frac{\bar{Q}_{\alpha,22}}{10^{-3}} \right)^2,$$

It is a coincidence that $|\Delta\Phi| \lesssim 1 \dots$ Unlike equilibrium tide

For white dwarf binaries, $|\Delta\Phi| \gg 1$ (LISA), lead to significant tidal heating

Challenge to nuclear physicists:

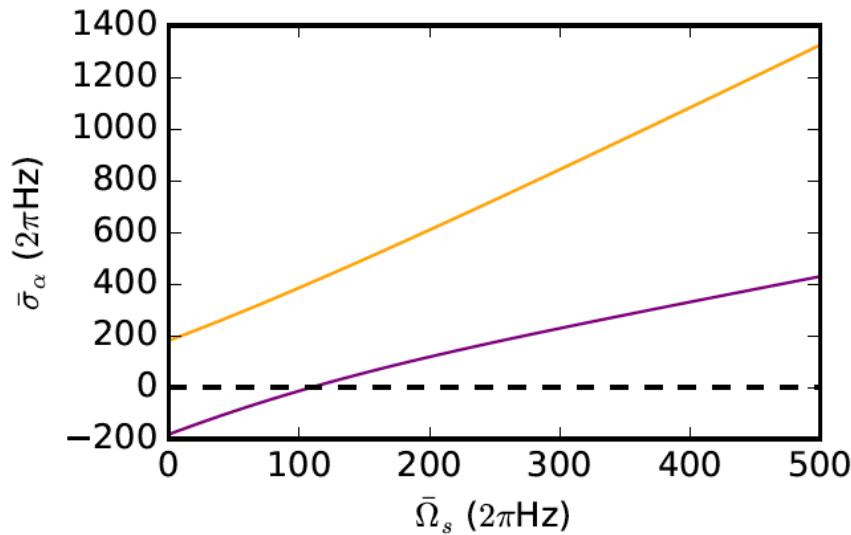
Survey $E_{\text{sym}}(n)$ (constrained by exps/obs), calculate g-modes and $\Delta\Phi$

With rotation: pure r-modes (inertial modes)

$$\Delta\Phi = \mp 0.0027 \left(\frac{R_{10}^8}{M_{1.4}^6} \right) \frac{2}{q(1+q)} \left(\frac{\epsilon_\alpha |\sigma_\alpha|}{\Omega_s^2} \right)^{-1} \\ \times \left(\frac{Q_{\alpha,2m}}{0.02 \hat{\Omega}_s^2} \right)^2 \left(\frac{f_s}{500 \text{ Hz}} \right)^2 \left| \mathcal{D}_{m\pm 2}^{(2)} \right|^2,$$

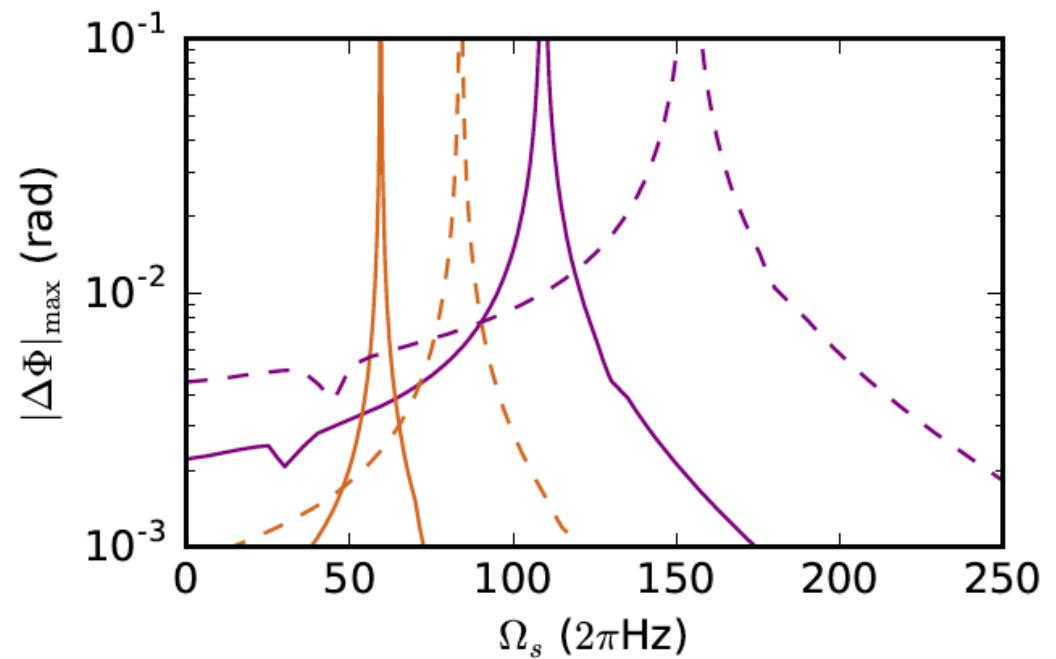
Usually negligible (unless R^8/M^6 is large)...

With rotation: Mixed modes (Inertial-Gravity Modes)



Rotation can reduce (retrograde) g-mode $|\sigma_\alpha|$, thus increase $|\Delta\Phi|$

$\gamma = 2$ $\Gamma - \gamma = 0.01$
 $\gamma \approx 3$ $\Gamma - \gamma = 0.02$

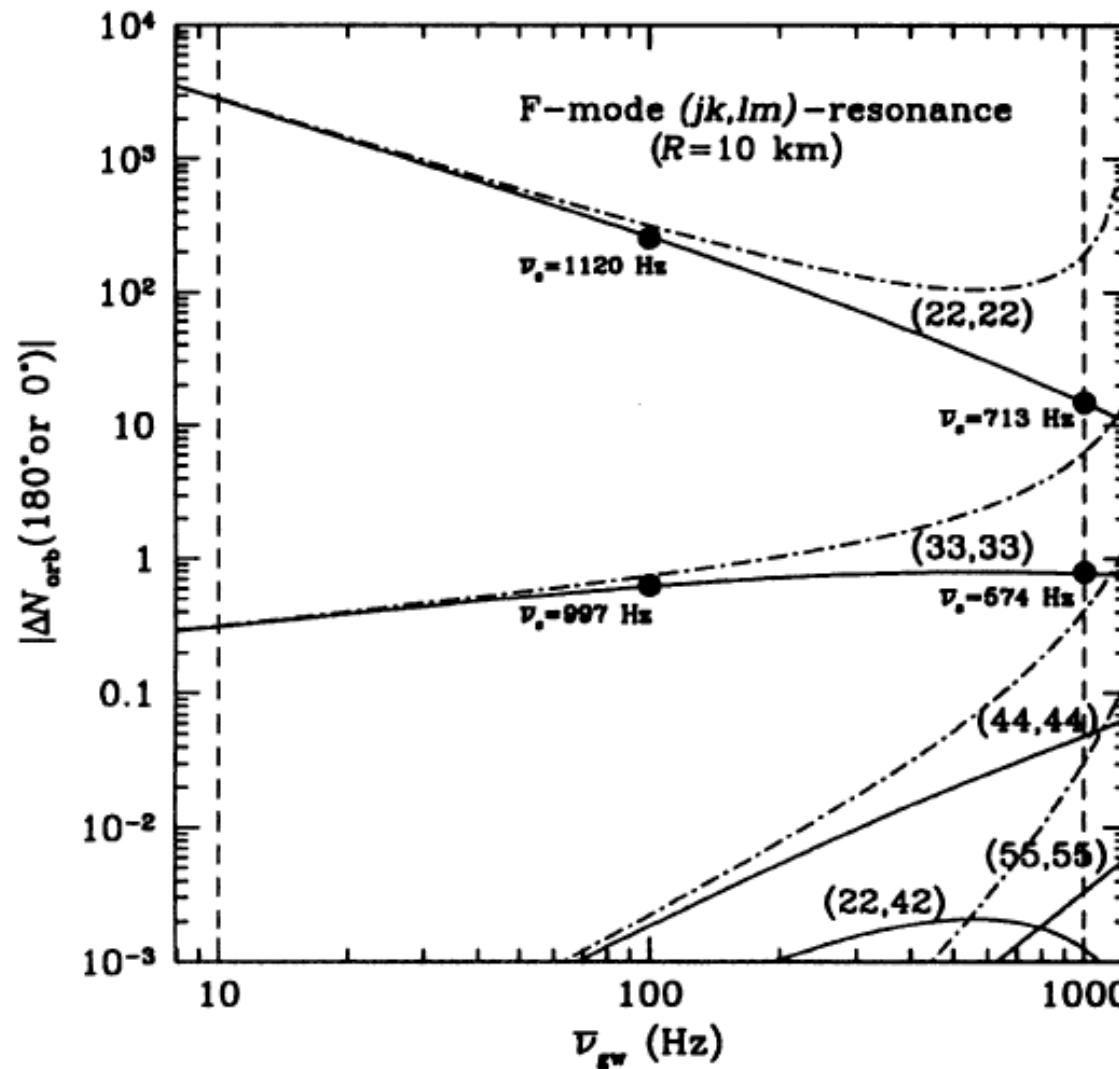


Xu & DL 2017

Resonant excitation of F-modes?

Very strong ... possible with **very rapid rotations**

Wynn Ho & DL 1999

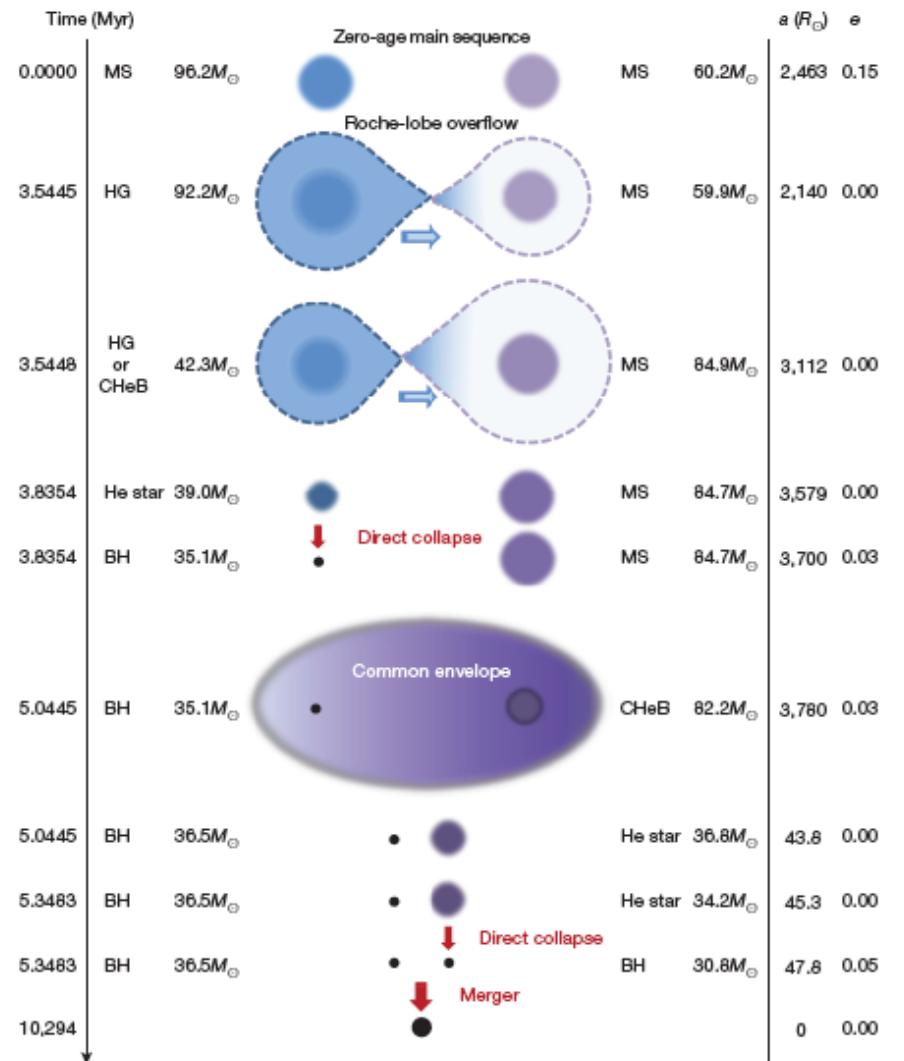


Physics allows it.
Does Nature provide it?

So far: Circular mergers

What about eccentric mergers?

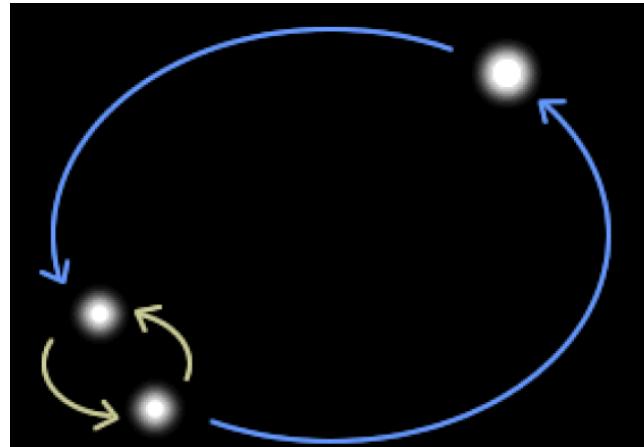
Standard Binary Evolution Channel → Circular merger



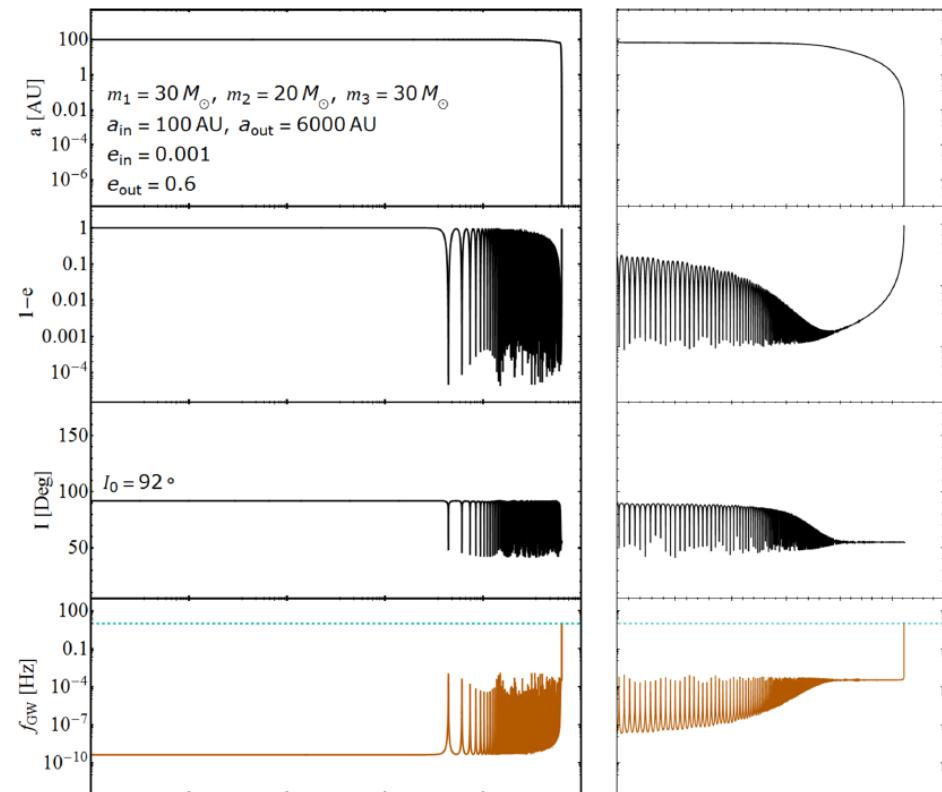
Belczynski +16

Possibility of eccentric mergers: Dynamical formation merging compact binaries

- Binary + single scatterings dense clusters
- Binary mergers induced by external companion (SMBH, stellar triple)

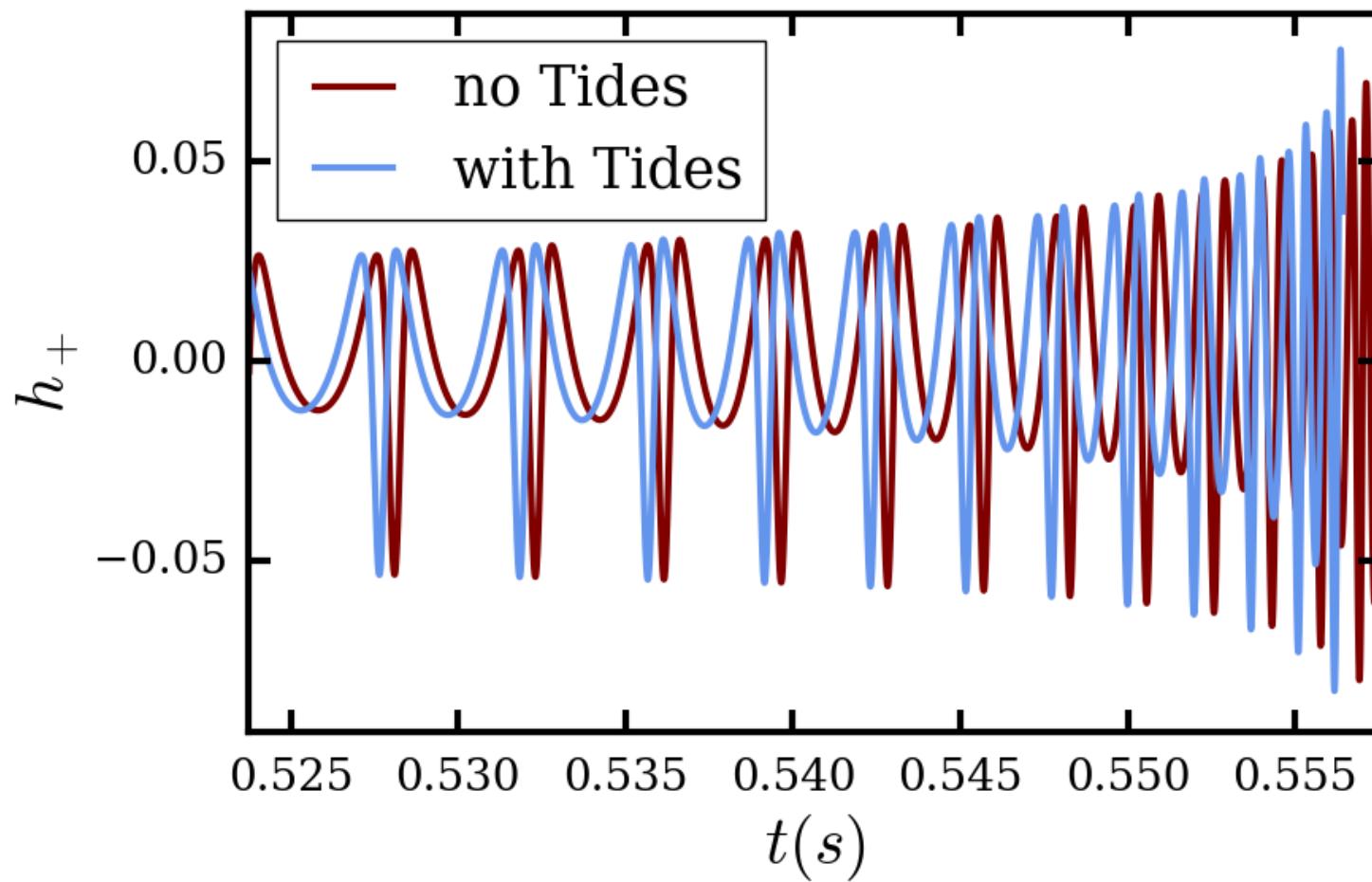


Bin Liu & DL 2017,18,19
also Silsbee & Tremaine 2017 etc



Tidal effect in eccentric mergers

Michelle Vick & DL 2019



Initial $e=0.9$, $r_p=5R$

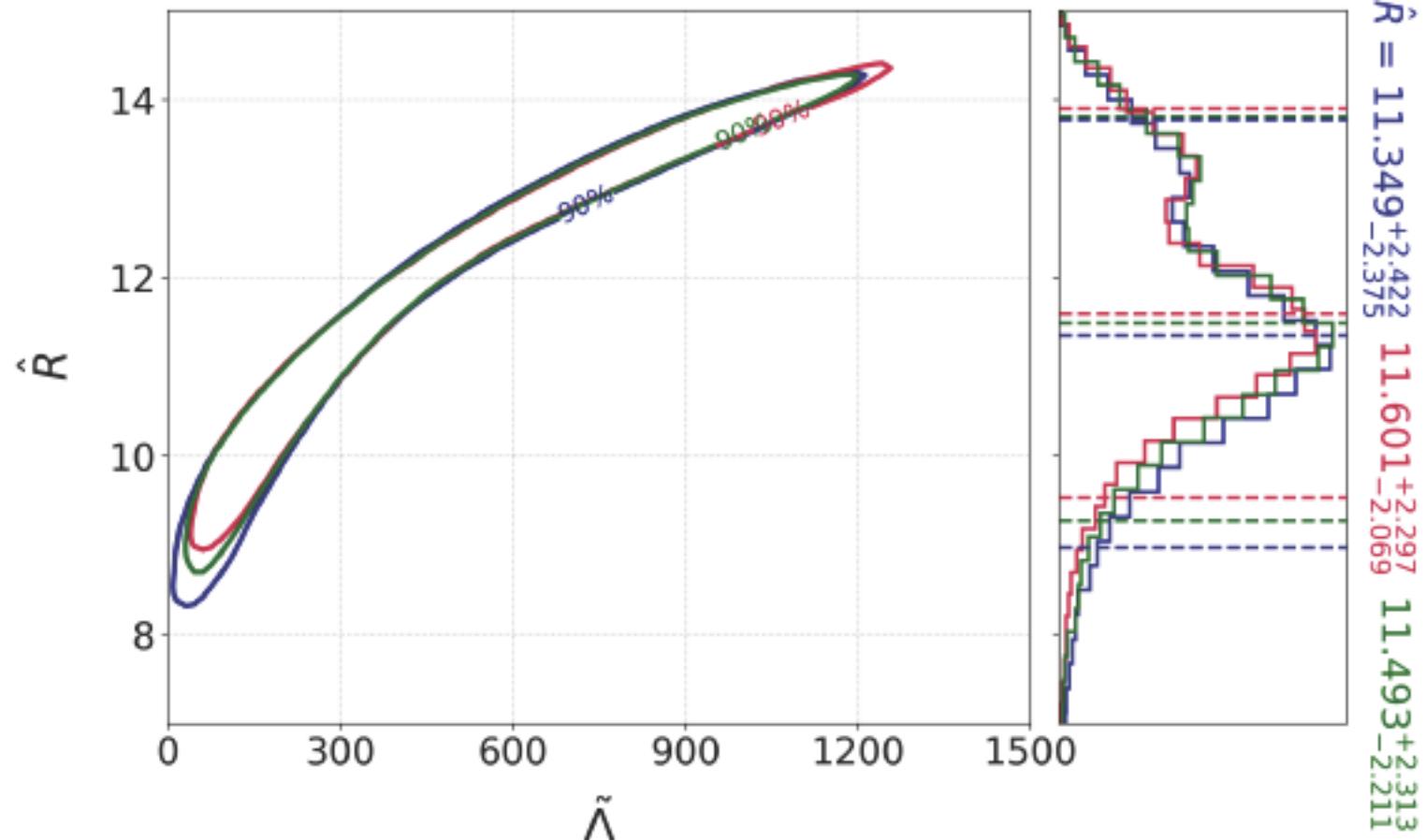
Summary

- Tidal distortion (non-resonant f-mode) constrains R-M
- Resonant tidal excitation of modes occurs at low frequency (< 100 Hz).
If strong, would provide a unique probe of EOS (e.g. symmetry energy)
 - g-modes: Important for low-M, large-R NSs
 - rotation enhances resonance.
 - f-mode resonance: Important for very rapid rotation
- Highly eccentric mergers ?

Note: Resonant tides are certainly important for WD binaries
(or WD-NS, WD-BH): LISA band
Also for 1-10 Hz band (Cosmic Explorer, Einstein Telescope)

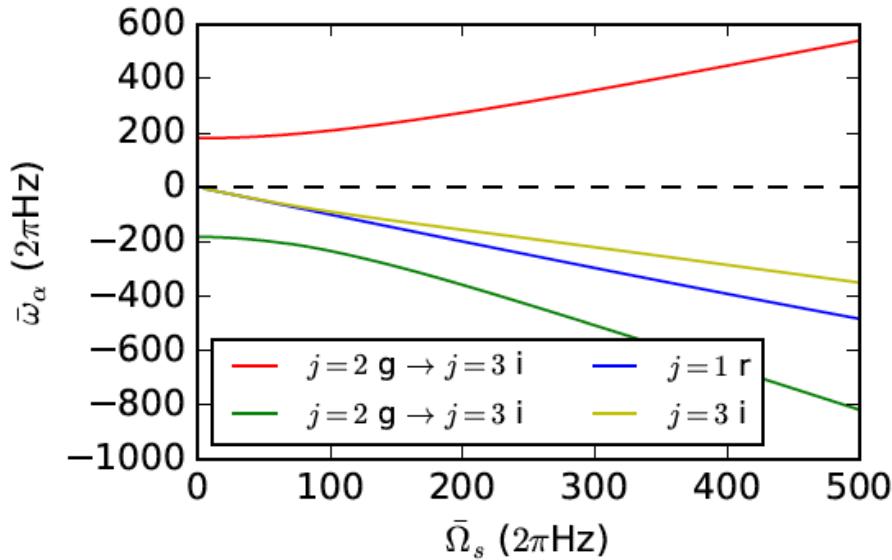
Thanks.

GW170817:



De, Finstad, Lattimer et al. 2018

With rotation: Many other inertial modes (modified by stratification)



Xu & DL 2017

Not as important...

