

## Rapid solutions of problems by quantum computation

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### Deutsch-Jozsa Problem

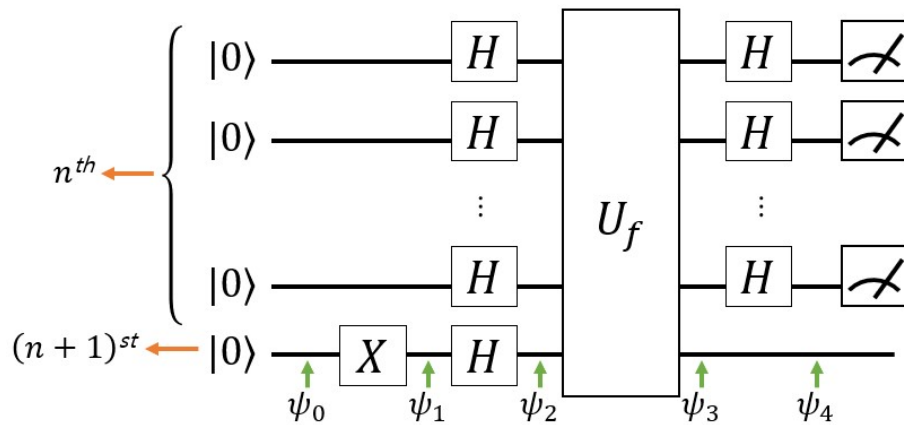
Given  $n$ -bit function  $f: \{0,1\}^n \rightarrow \{0,1\}$  which is guaranteed to either be *balanced* or *constant*, determine whether which is the case. *Balanced* means  $f(x) = 0$  for half of the  $x \in \{0,1\}^n$  and  $f(x) = 1$  for the remaining input domain, and *constant* means  $f$  is the same for all  $x \in \{0,1\}^n$ .

Classically, we can randomly choose  $k(\leq n)$  inputs  $x_1, \dots, x_k \in \{0,1\}^n$ , evaluate  $f(x_i)$  for  $i = 1, \dots, k$ , and answer "constant" if  $f(x_1) = \dots = f(x_k)$  and "balanced" otherwise. If we want to get a correct answer every time with this approaches,  $(2^{n-1} + 1)$ -queries are needed in the worst case. However, we claim that a single query is sufficient to determine the solution with certainty.

Please note that *this problem* is the  $n$ -bit generalization of the *Deutsch's problem*.

### Deutsch-Jozsa Algorithm

Deutsch-Jozsa algorithm can be implemented with the following  $(n + 1)$ -qubits quantum circuit (represented by  $(n + 1)$ -lines) initialized with, respectively,  $|0\rangle$ :



To analyze the state of the quantum circuit, we first need to understand what  $H^{\otimes n}|a\rangle$  equals for arbitrary  $a \in \{0,1\}$ . (Here,  $H^{\otimes n} = H \otimes \dots \otimes H$  ( $n$  times)). For this, we begin with a clean and formal way for the action of  $H$  on a single qubit. For  $a_1 \in \{0,1\}$ , we have

$$H|a_1\rangle = \frac{1}{\sqrt{2}}|0\rangle + \frac{1}{\sqrt{2}}(-1)^{a_1}|1\rangle = \frac{1}{\sqrt{2}}\sum_{b_1 \in \{0,1\}}(-1)^{a_1 b_1}|b_1\rangle.$$

If we instead had two qubits, starting in state  $|a\rangle$  for  $a = a_1 a_2 \in \{0,1\}^2$ , and applied Hadamard transforms to both, we would obtain

$$\begin{aligned} (H \otimes H)|a\rangle &= H|a_1\rangle \otimes H|a_2\rangle \\ &= \left(\frac{1}{\sqrt{2}}\sum_{b_1 \in \{0,1\}}(-1)^{a_1 b_1}|b_1\rangle\right) \otimes \left(\frac{1}{\sqrt{2}}\sum_{b_2 \in \{0,1\}}(-1)^{a_2 b_2}|b_2\rangle\right) = \frac{1}{(\sqrt{2})^2}\sum_{b \in \{0,1\}^2}(-1)^{a_1 b_1 + a_2 b_2}|b\rangle. \end{aligned}$$

We can generalize this to  $n$ -qubit states. Specifically, if we write the state  $|a\rangle = |a_1 \dots a_n\rangle$  as  $|a_1\rangle \otimes \dots \otimes |a_n\rangle$ , we have

$$\begin{aligned} H^{\otimes n}|a\rangle &= H|a_1\rangle \otimes \dots \otimes H|a_n\rangle \\ &= \left(\frac{1}{\sqrt{2}}\sum_{b_1 \in \{0,1\}}(-1)^{a_1 b_1}|b_1\rangle\right) \otimes \dots \otimes \left(\frac{1}{\sqrt{2}}\sum_{b_n \in \{0,1\}}(-1)^{a_n b_n}|b_n\rangle\right) = \frac{1}{(\sqrt{2})^n}\sum_{b \in \{0,1\}^n}(-1)^{a_1 b_1 + \dots + a_n b_n}|b\rangle. \end{aligned}$$

We also define  $a \cdot b = \sum_{i=1}^n a_i b_i \pmod{2}$ , (here, the mod 2 arises since the base is  $(-1)$ , so all we care about is if the exponent  $a \cdot b$  is even or odd) so that we may write

$$H^{\otimes n}|a\rangle = \frac{1}{(\sqrt{2})^n}\sum_{b \in \{0,1\}^n}(-1)^{a \cdot b}|b\rangle.$$

First, the algorithm starts with  $|\psi_0\rangle = |00 \dots 0\rangle = |0\rangle \otimes \dots \otimes |0\rangle$  ( $(n+1)$  times)  $= |0\rangle^{\otimes n} \otimes |0\rangle$

Analogously to Deutsch's algorithm, It is clear that

$$|\psi_1\rangle = (I^{\otimes n} \otimes X)|\psi_0\rangle = I|00 \dots 0\rangle \otimes X|0\rangle = |0\rangle^{\otimes n} \otimes |1\rangle,$$

$$|\psi_2\rangle = (H^{\otimes n} \otimes H)|\psi_1\rangle = H^{\otimes n}|00 \dots 0\rangle \otimes H|1\rangle$$

$$= \frac{1}{(\sqrt{2})^n}\sum_{x \in \{0,1\}^n}(-1)^{0 \cdot x_1 + \dots + 0 \cdot x_n}|x\rangle \otimes \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}}\right) = \frac{1}{(\sqrt{2})^n}\sum_{x \in \{0,1\}^n}|x\rangle \otimes \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}}\right).$$

$$\therefore H^{\otimes n}|a\rangle = \frac{1}{(\sqrt{2})^n}\sum_{b \in \{0,1\}^n}(-1)^{a \cdot b}|b\rangle \text{ and } H|1\rangle = \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}}\right)$$

Second, remembering that  $U_f|x\rangle|y\rangle = |x\rangle|y \oplus f(x)\rangle$ :

$$|\psi_3\rangle = U_f|\psi_2\rangle = \frac{1}{(\sqrt{2})^n}\sum_{x \in \{0,1\}^n}|x\rangle \otimes \left(\frac{|0 \oplus f(x)\rangle - |1 \oplus f(x)\rangle}{\sqrt{2}}\right) = \frac{1}{(\sqrt{2})^n}\sum_{x \in \{0,1\}^n}(-1)^{f(x)}|x\rangle \otimes \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}}\right).$$

$$\therefore |\psi_3\rangle = \begin{cases} \frac{1}{(\sqrt{2})^n}\sum_{x \in \{0,1\}^n}|x\rangle \otimes \left(\frac{|0 \oplus 0\rangle - |1 \oplus 0\rangle}{\sqrt{2}}\right) = \frac{1}{(\sqrt{2})^n}\sum_{x \in \{0,1\}^n}|x\rangle \otimes \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}}\right), & \text{if } f(x) = 0 \\ \frac{1}{(\sqrt{2})^n}\sum_{x \in \{0,1\}^n}|x\rangle \otimes \left(\frac{|0 \oplus 1\rangle - |1 \oplus 1\rangle}{\sqrt{2}}\right) = \frac{1}{(\sqrt{2})^n}\sum_{x \in \{0,1\}^n}|x\rangle \otimes \left(\frac{|1\rangle - |0\rangle}{\sqrt{2}}\right), & \text{if } f(x) = 1 \end{cases}$$

Lastly, we must apply the last set of Hadamard gates,

$$\begin{aligned} |\psi_4\rangle &= (H^{\otimes n} \otimes I)|\psi_3\rangle = H^{\otimes n}\left(\frac{1}{(\sqrt{2})^n}\sum_{x \in \{0,1\}^n}(-1)^{f(x)}|x\rangle\right) \otimes I\left(\frac{|0\rangle - |1\rangle}{\sqrt{2}}\right), \\ &= \left(\frac{1}{(\sqrt{2})^n}\sum_{x \in \{0,1\}^n}(-1)^{f(x)}\left(\frac{1}{(\sqrt{2})^n}\sum_{z \in \{0,1\}^n}(-1)^{x \cdot z}|z\rangle\right)\right) \otimes \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}}\right), \\ &= \left(\sum_{z \in \{0,1\}^n}\left(\frac{1}{2^n}\sum_{x \in \{0,1\}^n}(-1)^{f(x) + x \cdot z}\right)\right) \otimes \left(\frac{|0\rangle - |1\rangle}{\sqrt{2}}\right). \end{aligned}$$

**Case 1: balanced  $f$ .** Consider the amplitude on  $|z\rangle = |00 \dots 0\rangle$ :

$$\frac{1}{2^n} \sum_{x \in \{0,1\}^n} (-1)^{f(x) + x_1 \times 0 + \dots + x_n \times 0} = \frac{1}{2^n} \sum_{x \in \{0,1\}^n} (-1)^{f(x)} = \frac{1}{2^n} \times 0 = 0.$$

Since  $f$  is balanced, all the terms in this sum of  $(-1)^{f(x)}$  cancel, i.e. the sum equals 0. Thus, we conclude that the amplitude on  $|z\rangle = |00 \dots 0\rangle$  is 0, and so we will never see outcome  $|00 \dots 0\rangle$  in the final measurement.

**Case 2: Constant  $f$ .** Consider the amplitude on  $|z\rangle = |00 \dots 0\rangle$  and factor out  $(-1)^{f(x)}$  in  $|\psi_4\rangle$ :

$$|\psi_4\rangle = (-1)^{f(x)} \left( \sum_{z \in \{0,1\}^n} \left( \frac{1}{2^n} \sum_{x \in \{0,1\}^n} (-1)^{x \cdot z} |z\rangle \right) \right) \otimes \left( \frac{|0\rangle - |1\rangle}{\sqrt{2}} \right),$$

$$\frac{1}{2^n} \sum_{x \in \{0,1\}^n} (-1)^{x_1 \times 0 + \dots + x_n \times 0} = \frac{1}{2^n} \sum_{x \in \{0,1\}^n} (-1)^0 = \frac{1}{2^n} \sum_{x \in \{0,1\}^n} 1 = \frac{1}{2^n} \times 2^n = 1.$$

In other words, the state  $|0\rangle^{\otimes n} \otimes \left( \frac{|0\rangle - |1\rangle}{\sqrt{2}} \right)$  has amplitude 1. Since the quantum state is a unit vector, we can conclude that we must have  $|\psi_4\rangle = (-1)^{f(x)} |00 \dots 0\rangle \otimes \left( \frac{|0\rangle - |1\rangle}{\sqrt{2}} \right)$ , i.e. all the weight is on this one term. Thus, if  $f$  is constant, then measuring the first  $n$ -qubits yields outcome  $|00 \dots 0\rangle$  with certainty.

## Conclusion

By measuring the first  $n$ -qubits, we can conclude that  $f$  is balanced (any  $n$ -bit measurement outcome except  $|00 \dots 0\rangle$ ) or constant (always  $|00 \dots 0\rangle$  outcome) with just a single-query.