

ZAMS Project Report

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Introduction

The goal of this project is to calculate the radius and luminosity of low-mass stars and see how their structure and luminosity depend on p+p reaction rate. By finding the radius at which the surface luminosity is equal to the nuclear luminosity, the zero-aged main-sequence can be determined for low mass stars. We can do so by modifying the routines we developed in the white dwarf project.

Step 1

To start, we needed to revisit the last project's goal and make sure that we have found the listed scalings in the instructions. After reworking the code from project 2 into a correct fashion, we have found the the P_c should scale with a factor of 0.77 and p_c should scale with a factor of 5.99. Our values of δ , ξ , and η are also sufficiently small after a test to make sure they agree to two significant digits across a range of masses. The parameter values that work include: $\eta = 1.0e-10$, $\delta = 1.0e-8 \cdot M_{\text{sun}}$, and $\xi = 0.05$. These values will be further investigated in step 8 of the instructions. Also, by rerunning the equation of state with K multiplied by 0.5 and 2, we found that the equations for P_c and p_c remained satisfied. We did not write a script specific for this project to test these scenarios, as this code was crafted in project 2, but corrected that project 2 code and quickly performed these tests from that repository.

Steps 2 – 5; 8

For these steps, we needed to verify that the routines we finalized to be used for the latter parts of the project worked as intended. For the second and third steps we needed to compute the mean molecular weight for fully ionized plasma and write a routine that computes the density and temperature along an adiabat; the tests under *test_eos.py* for these two steps worked as intended. Since we are computing ZAMS that are at the start of their hydrogen

burning phase we can approximate their composition to contain hydrogen, helium and nitrogen with the corresponding mass fractions of $X(\text{H})=0.706$, $X(\text{He})=0.275$ and $X(\text{N})=0.019$; we used these exact values (they add to 1) to calculate the mean molecular weight. For step three, we needed to create a routine that would compute the heating rate of a low mass star which included in this calculation is a pp factor that gets multiplied by the pp rate; under *test_reactions.py*, the test for part four was successful. For step five, we wrote a routine that computes the central pressure, density, and temperature given a stellar mass and radius; this test under *test_reactions.py* was successful. We wrote a routine to interpolate the effective temperatures from Table 1 in the instructions and for the surface luminosity; these tests were successful under *test_zams.py*.

These tests all can be ran by deleting the necessary line in their respective test routine and then running *testing.py*. For notice, the routines in *structure.py* are fitted so they can produce correct results for the testing, while the updated routines for *central_thermal* and *Teff* are in *structure_for_main.py* and are used with unit inputs that fit better with the main files.

Step 6

For this part, we needed to copy four routines from the white dwarf project and modify them for use in this one. Besides updating the routines so they now fit this project, the main difference was requiring mass and radius as inputs instead of P_c for *central_values* and *integrate*; while perhaps these can be computed with P_c , this seemed like a necessary change considering that the new routine of *central_thermal* took in mass and radius as inputs. For example, in project two, *central_values* includes *get_scaled_rho* which input P , so *central_values* needed to input P_c , as well; however, since *central_thermal* inputs m and r and is necessary likewise for *central_values* in this project, *central_values* should take in m and r instead of P_c . Additionally, instruction specific modifications included adding a luminosity array and differential, including a new scale for stepsize, and a boundary condition for luminosity. As stated previously, the two routines for central thermal values and effective temperature are updated in *structure_for_main.py* for easier use with the main files.

Step 7

For this part, we needed to test the convergence of the integration by varying parameters η , ξ , and δ to ensure that our integration satisfied precision requirements. Our script under *test_convergence.py* involved two results: a return of arrays for the four integration return arrays (mass, radius, pressure, and luminosity) and a `mass_keep` array that returned each variation of η , ξ , and δ that satisfied the testing mass of $0.2 M_{\text{sun}}$. While this is essentially the same as the mass array, it helps to immediately where the convergence is happening, at least for the mass parameter. In the script, one line of the three at the bottom can be uncommented to test that parameter's convergence; the default that is uncommented is for δ .

From the project 2 solutions, we know that preferred values are $\eta = 1.0\text{e-}10$, $\delta = 1.0\text{e-}8 M_{\text{sun}}$, and $\xi = 0.05$. We used these values for this project, but we still needed to confirm their correctness. For δ , the output parameters began to converge starting at $\delta = 1.0\text{e-}7 M_{\text{sun}}$; this means it is reasonable to use $\delta = 1.0\text{e-}8 M_{\text{sun}}$ as it is not too small, but still corresponds to after the integration converged. For η , the output parameters began to converge starting at $\eta = 1.0\text{e-}6$; this means it is reasonable to use $\eta = 1.0\text{e-}10$ as it is not too small, but still corresponds to after the integration converged. For ξ , the output parameters began to converge starting at $\xi = 0.1$; this means it is reasonable to use $\xi = 0.05$ as it is not too small, but still corresponds to after the integration converged. As stated before, only one target mass of $m = 0.2 M_{\text{sun}}$ was necessary and the `pp_factor` used was 1.0.

Step 9

For this part, we needed to create a root finder routine for stellar radius. This involved writing a function that computing the difference in nuclear and surface luminosity. We then needed to create a radius finder routine that solved for this radius in the luminosity equation. This root finder requires an interval for the radius. An effort to try and make use of the virial relation that central temperature varies very little with mass ended in either too many iterations occurring or $f(a)$ and $f(b)$ being of equal signs, so instead the interval of $[0.001 R_{\text{sun}}, 4 R_{\text{sun}}]$ was utilized. Using prior knowledge about ZAMS stars we are analyzing, this interval is still sensible

because we know that the minimum radius must be at least $0.001 \cdot R_{\text{sun}}$ and the maximum of $4 \cdot R_{\text{sun}}$ should cover all masses in between the intended range of $[0.1 \cdot M_{\text{sun}}, 0.3 \cdot M_{\text{sun}}]$.

Step 10

a)

For this part, we needed to find the main-sequence radii, and hence L and T_{eff} , for several masses (25) between $0.1 M_{\text{sun}}$ and $0.3 M_{\text{sun}}$. This plot of $\log(L/L_{\text{sun}})$ against $\log(T_{\text{eff}}/K)$ required that the $\log(T_{\text{eff}}/K)$ axis (x-axis) display as inverted to correspond to a Hertzsprung-Russell diagram.

b)

For this part, we needed to plot $\log(T_c/K)$ against $\log(\rho_c/g \text{ cm}^{-3})$.

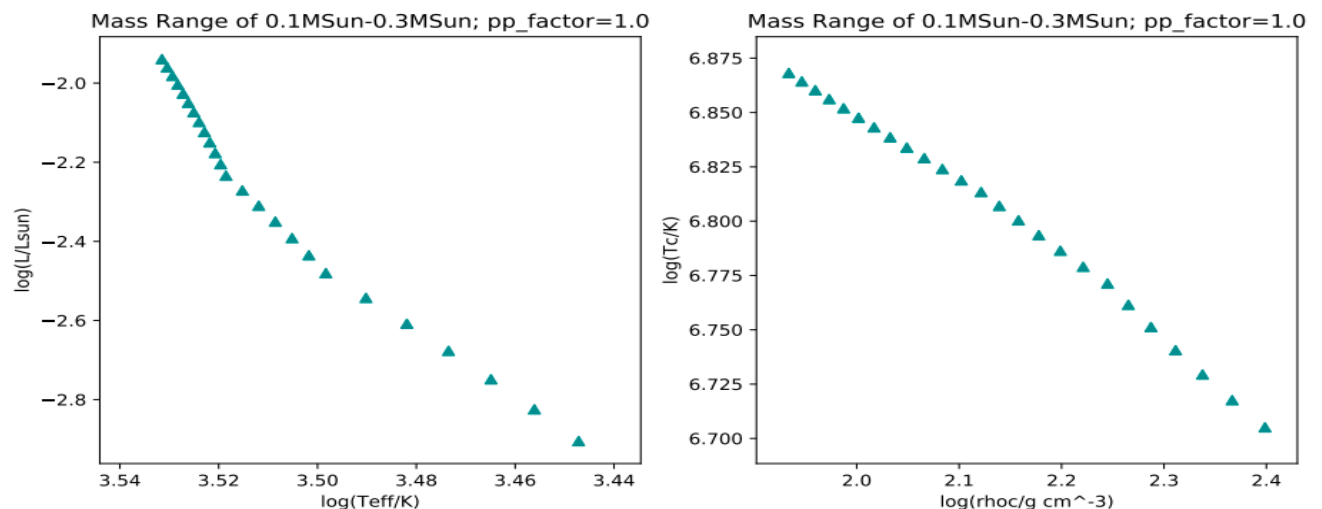


Figure 1a: Plots of $\log(L/L_{\text{sun}})$ against $\log(T_{\text{eff}}/K)$ and $\log(T_c/K)$ against $\log(\rho_c/g \text{ cm}^{-3})$.

Triangles represent 25 ZAMS stars between $0.1 M_{\text{sun}}$ and $0.3 M_{\text{sun}}$. The pp_factor used is 1.0

The above plots are for 10a and b. We can compare to these to figures **A** and **B** in this appendix section of this report. We can see with **A** that the $\log(L/L_{\text{sun}})$ against $\log(T_{\text{eff}}/K)$ closely resembles what is expected from the ZAMS stars in the mass range $0.1 M_{\text{sun}}$ to $0.3 M_{\text{sun}}$. At around (3.52, -2.25) on the above and **A**, we can see that gradient increases; also, the overall

path of the stars follows **A**, including specific points that the stars fall on. We can see with **B** that the $\log(T_c/K)$ against $\log(\rho/g \text{ cm}^{-3})$ closely resembles what is expected from the ZAMS stars in the mass range $0.1M_{\text{sun}}$ to $0.3M_{\text{sun}}$. For the above and **B**, we can see that at around (2.2, 6.75), the gradient decreases; the overall trend is very similar as well.

Step 11

For this part, we needed to plot $T(r)$, $T(m)$, $L(r)$, and $L(m)$ for a star with $M = 0.3M_{\text{sun}}$. For this part, we needed implement a different temperature equation as in part 10 as T is different than T_{eff} and of course T_c . In fact, the instruction manual provided the equation of $T = T_c(P/P_c)^{(1-1/\gamma)}$, where T_c and P_c are constants that depend on the star and γ is an absolute constant of $\gamma = 5/3$. This temperature is computed in the routine `get_rho_and_T`, however, we add it as a simple equation in our main execution file to compute T when mass is at $0.3 M_{\text{sun}}$.

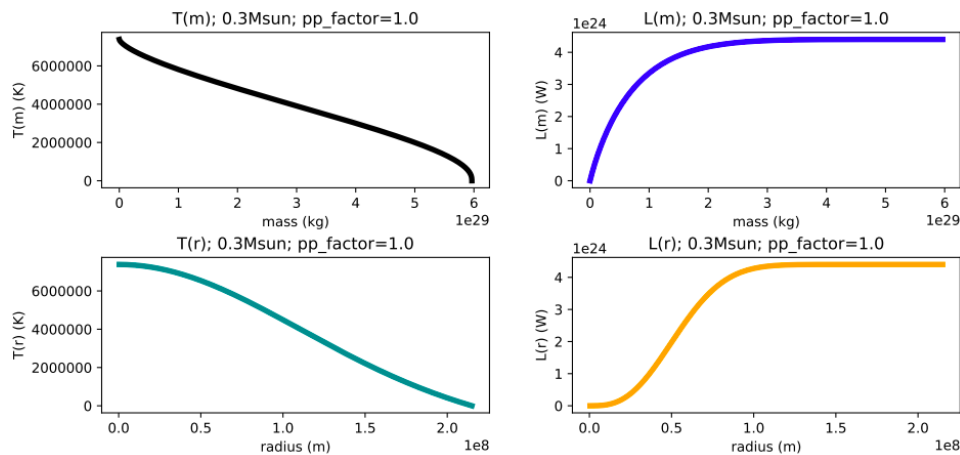


Figure 1b: Plots of T (K) against m (kg), T (K) against r (m), L (W) against m (kg), L (W) against r (m). Star used is of mass $0.3M_{\text{sun}}$. The pp_factor used is 1.0.

We can analyze each plot; while there are no comparison plots, we can reason why each behaves as they do as a good test of the content learned in the course. For $T(r)$, we can see that

that as the radius of the star increases, the temperature decreases; this is sensible considering that the central temperature will be much hotter than the surface as radius is smaller and density is larger and the kinetic energy of particles is greater which allows temperature to increase. For $T(m)$, we can see that as the mass of the star we integrate over increases, the temperature decreases. This is sensible considering that as we include more mass in the star, radius and mass increase. Since density is much more dependent on radius, it decreases as we collect more mass which makes temperature decrease as total mass increase. For $T(r)$, we can see a clear decrease in gradient around $0.5e8$ meters which means that location has reached large enough radius for the star to begin greatly declining in temperature. For $L(r)$, we can see that the luminosity greatly increases with radius and then reaches a constant value. From the surface luminosity equation it is sensible that luminosity increases with radius increasing, but the constant value it approaches could correspond to eventual surface value it settles upon. For $L(m)$, we can see that luminosity increases but with a declining gradient. The overall trend of increasing luminosity makes sense because as more mass is integrated over, so is more radius, where Luminosity greatly increases with. For $L(m)$ and $L(r)$, the luminosity seems to reach its constant level around the same location; this is sensible considering these plots are integrations over a star, so the luminosity should reach its constant value near the same integrated location; this is the case for the temperature as well.

We were also tasked to find what radius that $L(r)$ reaches 90% of its final value, as well as what fraction of mass is enclosed by this radius; these values can be seen as print statements from the main files. Including them here, as well, $L(r)$ reaches 90% of its final value at approximately a radius of $8.54e+7$ meters; the fraction of mass enclosed by this radius is approximately 0.275 of the stars mass. These values seem to correspond to the plot above and the radius size makes sense as it at least should be smaller than of course its own maximum radius, but also for a sanity check, a lot smaller than the sun's maximum radius; however, it should not be too small – say $8.54e+5$ meters.

Part 3)

For this part, we needed to suppose that the weak interactions were 10^5 which would increase E_{nucl} by the same factor. We needed to recompute the main-sequence stars with this increases reaction rate; we have included a `pp_factor` parameter in the necessary functions to make these easier. Additionally, instead a creating an unnecessarily long function in the main loop to change the `pp_factor`, we simply created separate main files to compute with the different `pp_factor`s.

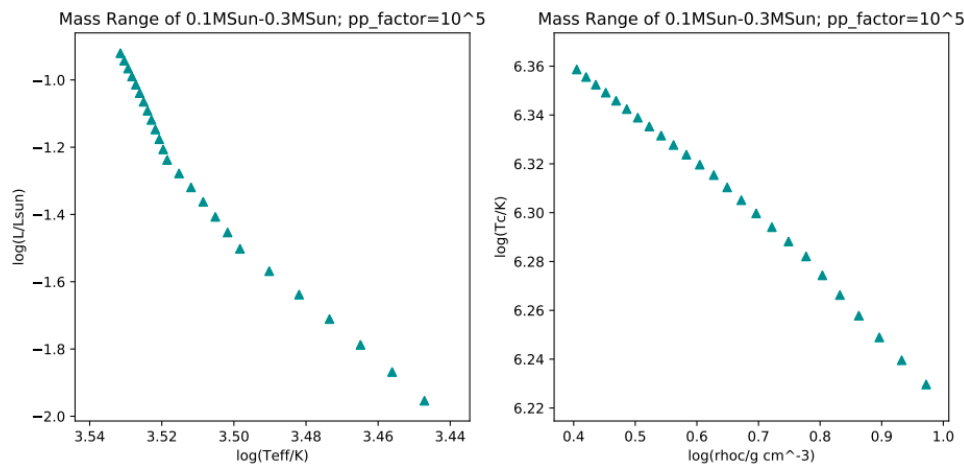


Figure 2a: Plots of $\log(L/L_{\text{sun}})$ against $\log(T_{\text{eff}}/K)$ and $\log(T_c/K)$ against $\log(\rho_{\text{hoc}}/g \text{ cm}^{-3})$. Triangles represent 25 ZAMS stars between 0.1Msun and 0.3Msun. The `pp_factor` used is 10^5

We can see that the stars vary in luminosity for $\log(L/L_{\text{sun}})$ against $\log(T_{\text{eff}}/K)$. This makes sense considering that T_{eff} does not depend on a change in `pp_factor`. Additionally, if the weak interaction is increased and therefore, E_{nucl} , this makes luminosity increase as the integrated luminosity depends on E_{nucl} ; if E_{nucl} increases, then luminosity will. For $\log(T_c/K)$ against $\log(\rho_{\text{hoc}}/g \text{ cm}^{-3})$, we can see that central temperature and density both decrease for all stars in the mass range when `pp_factor` is increased. Obviously, the weaker interaction doesn't make the star burn as warm centrally; if the temperature is not centrally as warm, then the central density will also decrease as the reaction slows.

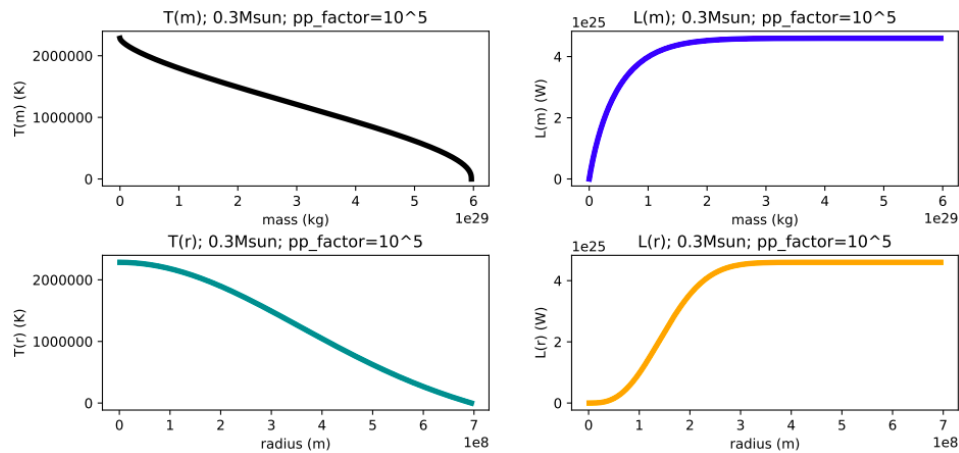


Figure 2b: Plots of T (K) against m (kg), T (K) against r (m), L (W) against m (kg), L (W) against r (m). Star used is of mass $0.3M_{\text{sun}}$. The pp_factor used is 10^5 .

Although not intended to do so in the instructions, we have included $T(r)$, $T(m)$, $L(r)$, and $L(m)$. We can see the major changes for a $0.3M_{\text{sun}}$ star includes a decrease as discussed previously in central temperature, and luminosity values increase by a factor of 10; the overall trends remain similar, as well. For this scenario, $L(r)$ at 90% corresponds to a radius $2.39e+8$ meters; this means that for a larger pp factor, the radius at which $L(r)$ reaches 90% of its final value greatly increases. However, because the final radius also greatly increases, this means that even less of the star's mass is enclosed by this radius; 0.194 of the star's mass is enclosed by this radius.

We can now speculate if the sun changed in a similar fashion to these lower-mass stars, how would surface temperatures on Earth change? We can first analyze the decrease in central temperature and temperature integrated over the star. In terms of analyzing surface temperatures on Earth, we want to consider luminosity of the sun, and although temperature varies luminosity greater than radius does, the change in temperature is not as great as the change in radius. Since the change in radius using this pp factor of the star is so greatly positive, this greatly increases the luminosity of the star. This makes sense when looking at the luminosity graphs and can see from the lot that luminosity increased by approximately a factor of ten for the main-sequence stars and the $0.3M_{\text{sun}}$ star. Since this is the case, if the sun

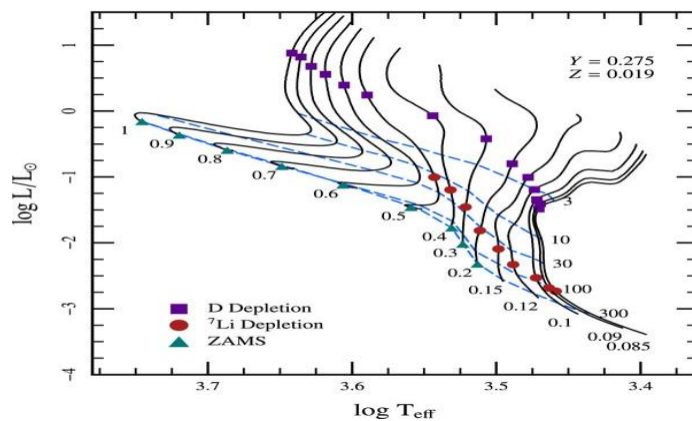
changed in a similar fashion with this increased factor, surface temperatures on Earth would greatly increase. This would likely cause life on Earth to not exist, as well as possibly further in planets be engulfed by the increased radius the sun would receive because of the increase in factor.

References

1. Bill Paxton, Lars Bildsten, Aaron Dotter, Falk Herwig, Pierre Lesaffre, and Frank Timmes. Modules for experiments in stellar astrophysics (MESA).
ApJS, 192:3, January 2011.

Appendix

A¹



B¹

