

Asteroid Astrometry from CCD Images

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Abstract

astrometry for asteroid name..etc

1. Introduction

(Some motivation for studying small bodies??)

2. Data Reduction

2.1. Systematic Corrections

The purpose of bias, dark, flat correction is to make the intensity of the image approximately linearly proportional to the number of astronomical photons arriving at the detector.

Bias subtraction. Bias frames are images taken with zero integration time and used to subtract the digital offset at zero level. Common practice include taking a scalar value representative of the bias frame dataset (i.e. mean or median) and subtracting that from the original image. For the simple averaging method, the more bias frames we have the better, since as the number of bias frames we have increase, the standard deviation of the mean falls off as $1/\sqrt{N}$, thereby decreasing the uncertainty on the obtained mean value. We applied the method of indiscriminate rejection described by Chromey (2010) which removes the largest pixel value of the bias frame, then take the median on the remaining values. This method

rejects the largest statistical outlier which is included when we simply take the mean for correction. Another method of bias correction is to set extra clock cycles so that the CCD reading continues for an additional 32 pixels after the physical image has been read out. Such use of overscan pixels can be advantageous over bias frames for bias correction in the case when the bias is changing over time since each image has its own bias correction taken immediately after the exposure. We corrected an image by taking the median of the 32 pixel. The median is chosen over the mean so that anomalies such as cosmic ray events or nearby radio activity are rejected. Even though we experimented with all the methods described above, since the goal of this lab is to conduct astrometric measurements using these images rather than detailed photometric study, we have simply chosen the more computationally trivial method of bias frame subtraction to apply to every image.

Dark and Flat Correction. Dark frames are long-exposure images taken with the shutters closed to calibrate for dark responses. We conduct dark correction by subtracting from a single dark frame with exposure time of 300 seconds. Since the CCD for Nickel Telescope

is cooled by the LN2 dewar to a relatively low temperature, the thermal production of electrons is much lower than in room temperature. Therefore dark frames will not significantly improve the data quality. We took the dark and flat frames for the I,R filters taken on October 16, 2013 and applied them to the other data on the other days since dark and flat frames were not taken for the observation at the other dates. Using the twilight sky as a source of uniform illumination, the flat field images calibrate the pixel-by-pixel variation of each pixel. The exposure time was short and adjusted for every frame to accommodate for the sky's changing brightness. Every pixel on the 1024-by-1024 image is processed by :

$$\frac{\text{image} - \text{dark}}{\text{flat}} \times \text{median}(\text{flat}) \quad (1)$$

As shown in the equation, we first dark-subtract from the image so that the intensity is linearly proportional to the number of photon counts then we divide by the flat in order to scale for the pixel-by-pixel variations. By plotting the pixel histogram of the image, we see that most of the pixel values are very small relative to our original image data at around 0.005 (@@@@UNITSSS). In order to preserve the comprehensible data ranges of the original intensity values after subtraction, the image is normalized by a scalar value that is representative of the magnitude of the flat field data. The median value is chosen because unlike the mean its value is not skewed due to the bad pixel and columns present on the Nickel Telescope.

2.2. Centroid-finding

In order to determine the positions of the stars on the image, we need to find the centroid position with the weight as the First, we set a threshold which was empirically adjusted to yield just enough stars for pattern recognition at around 5000 (UNITSS!!!!). This identifies Since bad columns, corners, and bor-

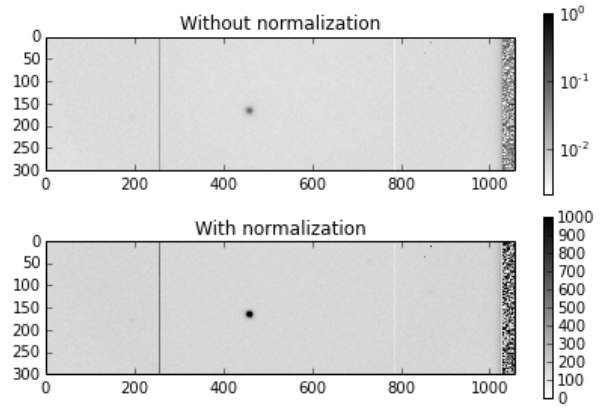


Figure 1: Without normalization, the image requires a logarithmic stretch to distinguish the asteroid and the background since it has to accommodate for the range of decimal values resulted from the division. With normalization, the image can be viewed on a linear scale with a lower intensity limit of 0 and upper limit of 1000.

ders often show up as completely dark pixels¹, it is mistaken as a whole row of stars, we have to manually mask the pixels by zeroing the intensity value at the selected pixel. Since this is a fairly trivial case where the dat-

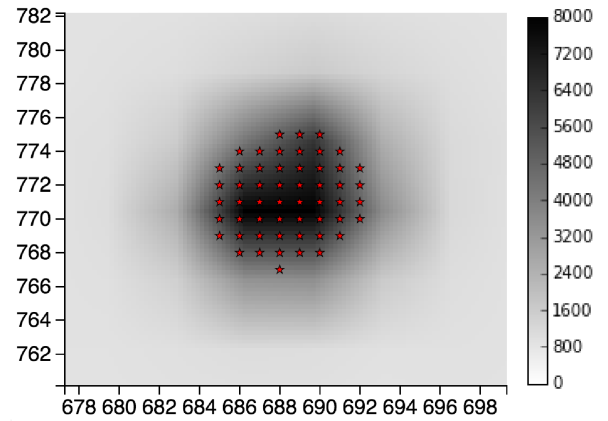


Figure 2: CENTROID MANY STUFF

apoints for each cluster is far apart, we use the more conceptually simplistic but computational intensive method of distance estimation instead of other clustering algorithm in

¹Since image color scheme is inverted, the darkest pixel mean highest intensity.

the literature. We loop through list of points above the threshold and for every p we compute its pixel distance from another point in the same list. If the separation distance between two datapoints is smaller than 20 pixel, then it is clustered into the same star. Each cluster represent a star. This threshold radius is selected r , most star have approx same aperture/ size because r . Finally, in order to get a representative value from cluster information, we calculate the centroid by Eq. 2 :

$$x = \frac{\sum_i x_i I_i}{\sum_i I_i} \quad (2)$$

Error on centroid.

2.3. Pixel to Standard Coordinate conversion

Since most stars used for (?) comparison are often very far away, it is convenient to treat them as projected position onto a sphere centered at the observer.

3. Astrometry

Using stellar positions to calibrate the image

3.0.1. Algorithm

3.1. Least Square

The purpose of using the least square method is to find a conversion matrix between the CCD-based pixel coordinate to standard coordinate. The standard coordinates can then be projected to celestial coordinates based on the Eq. (LABEL!!!): (TYPE EQUATION HERE)

3.1.1. USNO

intentionally adjust the Vmag lower cut than Vmag observed in NASA's small object dynamic so that it is able to identify the asteroid as a star.

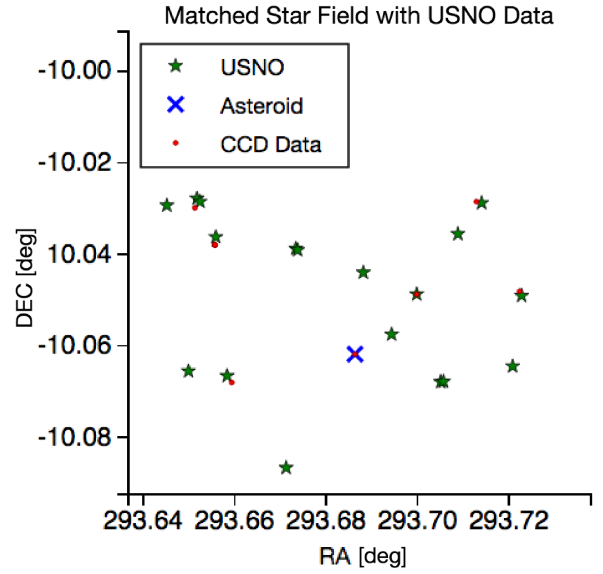


Figure 3: Matching starfield with USNO-B1 of stars with Vmag_i (—), October 16, 2013. The asteroid is easily identified as it shows up on the CCD image but not on the USNO-B1 catalog.

3.2. Ephemeris information

Position vector can be obtained from the HORIZONS Web-Interface heliocentric vector to Earth This could be verified with the vector from earth to asteroid. JPL horizon by default returns the ecliptic coordinate, Make sure to chose equatorial coordinate . This would affect the values for y and z components of the position (and velocity) vector returned since in the equatorial system the plane is around the equator

xy-plane: plane of the Earth's mean equator at the reference epoch x-axis : out along ascending node of instantaneous plane of the Earth's orbit and the Earth's mean equator at the reference epoch z-axis : along the Earth mean north pole at the reference epoch

3.3. Iterative solution

talk about ρ ...² From running this on (??)

²Since unit vectors are by normalized by its own magnitude, it is by definition dimensionless. Therefore, to construct physical units in the Earth-Asteroid vector, ρ must have units of AU.

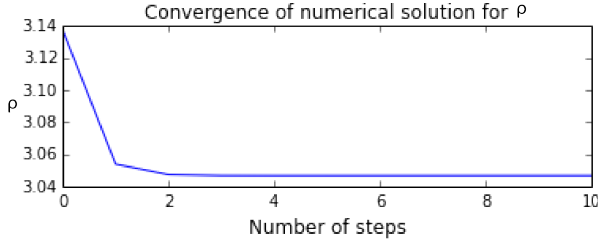


Figure 4: Convergence

1000 trials, the solution takes on average 5 steps to converge³. Each step approximately brings closer to the actual solution by an order of magnitude.

4. Error and Uncertainty

4.1. Centroid Error

In order to find the error on the value of ρ that scales the Earth-to-asteroid distance, we must first determine the error on the centroid values used to find ρ . By assuming that the error on each pixel is uncorrelated, we can apply the method of error propagation on Eq. 2 and obtain:

$$\sigma_{\langle x \rangle}^2 = \sum_i \left(\frac{\partial \langle x \rangle}{\partial I_i} \right)^2 \quad (3)$$

Computing $\frac{\partial \langle x \rangle}{\partial I_i}$ and plugging it back into Eq. 3:

$$\sigma_{\langle x \rangle}^2 = \frac{\sum_j I_j (x_j - \langle x \rangle)^2}{\left(\sum_i I_i \right)^2} \quad (4)$$

Using the list of clustered star position as described in Sec. 2.2. The same method is also used to calculate the error on $\langle y \rangle$. However this obtained results is in units of pixels, we must convert this error to units of degree as shown in Fig.5

³We define a solution to be converging when the difference between solutions at consecutive steps is less than 1.0×10^{-5}

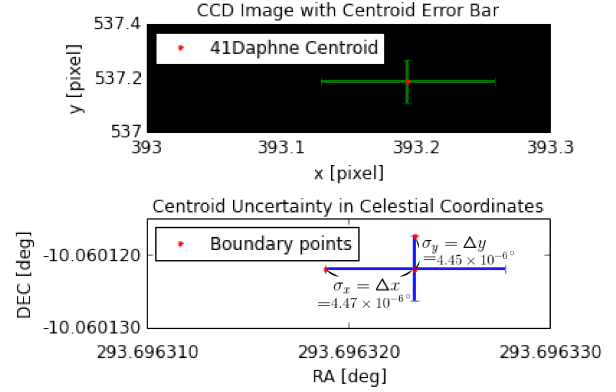


Figure 5: Since the pixel-to-celestial coordinate conversion procedure described in Sec. 3.1, works on coordinate points rather than value measurement (—). We select 2 boundary points that defines where the x and y error bar ends (larges — that substands — - maximum error —). Converting these two points and the asteroid's centroid location to celestial coordinates, then subtracting boundary points from the center position, we obtain the x and y error ranges in units of degrees. We find that the uncertainty on the centroid position is around 0.016 arcseconds.

4.2. Monte Carlo

Since we do not have an analytical solution to solve for ρ we can not compute the partial derivatives used in the error propagation equation. Instead we use the Monte Carlo method to approximate the error on our obtained value of ρ given the centroid error of the asteroid's celestial position. First, we construct two Gaussian noise distribution with standard deviation as the x and y centroid error that we obtained in Sec. 4.1. Then we draw a random value from each distribution to construct a noise vector for the x,y component with the noise on the z component set as zero.⁴ The noise vector is added to \vec{R} and this newly computed \vec{R} is used in the iterative algorithm to compute the solution for ρ as described in Sec.3.3 with other parameters kept the same. The Gaussian distribution is centered at zero

⁴We have no information about the error on z because it is not involved in the centroid-finding calculations. Since the CCD image is 2D, we can only compute the centroid error along x and y component.

so that the random noise may be positive or negative. This ensures that the computed ρ solutions are not skewed to a higher value by the addition of the noise vector. By performing the Monte Carlo procedure on 100000 samples, the standard deviation of the list of computed ρ values yields the approximate error on the value of ρ .

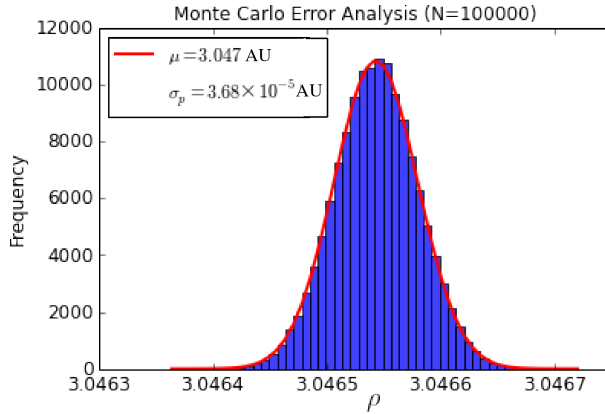


Figure 6: The standard deviation on this Gaussian distribution approximates the error on ρ to be 3.68×10^{-5} AU (≈ 5505 km).

4.3. Other possible error contribution

we only did monte carlo on x y error and not on the z component added noise vector For this excercise, we only examined errors on the asteroid. However, uncertainty in the measurement of stellar position can also affect our asteroid position result since stellar positions are taken into consideration in the least square procedure. This propagates to uncertainty on the plate constant values we obtain which is used in the calculation for conversion from α to δ coordinates. However this effect should be secondary to the ones we computed directly for the asteroid.

5. Conclusion

6. References

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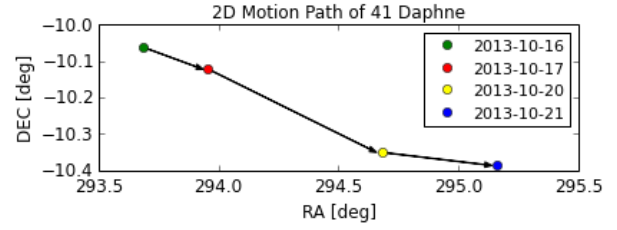


Figure 7: The minor deviation from the straight motion path could be explained by the asteroid's motion in the other directions. As shown in the JPL emphemeris velocity vectors, even though the V_y is much larger than V_x and V_z , V_x and V_z are still non-zero. Since the 3D orbital of the asteroid are suppressed onto the 2D celestial coordinates, — (?)

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