

1 Modal Expansion Fields With Separating Potential Fields

Any eigenmode basis for any waveguide structure is incomplete inside the region of any sources of EM excitation. Classic CMT ignores this fact. The expansion of the modal fields must be supplemented with orthogonal complementary fields E_b and H_b . The complementary fields are longitudinal fields, related to the longitudinal components of the bulk exciting currents J_{bz} and J_{bz} . [XREF: 3.4.24, 3.4.25]

```
(% i2) kill(all);
      load(vect);
      load(eigen);

done                                     (% o0)

                                           (% o1)

"/usr/share/maxima/5.45.1/share/vector/vect.mac"

                                           (% o2)

"/usr/share/maxima/5.45.1/share/matrix/eigen.mac"

(% i3) cmt_ch4_101:[E[a]=sum(A[k]*E[k],k,-N,N-1 ),H[a]=sum(A[k]*H[k],k,-N,N-1 )
];
```

$$\left[E_a = \sum_{k=-N}^{N-1} A_k E_k, H_a = \sum_{k=-N}^{N-1} A_k H_k \right] \quad (\% \text{ o3})$$

An unknown longitudinal dependence of $A_k(z)$ in the 2 equations below is due to the external sources and the orthogonal complementary fields E_b and H_b . The complete electromagnetic fields inside the source region have the following form: [XREF: 3.4.9, 3.4.10]

```
(% i4) cmt_ch4_421:[E(r[t],z)=E[a](r[t],z)+E[b](r[t],z), E(r[t],z)=sum(A[k](z)*E[k](r[t],z)+E[b](r[t],z),k,-
N,N-1) ];
```

$$\left[E(r_t, z) = E_b(r_t, z) + E_a(r_t, z), E(r_t, z) = \sum_{k=-N}^{N-1} E_k(r_t, z) A_k(z) + E_b(r_t, z) \right]$$

(% o4)

(% i5) cmt_ch4_422:[H(r[t],z)=H[a](r[t],z)+H[b](r[t],z), H(r[t],z)=sum(A[k](z)*H[k](r[t],z)+H[b](r[t],z),k,-N,N-1)];

$$\left[H(r_t, z) = H_b(r_t, z) + H_a(r_t, z), H(r_t, z) = \sum_{k=-N}^{N-1} H_k(r_t, z) A_k(z) + H_b(r_t, z) \right]$$

(% o5)

Total fields H and E can be represented as the sum of their curl (Ec,Hc) and potential (Ep,Hp) components, based on the Helmholtz decomposition theorem.

(% i6) cmt_ch4_422_A:[Ep=-grad(phi),Hp=-grad(psi)];

$$[Ep = -\text{grad}(\phi), Hp = -\text{grad}(\psi)]$$

(% o6)

Then the eigenfields of every kth mode can be represented as:

(% i7) cmt_ch4_423:[div(Ec[k])=0, div(Hc[k])=0,E[k]=Ec[k]+Ep[k],E[k]=Ec[k]-grad(phi[k]),H[k]=Hc[k]+Hp[k],H[k]=Hc[k]-grad(psi[k])];

$$[\text{div}(Ec_k) = 0, \text{div}(Hc_k) = 0, E_k = Ep_k + Ec_k, E_k = Ec_k - \text{grad}(\phi_k), H_k = Hp_k + Hc_k, H_k = Hc_k - \text{grad}(\psi_k)]$$

(% o7)

Now two sets of total eigenfields form the basis of a resulting Hilbert space, instead of just the (E[k],H[k]) eigenfields. The second set consists of the quasi-static eigenpotentials (psi[k] and phi[k]) and the curl eigenfields (Ec[k],Hc[k]). This increases the number of dimensions of the Hilbert space. This makes it possible to expand the complementary fields Eb and Hb in terms of the scalar basis potentials (psi[k], phi[k]). Apply the vector curl-field basis (Ec[k],Hc[k]) to expand the curl fields

(% i8) cmt_ch4_424:[Ec(r[t],z)=sum(A[k](z)*Ec[k](r[t],z),k,-N,N-1), Hc(r[t],z)=sum(A[k](z)*Hc[k](r[t],z),k,-N,N-1)];

$$\left[Ec(r_t, z) = \sum_{k=-N}^{N-1} Ec_k(r_t, z) A_k(z), Hc(r_t, z) = \sum_{k=-N}^{N-1} Hc_k(r_t, z) A_k(z) \right]$$

(% o8)

Apply the scalar-potential basis (phi(k),psi(k)) to expand the quasi-static potentials. Have to use fpsi and fphi because psi is part of the Gamma Functions package.

(% i9) cmt_ch4_425:[fphi(r[t],z)=sum(A[k](z)*fphi[k](r[t],z),k,-N,N-1),
fpsi(r[t],z)=sum(A[k](z)*fpsi[k](r[t],z),k,-N,N-1)];

$$\left[f\phi(r_t, z) = \sum_{k=-N}^{N-1} f\phi_k(r_t, z) A_k(z), f\psi(r_t, z) = \sum_{k=-N}^{N-1} f\psi_k(r_t, z) A_k(z) \right] \quad (\% \text{ o9})$$

In this case the expressions for complete fields inside the source region have the following form:

(% i10) cmt_ch4_426:[E=Ec-express(grad(phi)), E=sum(A[k]*(Ec[k]-grad(phi[k])) ,
k,-N,N-1) -e[z]*sum('diff(A[k],z)*phi[k],k,-N,N-1), E=sum(A[k]*E[k],k,-N,N-1)-
e[z]*sum('diff(A[k],z),k,-N,N-1)*psi[k]];

$$[E = \left[E_c - \frac{d}{dx} \phi, E_c - \frac{d}{dy} \phi, E_c - \frac{d}{dz} \phi \right], E = \left(\sum_{k=-N}^{N-1} A_k (E_{c_k} - \text{grad}(\phi_k)) \right) - \left(\sum_{k=-N}^{N-1} \phi_k \left(\frac{d}{dz} A_k \right) \right)] \quad (\% \text{ o10})$$

(% i11) cmt_ch4_427:[H=Hc-express(grad(psi)) ,H=sum(A[k]*(Hc[k]-grad(psi[k])) ,
k,-N,N-1) -e[z]*sum('diff(A[k],z)*psi[k],k,-N,N-1), H=sum(A[k]*H[k],k,-N,N-1)-
e[z]*sum('diff(A[k],z),k,-N,N-1)*psi[k]];

$$[H = \left[H_c - \frac{d}{dx} \psi, H_c - \frac{d}{dy} \psi, H_c - \frac{d}{dz} \psi \right], H = \left(\sum_{k=-N}^{N-1} A_k (H_{c_k} - \text{grad}(\psi_k)) \right) - \left(\sum_{k=-N}^{N-1} \psi_k \left(\frac{d}{dz} A_k \right) \right)] \quad (\% \text{ o11})$$

(% i12) cmt_ch4_428:[Eb=-e[z]*sum('diff(A[k],z)*phi[k],k,-N,N-1)];

$$[E_b = - \left(\sum_{k=-N}^{N-1} \phi_k \left(\frac{d}{dz} A_k \right) \right) e_z] \quad (\% \text{ o12})$$

(% i13) cmt_ch4_429:[Hb=-e[z]*sum('diff(A[k],z)*psi[k],k,-N,N-1)];

$$[H_b = - \left(\sum_{k=-N}^{N-1} \psi_k \left(\frac{d}{dz} A_k \right) \right) e_z] \quad (\% \text{ o13})$$

$A[k]$ is replaced by $\partial A[k]/\partial z$, which vanishes outside the source region. Outside the source region, $A[k]$ is constant, making its derivative zero. The Hilbert space spanned by the 2 sets of basis functions is closed w.r.t. any function that corresponds to any external source. The expressions for the curl fields and

the quasi-static potentials do not contain orthogonal complements. Then if the potential fields E_b and H_b are excluded, we can build an eigenmode basis that produces a modal expansion with no orthogonal complements. Use the scalar potentials instead of the potential field expressions, together with the curl fields E_c and H_c . The appropriate sets of mode quantities $\phi[k]$, $\psi[k]$ will then be a complete basis with no orthogonal complements. This does not apply to the polarization vector and magnetic field vector. Equations of motion and constitutive relations (tensors) contain the complete fields, not the curl and quasi-static parts. Polarization P and Magnetization M are derived from the tensors and equations of motion. Therefore P and M are required to have orthogonal complements P_b , M_b generated by E_b and H_b . $f\psi$ and $f\phi$ are maxima placeholders for ψ and ϕ . Total fields are given here:

```
(% i14) cmt_ch4_429_A:[E=Ec-grad(fphi),E=Ea+Eb, H=Hc-grad(fpsi), H=Ha+Hb,
P=Pa+Pb, M=Ma+Mb, Ecb=0, Hcb=0, fphi_b=0, fpsi_b=0];
```

```
[E = Ec-grad (fphi) , E = Eb+Ea , H = Hc-grad (fpsi) , H = Hb+Ha , P = Pb+Pa , M = Mb+Ma , Ecb = 0 ,
(% o14)
```

The Quasi-Static Approximation $\nabla \times E \approx 0$ and $E \approx E_p = -\nabla \phi$, but $H = H_c$ because $\nabla \cdot H = 0$ and $\nabla \times H \neq 0$. For this purpose I will introduce a new assertion using function notation. `approx(e1, e2, [tol])` means $e1$ is close to $e2$ but may vary from $e2$ by tol (tolerance). Maxima doesn't recognize the wxmaxima approximation symbol, so I defined something that works.

```
(% i15) approx(curl(E),0);
```

```
approx (curl(E) , 0) (% o15)
```

The Quasi-Magnetostatic Approximation Useful for MSW in magnetized Ferrites $\nabla \times H \approx 0$ and $H \approx H_p = -\nabla \psi$ but $E = E_c$ because $\nabla \cdot E = 0$ and $\nabla \times E \neq 0$. The consequence is that there are slow waves due to the existence of a specific potential field. The curl component (fast waves) is eliminated from the approximation expression. The basic task of Coupled Mode Theory is to find the modal excitation amplitudes $A_k(z)$ inside the region of influence of external sources. The external sources are assumed to be given when starting the first stage of the analysis.

4.3 Constitutive Relations and Dynamic Equations For Space-Dispersive Active Media

4.3.1 Piezoelastic Properties of a Medium

```
(% i16) cmt_ch04_431:[S_ij = (1/2) * ( 'diff(μ_i, r_j) + 'diff(μ_j, r_i) )];
```

$$S_{ij} = \frac{\frac{d}{dr_i} \mu_j + \frac{d}{dr_j} \mu_i}{2} \quad (\% \text{ o16})$$

```
(% i18) T_bar:[T_x, T_y, T_z];T_bar_Σ:[T_Σ_x,T_Σ_y,T_Σ_z];
```

```
[T_x, T_y, T_z] (% o17)
```

$$[T_{-x}, T_{-y}, T_{-z}] \quad (\% \text{ o18})$$

(% i19) cmt_ch04_432: [p_m * 'diff(U_i ,t) = 'diff(T_ij,r_j), p_m * 'diff(U,t) = express(div(T_bar))];

$$\left[\left(\frac{d}{dt} U_i \right) p_m = \frac{d}{dr_j} T_{ij}, \left(\frac{d}{dt} U \right) p_m = \frac{d}{dz} T_z + \frac{d}{dy} T_y + \frac{d}{dx} T_x \right] \quad (\% \text{ o19})$$

Values and Tensors Used For Piezoelectric Analysis
 Total Stress Tensor T_bar[Σ]
 = T_bar + T_bar_fr
 2nd rank stress tensor T_bar
 2nd rank susceptibility X_bar[S]
 2nd rank permittivity ε_bar[S]
 3rd rank piezoelectric stress e_bar_bar
 4th rank elastic stiffness c_bar_bar[E]
 internal friction stress T_bar[fr]
 viscosity tensor η_bar
 inverse relaxation time τ(-1)

(% i20) cmt_ch04_432_A: [T_bar_Σ, X_bar[s], ε_bar[s] = ε_0 * (I_bar + X_bar[s]), e_bar_bar, c_bar_bar[E], T_bar_fr] ;

$$[[T_{-x}, T_{-y}, T_{-z}], X_{bar_s}, \epsilon_{bar_s} = (X_{bar_s} + I_{bar}) \epsilon_0, e_{bar_bar}, c_{bar_bar_E}, T_{bar_fr}] \quad (\% \text{ o20})$$

(% i21) cmt_ch04_433: [P[k] = e[k,i,j] * S[i,j] + ε_0 * X[i,k,s] * E[i], P = e_bar_bar . S_bar + ε_0 * X_bar[s] . E];

$$[P_k = E_i X_{i,k,s} \epsilon_0 + S_{i,j} e_{k,i,j}, P = (X_{bar_s} . E) \epsilon_0 + e_{bar_bar} . S_{bar}] \quad (\% \text{ o21})$$

(% i22) cmt_ch04_434: [T[i,j] = c[i,j,k,l,E] * S[i,j] + ε[i,k] * E[i], T_bar = c_bar_bar[E] . S_bar + e_bar_bar . E];

$$[T_{i,j} = E_i \epsilon_{i,k} + S_{i,j} c_{i,j,k,l,E}, [T_x, T_y, T_z] = e_{bar_bar} . E + c_{bar_bar_E} . S_{bar}] \quad (\% \text{ o22})$$

(% i23) cmt_ch04_435: [D[k] = e[i,k,i,j] * S[i,j] + ε[i,k,s] * E[i], D = e_bar_bar . S_bar + ε_bar[s] . E];

$$[D_k = S_{i,j} e_{ik,i,j} + E_i \epsilon_{i,k,s}, D = \epsilon_{bar_s} . E + e_{bar_bar} . S_{bar}] \quad (\% \text{ o23})$$

(% i24) cmt_ch04_436: [T_fr[i,j] = η_[i,j,k,l] * 'diff(S[k,l], t), T_bar_fr = η_bar_bar . S_bar_dot];

$$\left[T_{fr\,i,j} = \eta_{-i,j,k,l} \left(\frac{d}{dt} S_{k,l} \right), T_{bar_fr} = \eta_{bar_bar} . S_{bar_dot} \right] \quad (\% \text{ o24})$$

(% i25) cmt_ch04_437:[F_fr[i] = - $\tau[i,j]^{-1} * \rho_m * U[i]$, F_fr = - $\tau^{-1} . \rho_m * U$];

$$\left[F_{fr i} = -\frac{U_i \rho_m}{\tau_{i,j}}, F_{fr} = -U \left(\frac{1}{\tau} . \rho_m \right) \right] \quad (\% \text{ o25})$$

Allowing for (4.3.6) and (4.3.7), re-write (4.3.1) and (4.3.2) as

(% i26) cmt_ch04_438:['diff(S[i,j], t) = (1/2) * ('diff(U[i],r[j]) + 'diff(U[j],r[i]))];

$$\left[\frac{d}{dt} S_{i,j} = \frac{\frac{d}{dr_i} U_j + \frac{d}{dr_j} U_i}{2} \right] \quad (\% \text{ o26})$$

(% i27) cmt_ch04_439:[$\rho_m * 'diff(U,t) = \text{express}(\text{div}(T_bar_ \Sigma)) + F_fr$];

$$\left[\left(\frac{d}{dt} U \right) \rho_m = \frac{d}{dz} T_{-z} + \frac{d}{dy} T_{-y} + \frac{d}{dx} T_{-x} + F_{fr} \right] \quad (\% \text{ o27})$$

For Pure Harmonic Processes

(% i28) cmt_ch04_4310:[U = U[1], U[1] = 'diff(u[1], t), $\rho_m = \rho_m[0] + \rho_m[1]$];

$$\left[U = U_1, U_1 = \frac{d}{dt} u_1, \rho_m = \rho_{m1} + \rho_{m0} \right] \quad (\% \text{ o28})$$

4.3.2 Ferrimagnetic Properties of a Medium
 External static magnetic field $H[0,e]$ Saturation Magnetization $M[0]$ Total magnetization vector M Effective Magnetic Field H_eff External DC Field H_0_e Maxwellian Field H Crystal anisotropy $H_c = -N_bar[c]$.
 MDemagnetizing Field $H[d] = N_bar[d]$. MExchange Field $H[ex] = \lambda[ex] *$
 2 MNet Tensor $N_bar = N_bar[c] + N_bar[d]$ Relaxation Term R

(% i30) H_eff:[Heff_x, Heff_y, Heff_z];M:[M_x, M_y, M_z];

$$[Heff_x, Heff_y, Heff_z] \quad (\% \text{ o29})$$

$$[M_x, M_y, M_z] \quad (\% \text{ o30})$$

(% i31) cmt_ch04_4311:[$\gamma = \text{abs}(e) / m[0]$, 'diff(M, t) = - $\gamma * \mu_0 * (\text{express}(M \sim H_eff)) + R$];

$$\left[\gamma = \frac{|e|}{m_0}, \frac{d}{dt} [M_x, M_y, M_z] = [R - (Heff_z M_y - Heff_y M_z) \gamma \mu_0, R - (Heff_x M_z - Heff_z M_x) \gamma \mu_0, R - (Heff_y M_x - \right. \quad (\% \text{ o31})$$

(% i32) cmt_ch04_4312:[H_eff = H[0,e] + H - N_bar . M + λ[ez] * express(laplacian(M))];

$$[\text{Heff}_x, \text{Heff}_y, \text{Heff}_z] = \left(\frac{d^2}{dz^2} [M_x, M_y, M_z] + \frac{d^2}{dy^2} [M_x, M_y, M_z] + \frac{d^2}{dx^2} [M_x, M_y, M_z] \right) \lambda_{ez} - N_{\text{bar}} \cdot [M_x, M_y, M_z] \quad (\% \text{ o32})$$

(% i33) cmt_ch04_4313:[α=ΔH/H[0],R = α*((M/M_0) + 'diff(M,t))];

$$\left[\alpha = \frac{H}{H_0}, R = \left(\frac{d}{dt} [M_x, M_y, M_z] + \left[\frac{M_x}{M_0}, \frac{M_y}{M_0}, \frac{M_z}{M_0} \right] \right) \alpha \right] \quad (\% \text{ o33})$$

(% i34) cmt_ch04_4314:[M = M_0 + M_1, abs(M_1) < abs(M_0)];

$$[M_x, M_y, M_z] = M_1 + M_0, |M_1| < |M_0| \quad (\% \text{ o34})$$

(% i35) cmt_ch04_4315:[H_0 = H[0,e] - N_bar . M_0, H_eff = H_0 + H_1 - N_bar . M_1 + λ[ez] * express(laplacian(M_1))];

$$\left[H_0 = H_{0,e} - N_{\text{bar}} \cdot M_0, [\text{Heff}_x, \text{Heff}_y, \text{Heff}_z] = \left(\frac{d^2}{dz^2} M_1 + \frac{d^2}{dy^2} M_1 + \frac{d^2}{dx^2} M_1 \right) \lambda_{ez} - N_{\text{bar}} \cdot M_1 + H_1 + H_0 \right] \quad (\% \text{ o35})$$

(4.3.3) Drifting Charge Carriers in a Medium (Plasmas)Hydrodynamic Force EquationEffective electron mass m_0 Free electron mass m_0 energy relaxation time $T_M = \epsilon/\sigma = \epsilon^*m/e^2n^*\tau$ Maxwellian relaxation time τ_e determines rate of electron perturbations T_M determines time scale of signal changes in the electric field and charge distribution $\tau_e \ll T_M$ This condition means the temperature keeps pace with signal changes in the electric field This provides a local relationship between T and E . This allows the momentum relaxation time τ to be considered a function of E E is found from measuring the field dependence of mobility $\mu(E) = (e/m)^*\tau(E)$ diffusion $D(E) = v_T^2 * \tau(E)$

(% i41) B:[B_x,B_y,B_z];v:[v_x, v_y, v_z];r_l:[r1_x, r1_y, r1_z]; r_0:[r0_x, r0_y, r0_z];
J_b1_e:[J_b1e_x, J_b1e_y, J_b1e_z]; J_b1_m:[J_b1m_x, J_b1m_y, J_b1m_z];

$$[B_x, B_y, B_z] \quad (\% \text{ o36})$$

$$[v_x, v_y, v_z] \quad (\% \text{ o37})$$

$$[r1_x, r1_y, r1_z] \quad (\% \text{ o38})$$

$$[r0_x, r0_y, r0_z] \quad (\% \text{ o39})$$

$$[J_b1e_x, J_b1e_y, J_b1e_z] \quad (\% \text{ o40})$$

$$[J_b1m_x, J_b1m_y, J_b1m_z] \quad (\% \text{ o41})$$

$$(\% \text{ i42}) \text{ cmt_ch04_4316:} [\text{'diff(v,t) + (v . express(div(v))) = (e/m) * (E + express(v} \\ \sim B) - \text{express(grad((n*k[B] * T)/(m*n)))- (v/\tau)}]);$$

$$[[v_x, v_y, v_z] \cdot \left(\frac{d}{dz} v_z + \frac{d}{dy} v_y + \frac{d}{dx} v_x \right) + \frac{d}{dt} [v_x, v_y, v_z] = [\frac{e \left(-B_y v_z + B_z v_y - \frac{d}{dx} \frac{k_{[B_x, B_y, B_z]} T}{m} + E \right)}{m} - \frac{v_x}{\tau}, \frac{e \left(\right)}{m} - \frac{v_y}{\tau}, \frac{e \left(\right)}{m} - \frac{v_z}{\tau}] \quad (\% \text{ o42})$$

$$(\% \text{ i43}) \text{ cmt_ch04_4316_A:} [\text{express(grad(p)) = (m*v_T^2)*express(grad(n))}];$$

$$\left[\left[\frac{d}{dx} p, \frac{d}{dy} p, \frac{d}{dz} p \right] = \left[m \left(\frac{d}{dx} n \right) v_T^2, m \left(\frac{d}{dy} n \right) v_T^2, m \left(\frac{d}{dz} n \right) v_T^2 \right] \right] \quad (\% \text{ o43})$$

$$(\% \text{ i44}) \mu(E) := (e/m) * \tau(E);$$

$$\mu(E) := \frac{e}{m} \tau(E) \quad (\% \text{ o44})$$

$$(\% \text{ i45}) D(E) := v_T^2 * \tau(E);$$

$$D(E) := v_T^2 \tau(E) \quad (\% \text{ o45})$$

For Plasmas in a magnetic field B the electron heating is the result of an electric field called the effective heating field (Appendix D.5) b takes into account an influence of the magnetic fields on the heating effect τ, μ , D now depend on E_h

(% i46) cmt_ch04_4317:[b = μ^*B , E_h = $\sqrt{E^2 + ((b_{\text{vec}} \cdot E)^2 / (1+b^2))}$] ;

$$\left[b = [B_x \mu, B_y \mu, B_z \mu], E_h = \sqrt{\frac{(b_{\text{vec}} \cdot E)^2}{b^2 + 1} + E^2} \right] \quad (\% \text{ o46})$$

Small Signal Analysis: All signal values with subscript 1 << those with subscript 0

(% i47) cmt_ch04_4317_B:[E=E_0 + E_1, B= B_0+B_1, E_h = E_h0 + E_h1,
 $\tau(E_h) = \tau(E_{h0}) + \text{'diff'}(\tau, E)$, E_h1= $\tau_0 + \tau_1$, $\tau_0 = \tau(E_{h0})$] ;

$$\left[E = E_1 + E_0, [B_x, B_y, B_z] = B_1 + B_0, E_h = E_{h1} + E_{h0}, \tau(E_h) = \frac{d}{dE} \tau + \tau(E_{h0}), E_{h1} = \tau_1 + \tau_0, \tau_0 = \tau(E_{h0}) \right] \quad (\% \text{ o47})$$

(% i48) cmt_ch04_4317_C:[(τ_1/τ_0) = $\text{'diff'}(\log(\tau), \log(E)) * (E_{h1}/E_{h0})$,
 $\tau_1/\tau_0 = (\kappa_0 - 1) * (E_{h1}/E_{h0})$] ;

$$\left[\tau - \frac{1}{\tau_0} = \frac{E_{h1} \left(\frac{d}{d \log(E)} \log(\tau) \right)}{E_{h0}}, \frac{\tau_1}{\tau_0} = \frac{E_{h1} (\kappa_0 - 1)}{E_{h0}} \right] \quad (\% \text{ o48})$$

(% i49) cmt_ch04_4318:[$\tau_1/\tau_0 = (\kappa_0 - 1) * (F_0/E_0) \cdot ((E_1 + v_0 \sim B_1)/E_0)$] ;

$$\left[\frac{\tau_1}{\tau_0} = \left(\frac{F_0}{E_0} \cdot \frac{E_1 - B_1 \sim v_0}{E_0} \right) (\kappa_0 - 1) \right] \quad (\% \text{ o49})$$

(% i50) cmt_ch04_4319:[b_0 = $\mu_e B_0$, F_0= $((1 + b_0^2) * (E_0 + (b_0 \cdot E_0) * b_0)) / ((1 + \kappa_0 * b_0^2) + ((1 + b_0^2) + (1 - \kappa_0)) * (b_0 \cdot E_0)^2 / E_0^2)$] ;

$$\left[b_0 = B_0 \mu_e, F_0 = \frac{(b_0^2 + 1) (b_0 (b_0 \cdot E_0) + E_0)}{b_0^2 \kappa_0 + \frac{(b_0 \cdot E_0)^2 (-\kappa_0 + b_0^2 + 2)}{E_0^2} + 1} \right] \quad (\% \text{ o50})$$

(% i51) cmt_ch04_4320:[E=E[k,0], $\mu_d = \text{'diff'}(\mu(E) * E, E)$, $\mu_e = \mu(E[h,0])$, $\mu_e = (e/m) * \tau(E[h,0])$, $\mu_e = (e/m) * \tau_0$] ;

$$\left[E = E_{k,0}, \mu_d = \frac{d}{dE} \frac{E \tau(E) e}{m}, \mu_e = \frac{e \tau(E_{h,0})}{m}, \mu_e = \frac{e \tau(E_{h,0})}{m}, \mu_e = \frac{e \tau_0}{m} \right] \quad (\% \text{ o51})$$

Electron displacement vector, a function of the unperturbed position vector \mathbf{r}_0 . The trajectory of liquid (virtual) particle motion is replaced by the E field description. Now deal with vector field of electron displacement $\mathbf{r}_1(\mathbf{r}_0, t)$. \mathbf{r}_1 is identical to the field of the lattice particle displacement $\mu(\mathbf{r}_0, t)$. $\mu(\mathbf{r}_0, t)$ is a term from elasticity theory

(% i52) cmt_ch04_4321:[$\mathbf{r}_1(\mathbf{r}_0, t) = \mathbf{r}(t) - \mathbf{r}_0(t)$];

$$[[r_{1x}, r_{1y}, r_{1z}][r_{0x}, r_{0y}, r_{0z}], t] = \mathbf{r}(t) - [r_{0x}, r_{0y}, r_{0z}](t) \quad (\% \text{ o52})$$

Total instantaneous velocity $\mathbf{v}(\mathbf{r}, t)$ of a group of charges satisfying (4.3.16) is:

(% i53) cmt_ch04_4322:[$\mathbf{v}(\mathbf{r}_{\text{vec}}, t) = \mathbf{v}_0(\mathbf{r}_{\text{vec}}) + \mathbf{u}_1(\mathbf{r}_{\text{vec}}, t)$, $\mathbf{v}(\mathbf{r}_{\text{vec}}, t) = \mathbf{v}_0(\mathbf{r}_0) + \mathbf{v}_1(\mathbf{r}_0, t)$];

$$[[v_x, v_y, v_z](r_{\text{vec}}, t) = u_1(r_{\text{vec}}, t) + v_0(r_{\text{vec}}), [v_x, v_y, v_z](r_{\text{vec}}, t) = v_1([r_{0x}, r_{0y}, r_{0z}], t) + v_0([r_{0x}, r_{0y}, r_{0z}]] \quad (\% \text{ o53})$$

Barybin uses $(\mathbf{r} \cdot \mathbf{v}_0)$ but maxima won't accept that operator

(% i54) cmt_ch04_4322_A:[$\mathbf{v}_1 = \mathbf{u}_1 + (\mathbf{r}_1 \cdot \text{express}(\text{grad}(\mathbf{v}_0)))$];

$$\left[v_1 = r_{1z} \left(\frac{d}{dz} v_0 \right) + r_{1y} \left(\frac{d}{dy} v_0 \right) + r_{1x} \left(\frac{d}{dx} v_0 \right) + u_1 \right] \quad (\% \text{ o54})$$

Polarization vector \mathbf{v}_1 obeys the equation of motion obtained from (4.3.22) as follows: Again Barybin uses $(\mathbf{r}_1 \cdot \mathbf{v}_0)$ for example, but maxima won't accept this directly. Will have to use apply/define/makefun/buildq and macros to make it work. Using \mathbf{E}_1 as a small signal field vector. Using $\mathbf{E}[1, p]$ as a symbolic placeholder for $\mathbf{E}[1, +]$ for entry into maxima

(% i55) cmt_ch04_4323:[$\tau_0 = \tau(\mathbf{E}_{h0})$, $\text{'diff}(\mathbf{v}_1, t) + (\mathbf{v}_0 \cdot \text{express}(\text{grad}(\mathbf{v}_1))) = (\mathbf{e}/m) * (\mathbf{E}_1 + \mathbf{r}_1 \cdot \text{express}(\text{grad}(\mathbf{E}_0))) + \mathbf{v}_1 \sim \mathbf{B}_0 + \mathbf{v}_0 \sim \mathbf{B}_1 + \mathbf{v}_0 \sim (\mathbf{r}_1 \cdot \text{express}(\text{grad}(\mathbf{B}_0))) + (\mathbf{v}_T^2 / \rho_0) * (\rho_0 * \text{express}(\text{grad}(\text{express}(\text{div}(\mathbf{r}_1)))) + \text{express}(\text{grad}(\mathbf{r}_1)) \cdot \text{express}(\text{grad}(\rho_0))) - (\mathbf{v}_1 / \tau_0) + (\mathbf{v}_0 / \tau_0) * ((\tau_1 + (\mathbf{r}_1 \cdot \text{express}(\text{grad}(\tau_0)))) / \tau_0)$];

$$[\tau_0 = \tau(E_{h0}), v_0 \cdot \left[\frac{d}{dx} v_1, \frac{d}{dy} v_1, \frac{d}{dz} v_1 \right] + \frac{d}{dt} v_1 = \frac{v_0 \left(\tau_1 + r_{1z} \left(\frac{d}{dz} \tau_0 \right) + r_{1y} \left(\frac{d}{dy} \tau_0 \right) + r_{1x} \left(\frac{d}{dx} \tau_0 \right) \right)}{\tau_0^2} - \frac{v_1}{\tau_0} + (v_T^2 (\quad (\% \text{ o55})$$

(% i56) cmt_ch04_4324:[$\mathbf{v}_1 = \text{'diff}(\mathbf{r}_1, t) + (\mathbf{v}_0 \cdot \text{express}(\text{grad}(\mathbf{r}_1)))$];

$$\left[v_1 = v_0 \cdot \left[\frac{d}{dx} [r_{1x}, r_{1y}, r_{1z}], \frac{d}{dy} [r_{1x}, r_{1y}, r_{1z}], \frac{d}{dz} [r_{1x}, r_{1y}, r_{1z}] \right] + \frac{d}{dt} [r_{1x}, r_{1y}, r_{1z}] \right]$$

(% o56)

(% i67) J_1: [J1_x, J1_y,J1_z];p_1:[p1_x, p1_y, p1_z];v_0:[v0_x,v0_y,v0_z];E_1:
 [E1_x, E1_y,E1_z];D_1: [D1_x, D1_y,D1_z];
 H_1: [H1_x, H1_y,H1_z]; B_1: [B1_x, B1_y,B1_z]; M_1: [M1_x,
 M1_y,M1_z]; P_1: [P1_x, P1_y,P1_z];
 H_2_p:[H2p_x, H2p_y, H2p_z];
 E_2:[E2_x, E2_y, E2_z];

$[J1_x, J1_y, J1_z]$ (% o57)

$[p1_x, p1_y, p1_z]$ (% o58)

$[v0_x, v0_y, v0_z]$ (% o59)

$[E1_x, E1_y, E1_z]$ (% o60)

$[D1_x, D1_y, D1_z]$ (% o61)

$[H1_x, H1_y, H1_z]$ (% o62)

$[B1_x, B1_y, B1_z]$ (% o63)

$[M1_x, M1_y, M1_z]$ (% o64)

$[P1_x, P1_y, P1_z]$ (% o65)

$$[H2p_x, H2p_y, H2p_z] \quad (\% \text{ o66})$$

$$[E2_x, E2_y, E2_z] \quad (\% \text{ o67})$$

$$(\% \text{ i68}) \text{ cmt_ch04_4325:}[\text{'diff}(\rho_1, t) + \text{express}(\text{div}(\mathbf{J_1}) = 0)];$$

$$\left[\frac{d}{dt} \rho_1 + \frac{d}{dz} J1_z + \frac{d}{dy} J1_y + \frac{d}{dx} J1_x = \frac{d}{dt} \rho_1 \right] \quad (\% \text{ o68})$$

$$(\% \text{ i69}) \text{ cmt_ch04_4326:}[\rho_1 = -\text{express}(\text{div}(\mathbf{p_1}))];$$

$$\left[\rho_1 = -\frac{d}{dz} p1_z - \frac{d}{dy} p1_y - \frac{d}{dx} p1_x \right] \quad (\% \text{ o69})$$

$$(\% \text{ i70}) \text{ cmt_ch04_4327:}[\mathbf{J_1} = \text{'diff}(\mathbf{p_1}, t) + \text{express}(\text{curl}(\mathbf{p_1} \sim \mathbf{v_0}))];$$

$$[[J1_x, J1_y, J1_z] = [\frac{d}{dy} (p1_x v0_y - p1_y v0_x) - \frac{d}{dz} (p1_z v0_x - p1_x v0_z), \frac{d}{dz} (p1_y v0_z - p1_z v0_y) - \frac{d}{dx} (p1_x v0_y - p1_y v0_x), \frac{d}{dx} (p1_z v0_x - p1_x v0_z) - \frac{d}{dy} (p1_y v0_x - p1_x v0_y)]] \quad (\% \text{ o70})$$

4.3.4 Electrodynamic Formulations For Active Polarized Media

$$(\% \text{ i71}) \text{ cmt_ch04_4328:}[\text{express}(\text{curl}(\mathbf{E_1})) = \text{'diff}(\mathbf{B_1}, t)];$$

$$\left[\left[\frac{d}{dy} E1_z - \frac{d}{dz} E1_y, \frac{d}{dz} E1_x - \frac{d}{dx} E1_z, \frac{d}{dx} E1_y - \frac{d}{dy} E1_x \right] = -\frac{d}{dt} [B1_x, B1_y, B1_z] \right] \quad (\% \text{ o71})$$

$$(\% \text{ i72}) \text{ cmt_ch04_4329:}[\text{express}(\text{curl}(\mathbf{H_1})) = \text{'diff}(\mathbf{D_1}, t) + \mathbf{J1}];$$

$$\left[\left[\frac{d}{dy} H1_z - \frac{d}{dz} H1_y, \frac{d}{dz} H1_x - \frac{d}{dx} H1_z, \frac{d}{dx} H1_y - \frac{d}{dy} H1_x \right] = \mathbf{J1} + \frac{d}{dt} [D1_x, D1_y, D1_z] \right] \quad (\% \text{ o72})$$

$$(\% \text{ i73}) \text{ cmt_ch04_4330:}[\text{express}(\text{div}(\mathbf{D_1})) = \rho_1];$$

$$\left[\frac{d}{dz} D1_z + \frac{d}{dy} D1_y + \frac{d}{dx} D1_x = \rho_1 \right] \quad (\% \text{ o73})$$

(% i74) cmt_ch04_4331:express(div(B_1)) = 0;

$$\frac{d}{dz}B1_z + \frac{d}{dy}B1_y + \frac{d}{dx}B1_x = 0 \quad (\% \text{ o74})$$

(% i75) cmt_ch04_4332:[D_1 = (ε_0 * E_1) + P_1];

$$[D1_x, D1_y, D1_z] = [E1_x \epsilon_0 + P1_x, E1_y \epsilon_0 + P1_y, E1_z \epsilon_0 + P1_z] \quad (\% \text{ o75})$$

(% i76) cmt_ch04_4333:[H_1 = ((1/μ_0) * B_1) - M_1];

$$\left[H1_x, H1_y, H1_z \right] = \left[\frac{B1_x}{\mu_0} - M1_x, \frac{B1_y}{\mu_0} - M1_y, \frac{B1_z}{\mu_0} - M1_z \right] \quad (\% \text{ o76})$$

(% i78) n[s,p]:[n_sp_x, n_sp_y, n_sp_z];n[s,m]:[n_sm_x, n_sm_y, n_sm_z];

$$[n_sp_x, n_sp_y, n_sp_z] \quad (\% \text{ o77})$$

$$[n_sm_x, n_sm_y, n_sm_z] \quad (\% \text{ o78})$$

The field vectors at the boundaries are also plus or minus (outward or inward)

(% i80) E_1_p: [E1_p_x, E1_p_y, E1_p_z];E_1_m: [E1_m_x, E1_m_y, E1_m_z];

$$[E1_p_x, E1_p_y, E1_p_z] \quad (\% \text{ o79})$$

$$[E1_m_x, E1_m_y, E1_m_z] \quad (\% \text{ o80})$$

(% i82) H_1_p: [H1_p_x, H1_p_y, H1_p_z];H_1_m: [H1_m_x, H1_m_y, H1_m_z];

$$[H1_p_x, H1_p_y, H1_p_z] \quad (\% \text{ o81})$$

$$[H1_m_x, H1_m_y, H1_m_z] \quad (\% \text{ o82})$$

(% i86) D_1_p: [D1_p_x, D1_p_y, D1_p_z];D_1_m: [D1_m_x, D1_m_y, D1_m_z];H_1P:[H_1P_x, H_1P_y, H_1P_z];D_1P:[D_1P_x, D_1P_y, D_1P_z];

$$[D1_p_x, D1_p_y, D1_p_z] \quad (\% \text{ o83})$$

$$[D1_m_x, D1_m_y, D1_m_z] \quad (\% \text{ o84})$$

$$[H_1P_x, H_1P_y, H_1P_z] \quad (\% \text{ o85})$$

$$[D_1P_x, D_1P_y, D_1P_z] \quad (\% \text{ o86})$$

$$(\% \text{ i93}) \text{ B_1_p: } [B1_p_x, B1_p_y, B1_p_z]; B_1_m: [B1_m_x, B1_m_y, B1_m_z]; M_net: [Mnet_x, Mnet_y, Mnet_z]; P_net: [Pnet_x, Pnet_y, Pnet_z]; v_o: [v0_x, v0_y, v0_z]; B_0: [B0_x, B0_y, B0_z]; r: [r_x, r_y, r_z];$$

$$[B1_p_x, B1_p_y, B1_p_z] \quad (\% \text{ o87})$$

$$[B1_m_x, B1_m_y, B1_m_z] \quad (\% \text{ o88})$$

$$[Mnet_x, Mnet_y, Mnet_z] \quad (\% \text{ o89})$$

$$[Pnet_x, Pnet_y, Pnet_z] \quad (\% \text{ o90})$$

$$[v0_x, v0_y, v0_z] \quad (\% \text{ o91})$$

$$[B0_x, B0_y, B0_z] \quad (\% \text{ o92})$$

$$[r_x, r_y, r_z] \quad (\% \text{ o93})$$

$$(\% \text{ i94}) \text{ cmt_ch04_4334: } [\text{express}(n[s,p] \sim E_1_p) + \text{express}(n[s,m] \sim E_1_m) = 0];$$

$$[[-E1_p_y n_sp_z + E1_p_z n_sp_y - E1_m_y n_sm_z + E1_m_z n_sm_y, E1_p_x n_sp_z - E1_p_z n_sp_x + E1_m_x n_sm_z -$$

(% o94)

(% i95) cmt_ch04_4335:[express(n[s,p] ~ H_1_p) + express(n[s,m] ~ H_1_m) =
J[s,eq]];

[[-H1_p_y n_sp_z + H1_p_z n_sp_y - H1_m_y n_sm_z + H1_m_z n_sm_y, H1_p_x n_sp_z - H1_p_z n_sp_x + H1_m_x n_sm_z
(% o95)

(% i96) cmt_ch04_4336:[express(n[s,p] ~ D_1_p) + express(n[s,m] ~ D_1_m)
=ρ[s,eq]];

[[-D1_p_y n_sp_z + D1_p_z n_sp_y - D1_m_y n_sm_z + D1_m_z n_sm_y, D1_p_x n_sp_z - D1_p_z n_sp_x + D1_m_x n_sm_z
(% o96)

(% i97) cmt_ch04_4337:[express(n[s,p] ~ B_1_p) + express(n[s,m] ~ B_1_m) =
0];

[[-B1_p_y n_sp_z + B1_p_z n_sp_y - B1_m_y n_sm_z + B1_m_z n_sm_y, B1_p_x n_sp_z - B1_p_z n_sp_x + B1_m_x n_sm_z
(% o97)

(% i98) cmt_ch04_4338:[p_eq[s] = p_0*(n . r_1) , p_eq[s]= n . p_1];

$[p_{eq_s} = (n. [r_{1x}, r_{1y}, r_{1z}]) p_0, p_{eq_s} = n. [p_{1x}, p_{1y}, p_{1z}]]$ (% o98)

(% i99) cmt_ch04_4339:[J_eq[s] = p_eq[s] * v_0 , J_eq[s] = (n . p-1) * v_0];

$[J_{eq_s} = [p_{eq_s} v_{0x}, p_{eq_s} v_{0y}, p_{eq_s} v_{0z}], J_{eq_s} = (n.p - 1) [v_{0x}, v_{0y}, v_{0z}]]$ (% o99)

(% i100) cmt_ch04_4340:[P_net = P_1 + p_1, M_net = M_1 + m_1, M__net =
approx(M_1 + (P_1 ~ v_0))];

[[Pnet_x, Pnet_y, Pnet_z] = [p1_x + P1_x, p1_y + P1_y, p1_z + P1_z] , [Mnet_x, Mnet_y, Mnet_z] = [m1 + M1_x, m1 + M1_y, m1 + M1_z]
(% o100)

(% i101) cmt_ch04_4341:ρ_P = -express(div(P_net)) , J+P = 'diff(P_net,t), J_M =
express(curl(M_net));

$\rho_P = -\frac{d}{dz}P_{net_z} - \frac{d}{dy}P_{net_y} - \frac{d}{dx}P_{net_x}$ (% o101)

Chu Formulation Of Maxwell's Equations

(% cmt_ch04_4342:[express(curl(E_1)) = 'diff(B_1,t)];
i102)

$$\left[\left[\frac{d}{dy} E_{1z} - \frac{d}{dz} E_{1y}, \frac{d}{dz} E_{1x} - \frac{d}{dx} E_{1z}, \frac{d}{dx} E_{1y} - \frac{d}{dy} E_{1x} \right] = \frac{d}{dt} [B_{1x}, B_{1y}, B_{1z}] \right] \quad (\% \text{ o102})$$

(% cmt_ch04_4343:[express(curl((1/μ_0) * B_1)) = ε_0 * 'diff(E_1,t) + (J_P
i103) + J_M)];

$$\left[\left[\frac{d}{dy} \frac{B_{1z}}{\mu_0} - \frac{d}{dz} \frac{B_{1y}}{\mu_0}, \frac{d}{dz} \frac{B_{1x}}{\mu_0} - \frac{d}{dx} \frac{B_{1z}}{\mu_0}, \frac{d}{dx} \frac{B_{1y}}{\mu_0} - \frac{d}{dy} \frac{B_{1x}}{\mu_0} \right] = \left(\frac{d}{dt} [E_{1x}, E_{1y}, E_{1z}] \right) \epsilon_0 + J_P + J_M \right] \quad (\% \text{ o103})$$

(% cmt_ch04_4344:[express(div(E_1)) = (1/ε_0) * ρ_P];
i104)

$$\left[\frac{d}{dz} E_{1z} + \frac{d}{dy} E_{1y} + \frac{d}{dx} E_{1x} = \frac{\rho_P}{\epsilon_0} \right] \quad (\% \text{ o104})$$

(% cmt_ch04_4345:[express(curl(B_1)) = 0];
i105)

$$\left[\left[\frac{d}{dy} B_{1z} - \frac{d}{dz} B_{1y}, \frac{d}{dz} B_{1x} - \frac{d}{dx} B_{1z}, \frac{d}{dx} B_{1y} - \frac{d}{dy} B_{1x} \right] = 0 \right] \quad (\% \text{ o105})$$

Minkowski Formulation Of Maxwell's Equations

(% cmt_ch04_4346:[D_1_P = ε_0 * E_1 + P_net];
i106)

$$[D_{1P} = [E_{1x}\epsilon_0 + P_{net_x}, E_{1y}\epsilon_0 + P_{net_y}, E_{1z}\epsilon_0 + P_{net_z}]] \quad (\% \text{ o106})$$

(% cmt_ch04_4347:[H_1_P = (1/μ_0) * B_1 - M_net];
i107)

$$[H_{1P} = \left[\frac{B_{1x}}{\mu_0} - M_{net_x}, \frac{B_{1y}}{\mu_0} - M_{net_y}, \frac{B_{1z}}{\mu_0} - M_{net_z} \right]] \quad (\% \text{ o107})$$

(% cmt_ch04_4348:[express(curl(E_1)) = 'diff(B_1,t)];
i108)

$$\left[\left[\frac{d}{dy} E_{1z} - \frac{d}{dz} E_{1y}, \frac{d}{dz} E_{1x} - \frac{d}{dx} E_{1z}, \frac{d}{dx} E_{1y} - \frac{d}{dy} E_{1x} \right] = \frac{d}{dt} [B_{1x}, B_{1y}, B_{1z}] \right]$$

(% o108)

(% cmt_ch04_4349:[express(curl(H_1P)) = 'diff(D_1_P,t)];
i109)

$$\left[\left[\frac{d}{dy} H_{1P_z} - \frac{d}{dz} H_{1P_y}, \frac{d}{dz} H_{1P_x} - \frac{d}{dx} H_{1P_z}, \frac{d}{dx} H_{1P_y} - \frac{d}{dy} H_{1P_x} \right] = \frac{d}{dt} D_{1P} \right]$$

(% o109)

(% cmt_ch04_4350:[express[div(D_1P)]= 0];
i110)

$$\left[\frac{d}{dz} D_{1P_z} + \frac{d}{dy} D_{1P_y} + \frac{d}{dx} D_{1P_x} = 0 \right]$$

(% o110)

(% cmt_ch04_4351:[express(div(B_1))=0];
i111)

$$\left[\frac{d}{dz} B_{1z} + \frac{d}{dy} B_{1y} + \frac{d}{dx} B_{1x} = 0 \right]$$

(% o111)

(% cmt_ch04_4352:[D_1P = D_1 + p_1, H_1P = H_1 + express(v_0 ~ p_1)];
i112)

$$[[D_{1P_x}, D_{1P_y}, D_{1P_z}] = [p_{1x} + D_{1x}, p_{1y} + D_{1y}, p_{1z} + D_{1z}], [H_{1P_x}, H_{1P_y}, H_{1P_z}] = [-p_{1y} v_{0z} + D_{1x}, -p_{1x} v_{0z} + D_{1y}, -p_{1x} v_{0y} + D_{1z}]]$$

(% o112)

Use express() here once the E and n vectors are defined properly

(% cmt_ch04_4353:[n[s,p] ~ E[1,p] + n[s,m] ~ E[1,m] = 0] ;
i113)

$$[-E_{1,p} \sim [n_{sp_x}, n_{sp_y}, n_{sp_z}] - E_{1,m} \sim [n_{sm_x}, n_{sm_y}, n_{sm_z}] = 0]$$

(% o113)

(% cmt_ch04_4354:[n[s,p] ~ H[1, p,p] + n[s,m] ~ H[1,p,m] = 0] ;
i114)

$$[-H_{1,p,p} \sim [n_{sp_x}, n_{sp_y}, n_{sp_z}] - H_{1,p,m} \sim [n_{sm_x}, n_{sm_y}, n_{sm_z}] = 0]$$

(% o114)

```
(%      cmt_ch04_4355:[n[s,p] ~ D[1,p,p] + n[s,m] ~ D[1,p,m] = 0] ;
i115)
```

$$[-D_{1,p,p} \sim [n_{sp_x}, n_{sp_y}, n_{sp_z}] - D_{1,p,m} \sim [n_{sm_x}, n_{sm_y}, n_{sm_z}] = 0]$$

(% o115)

Removing B to allow for use of symbolic placeholder

```
(%      kill(B);
i116)
```

done (% o116)

```
(%      cmt_ch04_4356:[n[s,p] ~ B[1,p,p] + n[s,m] ~ B[1,p,m] = 0] ;
i117)
```

$$[-B_{1,p,p} \sim [n_{sp_x}, n_{sp_y}, n_{sp_z}] - B_{1,p,m} \sim [n_{sm_x}, n_{sm_y}, n_{sm_z}] = 0]$$

(% o117)

```
(%      B:[B_x, B_y, B_z];
i118)
```

$$[B_x, B_y, B_z]$$

(% o118)

4.4 General Power-Energy Relations for Space-Dispersive Active Media4.4.1
Generalized Poynting's Theorem For SDAM

```
(%      E:[E_x,E_y, E_z];H:[H_x,H_y,H_z];D:[D_x,D_y, D_z];S:[X_x,S_y,
i123) S_z];p_1:[p1_x, p1_y, p1_z];
```

$$[E_x, E_y, E_z]$$

(% o119)

$$[H_x, H_y, H_z]$$

(% o120)

$$[D_x, D_y, D_z]$$

(% o121)

$$[X_x, S_y, S_z]$$

(% o122)

$$[p1_x, p1_y, p1_z] \quad (\% \text{ o123})$$

(% cmt_ch04_441:[express(curl(E) = - 'diff(B,t))];
i124)

$$\left[\left[\frac{d}{dy} E_z - \frac{d}{dz} E_y, \frac{d}{dz} E_x - \frac{d}{dx} E_z, \frac{d}{dx} E_y - \frac{d}{dy} E_x \right] = -\frac{d}{dt} [B_x, B_y, B_z] \right] \quad (\% \text{ o124})$$

(% cmt_ch04_442:[express(curl(H)) = 'diff(D,t) + J];
i125)

$$\left[\left[\frac{d}{dy} H_z - \frac{d}{dz} H_y, \frac{d}{dz} H_x - \frac{d}{dx} H_z, \frac{d}{dx} H_y - \frac{d}{dy} H_x \right] = J + \frac{d}{dt} [D_x, D_y, D_z] \right] \quad (\% \text{ o125})$$

Scalar Multiply 4.4.1 by H, 4.4.2 by -E and add results.This gives the instantaneous Poynting theorem

(% cmt_ch04_443:['diff(((E . D)/2) + ((H.B)/2),t) + express(div(E ~ H)) =
i126) -I_P - I_m - I_J];

$$\left[\frac{d}{dt} \left(\frac{B_z H_z + B_y H_y + B_x H_x}{2} + \frac{D_z E_z + D_y E_y + D_x E_x}{2} \right) + \frac{d}{dx} (E_y H_z - E_z H_y) + \frac{d}{dy} (E_z H_x - E_x H_z) + \frac{d}{dz} (E_x H_y - E_y H_x) \right] \quad (\% \text{ o126})$$

where

(% cmt_ch04_444:[(1/2) * ((E . 'diff(D,t)) - (D . 'diff(E,t))) = (1/2) * ((E .
i127) 'diff(P,t)) - (P . 'diff(E,t)))];

$$\left[\frac{[E_x, E_y, E_z] \cdot \frac{d}{dt} [D_x, D_y, D_z] - [D_x, D_y, D_z] \cdot \frac{d}{dt} [E_x, E_y, E_z]}{2} = \frac{[E_x, E_y, E_z] \cdot \frac{d}{dt} P - P \cdot \frac{d}{dt} [E_x, E_y, E_z]}{2} \right] \quad (\% \text{ o127})$$

(% cmt_ch04_445:[I_M = (1/2)*((H . 'diff(B,t)) - (B . 'diff(H,t))), I_M = (1/2) *
i128) (H . 'diff(μ_0*M,t)) - μ_0*M . 'diff(H,t))];

$$\left[I_M = \left(\frac{I_M}{2} = \frac{\frac{[H_x, H_y, H_z] \cdot \text{del}([M_x, M_y, M_z] \cdot t) \mu_0}{2} - ([M_x, M_y, M_z] \cdot \frac{d}{dt} [H_x, H_y, H_z]) \mu_0}{2} \right) \right]$$

(% o128)

(% cmt_ch04_446:[I_J = J . E];
i129)

$$[I_J = J. [E_x, E_y, E_z]] \quad (\% \text{ o129})$$

(% cmt_ch04_4477:[D = \epsilon_0 * E + P , B = \mu_0 * (H+M)];
i130)

$$[[D_x, D_y, D_z] = [E_x \epsilon_0 + P, E_y \epsilon_0 + P, E_z \epsilon_0 + P], [B_x, B_y, B_z] = [(M_x + H_x) \mu_0, (M_y + H_y) \mu_0, (M_z + H_z) \mu_0]] \quad (\% \text{ o130})$$

(% cmt_ch04_448:[w_em = (E . D) / 2 + (H . B) /2];
i131)

$$\left[w_{\text{em}} = \frac{B_z H_z + B_y H_y + B_x H_x}{2} + \frac{D_z E_z + D_y E_y + D_x E_x}{2} \right] \quad (\% \text{ o131})$$

(% cmt_ch04_449:[S_em = express(E ~ H)];
i132)

$$[S_{\text{em}} = [E_y H_z - E_z H_y, E_z H_x - E_x H_z, E_x H_y - E_y H_x]] \quad (\% \text{ o132})$$

(% cmt_ch04_4410:['diff(w,t) + express(div(S) + q = 0)];
i133)

$$\left[\frac{d}{dt} w + q + \frac{d}{dx} X_x + \frac{d}{dz} S_z + \frac{d}{dy} S_y = \frac{d}{dt} w \right] \quad (\% \text{ o133})$$

(% cmt_ch04_4411:[div(avg_t(S) + avg_t(q) = 0)];
i134)

$$[\text{div}(\text{avg}_t(q) + \text{avg}_t([X_x, S_y, S_z])) = 0] \quad (\% \text{ o134})$$

(% cmt_ch04_4412:[w_add = w_el + w_fm + w_pl];
i135)

$$[w_{\text{add}} = w_{\text{pl}} + w_{\text{fm}} + w_{\text{el}}] \quad (\% \text{ o135})$$

(% cmt_ch04_4413:[S_add = S_el + S_fm + S_pl];
i136)

$$[S_{\text{add}} = S_{\text{pl}} + S_{\text{fm}} + S_{\text{el}}] \quad (\% \text{ o136})$$

```
(%      cmt_ch04_4414:[q_add = q_el + q_fm + q_pl];
```

```
i137)
```

$$[q_{\text{add}} = q_{\text{pl}} + q_{\text{el}} + q_{\text{fm}}] \quad (\% \text{ o137})$$

4.4.1.1 Contribution From Piezoelectric Properties of a Medium

```
(%      cmt_ch04_4415:[I_P=(1/2) * (E . 'diff(P,t) - P . 'diff(E,t) ), I_P = 'diff(w_el,t)
```

```
i138)      + div(S_el) + q_el ];
```

$$\left[I_P = \frac{[E_x, E_y, E_z] \cdot \frac{d}{dt} P - P \cdot \frac{d}{dt} [E_x, E_y, E_z]}{2}, I_P = \text{del}(w_{\text{el}}.t) + q_{\text{el}} + \text{div}(S_{\text{el}}) \right] \quad (\% \text{ o138})$$

The author uses a colon : instead of dot between T_bar and S_bar. T_bar and S_bar are tensors in this context. The : operator is a double dot product given by $\sum_j \sum_i (a_i \cdot d_j) * (b_i \cdot c_j)$ or $\sum_j \sum_i (a_i \cdot c_j) * (b_i \cdot d_j)$. Load the following package to get mattrace (trace) of a matrix. dd(A,B) gives the proper result for a double dot product (:). TODO: RETROFIT CHAPTERS 2-3 WITH THIS DEFINITION

```
(%      load ("nchrpl");
```

```
i139)
```

(% o139)

```
"/usr/share/maxima/5.45.1/share/matrix/nchrpl.mac"
```

```
(%      dd(A,B):=mattrace(A * transpose(B));
```

```
i140)
```

$$\text{dd}(A, B) := \text{mattrace}(A \text{ transpose}(B)) \quad (\% \text{ o140})$$

```
(%      T_bar:matrix([T_11, T_12, T_13],[T_21,T_22, T_23],[T_31, T_32, T_33]);
```

```
i141)
```

$$\begin{pmatrix} T_{11} & T_{12} & T_{13} \\ T_{21} & T_{22} & T_{23} \\ T_{31} & T_{32} & T_{33} \end{pmatrix} \quad (\% \text{ o141})$$

```
(%      S_bar:matrix([S_11, S_12, S_13],[S_21,S_22, S_23],[S_31, S_32, S_33]);
```

```
i142)
```

$$\begin{pmatrix} S_{11} & S_{12} & S_{13} \\ S_{21} & S_{22} & S_{23} \\ S_{31} & S_{32} & S_{33} \end{pmatrix} \quad (\% \text{ o142})$$

(% $\eta_bar:matrix([\eta_11, \eta_12, \eta_13],[\eta_21,\eta_22, \eta_23],[\eta_31, \eta_32, \eta_33]);$
i143)

$$\begin{pmatrix} \eta_{11} & \eta_{12} & \eta_{13} \\ \eta_{21} & \eta_{22} & \eta_{23} \\ \eta_{31} & \eta_{32} & \eta_{33} \end{pmatrix} \quad (\% \text{ o143})$$

(% $S_bar_dot:matrix(['diff(S_11,t)', 'diff(S_12,t)', 'diff(S_13,t)'], ['diff(S_21,t)', 'diff(S_22,t)', 'diff(S_23,t)'], ['diff(S_31,t)', 'diff(S_32,t)', 'diff(S_33,t)']);$
i144)

$$\begin{pmatrix} \frac{d}{dt}S_{11} & \frac{d}{dt}S_{12} & \frac{d}{dt}S_{13} \\ \frac{d}{dt}S_{21} & \frac{d}{dt}S_{22} & \frac{d}{dt}S_{23} \\ \frac{d}{dt}S_{31} & \frac{d}{dt}S_{32} & \frac{d}{dt}S_{33} \end{pmatrix} \quad (\% \text{ o144})$$

(% $cmt_ch04_4416:[w_el = (1/2) * p_m * U^2 + (1/2) * dd(T_bar, S_bar)];$
i145)

$$\left[w_{el} = \frac{U^2 p_m}{2} + \frac{S_{33}T_{33} + S_{22}T_{22} + S_{11}T_{11}}{2} \right] \quad (\% \text{ o145})$$

Assuming dot over bar means ordinary derivative of matrix elements with
timeTODO: Check on matrix differentiation rules for this

(% $cmt_ch04_4417:[S_el = (1/2) * p_m * U^2 * Y - T_bar^2 . U];$
i146)

$$\left[S_{el} = \frac{U^2 Y p_m}{2} - \begin{pmatrix} T_{11}^2 & T_{12}^2 & T_{13}^2 \\ T_{21}^2 & T_{22}^2 & T_{23}^2 \\ T_{31}^2 & T_{32}^2 & T_{33}^2 \end{pmatrix} . U \right] \quad (\% \text{ o146})$$

(% $cmt_ch04_4418:[q_el = dd(dd(S_bar_dot, \eta_bar), S_bar_dot)];$
i147)

$$[q_{el} = \left(\frac{d}{dt}S_{33} \right) \left(\left(\frac{d}{dt}S_{33} \right) \eta_{33} + \left(\frac{d}{dt}S_{22} \right) \eta_{22} + \left(\frac{d}{dt}S_{11} \right) \eta_{11} \right) + \left(\frac{d}{dt}S_{22} \right) \left(\left(\frac{d}{dt}S_{33} \right) \eta_{33} + \left(\frac{d}{dt}S_{22} \right) \eta_{22} + \left(\frac{d}{dt}S_{11} \right) \eta_{11} \right) + \left(\frac{d}{dt}S_{11} \right) \left(\left(\frac{d}{dt}S_{33} \right) \eta_{33} + \left(\frac{d}{dt}S_{22} \right) \eta_{22} + \left(\frac{d}{dt}S_{11} \right) \eta_{11} \right) \quad (\% \text{ o147})$$

(% $cmt_ch04_4419:[w_el = (1/2) * p_m0 * U_1^2 + (1/2) * dd(T_bar[1], S_bar[1]);$
i148)

$$\left[w_{el} = \frac{U_1^2 p_{m0}}{2} + \frac{S_{11}T_{11}}{2} \right] \quad (\% \text{ o148})$$

(% $cmt_ch04_4420:[S_el = -T[1,\Sigma] . U_1];$
i149)

$$[S_{el} = -T_{1,\Sigma} . U_1] \quad (\% \text{ o149})$$

```
(% cmt_ch04_4421:[q_el = dd(dd(S_bar_dot[1], eta_bar), S_bar_dot) + p_m0 *
i150) U_1 . tau_bar[-1] . U_1];
```

$$[q_{el} = \left(\frac{d}{dt} S_{33}\right) \left(\left(\frac{d}{dt} S_{13}\right) \eta_{33} + \left(\frac{d}{dt} S_{12}\right) \eta_{22} + \left(\frac{d}{dt} S_{11}\right) \eta_{11} \right) + \left(\frac{d}{dt} S_{22}\right) \left(\left(\frac{d}{dt} S_{13}\right) \eta_{33} + \left(\frac{d}{dt} S_{12}\right) \eta_{22} + \left(\frac{d}{dt} S_{11}\right) \eta_{11} \right)$$

(% o150)

4.4.1.2 Contribution From Ferrimagnetic Properties of a Medium

```
(% S_fm:[S_fm_x, S_fm_y, S_fm_z];  
i151)
```

$$[S_fm_x, S_fm_y, S_fm_z] \quad (\% \text{ o151})$$

```
(% cmt_ch04_4422:[I_M = (1/2) * (H . 'diff(mu_0*M,t) - mu_0*M . 'ddiff(H,t)),
i152) I_M = 'diff(w_fm,t) + express(div(S_fm)) + q_fm];
```

$$[I_M = \frac{[H_x, H_y, H_z] \cdot \frac{d}{dt} [M_x \mu_0, M_y \mu_0, M_z \mu_0] - ([M_x, M_y, M_z] \cdot \text{ddiff}([H_x, H_y, H_z], t)) \mu_0}{2}, I_M = \frac{d}{dt} w_{\text{fm}} + q_{\text{fm}}]$$

```
(% cmt_ch04_4423:[w_fm = w_z + w_md + w_an+ w_ex];
i153)
```

$$[w_{\text{fm}} = w_z + w_{\text{md}} + w_{\text{ex}} + w_{\text{an}}] \quad (\% \text{ o153})$$

```
(% cmt_ch04_4424:[w_z=-μ_0*M . H_0[e]];
i154)
```

$$[w_z = -([M_x, M_y, M_z] \cdot H_{0e}) \mu_0] \quad (\% \text{ o154})$$

```
(% cmt_ch04_4435:[w_md=(-μ_0/2) * M . H];
i155)
```

$$\left[w_{\text{md}} = -\frac{(H_z M_z + H_y M_y + H_x M_x) \mu_0}{2} \right] \quad (\% \text{ o155})$$

(% cmt_ch04_4426:[w_an(-μ_0/2) * M . H_an, w_an = (μ_0/2)*M . N_bar .
i156) M];

$$\left[([M_x, M_y, M_z] \cdot H_{\text{an}}) w_{\text{an}} \left(-\frac{\mu_0}{2} \right), w_{\text{an}} = \frac{([M_x, M_y, M_z] \cdot N_{\text{bar}} \cdot [M_x, M_y, M_z]) \mu_0}{2} \right]$$

(% o156)

(% cmt_ch04_4427:[w_ez=(-μ_0/2) * M . H_ez, H_ex = λ_ex * ex-
i157) press(laplacian(M)),w_an = -λ_ez * (μ_0/2) * express(laplacian(M))];

$$[w_{ez} = -\frac{([M_x, M_y, M_z] \cdot H_{ez}) \mu_0}{2}, H_{ex} = \left(\frac{d^2}{dz^2} [M_x, M_y, M_z] + \frac{d^2}{dy^2} [M_x, M_y, M_z] + \frac{d^2}{dx^2} [M_x, M_y, M_z] \right) \lambda_{ez}]$$

(% o157)

(% cmt_ch04_4428:[S_fm= λ_ex * (μ_0/2) * (M . 'diff(express(grad(M)) ,t) -
i158) express(grad(M)) . 'diff(M,t))] ;

$$[[S_{fm_x}, S_{fm_y}, S_{fm_z}] = (([M_x, M_y, M_z] \cdot \frac{d}{dt} \left[\frac{d}{dx} [M_x, M_y, M_z], \frac{d}{dy} [M_x, M_y, M_z], \frac{d}{dz} [M_x, M_y, M_z] \right) - [M_x, M_y, M_z] \cdot \frac{d}{dt} [M_x, M_y, M_z])]$$

(% o158)

(% cmt_ch04_4429:[q_fm = α* (μ_0/w_M) * ('diff(M,t) . 'diff(M,t))];
i159)

$$\left[q_{fm} = \frac{\left(\frac{d}{dt} [M_x, M_y, M_z] \cdot \frac{d}{dt} [M_x, M_y, M_z] \right) \alpha \mu_0}{w_M} \right]$$

(% o159)

(% cmt_ch04_30:[w_z = μ_0*M . H_0, w_z=-μ_0*H_0*M_z, w_z=-
i160) μ_0*H_0*sqrt(M_0^2 - M_1^2), w_z=-μ_0*H_0*M_0 + (μ_0/2)
*(H_0/M_0) * (M_1 . M_2)];

$$[w_z = ([M_x, M_y, M_z] \cdot H_0) \mu_0, w_z = -H_0 M_z \mu_0, w_z = \left[-H_0 \sqrt{M_0^2 - M_{1x}^2} \mu_0, -H_0 \sqrt{M_0^2 - M_{1y}^2} \mu_0, -H_0 \sqrt{M_0^2 - M_{1z}^2} \mu_0 \right]]$$

(% o160)

(% cmt_ch04_4431:[w_fm = (μ_0/2) * (((H_0/M_0) * M_1 - H_1) . M_1 +
i161) M_1 . N_bar . M_1 - λ_ex * express(laplacian(M_1) . M_1))];

$$[w_{fm} = ((-\left(\left(\frac{d^2}{dz^2} [M_{1x}, M_{1y}, M_{1z}] + \frac{d^2}{dy^2} [M_{1x}, M_{1y}, M_{1z}] + \frac{d^2}{dx^2} [M_{1x}, M_{1y}, M_{1z}] \right) \cdot [M_{1x}, M_{1y}, M_{1z}] \right) \lambda_{ex})]$$

(% o161)

(% cmt_ch04_4432:[S_fm = λ_ex * (μ_0/2) * (M_1 . diff(express(grad(M1)),t)
i162) - express(grad(M_1)) . diff(M_1,t))];

$$[[S_{fm_x}, S_{fm_y}, S_{fm_z}] = 0]$$

(% o162)

(% cmt_ch04_4433:[q_fm = $\alpha^*(\mu_0/w_M) * ('diff(M_1,t) . 'diff(M_1,t))$];
i163)

$$\left[q_{fm} = \frac{\left(\frac{d}{dt} [M1_x, M1_y, M1_z] \cdot \frac{d}{dt} [M1_x, M1_y, M1_z] \right) \alpha \mu_0}{w_M} \right] \quad (\% \text{ o163})$$

4.4.1.3 Contribution Of Drifting Charge Carriers in a Medium

(% cmt_ch04_4434:[I_J = J_1 . E_1 , I_J = (1/2) * 'diff((E_1 . p_1 - B_1 .
i164) express(p_1 ~ v_0)) - express(div(express(E_1 ~ express(p_1 ~ v_0))))
,t)] ;

$$[I_J = E1_z J1_z + E1_y J1_y + E1_x J1_x , I_J = \text{diff} \setminus \left(\left(-\frac{d}{dy} (E1_z (p1_y v0_z - p1_z v0_y) - E1_x (p1_x v0_y - p1_y v0_x)) - \frac{d}{dz} (E1_z (p1_x v0_z - p1_x v0_y) - E1_y (p1_y v0_z - p1_z v0_x)) \right) \right)] ; \quad (\% \text{ o164})$$

(% cmt_ch04_4435: ['diff(((E_1 . D_1) / 2) + ((H_1 . B_1) / 2),t), 'diff(((E_1
i165) . D_1_p)/2) + ((H_1_p . B_1) / 2) , t)];

$$\left[\frac{d}{dt} \left(\frac{B1_z H1_z + B1_y H1_y + B1_x H1_x}{2} + \frac{D1_z E1_z + D1_y E1_y + D1_x E1_x}{2} \right), \frac{d}{dt} \left(\frac{B1_z H1_p + B1_y H1_p + B1_x H1_p}{2} \right) \right] ; \quad (\% \text{ o165})$$

(% cmt_ch04_4436:[express(div(express(E_1 ~ H_1))), express(div(express(E_1
i166) ~ H_1))] ;

$$\left[\frac{d}{dx} (E1_y H1_z - E1_z H1_y) + \frac{d}{dy} (E1_z H1_x - E1_x H1_z) + \frac{d}{dz} (E1_x H1_y - E1_y H1_x), \frac{d}{dx} (E1_y H1_z - E1_z H1_y) + \frac{d}{dy} (E1_z H1_x - E1_x H1_z) + \frac{d}{dz} (E1_x H1_y - E1_y H1_x) \right] ; \quad (\% \text{ o166})$$

(% cmt_ch04_4437:[w_pl = W_ek, w_pl = (m/(2*c)) * (v_1_p . 'diff(p_1,t) -
i167) p_1 . 'diff(v_1_p,t))] ;

$$\left[w_{pl} = W_{ek}, w_{pl} = \frac{m (v_{1p} \cdot \frac{d}{dt} [p1_x, p1_y, p1_z] - [p1_x, p1_y, p1_z] \cdot \frac{d}{dt} v_{1p})}{2c} \right] \quad (\% \text{ o167})$$

(% cmt_ch04_4438:[S_pl = S_ek + Sth, S_pl = (m/(2*e)) * v_0 * (v_1_p .
i168) 'diff(p_1,t) - p_1 . 'diff(v_1_p,t)) + (m/(2*e)) * (v_1^2/p_0) * p_1 * 'diff(
express(div(p_1)),t) - express(div(p_1))*diff(p_1,t)] ;

$$[S_{pl} = S_{th} + S_{ek}, S_{pl} = \left[\frac{mv0_x}{2e}, \frac{mv0_y}{2e}, \frac{mv0_z}{2e} \right] \left(v_{1p} \cdot \frac{d}{dt} [p1_x, p1_y, p1_z] - [p1_x, p1_y, p1_z] \cdot \frac{d}{dt} v_{1p} \right) + \left(\frac{d}{dt} [p1_x, p1_y, p1_z] \cdot v_{1p} \right)] ;$$

(% o168)

(% cmt_ch04_4339:[q_pl = (1/tau_0) * (m/(2*e)) * ((v_1 . 'diff(p_1,t)) - p_1 .
i169) 'diff(v_1,t)) - (f_1 *'diff(p_1,t) - p_1 * 'diff(f_1,t)) . v_0, f_1 = (tau_1 +
r_1 . express(grad(tau_0))) /tau_0];

$$[q_{pl} = \frac{m \left(v_1 \cdot \frac{d}{dt} [p_{1x}, p_{1y}, p_{1z}] - [p_{1x}, p_{1y}, p_{1z}] \cdot \frac{d}{dt} v_1 \right) - \left(f_1 \left(\frac{d}{dt} [p_{1x}, p_{1y}, p_{1z}] \right) \right) + \left[- \left(\frac{d}{dt} f_1 \right) p_{1x}, - \left(\frac{d}{dt} f_1 \right) p_{1y}, - \left(\frac{d}{dt} f_1 \right) p_{1z} \right]}{2e\tau_0}$$

(% o169)

(% cmt_ch04_4440:[w_L = -(e/(2*m)) * B_0, v_1_p = v_1 - (e/(2*m)) * ex-
i170) press(r ~ B_0) , v_1_p = v_1 + express(r_1 ~ ((e/(2*m)) * B_0))];

$$[w_L = \left[-\frac{B_0 x e}{2m}, -\frac{B_0 y e}{2m}, -\frac{B_0 z e}{2m} \right], v_{1p} = \left[v_1 - \frac{e (B_0 z r_y - B_0 y r_z)}{2m}, v_1 - \frac{e (B_0 x r_z - B_0 z r_x)}{2m}, v_1 - \frac{e (B_0 y r_x - B_0 x r_y)}{2m} \right]$$

(% o170)

4.4.1.4 Small-signal power theorem for generalized space-dispersive active media

(% cmt_ch04_4441:['diff(w_em + w_el + efm + w_pl,t) + express(div(S_em
i171) + S_el + S_fm + S_pl)) + (q_el + q_fm + q_pl) = 0] ;

$$\left[\frac{d}{dt} (w_{pl} + w_{em} + w_{el} + efm) + q_{pl} + q_{fm} + q_{el} + \frac{d}{dz} (S_{pl} + S_{fmz} + S_{em} + S_{el}) + \frac{d}{dy} (S_{pl} + S_{fmy} + S_{em} + S_{el}) + \frac{d}{dx} (S_{pl} + S_{fx} + S_{em} + S_{el}) \right]$$

(% o171)

(% cmt_ch04_4442:[w_em = ((E_1 . D_1_p) / 2) + ((H_1_p . B_1) / 2)];
i172)

$$\left[w_{em} = \frac{B_{1z} H_{1p_z} + B_{1y} H_{1p_y} + B_{1x} H_{1p_x}}{2} + \frac{D_{1p_z} E_{1z} + D_{1p_y} E_{1y} + D_{1p_x} E_{1x}}{2} \right]$$

(% o172)

(% cmt_ch04_4443:[S_em = express(E ~ H_1_p)];
i173)

$$[S_{em} = [E_y H_{1p_z} - E_z H_{1p_y}, E_z H_{1p_x} - E_x H_{1p_z}, E_x H_{1p_y} - E_y H_{1p_x}]]$$

(% o173)

Left out express(div()) here because avg_t is not yet defined avg_t(F, n) will be a time-averaged value of F over n time intervals)

(% cmt_ch04_4444:[div(avg_t(S_em) + avg_t(S_el) + avg_t(S_fm) +
i174) avg_t(S_pl)) + (avg_t(q_el) + avg_t(q_fm) + avg_t(q_pl))) = 0];

$$[\text{div}(\text{avg}_t(q_{pl}) + \text{avg}_t(q_{fm}) + \text{avg}_t(q_{el}) + \text{avg}_t(S_{pl}) + \text{avg}_t([S_{fm_x}, S_{fm_y}, S_{fm_z}])) + \text{avg}_t(S_{em}) + \text{avg}_t(S_{el})] = 0 \quad (\% \text{ o174})$$

Elastic Properties of a Medium
Elastic properties of a medium contribute these terms to the time-averaged power relation (4.4.4.4)
The average elastic energy flux density

(% cmt_ch04_4445:[avg_t(S_el) = (1/2) * Re(-E_1_\Sigma . conjugate(U_1)),
i175) avg_t(S_el) = (1/2) * Re(V_1_a . conjugate(J_1_a))];

$$\left[\text{avg}_t(S_{el}) = \frac{\text{Re}(-E_1_\Sigma . U_1)}{2}, \text{avg}_t(S_{el}) = \frac{\text{Re}(V_{1a} . J_{1a})}{2} \right] \quad (\% \text{ o175})$$

The average elastic power loss density
Author uses : for double-dot or tensor (term-by-term) multiplication
NOTE: For now I use dot product but this needs to be addressed properly
Re can be defined to call realpart()

(% Re(ev):="realpart(ev);
i176)

$$\text{Re}(ev) := \text{realpart}(ev) \quad (\% \text{ o176})$$

(% cmt_ch04_4446:[
i177) avg_t(q_el,n) = (1/2) * Re(conjugate(S_1_dot_bar) . \eta_bar . S_1_dot_bar
+ p_m0 * conjugate(U1) . \tau^(-1) . U_1),
avg_t(q_el,n) = (\omega^2/2) * conjugate(S_1) . \eta_bar . S_1 + (p_m0/2) * conjugate(U_1) . \tau^(-1) . U_1
];

$$[\text{avg}_t(q_{el}, n) = \frac{\text{realpart}(U_1 . \frac{1}{\tau} . U_1) p_{m0} + \text{realpart}\left(S_{1_dot_bar} . \begin{pmatrix} \eta_{11} & \eta_{12} & \eta_{13} \\ \eta_{21} & \eta_{22} & \eta_{23} \\ \eta_{31} & \eta_{32} & \eta_{33} \end{pmatrix} . S_{1_dot_bar}\right)}{2}, \text{avg}_t(q_{el}, n) = \frac{\omega^2}{2} \text{realpart}(S_1 . \eta_{bar} . S_1) + \frac{p_{m0}}{2} \text{realpart}(U_1 . \tau^{-1} . U_1)] \quad (\% \text{ o177})$$

(% cmt_ch04_4447:[J_1_a = %i*\omega*u_1, J_1_a = U_1];
i178)

$$[J_{1a} = i\omega u_1, J_{1a} = U_1] \quad (\% \text{ o178})$$

(% cmt_ch04_4448:[V_1_a = -T_1_\Sigma];
i179)

$$[V_{1a} = -T_{1_}] \quad (\% \text{ o179})$$

(% cmt_ch04_4449:[avg_t(S_fm,n) = (1/2) * Re(%i*ω*μ_0*λ_ex * ex-
i180) press(grad(M_1)) . conjugate(M_1)) , avg_t(S_fm,n) = (1/2) * Re(V_1_m
 . conjugate(J_1_m))];

$$[\text{avg}_t([S_{\text{fm}_x}, S_{\text{fm}_y}, S_{\text{fm}_z}], n) = \frac{\left(-\left(\frac{d}{dz}[0, 0, 0]\right)M1_z - \left(\frac{d}{dy}[0, 0, 0]\right)M1_y - \left(\frac{d}{dx}[0, 0, 0]\right)M1_x\right)\lambda_{\text{ex}}\mu_0}{2} \quad (\% \text{ o180})$$

(% cmt_ch04_4450:[avg_t(q_fm,n) = (1/2) * Re(α * (μ_0/ω_M) * ('diff(M_1,t)
i181) . 'diff(conjugate(M_1),t))) , avg_t(q_fm,n) = ν_M * (μ_0/2) * (ω/ω_M)^2 *
abs(M_1)^2];

$$[\text{avg}_t(q_{\text{fm}}, n) = \frac{\text{realpart}\left(\frac{d}{dt}[M1_x, M1_y, M1_z] \cdot \frac{d}{dt}[M1_x, M1_y, M1_z]\right)\alpha\mu_0}{2\omega_M}, \text{avg}_t(q_{\text{fm}}, n) = \left[\frac{M1_x^2\mu_0\nu_M\omega^2}{2\omega_M^2}, \frac{M1_y^2\mu_0\nu_M\omega^2}{2\omega_M^2}, \frac{M1_z^2\mu_0\nu_M\omega^2}{2\omega_M^2}\right] \quad (\% \text{ o181})$$

(% cmt_ch04_4451:[J_1_m = %i*ω*μ_00*M_1];
i182)

$$[J_{1_m} = [\%iM1_x\mu_{00}\omega, \%iM1_y\mu_{00}\omega, \%iM1_z\mu_{00}\omega]] \quad (\% \text{ o182})$$

(% cmt_ch04_4452:[V_1_m = -λ_ex * express(grad(M_1))];
i183)

$$[V_{1_m} = \left[-\left(\frac{d}{dx}[M1_x, M1_y, M1_z]\right)\lambda_{\text{ex}}, -\left(\frac{d}{dy}[M1_x, M1_y, M1_z]\right)\lambda_{\text{ex}}, -\left(\frac{d}{dz}[M1_x, M1_y, M1_z]\right)\lambda_{\text{ex}}\right] \quad (\% \text{ o183})$$

Plasma Properties of a MediumPlasma properties contribute these terms to the
time-average power relation (4.4.44)(i) The average plasma energy flow density

(% cmt_ch04_4453:[avg_t(S_pl,n) = avg_t(S_ek,n) + avg_t(S_th,n),
i184) avg_t(S_pl,n) = (1/2)* Re((m/e) * v_0*v_1_p . conjugate(%i*ω*p_1)) +
(1/2) * Re(-(m/e) * (v_τ^2/p_0) * express(div(p_1)) * conjugate(i%ω*p_1))
,
avg_t(S_pl,n) = (1/2) * Re(V_1_ek . conjugate(J_1_e)) + (1/2) *
Re(V_1_th * conjugate(J_1_e)) ,
avg_t(S_pl,n) = (1/2) * Re(V_1_e . conjugate(J_1_e))];

$$[\text{avg}_t(S_{\text{pl}}, n) = \text{avg}_t(S_{\text{th}}, n) + \text{avg}_t(S_{\text{ek}}, n), \text{avg}_t(S_{\text{pl}}, n) = (\text{realpart}(v_{1_p} \cdot [-\%ip1_x\omega, -\%ip1_y\omega, -\%ip1_z\omega])$$

(% o184)

(ii) The average plasma power loss density where $\mu_e = (e/m)$ is the static electron mobility and $j_1 = (\tau_1 + r_1 \cdot \text{express}(\text{grad}(\tau_0))) / \tau_0$ and $j_1 = (\kappa_0 - 1) * (E_0 \cdot v / E_0) \cdot ((E_1 + \text{express}(v_0 \sim B-1)) / E_0)$

(% cmt_ch04_4454:[avg_t(q_pl,n) = (1/2) * Re((1/\tau_0) * (m/e) * (v_1 - j_1 *
i185) v_0) . conjugate(%i*\omega*p_1)), avg_t(q_pl,n) = (1/2) * Re((1/\mu_e) * (v_1
- j_1 * v_0) . conjugate(J_1_e))];

$$\left[\text{avg}_t(q_{pl}, n) = 0, \text{avg}_t(q_{pl}, n) = \frac{\text{realpart}([v_1 - j_1 v_{0x}, v_1 - j_1 v_{0y}, v_1 - j_1 v_{0z}] \cdot J_{-1_e})}{2\mu_e} \right]$$

(% o185)

Two new quantities in (4.4.53) A vector of the electronic polarization current density $j_{-1_e} = \%i*\omega*p_{-1}$

(% cmt_ch04_4455:[j_{-1_e} = \%i*\omega*p_{-1}];
i186)

$$[j_{-1_e} = [\%ip_{1_x}\omega, \%ip_{1_y}\omega, \%ip_{1_z}\omega]]$$

(% o186)

A tensor of the effective electronic potential $V_{-1_e} = V_{-1_{ek}} + V_{-1_{th}} *$
 $I_{\text{bar}} = (m/e) * (v_0 * v_{-1_p} + (v_r^2/p_0) * p_{-1} * I_{\text{bar}})$

(% cmt_ch04_4456:[V_{-1_e} = V_{-1_{ek}} + V_{-1_{th}} * I_{\text{bar}}, V_{-1_e} = (m/e) *
i187) (v_0*v_{-1_p} + (v_r^2/p_0) * p_{-1} * I_{\text{bar}})];

$$[V_{-1_e} = I_{\text{bar}}V_{-1_{th}} + V_{-1_{ek}}, V_{-1_e} = \left[\frac{m \left(\frac{I_{\text{bar}} p_{1_x} v_r^2}{p_0} + v_{0x} v_{-1_p} \right)}{e}, \frac{m \left(\frac{I_{\text{bar}} p_{1_y} v_r^2}{p_0} + v_{0y} v_{-1_p} \right)}{e}, \frac{m \left(\frac{I_{\text{bar}} p_{1_z} v_r^2}{p_0} + v_{0z} v_{-1_p} \right)}{e} \right]$$

(% o187)

which involves the tensor of electrostatic potential

(% cmt_ch04_4457:[V_{-1_{ek}} = (m/e) * v_0 * v_{-1_p}];
i188)

$$[V_{-1_{ek}} = \left[\frac{mv_{0x} v_{-1_p}}{e}, \frac{mv_{0y} v_{-1_p}}{e}, \frac{mv_{0z} v_{-1_p}}{e} \right]]$$

(% o188)

and a scalar of the thermal potential (with unit dyadic I_{bar})

(% cmt_ch04_4458:[V_{-1_{th}} = -(m/e) * (v_{\tau}^2/p_0) * (\text{express}(\text{div}(p_1)) ,
i189) V_{-1_{th}} = ((k_B * T)/e) * (p_1/p_0)];

$$[V_{-1_{th}} = \left(\left[-\frac{V_{-1_{th}} m p_{1_x} v_{\tau}^2}{e p_0^2}, -\frac{V_{-1_{th}} m p_{1_y} v_{\tau}^2}{e p_0^2}, -\frac{V_{-1_{th}} m p_{1_z} v_{\tau}^2}{e p_0^2} \right] = \left[-\frac{T k_B m p_{1_x} v_{\tau}^2}{e^2 p_0^2}, -\frac{T k_B m p_{1_y} v_{\tau}^2}{e^2 p_0^2}, -\frac{T k_B m p_{1_z} v_{\tau}^2}{e^2 p_0^2} \right] \right)$$

(% o189)

EM contribution to time-average power relation (4.4.44) according to generalized Poynting vecdtr (4.4.43) is:

(% cmt_ch04_4459:[avg_t(S_em,n) = (1/2) * Re(express(E_1 ~ conju-
i190) gate(H_1_p))]);

$$\left[\text{avg}_t(S_{\text{em}}, n) = \left[\frac{E_{1y}H_{1pz} - E_{1z}H_{1py}}{2}, \frac{E_{1z}H_{1px} - E_{1x}H_{1pz}}{2}, \frac{E_{1x}H_{1py} - E_{1y}H_{1px}}{2} \right] \right]$$

(% o190)

Use the Helmholtz Decomposition Theorem to represent the EM field intensity vectors as sum of curl and potential components

(% cmt_ch04_4460:[E_1 = E_c1 - express(grad(phi_1)), H_1 = H_c1 -
i191) express(grad(psi_1))];

$$[E_{1x}, E_{1y}, E_{1z}] = \left[E_{c1} - \frac{d}{dx}\phi_1, E_{c1} - \frac{d}{dy}\phi_1, E_{c1} - \frac{d}{dz}\phi_1 \right], [H_{1x}, H_{1y}, H_{1z}] = \left[H_{c1} - \frac{d}{dx}\psi_1, H_{c1} - \frac{d}{dy}\psi_1, H_{c1} - \frac{d}{dz}\psi_1 \right]$$

(% o191)

Modify S_em by replacing H_1 with H_1_p to get S_em = express(E_1 ~ H_1_p) From (4.3.48 and 4.3.61) it follows that:

(% cmt_ch04_4461:[H_1_p = H_e1_p - express(grad(psi_1)), H_c1_p = H_c1 -
i192) express(p_1 ~ v_0)];

$$[H_{1px}, H_{1py}, H_{1pz}] = \left[H_{e1p} - \frac{d}{dx}\psi_1, H_{e1p} - \frac{d}{dy}\psi_1, H_{e1p} - \frac{d}{dz}\psi_1 \right], H_{c1p} = [-p_{1y}v_{0z} + p_{1z}v_{0y}]$$

(% o192)

Use Maxwell's equations in dielectric form (4.43.48) and (4.3.61) along with (4.4.60) and (4.4.61) Also use the following vector identities

(% A:[A_x,A_y,A_z];Phi:[Phi-x, Phi_y, Phi_z];
i194)

$$[A_x, A_y, A_z]$$

(% o193)

$$[\Phi - x, \Phi_y, \Phi_z]$$

(% o194)

```
(%      cmt_ch04_4461_A:[
i195)   express(curl(Φ*A) ) = express(Φ*A ~ A) + express(express(grad(Φ)) ~ A) ,
        express(div( express(curl(Φ*A)))) = 0,
        express(curl(curl(Φ))) = 0
        ] ;
```

$$\left[\left[\frac{d}{dy} (A_z \Phi_z) - \frac{d}{dz} (A_y \Phi_y), \frac{d}{dz} (A_x (\Phi - x)) - \frac{d}{dx} (A_z \Phi_z), \frac{d}{dx} (A_y \Phi_y) - \frac{d}{dy} (A_x (\Phi - x)) \right] \right] = \left[A_z \left(\frac{d}{dy} [\Phi - x, \Phi - y, \Phi - z] \right) \right]$$

(% o195)

This gives rise to the relation:TODO: CREATE A VECTOR-VALUED VER-
SION OF avg_tTODO: USE express to evaluate curls, divergence, grad and
cross-product

```
(%      cmt_ch04_4462:[(div(avg_t(S_em,n))) = avg_t(S_add,n) and avg_t(q_b,n)
i196)   = (div( (E_1 ~ h_1_P))) ,
        div(avg_t(S_em,n) ) = div(avg_t( (E_c1 ~ H_c1_p) + ϕ_1 *diff(D_1_p,t)
        + ψ_1 *diff(B_1,t) ,n) ) ] ;
```

$$[\text{false}, \text{div}(\text{avg}_t(S_{em}, n)) = \text{div} \left(\text{avg}_t \left(\left(\frac{d}{dt} [B1_x, B1_y, B1_z] \right) \psi_1 + \phi + E_{c1} \sim H_{c1_p} - \frac{d}{dt} [D1_{px}, D1_{py}, D1_{pz}] \right) \right]$$

(% o196)

Substitute (4.4.62) into (44.44) to give:

```
(%      cmt_ch04_4463:[ div( avg_t(S_el,n) + avg_t(S_add)) + avg_t(q_b,n) = 0 ]
i197)   ;
```

$$[\text{avg}_t(q_b, n) + \text{div}(\text{avg}_t(S_{el}, n) + \text{avg}_t(S_{add})) = 0]$$

(% o197)

where avg_t(S_add,n) and avg_t(q_b,n) are time-averaged values of (4.4.13)
and (4.4.14)The constituents of avg_t(S_add,n) and avg_t(q_b,n) are obtained
fromeqns(4.4.45),(4.4.46), (4.4.49),(4.4.50),(4.4.53),(4.4.54)Now avg_t(S_add,n)
and avg_t(q_b,n) can be given as follows:Not using express for div curl grad
cross-product (~) for nowUse express when vector quantities are properly de-
fined

```
(%      cmt_ch04_4464:[
i198)   avg_t(S_add, n) = avg_t(S_el,n) + avg_t(S_fm,n) + avg_t(S_pl,n) ,
        avg_t(S_add, n) = (1/2) * Re( V_1_a . conjugate(J_1_a) + V_1_m . con-
        jugate(J_1_m) + V_1_e . conjugate(J_1_e))
        ];
```

$$[\text{avg}_t(S_{add}, n) = \text{avg}_t(S_{pl}, n) + \text{avg}_t([S_{fm_x}, S_{fm_y}, S_{fm_z}], n) + \text{avg}_t(S_{el}, n), \text{avg}_t(S_{add}, n) = \frac{\text{realpart}(V_1_a \cdot \text{conjugate}(J_1_a) + V_1_m \cdot \text{conjugate}(J_1_m) + V_1_e \cdot \text{conjugate}(J_1_e))}{2}]$$

(% o198)

```
(% cmt_ch04_4465:[
i199) avg_t(q_b,n) = avg_t(q_add,n),
      avg_t(q_b,n) = avg_t(q_el,n) + avg_t(q_fm,n) + avg_t(q_pl,n),
      avg_t(q_b,n) = (ω^2/2 * conjugate(S_1_bar) . η_bar . S_1_bar)
      + (p_m0/2) * conjugate(U_1) . τ^(-1) . U_1
      + v_M * (μ_0/2) * (ω/ω_M)^2 * abs(M_1)^2 + (1/2) * Re((1/μ_e) * (v_1 -
      j_1*v_0) . conjugate(J_1_e) )
      ];
```

$$[\text{avg}_t(q_b, n) = \text{avg}_t(q_{\text{add}}, n), \text{avg}_t(q_b, n) = \text{avg}_t(q_{\text{pl}}, n) + \text{avg}_t(q_{\text{fm}}, n) + \text{avg}_t(q_{\text{el}}, n), \text{avg}_t(q_b, n) = \frac{\left(S_{1_ba} \right)}{(\% o199)}$$

```
(% cmt_ch04_4466:[avg_t(S_em,n) = (1/2) * Re( (E_c1 ~ conjugate(H_c1_p)
i200) + φ_1 * conjugate(%i*ω*D_1_p)) + ψ_1 * conjugate(%i*ω*B_1)),
      avg_t(S_em,n) = avg_t(S_em_c,n) + avg_t(S_es,n) + avg_t(S_ms,n) ];
```

$$[\text{avg}_t(S_{\text{em}}, n) = \left[\frac{\text{realpart}(E_{c1} \sim H_{c1_p})}{2}, \frac{\text{realpart}(E_{c1} \sim H_{c1_p})}{2}, \frac{\text{realpart}(E_{c1} \sim H_{c1_p})}{2} \right], \text{avg}_t(S_{\text{em}}, n) = (\% o200)$$

```
(% cmt_ch04_4467:[avg_t(S_em_c,n) = (1/2) * Re( E_c1 ~ conju-
i201) gate(H_c1_p))];
```

$$\left[\text{avg}_t(S_{\text{em}_c}, n) = \frac{\text{realpart}(E_{c1} \sim H_{c1_p})}{2} \right] (\% o201)$$

```
(% cmt_ch04_4468:[avg_t(S_es,n) = (1/2) * Re( φ_1 *
i202) conjugate(%i*ω*D_1_p))];
```

$$[\text{avg}_t(S_{\text{es}}, n) = [0, 0, 0]] (\% o202)$$

```
(% cmt_ch04_4469:[avg_t(S_ms,n) = (1/2) * Re(ψ_1*conjugate(%i*ω*B_1)) ] ;
i203)
```

$$[\text{avg}_t(S_{\text{ms}}, n) = [0, 0, 0]] (\% o203)$$

4.4.2 Modal Transmission and Dissipation Of Power

4.4.2.1 General Power Relations

(% kill(S);
i204)

done (% o204)

(% cmt_ch04_4469A:[div(avg_t(S_em,n) + avg_t(S_add,n))] ;
i205)

$[\text{div}(\text{avg}_t(S_{\text{em}},n) + \text{avg}_t(S_{\text{add}},n))]$ (% o205)

(% cmt_ch04_4470:['integrate(div(avg_t(S_em,n) + avg_t(S_add,n)) ,S) =
i206) 'diff('integrate(e_z . avg_t(S_em,n),S) ,z) + 'diff('integrate(e_z .
avg_t(S_add,n),S),z)
- 'sum('integrate(n_i_p . avg_t(S_em_p,n) + n_i_m . avg_t(S_em_m,n)
,L_i),i,1,N)
- 'sum('integrate(n_i_p . avg_t(S_add_p,n) + n_i_m . avg_t(S_add_m,n)
,L_i),i,1,N)
];

$[S \text{div}(\text{avg}_t(S_{\text{em}},n) + \text{avg}_t(S_{\text{add}},n)) = -L_i N (n_{i_p} \cdot \text{avg}_t(S_{\text{em}_p},n) + n_{i_m} \cdot \text{avg}_t(S_{\text{em}_m},n)) - L_i N (n_{i_p}$
(% o206)

(% cmt_ch04_4471:[Q_em_s = 'sum('integrate(n_i_p . avg_t(S_em_p,n) +
i207) n_i_m . avgs_t(S_em_m,n),S),i,1,N];

$[Q_{\text{em}_s} = N S (n_{i_p} \cdot \text{avg}_t(S_{\text{em}_p},n) + n_{i_m} \cdot \text{avgs}_t(S_{\text{em}_m},n))]$ (% o207)

(% cmt_ch04_4472:[Q_add_s = 'sum('integrate(n_i_p . avg_t(S_add_p,n) +
i208) n_i_m . avgs_t(S_add_m,n),S),i,1,N];

$[Q_{\text{add}_s} = N S (n_{i_p} \cdot \text{avg}_t(S_{\text{add}_p},n) + n_{i_m} \cdot \text{avgs}_t(S_{\text{add}_m},n))]$
(% o208)

Contribution by a metallic boundary

(% cmt_ch04_4473:[Q_em_s = (1/2) * 'integrate(R_s*(H_tau . conjugate(H_tau)),
i209) L)];

$\left[Q_{\text{em}_s} = \frac{(H_\tau \cdot H_\tau) L R_s}{2} \right]$ (% o209)

(% cmt_ch04_4474:['diff(P(z),z)~ +~ Q(z)~ = 0];
i210)

$$\left[\frac{d}{dz} P(z) + Q(z) = 0 \right] \quad (\% \text{ o210})$$

(% cmt_ch04_4475:[P~ = P_EM +~ P_PM, P~ = 'integrate(avg_t(S_em,n)~ .
i211) e_z, S) +~ 'integrate(avg_t(S_add,n)~ . e_z, S)];

$$[P = P_{PM} + P_{EM}, P = S (\text{avg}_t(S_{em}, n) \cdot e_z) + S (\text{avg}_t(S_{add}, n) \cdot e_z)] \quad (\% \text{ o211})$$

(% cmt_ch04_4476:[
i212) P_EM = 'integrate(avg_t(S_EM,n)~ .e_z, S) ,
P_EM = 'integrate(avg_t(S_em,n)~ . e_z, S),
P_EM = (1/2)~ *~ Re('integrate((E_c1 ~ conjugate(H_c1_p) +~ \phi_1
~ conjugate(%i\omega*D_1_P) +~ \psi_1 *~ conjugate(%i*\omega*B_1)). e_z,S))
];

$$[P_{EM} = S (\text{avg}_t(S_{EM}, n) \cdot e_z), P_{EM} = S (\text{avg}_t(S_{em}, n) \cdot e_z), P_{EM} = (S)/2] \quad (\% \text{ o212})$$

(% E_c1:[Ec1_x, Ec1_y, Ec1_z];
i214) H_c1_p:[Hc1p_x, Hc1p_y, Hc1p_z];

$$[Ec1_x, Ec1_y, Ec1_z] \quad (\% \text{ o213})$$

$$[Hc1p_x, Hc1p_y, Hc1p_z] \quad (\% \text{ o214})$$

(% cmt_ch04_4477:[P_em = 'integrate(~ avt_t(S_em_c, n)~ . e_z, S), P_em
i215) = (1/2)~ *~ Re('integrate(express(E_c1 ~ conjugate(H_c1_p))~ . e_z, S))
];

$$[P_{em} = S (\text{avt}_t(S_{em_c}, n) \cdot e_z), P_{em} = \frac{\text{realpart}([Ec1_y Hc1p_z - Ec1_z Hc1p_y, Ec1_z Hc1p_x - Ec1_x Hc1p_z, Ec1_x Hc1p_y - Ec1_y Hc1p_z])}{2}] \quad (\% \text{ o215})$$

(% cmt_ch04_4478:[P_es = 'integrate(avg_t(S_es,n)~ . ez, S) , P_es = (1/2)~
i216) *~ Re('integrate(conjugate(\psi_1(%i*\omega*D_1_P)) . e_z, S))];

$$\left[P_{es} = S (\text{avg}_t(S_{es}, n) \cdot e_z), P_{es} = \frac{S \text{realpart}(\overline{\psi_1(\%i\omega D_1 P)} \cdot e_z)}{2} \right] \quad (\% \text{ o216})$$

(% cmt_ch04_4479:[P_ms = 'integrate(avg_t(S_ms,n)~ . e_z, S), P_ms =
i217) (1/2)~ *~ Re(~ 'integrate(conjugate(ψ_1(%i*ω*B_1)) . e_z,S))] ;

$$\left[P_{\text{ms}} = S (\text{avg}_t (S_{\text{ms}}, n) . e_z), P_{\text{ms}} = \frac{S \text{realpart} \left(\overline{\psi_1 ([\%iB1_x\omega, \%iB1_y\omega, \%iB1_z\omega])} . e_z \right)}{2} \right]$$

(% o217)

(% cmt_ch04_4480:[
i218) P_PM = 'integrate(avg_t(S_add,n)~ . e_z, S),
P_PM= (1/2)~ *~ Re(~ 'integrate((V_1_a . conjugate(J_1_a)~ +~
V_1_m . conjugate(J_1_m)~ +~ V_1_e . conjugate(J_1_e)) . e_z,S))
];

$$\left[P_{\text{PM}} = S (\text{avg}_t (S_{\text{add}}, n) . e_z), P_{\text{PM}} = \frac{S \text{realpart} ((V_{1_m} . J_{1_m} + V_{1_e} . J_{1_e} + V_{1_a} . J_{1_a}) . e_z)}{2} \right]$$

(% o218)

4480 involves 3 contributions

(% cmt_ch04_4481:[P_el = 'integrate(avg_t(S_el,n)~ . e_z, s), P_el = (1/2)~
i219) *~ Re(~ 'integrate((V_1_a . conjugate(J_1_a))~ . e_z,S))] ;

$$\left[P_{\text{el}} = \text{integrate} (\text{avg}_t (S_{\text{el}}, n) . e_z, s), P_{\text{el}} = \frac{S \text{realpart} ((V_{1_a} . J_{1_a}) . e_z)}{2} \right]$$

(% o219)

(% cmt_ch04_4482:[P_fm = 'integrate(avg_t(~ S_fm,n)~ . e-z, S), P_fm =
i220) (1/2)~ *~ Re(~ 'integrate ((V_1_m . conjugate(J_1_m))~ . e_z, S))];

$$\left[P_{\text{fm}} = S (\text{avg}_t ([S_{\text{fm}_x}, S_{\text{fm}_y}, S_{\text{fm}_z}], n) . e - z), P_{\text{fm}} = \frac{S \text{realpart} ((V_{1_m} . J_{1_m}) . e_z)}{2} \right]$$

(% o220)

(% cmt_ch04_4483: [
i221) P_pl = 'integrate(avg_t(S_pl,n)~ . e_z, S),
P_pl = (1/2)~ *~ Re(~ 'integrate((V_1_e . conjguate(J_1_e)) . e_z,S))
];

$$\left[P_{\text{pl}} = S (\text{avg}_t (S_{\text{pl}}, n) . e_z), P_{\text{pl}} = \frac{S \text{realpart} ((V_{1_e} . \text{conjguate} (J_{1_e})) . e_z)}{2} \right]$$

(% o221)

(% cmt_ch04_4484:[Q = Q_b +~ Q_s , Q = integrate(avg_t(q_b,n)~ ,S) +~
i222) integrate(avg_t(q_add_s,n),L)];

$$[Q = Q_s + Q_b, Q = Savg_t(q_b, n) + Lavg_t(q_add_s, n)] \quad (\% \text{ o222})$$

(% cmt_ch04_4485:[
i223) Q_b = 'integrate(avg_t(q_b,n),S)~ ,
Q_b = 'integrate(avg_t(q_add,n),S) ,
Q_b = 'integrate(avg_t(q_el,n)~ +~ avg_t(q_fm,n)~ +~
avg_t(q_pl,n),S),
Q_b=Q_el_b +~ Q_fm_b +~ Q_pl_b
];

$$[Q_b = Savg_t(q_b, n), Q_b = Savg_t(q_{add}, n), Q_b = S (avg_t(q_{pl}, n) + avg_t(q_{fm}, n) + avg_t(q_{el}, n)), Q_b = Q_{pl_b} + \dots] \quad (\% \text{ o223})$$

Note : (colon) double-dot notation which is a term-by-term tensor product is being replaced temporarily by standard multiplication. This is easy to implement in octave and matlab with the dotted operations. Bulk losses from 4485 include contributions from:

(% cmt_ch04_4486:[
i224) Q_el_b = 'integrate(avg_t(q_el,n),s),
Q_el_b = (ω^2/2)~*~ 'integrate(conjugate(S_1_bar)*~ η_bar *~ S_1_bar
,S)
+~
(ρ_m0~ /2) *~ 'integrate(conjugate(U_1)~ . (1/τ_bar). U_1,S)] ;

$$[Q_{el_b} = avg_t(q_{el}, n) s, Q_{el_b} = \begin{pmatrix} \frac{SS_{1_bar}^2 \eta_{11} \omega^2}{2} + \frac{S(U_1 \cdot \frac{1}{\tau_bar} \cdot U_1) \rho_{m0}}{2} & \frac{SS_{1_bar}^2 \eta_{12} \omega^2}{2} + \frac{S(U_1 \cdot \frac{1}{\tau_bar} \cdot U_1) \rho_{m0}}{2} \\ \frac{SS_{1_bar}^2 \eta_{21} \omega^2}{2} + \frac{S(U_1 \cdot \frac{1}{\tau_bar} \cdot U_1) \rho_{m0}}{2} & \frac{SS_{1_bar}^2 \eta_{22} \omega^2}{2} + \frac{S(U_1 \cdot \frac{1}{\tau_bar} \cdot U_1) \rho_{m0}}{2} \\ \frac{SS_{1_bar}^2 \eta_{31} \omega^2}{2} + \frac{S(U_1 \cdot \frac{1}{\tau_bar} \cdot U_1) \rho_{m0}}{2} & \frac{SS_{1_bar}^2 \eta_{32} \omega^2}{2} + \frac{S(U_1 \cdot \frac{1}{\tau_bar} \cdot U_1) \rho_{m0}}{2} \end{pmatrix}] \quad (\% \text{ o224})$$

(% cmt_ch04_4487:[
i225) Q_fm_b = 'integrate(avg_t(q_fm,n),S) , Q_fm_b = ν_M~
~ (μ_0/2)~~ (ω/ω_M)^2~*~ 'integrate(abs(M_1)^2,S)

];

$$[Q_{fm_b} = Savg_t(q_{fm}, n), Q_{fm_b} = \left[\frac{M1_x^2 S \mu_0 \nu_M \omega^2}{2 \omega_M^2}, \frac{M1_y^2 S \mu_0 \nu_M \omega^2}{2 \omega_M^2}, \frac{M1_z^2 S \mu_0 \nu_M \omega^2}{2 \omega_M^2} \right]]$$

(% o225)

```
(% cmt_ch04_4488:[
i226) Q_pl_b = 'integrate( avg_t(q_pl, n)~, S),
Q_pl_b = (1/2)~ *~ Re( 'integrate( (1/μ_e)~ *~ (v_1 - j_1 *~ v_0)~ .
conjugate(J_1) ,S))

];
```

$$\left[Q_{pl_b} = S \text{avg}_t(q_{pl}, n), Q_{pl_b} = \frac{S (J_{1z} (v_1 - j_1 v_{0z}) + J_{1y} (v_1 - j_1 v_{0y}) + J_{1x} (v_1 - j_1 v_{0x}))}{2\mu_e} \right]$$

(% o226)

```
(% cmt_ch04_4489:[
i227) Q_s = 'integrate(avg_t(q_add_s, n)~, L),
Q_s = 'integrate(n . avg_t(S_add, n), L),
Q_s = integrate(n . (~ avg_t(S_el,n)~ +~ avg_t(S_fm,n)~ +~
avg_t(S_pl,n) ) ,L) ,
Q_s = Q_el_s +~ Q_fm_s +~ Q_pl_s
];
```

$$[Q_s = L \text{avg}_t(q_{add_s}, n), Q_s = L (n. \text{avg}_t(S_{add}, n)), Q_s = L (n. (\text{avg}_t(S_{pl}, n) + \text{avg}_t([S_{fm_x}, S_{fm_y}, S_{fm_z}]$$

(% o227)

4489 Includes 3 contributions Q_el_S from elastic media, Q_fm_s from surfaces of ferrimagnetic media, Q_pl_s from surfaces of charge carrier streams in plasmas

```
(% cmt_ch04_4490:[
i228) Q_el_s = integrate(avg_t(q_el_s,n),L),
Q_el_s = integrate(n . avg_t(S_el,n)~, L) ,
Q_el_s = (1/2)~ *~ Re( integrate( (n . V_bar_1_a) . conjugate(J_1_a) ,L
) ) ,
Q_el_s = (1/2)~ *~ Re( integrate( - (n . T_1_Σ) . conjugate(U_1) ,L ) )
];
```

$$[Q_{el_s} = L \text{avg}_t(q_{el_s}, n), Q_{el_s} = L (n. \text{avg}_t(S_{el}, n)), Q_{el_s} = \frac{L \text{realpart}((n. V_{\bar{1}a}). J_{1a})}{2}, Q_{el_s} =$$

(% o228)

```
(% cmt_ch04_4491:[
i229) Q_fm_s = integrate(avg_t(q_fm_s,n),L),
      Q_fm_s = integrate(n . avg_t(S_fm,n) ,L) ,
      Q_fm_s = (1/2) * Re( integrate( (n . V_bar_1_m) . conjugate(J_1_m) ,L
      ) ) ,
      Q_fm_s = (1/2) * Re( integrate( %i*omega*mu_0*lambda_ex(n . express(grad(M_1))) .
      conjugate(M_1)~ ,L ) )
];
```

$$[Q_{fm_s} = L \text{avg}_t(q_{fm_s}, n), Q_{fm_s} = L (n \cdot \text{avg}_t([S_{fm_x}, S_{fm_y}, S_{fm_z}], n)), Q_{fm_s} = \frac{L \text{realpart}((n \cdot V_{\bar{1}_m}) \cdot J_{1_m})}{2}$$

(% o229)

```
(% cmt_ch04_4492:[
i230) Q_pl_s = integrate(avg_t(q_pl_s,n),L),
      Q_pl_s = integrate(n . avg_t(S_pl,n) ,L) ,
      Q_pl_s = (1/2) * Re( integrate( (n . V_bar_1_e) . conjugate(J_1_e) ,L ) )
      ,
      Q_pl_s = (1/2) * Re( integrate( %i*omega*(m/e)~ *~ (v_r^2/p_0) *~ (n .
      p_v_1)~ *~ conjugate(p_1)~ ,L ) )
];
```

$$[Q_{pl_s} = L \text{avg}_t(q_{pl_s}, n), Q_{pl_s} = L (n \cdot \text{avg}_t(S_{pl}, n)), Q_{pl_s} = \frac{L \text{realpart}((n \cdot V_{\bar{1}_e}) \cdot J_{1_e})}{2}, Q_{pl_s} =$$

(% o230)

For free solid surface , For rigid solid surface

```
(% cmt_ch04_4493: n . T_bar_1_Sigma = 0;
i232) cmt_ch04_4494: U_1 = 0
      ;
```

$$n \cdot T_{\bar{1}_\Sigma} = 0 \quad (\% \text{ o231})$$

$$U_1 = 0 \quad (\% \text{ o232})$$

```
(% cmt_ch04_4494_A:~ [
i233) n . V_bar_1_a = 0 ,
      J_1_a = 0
      ];
```

$$[n \cdot V_{\bar{1}_a} = 0, J_{1_a} = 0] \quad (\% \text{ o233})$$

Boundary Conditions On Ferrite Surface For Calculation of Ferrimagnetic Loss-
esTextbook $\rightarrow \lambda_{\text{ex}}(\mathbf{n} \cdot \mathbf{M}_1)$. What exactly does this mean? TODO: $\sim I$ have
split it into $\mathbf{n} \cdot \text{div}(\mathbf{M}_1)$ but this needs to be verified

```
(%      cmt_ch04_4494_B:[
i234)   n . V_bar_1_m = -lambda_ex(n . div(M_1) ),
        j_1_M = %I* omega*~ mu_0*~ m_1
        ];
```

$$[n.V_{\text{bar}_1}_m = -\lambda_{\text{ex}}(n \cdot \text{div}([M_{1x}, M_{1y}, M_{1z}]))], j_{1M} = \%Im_1\mu_0\omega]$$

(% o234)

For free surface spins

```
(%      cmt_ch04_4495: rho_1 = 0;
i235)
```

$$\rho_1 = 0$$

(% o235)

For rigid surface spins

```
(%      cmt_ch04_4496: n . p_v_1 = 0;
i236)
```

$$n.p_v_1 = 0$$

(% o236)

4.4.97 and 4.4.98 ensure the following conditions on the unperturbed carrier
stream boundaryThe result is that the contour integral in (4.4.92) vanishes so
 $Q_{\text{pl}_s} = 0$ This shows that there are no plasma surface losses

```
(%      cmt_ch04_4498_A:[
i237)   n . V_bar_1_e = n *~ V_1_th, n . V_bar_1_e = 0,
        n . J_1_e = 0 , n . v_0 = 0
        ];
```

$$[n.V_{\text{bar}_1}_e = V_{1\text{th}}n, n.V_{\text{bar}_1}_e = 0, n.J_{1e} = 0, n.[v_{0x}, v_{0y}, v_{0z}] = 0]$$

(% o237)

```
(%      cmt_ch04_4499:[
i238)   a_k(z)~ = A_k *~ %e^(-gamma_k*z) , a_k(z) = A_k *~ %e^(-alpha_k*z) *~ %e
        ^(-%i*beta_k*z)
        ];
```

$$\left[a_k(z) = A_k e^{-z\gamma_k}, a_k(z) = A_k e^{-iz\beta_k - z\alpha_k} \right]$$

(% o238)

```
(% cmt_ch04_44100:[
i239) Φ_1(r_t,z) = 'sum(A_k *~ Φ_k(r_t,z) , k,1,N),
Φ_1(r_t,z)~ = 'sum (A_k *~ Φ_k_bar(r_t)~ *~ %e^(-γ_k*z), k,1,N),
Φ_1(r_t,z)~ = 'sum(a_k(z)~ *~ Φ_k_bar(r_t), k, 1,N)
];
```

$$\left[\Phi_1(r_t, z) = A_k N \Phi_k(r_t, z), \Phi_1(r_t, z) = A_k N _k_bar(r_t) \% e^{-z\gamma_k}, \Phi_1(r_t, z) = N _k_bar(r_t) a_k(z) \right]$$

(% o239)

For the total power flow P defined in (4.4.75)~

```
(% cmt_ch04_44101:[
i240) P(z) = P_EM(z) + P_PM(z) ,
P(z) = sum(sum(P_kl(z), L, 1,N),K,1,N) , (1/2)*Re(sum(sum(N_kl * con-
jugate(a_k(z)) * a_l(z),l, 1,m),k,1,n)) ,
P(z) = (1/4) * sum(sum( N_kl * conjugate(a_k(z)) * conjugate(a_l(z)),l,1,m),
k,1,n) ,
P(z) = (1/4) * sum(N_k*abs(a_k(z))^2,k,1,n) + (1/2)*Re(sum(sum(N_kl *
conjugate(a_k(z)) * a_l(z),l, 1,m),k,1,n)) ,
P(z) = sum(P_k(z),k,1,n) + sum(sum(Q_kl_pair(z) ,l,1,m),k,1,n)
];
```

$$[P(z) = P_{PM}(z) + P_{EM}(z), P(z) = N^2 P_{kl}(z), \frac{N_{kl} m n (\text{realpart}(a_k(z)) \text{realpart}(a_l(z)) + \text{imagpart}(a_k(z)) \text{imagpart}(a_l(z)))}{2}]$$

(% o240)

For the total power loss Q defined in (4.4.84)~

```
(% cmt_ch04_44102: [
i241) Q(z)~ = Q_b(z)~ +~ Q_s(z)~ ,
Q(z)~ = sum(sum( P_kl(z)~ , l,1,m),k,1,n),
Q(z)~ = (1/4)~ *~ sum(sum(M_kl *~ conjugate(a_k(z))~ *~
conjugate(a_l(z)) , l,1,m),k,1,n) ,
Q(z)~ = (1/4)~ *~ sum(~ sum(~ M_k *~ abs(a_k(z))^2,l,1,m),k,1,n)
+~ (1/2)~ *~ Re( sum(sum(~ M_kl *~ conjugate(a_k(z)~ )~ *~ a_l(z),
l,1,m),k,1,n)) ,
Q(z)~ = sum(Q_k(z), k,1,n)~ +~ sum(sum(Q_kl_pair(z),l,1,m),k,1,n)
];
```

$$[Q(z) = Q_s(z) + Q_b(z), Q(z) = m n P_{kl}(z), Q(z) = \frac{M_{kl} m n \overline{a_k(z)} a_l(z)}{4}, Q(z) = \frac{M_{kl} m n (\text{realpart}(a_k(z)) \text{realpart}(a_l(z)) + \text{imagpart}(a_k(z)) \text{imagpart}(a_l(z)))}{4}]$$

(% o241)


```
(%      cmt_ch04_44103:[
i242)   N_kl = conjugate(N_lk), M_kl = conjugate(M_lk)
];
```

$$[N_{kl} = N_{lk}, M_{kl} = M_{lk}] \quad (\% \text{ o242})$$

```
(%      cmt_ch04_44104:[
i243)   N_kl= N_kl_EM +~ N_kl_PM
];
```

$$[N_{kl} = N_{kl_PM} + N_{kl_EM}] \quad (\% \text{ o243})$$

```
(%      cmt_ch04_44105:[
i244)   N_kl_EM = integrate(
(
conjugate(E_bar_k)~ ~ H_bar_l_P)
+~
E_bar_l ~ conjugate(H_bar_ck_p)~
,S,1,N)~
+~
(~
conjugate(phi_hat_k(%i*omega*D_hat_l_p) )
+~
conjugate(phi_hat_l(%i*omega*D_hat_k_p))
)
+
(
conjugate(psi_hat_k(%i*omega*B_hat_l))
+~
conjugate(psi_hat_l(%i*omega*B_hat_k) )
. e_z, S)
];
```

$$[N_{kl_EM} = \overline{\phi_{\text{hat}_k}(\%iD_{\text{hat}_l_p\omega})} + \overline{\phi_{\text{hat}_l}(\%iD_{\text{hat}_k_p\omega})} + S + (E_{\text{bar}_l} \sim H_{\text{bar_ck}_p} + E_{\text{bar}_k} \sim H_{\text{bar_cl}})] \quad (\% \text{ o244})$$

```
(% cmt_ch04_44106:[
i245) N_kl_EM = 'integrate(
(
(conjugate(E_ck_hat)~ ~ Hn_l_p_hat) +~ (E_l_hat ~ conju-
gate(H_k_p_hat)))
+
(conjugate(phi_k_hat(%I*omega*D_l_p_hat))~ +~
conjugate(phi_l_hat(%i*omega*D_k_p_hat)))~
+
(conjugate(psi_k_hat(%i*omega*B_l_hat))~ +~
conjugate(psi_l_hat(%i*omega*B_k_hat))
)
. e_z, S )
];
```

$$[N_{kl_EM} = S \left(\left(\overline{\psi_{k_hat}(\%iB_{l_hat}\omega)} + \overline{\psi_{l_hat}(\%iB_{k_hat}\omega)} \right) . e_z + \overline{\phi_{k_hat}(\%ID_{l_p_hat}\omega)} + \overline{\phi_{l_hat}(\%ID_{k_p_hat}\omega)} \right) \quad (\% o245)$$

```
(% cmt_ch04_44107:[
i246) N_kl_PM = 'integrate( ( ((conjugate(V_k_a_bar_hat) . J_l_a_hat) +
(V_l_a_bar_hat . conjugate(J_k_a_hat)))
+
(conjugate(V_m_k_bar_hat) . J_l_m_hat)~ + (V_l_m_bar_hat~ . con-
jugate(J_k_m_hat))
+
(conjugate(V_k_e_bar_hat)~ . J_l_e_hat)~ +~ (V_l_e_bar_hat .
conjugate(J_k_e_hat)))~ . e_z, S )
];
```

$$[N_{kl_PM} = S \left((V_{m_k_bar_hat}.J_{l_m_hat} + V_{l_m_bar_hat}.J_{k_m_hat} + V_{l_e_bar_hat}.J_{k_e_hat}) \right) \quad (\% o246)$$

```
(% cmt_ch05_44108:[
i247) M_kl = M_kl_b +~ M_kl_s
];
```

$$[M_{kl} = M_{kl_s} + M_{kl_b}] \quad (\% o247)$$

Here he uses colon (:) for the term-by-term tensor product. Maxima will not accept that syntax, will create function `tbyt` or `prdt` to do this. For now using dot product or multiplication, but this has to be changed later. INCOMPLETE

cmt_ch04_44109~

```
(% cmt_ch04_44109:[
i248) M_kl_b = 2 *~ 'integrate( (ω^2 .~ ( ~ conjugate(S_k_bar_hat) . η_bar .
S_l_bar_hat)) ,S )
];
```

$$\left[M_{kl_b} = 2S \left(\omega^2 . S_{k_bar_hat} . \begin{pmatrix} \eta_{11} & \eta_{12} & \eta_{13} \\ \eta_{21} & \eta_{22} & \eta_{23} \\ \eta_{31} & \eta_{32} & \eta_{33} \end{pmatrix} . S_{l_bar_hat} \right) \right]$$

(% o248)

```
(% cmt_ch04_44110:[
i249) M_kl_s = - 'integrate( ( (n . conjugate(T_hat_k_Σ)) . U_hat_t +~ (n .
T_hat_l_Σ)~ . conjugate(U_hat_k))
+
λ_ex *~ ( ~ (n . grad(conjugate(M_hat_k)))~ . (%i*ω*μ_0 *~ M_hat_t)~
+~ (n . grad(M_hat_t) ) . conjugate(%i*ω*μ_0*M_hat_k) )
+
(m/e)~ *~ (v_r ^2 / ρ_0)~ *~ ( ~ (n . conjugate(p_hat_k))~ . ~
(%i*ω*ρ_hat_l)~ +~ (n . ρ_hat_l)~ . conjugate(%i*ω*ρ_bar_k))
) ,S)~
];
```

$$[M_{kl_s} = -S \left(\frac{m \left(\%i \left((n.p_hat_k) . \rho_hat_l \omega \right) - \%i \left((n.\rho_hat_l) . \rho_bar_k \omega \right) \right) v_r^2}{e \rho_0} + (\%i \left((n. grad (M_hat_k)) . M_h \right. \right.$$

(% o249)

→

```
(% cmt_ch04_44111:[ P_k(z) = P_kk(z), P_kk(z) = 1/4 *~ N_kk . conju-
i250) gate(a_k(z)) *~ a_k(z) ];
```

$$\left[P_k(z) = P_{kk}(z), P_{kk}(z) = \frac{\left(N_{kk} . \overline{a_k(z)} \right) a_k(z)}{4} \right]$$

(% o250)

```
(% cmt_ch04_44112:[ Q_k(z)~ = Q_kk(z), Q_kk(z)~ = 1/4 * M_kk . conju-
i251) gate(a_k(z)) *~ a_k(z)];
```

$$\left[Q_k(z) = Q_{kk}(z), Q_{kk}(z) = \frac{\left(M_{kk} . \overline{a_k(z)} \right) a_k(z)}{4} \right]$$

(% o251)

```
(%      cmt_ch04_44113:[P_kl(z) = conjugate(P_lk(z))~ , P_kl(z) = (1/4)~ *~
i252)  N_kl *~ conjugate(a_k(z)~ ) *~ a_l(z) ]
;
```

$$\left[P_{kl}(z) = \overline{P_{lk}(z)}, P_{kl}(z) = \frac{N_{kl}a_l(z)\overline{a_k(z)}}{4} \right] \quad (\% \text{ o252})$$

```
(%      cmt_ch04_44114:[Q_kl(z) = conjugate(Q_lk(z)) ,Q_kl(z) = (1/4) * M_kl *
i253)  conjugate(a_k(z) ) * a_l(z) ];
```

$$\left[Q_{kl}(z) = \overline{Q_{lk}(z)}, Q_{kl}(z) = \frac{M_{kl}a_l(z)\overline{a_k(z)}}{4} \right] \quad (\% \text{ o253})$$

Example Using tex - Output Latex From Derived Expressions
 Latex strings can be created using maxima and the wxmaxima notebook gets saved as pdf-latex. Then the exported notebook with latex can be processed by the pdflatex command to generate documents. You can also call maxima from SageMath or Maxima-Jupyter and export those notebooks in pdf or latex format. SageTeX also works well and can be integrated into TexMaker and TexStudio for a full publishing environment

```
(%      tex( [Q_kl(z) = conjugate(Q_lk(z))],false );
i256)  tex([Q_kl(z) = conjugate(Q_lk(z))],false ) ;
      tex([Q_kl(z) = (1/4) * M_kl * conjugate(a_k(z) ) * a_l(z) ],false ) ;
                                              (% o254)
```

”\$ \$

$\backslash left[\{ \backslash itQ_kl \} \backslash left(z \backslash right) = \{ \backslash itQ_lk \} \backslash left(z \backslash right) \wedge \backslash star \backslash right] \$ $$

(% o255)

”\$ \$ \backslash left[\{ \backslash itQ_kl \} \backslash left(z \backslash right) = \{ \backslash itQ_lk \} \backslash left(z \backslash right) \wedge \backslash star \backslash right] \\$ \$

(% o256)

”\$ \$ \backslash left[\{ \backslash itQ_kl \} \backslash left(z \backslash right) = \{ \{ \backslash itM_kl \} \backslash , \{ \backslash ita_l \} \backslash left(z \backslash right) \backslash , \{ \backslash ita_k \} \backslash left(z \backslash right) \wedge \backslash star \backslash right] \\$ \$

```
(%      cmt_ch04_44115:[ P_kl_pair(z) = P_kl(z)~ +~ P_lk(z) , P_kl_pair(z) =
i257)  (1/2)~ *~ Re(N_kl *~ conjugate(a_k(z)) *~ a_l(z)) ]
;
```

$$[P_kl_pair(z) = P_{lk}(z) + P_{kl}(z), P_kl_pair(z) = \frac{N_{kl}(\text{realpart}(a_k(z))\text{realpart}(a_l(z)) + \text{imagpart}(a_k(z))\text{imagpart}(a_l(z)))}{2}$$

(% o257)

(% cmt_ch04_4416:[Q_kl_pair(z) = Q_kl(z) +~ Q_lk(z) , Q_kl_pair(z) = 1/2
i258) * Re(M_kk . conjugate(a_k(z)) * a_k(z)~)] ;

$$[Q_{kl_pair}(z) = Q_{lk}(z) + Q_{kl}(z), Q_{kl_pair}(z) = \frac{\text{realpart}\left(M_{kk} \cdot \overline{a_k(z)}\right) \text{realpart}(a_k(z)) - \text{imagpart}\left(M_{kk} \cdot \overline{a_k(z)}\right) \text{imagpart}(a_k(z))}{2}]$$

(% o258)

4.5 Development of SDAM Excitation Theory For SDAM Waveguides
4.5.1 Generalized Reciprocity Theorem In Complex Conjugate Form
In the region of exciting sources in the bulk material Maxwell's equations take this form:

(% cmt_ch04_451:[express(curl(E_1)) = -%i*omega*B_1 - J_b1_m];
i259)

$$\left[\frac{d}{dy} E_{1z} - \frac{d}{dz} E_{1y}, \frac{d}{dz} E_{1x} - \frac{d}{dx} E_{1z}, \frac{d}{dx} E_{1y} - \frac{d}{dy} E_{1x} \right] = [-\%i B_{1x} \omega - J_{b1m_x}, -\%i B_{1y} \omega - J_{b1m_y}, -\%i B_{1z} \omega - J_{b1m_z}]$$

(% o259)

(% cmt_ch04_452:[express(curl(~ H_1_p)) ~ = %i*omega*D_1_p +~ J_b1_e];
i260)

$$\left[\frac{d}{dy} H_{1p_z} - \frac{d}{dz} H_{1p_y}, \frac{d}{dz} H_{1p_x} - \frac{d}{dx} H_{1p_z}, \frac{d}{dx} H_{1p_y} - \frac{d}{dy} H_{1p_x} \right] = [\%i D_{1p_x} \omega + J_{b1e_x}, \%i D_{1p_y} \omega + J_{b1e_y}, \%i D_{1p_z} \omega + J_{b1e_z}]$$

(% o260)

(% cmt_ch04_453:[express(div(D_1_p)) ~ = rho_b1_e];
i261)

$$\left[\frac{d}{dz} D_{1p_z} + \frac{d}{dy} D_{1p_y} + \frac{d}{dx} D_{1p_x} = \rho_{b1e} \right]$$

(% o261)

(% cmt_ch04_454:[express(div(B_1)) ~ = rho_b1_m];
i262)

$$\left[\frac{d}{dz} B_{1z} + \frac{d}{dy} B_{1y} + \frac{d}{dx} B_{1x} = \rho_{b1m} \right]$$

(% o262)

The charge and current densities (electric and magnetic) are related to each other by:

```
(%      cmt_ch04_455:[
i263)   %i*ω*ρ_b1_e +~ express(div(J_b1_e)) = 0,
        %i*ω*ρ_b1_m +~ express(div(J_b1_m))~ = 0
        ];
```

$$[\%i\rho_{b1_e}\omega + \frac{d}{dz}J_{b1_ez} + \frac{d}{dy}J_{b1_ey} + \frac{d}{dx}J_{b1_ex} = 0, \%i\rho_{b1_m}\omega + \frac{d}{dz}J_{b1_mz} + \frac{d}{dy}J_{b1_my} + \frac{d}{dx}J_{b1_mx} = 0]$$

(% o263)

The field-intensity vectors E_1 , $H_1_p = H_1 - p_1 \times v_0$ are linked with the flux-density vectors D_1_p and B_1 by:

```
(%      cmt_ch05_456:[ D_1_p = ε_0 *~ E_1 +~ P_net, D_1_p =  ε_0 *~ E_1
i264)   +~ (P_1 +~ p_1)];
```

$$[[D1_p_x, D1_p_y, D1_p_z] = [E1_x\epsilon_0 + Pnet_x, E1_y\epsilon_0 + Pnet_y, E1_z\epsilon_0 + Pnet_z], [D1_p_x, D1_p_y, D1_p_z] = [E$$

(% o264)

->

```
(%      cmt_ch05_457:[B+1 = μ_0 *~ (H_1_p +~ M_net)~ , B_1 = μ_0~ *~
i265)   (H_1 +~ M_1)];
```

$$[[B_x + 1, B_y + 1, B_z + 1] = [(Mnet_x + H1_p_x) \mu_0, (Mnet_y + H1_p_y) \mu_0, (Mnet_z + H1_p_z) \mu_0], [B1_x, B1_y, B$$

(% o265)

where P_1 , M_1 and p_1 describe physical properties of the medium being analyzed. To obtain the conjugate reciprocity theorem it is necessary to consider another (small-signal)~system denoted with subscript 2. The dynamic equations of this second system all take the complex conjugate form initially used for the BAM analysis. The conventional procedure applied to combine the two systems of equations gives equation (3.5.27). Equation (3.5.27)~is re-written here using equations (4.5.6) and (4.5.7) as follows:

```
(%      cmt_ch04_458:[
i266)   express(div( express(E_1 ~ conjugate(H_2_p)) + express( conjugate(E_2)
        ~ H_1_p))) =
        -%i*ω* (P_1 . conjugate(E_2))~ - (conjugate(P_2)~ . E_1)
        - %i*ω* (~ (M_1 . conjugate(H_2))~ - (conjugate(M_2) . H_1) )
        -%i *ω*~ ((~ p_1 . conjugate(E_2_primed)) - (conjugate(p_2) .
        E_1_primed))
        - (~ (J_b1_e . conjugate(E_2))~ +~ (conjugate(J_b2_e) . E_1) )
        - ~ (J_b1_m . conjugate(H_2_p))~ +~ (conjugate(J_b2_m)~ . H_1_p)
        ];
```

$$[\frac{d}{dx} (E1_y H2p_z - E1_z H2p_y + E2_y H1_p_z - E2_z H1_p_y) + \frac{d}{dy} (-E1_x H2p_z + E1_z H2p_x - E2_x H1_p_z + E2_z H1_p_x$$

(% o266)

where the electric fields $E_{1_primed} = E_1 + \sim v_0 \times B_1$ and $E_{2_primed} = E_2 + \sim v_0 \times B_2$ acmt_re measured relative to anobserver moving with the non-relativistic velocity v_0

(%
i268) S_12:[S_12_x, S_12_y, S_12_z];
H_2_p:[H_2p_x, H_2p_y, H_2p_z];

$[S_{12_x}, S_{12_y}, S_{12_z}]$ (% o267)

$[H_{2p_x}, H_{2p_y}, H_{2p_z}]$ (% o268)

(%
i269) cmt_ch04_459:[express(div(S_12))~ +~ q_12_b = r_12_b];

$\left[q_{12_b} + \frac{d}{dz} S_{12_z} + \frac{d}{dy} S_{12_y} + \frac{d}{dx} S_{12_x} = r_{12_b} \right]$ (% o269)

(%
i270) cmt_ch04_4510:[S_12 = S_12_EM +~ S_12_PM];

$[S_{12_x}, S_{12_y}, S_{12_z}] = S_{12_PM} + S_{12_EM}$ (% o270)

(%
i271) cmt_ch04_4511:[
S_12_EM~ = express(E_1 ~ conjugate(H_2_p))~ +~
express(conjugate(E_2)~ ~ H_1_p)
];

$[S_{12_EM} = [E_{1_y} H_{2p_z} - E_{1_z} H_{2p_y} + E_{2_y} H_{1_p_z} - E_{2_z} H_{1_p_y}, -E_{1_x} H_{2p_z} + E_{1_z} H_{2p_x} - E_{2_x} H_{1_p_z} + E_{2_z} H_{1_p_x}]]$
(% o271)

(%
i272) cmt_ch04_4512:[
S_12_PM~ = (V_1_a . conjugate(J_2_a))~ +~ (conjugate(V_2_a)~ .
J_1_a)
];

$[S_{12_PM} = V_{2_a} \cdot J_{1_a} + V_{1_a} \cdot J_{2_a}]$ (% o272)

```
(% cmt_ch04_4512:[
i273) S_12_PM = ( (V_bar_1_a . conjugate(J_2_a)) + (conjugate(V_bar_2_a)
. J_1_a))
+
( ( V_bar_1_m . conjugate(J_2_m)) +~ (conjugate(V_bar_2_m)~ .
J_1_m) )
+
(( v_bar_1_e . conjugate(J_2_e))~ +~ (conjugate(V_bar_2_e)~ . J_1_e))

];
```

```
[S_12_PM = v_bar_1_e.J_2_e+V_bar_2_m.J_1_m+V_bar_2_e.J_1_e+V_bar_2_a.J_1_a+V_bar_1_m.J_2_m+V_
(% o273)
```

```
(% cmt_ch04_4513:[
i274) q_12_b = 2 *~ (ω^2 *~ S_bar_1 .η_bar . conjugate(S_2))~ +~ (ρ_m0
*~ U_1 . (1/τ)~ . conjugate(U_2))
+
2*ν_M~ *~ μ_0 *~ (ω/ω_M)^2 *~ (M_1 . conjugate(M_2))
+
(1/μ_0)~ *~ (~ (~ v_1 - G_1*v_0 )~ . conjugate(J_2_e) +~
(conjugate(v_2)~ - conjugate(G_2)*v_0)~ . J_1_e )
];
```

```
[q_12_b = \frac{2 ([M1_x, M1_y, M1_z] . M_2) \mu_0 \nu_M \omega^2}{\omega_M^2} + 2 \left( S_{\text{bar}_1} . \begin{pmatrix} \eta_{11} & \eta_{12} & \eta_{13} \\ \eta_{21} & \eta_{22} & \eta_{23} \\ \eta_{31} & \eta_{32} & \eta_{33} \end{pmatrix} . S_2 \right) \omega^2 + \left( U_1 . \frac{1}{\tau} . U_2 \right) \rho_{\text{m}0} + \frac{[v_2 -
(% o274)
```

```
(% cmt_ch04_4514:[
i275) r_12_b = - ((J_b1_e . conjugate(E_2) ) +~ (conjugate(J_b2_e)~ . E_1))
-
((J_b1_m . conjugate(H_2_p)) +~ (conjugate(J_b2_m)~ . H_1_p))

];
```

```
[r_12_b = -J_b2_m . [H1_p_x, H1_p_y, H1_p_z] - J_b2_e . [E1_x, E1_y, E1_z] - H_2_p_z J_b1_m_z - H_2_p_y J_b1_m_y - H_2_p_x J_b1_m_x -
(% o275)
```



```
(% cmt_ch04_4515:[
i276) express(div(      express(E_1      ~      ~      conjugate(H_2_p)~      +~
      express(conjugate(E_2)~ ~ ~ H_1_p))))
=
((phi_1*~      conjugate(%i*omega*D_2_p))      +~      (conjugate(phi_2)~      *~
(%i*omega*D_1_p))) + ((psi_1* conjugate(%i*omega*B_2)) +~ (conjugate(psi_2)~
*~ (%i*omega*B_1)))
+~
((conjugate(J_b1_e)~      .      express(grad(conjugate(phi_2))))      +~
(conjugate(J_b2_e)~      .      express(grad(phi_1)))      )      +      ((J_b1_m      .
express(grad(conjugate(psi_2))))      +~      (conjugate(J_b2_m)~      .
express(grad(psi_1))) )
-
((~ (%i*omega*rho_b1_e) *~ conjugate(phi_2) ) ~ +~ (conjugate(%i*omega*rho_b2_e)
*~ phi_1)) - ( (%i*omega*rho_b1_m)~      *~      conjugate(psi_2))      +~
(conjugate(%i*omega*rho_b2_m) *~ psi_1)
];
```

$$\left[\frac{d}{dx} (E_{1y}H_{2p_z} - E_{1z}H_{2p_y} + E_{2y}H_{1p_z} - E_{2z}H_{1p_y}) + \frac{d}{dy} (-E_{1x}H_{2p_z} + E_{1z}H_{2p_x} - E_{2x}H_{1p_z} + E_{2z}H_{1p_x}) \right] \quad (\% \text{ o276})$$

```
(% E_c1:[E_c1_x, E_c1_y, E_c1_z];
i280) E_c2:[E_c2_x, E_c2_y, E_c2_z];
      H_c2_p:[H_c2p_x, H_c2p_y, H_c2p_z];
      H_c1_p:[H_c1p_x, H_c1p_y, H_c1p_z];
```

$$[E_{c1_x}, E_{c1_y}, E_{c1_z}] \quad (\% \text{ o277})$$

$$[E_{c2_x}, E_{c2_y}, E_{c2_z}] \quad (\% \text{ o278})$$

$$[H_{c2p_x}, H_{c2p_y}, H_{c2p_z}] \quad (\% \text{ o279})$$

$$[H_{c1p_x}, H_{c1p_y}, H_{c1p_z}] \quad (\% \text{ o280})$$

```
(% cmt_ch04_4516:[
i281) S_12_EM = (express(E_c1 ~ conjugate(H_c2_p))) +~
express(conjugate(E_c2)~ ~ ~ H_c1_p)
+
((phi_1*conjugate(%i*omega*D_2_p)) +~ conjugate(phi_2) *~ (%i*omega*D_1_p))
+
( psi_1*conjugate(%i*omega*B_2)) +~ (conjugate(psi_2)~ *~ (%i*omega*B_1))
];
```

$$[S_{12_EM} = [\%iB_{1x}\psi_2\omega - \%iB_{2y}\psi_1\omega + \%iD_{1_p_x}\phi_2\omega - \%iD_{2_p}\phi_1\omega + E_{c1_y}H_{c2_p_z} - E_{c1_z}H_{c2_p_y} + E_{c2_y}H_{c1_p_z} - E_{c2_z}H_{c1_p_y}] + \text{conjugate terms}]$$

(% o281)

```
(% cmt_ch04_4517:[
i282) r_12_b = - (J_b1_e . conjugate(E_c2)) +~ (conjugate(J_b2_e)~ . E_c1)
- (J_b1_m . conjugate(H_c2_p)) +~ (conjugate(J_b2_m)~ . H_c1_p)
+
((%i*omega*rho_b1_e)~ *~ conjugate(phi_2)) +~ (conjugate(%i*omega*rho_b2_e)~ *
phi_1)~
+
((%i*omega*rho_b1_m) * conjugate(psi_2)) + (conjugate(%i*omega*rho_b2_m) * psi_1)
];
```

$$[r_{12_b} = \%i\rho_{b1_m}\psi_2\omega - \%i\rho_{b2_m}\psi_1\omega + \%i\rho_{b1_e}\phi_2\omega - \%i\rho_{b2_e}\phi_1\omega + J_{b2_m} \cdot [H_{c1_p_x}, H_{c1_p_y}, H_{c1_p_z}] + \text{conjugate terms}]$$

(% o282)

4.5.2 Quasi-Orthogonality and Orthogonality of Modes In SDAM Waveguides

4.5.2.1 Mode quasi-orthogonality relations for lossy waveguides

```
(% cmt_ch04_4518:[
i283) 'diff(P_kl(z),z)~ +~ Q_kl(z)~ = 0
];
```

$$\left[\frac{d}{dz} P_{kl}(z) + Q_{kl}(z) = 0 \right]$$

(% o283)

```
(% cmt_ch04_4519:[
i284) P_kl(z)~ = (1/4)~ *~ 'integrate( conjugate(S_kl_EM (r_t,z)) . e_z,S, 1,N)
+
(1/4) * 'integrate( conjugate(S_kl_PM (r_t,z)) . e_z, S, 1, N),
P_kl(z)~ = P_kl_EM +~ P_kl_PM,
P_kl(z) = (1/4) * (N_kl_EM + N_kl_PM) * conjugate(A_k) * A_l * exp(-
(conjugate(γ_k) + γ_l) * z),
P_kl(z)~ = (1/4)~ *~ N_kl*~ conjugate(a_k(z))~ *a_l(z)~

];
```

$$[P_{kl}(z) = \frac{(N-1) \left(\overline{S_{kl_PM}(r_t, z)} \cdot e_z \right)}{4} + \frac{(N-1) \left(\overline{S_{kl_EM}(r_t, z)} \cdot e_z \right)}{4}, P_{kl}(z) = P_{kl_PM} + P_{kl_EM},$$

(% o284)

```
(% cmt_ch04_520:[
i285) Q_kl(z) = (1/4) * 'integrate( conjugate(q_kl_b(r_t,z)) , S, 1,N)
+
(1/4) * 'integrate( conjugate(q_kl_s(r_t,z)) , S, 1, N),
Q_kl(z)~ = Q_kl_b +~ Q_kl_s,
Q_kl(z)~ = (1/4) * (M_kl_b + M_kl_s) * conjugate(A_k) * A_l * exp(-
(conjugate(γ_k) + γ_l) * z),
Q_kl(z)~ = (1/4) * M_kl* conjugate(a_k(z)) *a_l(z)
];
```

$$[Q_{kl}(z) = \frac{(N-1) \overline{q_{kl_s}(r_t, z)}}{4} + \frac{(N-1) \overline{q_{kl_b}(r_t, z)}}{4}, Q_{kl}(z) = Q_{kl_s} + Q_{kl_b}, Q_{kl}(z) = \frac{A_k A_l (M_{kl_s} + M_{kl_b})}{4}$$

(% o285)

```
(% cmt_ch04_520_A:[
i286) 'diff(N_kl * exp(-(conjugate(γ_k) * γ_l) * z),z) + M_kl * exp(-
(conjugate(γ_k) + γ_l) * z) = 0
];
```

$$\left[\frac{d}{dz} \left(N_{kl} e^{-z\gamma_k \gamma_l} \right) + M_{kl} e^{z(-\gamma_l - \gamma_k)} = 0 \right]$$

(% o286)

```
(% cmt_ch04_521:[
i287) (conjugate(γ_k)~ +~ γ_l)~ *~ N_kl = M_kl, (conjugate(γ_k) +~ γ_l)~
*~ P_kl = Q_kl
];
```

$$[N_{kl}(\gamma_l + \gamma_k) = M_{kl}, P_{kl}(\gamma_l + \gamma_k) = Q_{kl}]$$

(% o287)

```
(%      cmt_ch04_4522:[
i288)  (conjugate( $\gamma_k$ ) +  $\gamma_l$ ) *~ (N_kl_EM +~ N_kl_PM)~ = M_kl_b +~
      M_kl_s
      ];
```

$$[(N_{kl_PM} + N_{kl_EM})(\gamma_l + \gamma_k) = M_{kl_s} + M_{kl_b}] \quad (\% \text{ o288})$$

4.5.2.2 Mode Orthogonality Relations For Lossless Waveguides

```
(%      cmt_ch04_4523:[
i289)  (conjugate( $\gamma_k$ ) +  $\gamma_l$ ) * N_kl = (conjugate( $\gamma_k$ ) +  $\gamma_l$ ) * (N_kl_EM +
      N_kl_PM),
      (conjugate( $\gamma_k$ ) +  $\gamma_l$ ) * N_kl = 0
      ];
```

$$[N_{kl}(\gamma_l + \gamma_k) = (N_{kl_PM} + N_{kl_EM})(\gamma_l + \gamma_k), N_{kl}(\gamma_l + \gamma_k) = 0] \quad (\% \text{ o289})$$

```
(%      cmt_ch04_4524:[
i290)  ( $\beta_k - \beta_l$ )~ *~ N_kl = ( $\beta_k - \beta_l$ )~ *~ (N_kl_EM + N_kl_PM),
      ( $\beta_k - \beta_l$ ) * N_kl = 0
      ];
```

$$[N_{kl}(\beta_k - \beta_l) = (N_{kl_PM} + N_{kl_EM})(\beta_k - \beta_l), N_{kl}(\beta_k - \beta_l) = 0] \quad (\% \text{ o290})$$

```
(%      cmt_ch04_4525:[
i291)  N_kl = n_KK *~  $\delta_{kl}$ ,
      if l != k then N_kl = 0,
      if l = k then N_kk = N_k
      ];
```

$$[N_{kl} = n_{KK}\delta_{kl}, \text{false}, \text{false}] \quad (\% \text{ o291})$$

```
(%      E_ck_hat:[Eckh_x, Eckh_y, Eckh_z];
i293)  H_ck_p_hat:[H_ckph_x, H_ckph_y, H_ckph_z];
```

$$[Eckh_x, Eckh_y, Eckh_z] \quad (\% \text{ o292})$$

$$[H_{ckph_x}, H_{ckph_y}, H_{ckph_z}] \quad (\% \text{ o293})$$

```
(% cmt_ch04_4526:[
i294) N_k = N_k_EM +~ N_k_PM,
      N_k = 2 *~ Re(
      integrate(
      (express(conjugate(E_ck_hat)~ ~ H_ck_p_hat)
      +~ conjugate(phi_k_hat)~ *~ (%i*omega*D_k_p_hat)
      +~ psi_k_hat *~ (%i*omega*B_k_hat))~ . e_z, s,1,N))
      ];
```

$$[N_k = N_{k_PM} + N_{k_EM}, N_k = 2(N-1) \text{realpart}([\%i B_{k_hat} \psi_{k_hat} \omega + \%i D_{k_p_hat} \phi_{k_hat} \omega +$$

(% o294)

```
(% cmt_ch04_4527:[
i295) V_k_a_bh = -T_k_Sigma_bh,
      J_k_a_hat = %i*omega*u_k_hat,
      J_k_a_hat = U_k_hat
      ];
```

$$[V_{k_a_bh} = -T_{k_bh}, J_{k_a_hat} = \%i u_{k_hat} \omega, J_{k_a_hat} = U_{k_hat}]$$

(% o295)

```
(% cmt_ch04_4528:[
i296) V_k_m_hat = -lambda_cx *~ M_k_hat,
      J_k_m_hat = %i*omega*mu_0 *~ M_k_hat
      ];
```

$$[V_{k_m_hat} = -\lambda_{cx} M_{k_hat}, J_{k_m_hat} = \%i M_{k_hat} \mu_0 \omega]$$

(% o296)

```
(% cmt_ch04_4529:[
i297) V_k_e_bh = (m/e)~ *~ (v_0*v_k_p_hat +~ (v_T^2 / o_0)~ *~ p_k_hat
      * I_bar),
      J_k_e_hat = %i*omega*p_k_hat
      ];
```

$$[V_{k_e_bh} = [\frac{m \left(v_{0x} v_{k_p_hat} + \frac{I_{bar} p_{k_hat} v_T^2}{o_0} \right)}{e}, \frac{m \left(v_{0y} v_{k_p_hat} + \frac{I_{bar} p_{k_hat} v_T^2}{o_0} \right)}{e}, \frac{m \left(v_{0z} v_{k_p_hat} + \frac{I_{bar} p_{k_hat} v_T^2}{o_0} \right)}{e}]$$

(% o297)

```
(% cmt_ch04_4530:[
i298) P_k = (1/4)~ *~ N_k *~ abs(a_k)^2];
```

$$\left[P_k = \frac{N_k a_k^2}{4} \right]$$

(% o298)

```
(%      cmt_ch04_4531:[
i299)  P_kk = conjugate(P_kk), P_kk = (1/4)~ *~ N_kk *~ conjugate(a_k)~ *
      a_k
      ];
```

$$\left[P_{kk} = P_{kk}, P_{kk} = \frac{N_{kk} a_k^2}{4} \right] \quad (\% \text{ o299})$$

```
(%      cmt_ch04_4532:[
i300)  N_kl = N_kk_tilde *~ δ_k_tilde_l,
      if l ≠ k then N_kl = 0,
      if l = k then N_kk_tilde = N_k
      ];
```

$$[N_{kl} = N_{kk_tilde} \delta_{kl}, N_{kl} = 0, \text{false}] \quad (\% \text{ o300})$$

```
(%      H_ck_tilde_p_hat:[H_cktph_x, H_cktph_y, H_cktph_z];
i301)
```

$$[H_{cktph_x}, H_{cktph_y}, H_{cktph_z}] \quad (\% \text{ o301})$$

```
(%
i302) cmt_ch04_4533:[
N_k = N_kk_EM +~ N_kk_PM,
N_k = 'integrate(
(
express(conjugate(E_ck_hat) ~ H_ck_tilde_p_hat) +~ express(E_ck_hat
~ conjugate(H_ck_p_hat))
+
(conjugate(phi_hat) *( %i*omega*D_k_tilde_p_hat)~ +~ phi_k_tilde_hat~ *~
conjugate(%i*omega*D_k_p_hat))
+
(conjugate(psi_k_hat)~ *~ (%i*omega*B_k_tilde_hat) +~ psi_k *~
conjugate(%i*omega*B_k_hat )) ). e_z,

S, 1, N)
+
'integrate(
(
(conjugate( V_k_a_bh) . J_k_tilde_a_hat~ +~ V_k_a_bh . conju-
gate(J_k_a_hat))
+
(conjugate(V_k_m_bh) . J_k_tilde_m_hat +~ V_k_tilde_m_bh . conju-
gate(J_k_m_hat))
+
(conjugate(V_k_e_bh)~ . J_k_tilde_e_hat +~ V_k_tilde_e_bh .
conjugate(J_k_e_hat))~
) . e_z, S, 1,N
)
]
;
```

$$[N_k = N_{kk_PM} + N_{kk_EM}, N_k = (N - 1) ([\%i B_k_tilde_hat \psi_k_hat \omega - \%i B_k_hat \psi_k \omega - \%i D_k_p_hat \psi_k_hat \omega])$$

(% o302)

```
(%
i303) cmt_ch04_4534:[
P_kk_pair = P_k_k_tilde +~ P_k_tilde_k~ ,
P_kk_pair = (1/2)~ *~ Re( ~ N_k_k_tilde *~ conjugate(a_k)~ *~
a_k_tilde)~ ,
P_kk_pair = (1/2 )~ *~ Re( N_k *~ conjugate(a_k)~ *~ a_k_tilde)
];
```

$$[P_{kk_pair} = P_{k_tilde_k} + P_{k_k_tilde}, P_{kk_pair} = \frac{N_{k_k_tilde} a_k a_{k_tilde}}{2}, P_{kk_pair} = \frac{N_k a_k a_{k_tilde}}{2}]$$

(% o303)

```
(%      cmt_ch04_4535:[
i304)   Φ_k(r_t,z)~ = Φ_k_bar(r_t) *~ exp(-γ_k*z)
];
```

$$\left[\Phi_k(r_t, z) = _k_bar(r_t) \% e^{-z\gamma_k} \right] \quad (\% \text{ o304})$$

```
(%      cmt_ch04_4536:[
i305)   Φ_1(r_t,z)~ = Φ_a(r_t,z)~ +~ Φ_b(r_t,z)~ ,
        Φ_1(r_t,z)~ = 'sum(A_l(z) *~ Φ_1(r_t,z)~ + Φ_b(r_t,z),l,1,N),
        Φ_1(r_t,z) = A_l(z)~ *~ Φ_1_hat(r_t)~ *~ exp(-γ_k*z) +~ Φ_b(r_t,z),
        Φ_1(r_t,z) = 'sum(a_l(z) *~ Φ_1_hat(r_t)~ +~ Φ_b(r_t,z), l,1,N)
];
```

$$\left[\Phi_1(r_t, z) = \Phi_b(r_t, z) + \Phi_a(r_t, z), \Phi_1(r_t, z) = N(\Phi_l(r_t, z) A_l(z) + \Phi_b(r_t, z)), \Phi_1(r_t, z) = _l_hat(r_t) A_l(z) \right] \quad (\% \text{ o305})$$

```
(%      n_s_plus:~ [n_sp_x, n_sp_y, n_sp_z];
i319)   E_c1_plus: [E_c1p_x, E_c1p_y, E_c1p_z];
        n_s_minus: [n_sm_x, n_sm_y, n_sm_z];
        E_c1_minus: [E_c1m_x, E_c1m_y, E_c1m_z];
        H_c1_p_plus: [H_c1pp_x, H_c1pp_y, H_c1pp_z];
        H_c1_p_minus: [H_c1pm_x, H_c1pm_y, H_c1pm_z];
        D_1_p_plus: [D_1pp_x, D_1pp_y, D_1pp_z];
        D_1_p_minus: [D_1pm_x, D_1pm_y, D_1pm_z];
        B_1_plus: [B_1p_x, B_1p_y, B_1p_z];
        B_1_minus: [B_1m_x, B_1m_y, B_1m_z];
        φ_1_plus: [φ_1p_x, φ_1p_y, φ_1p_z];
        φ_1_minus: [φ_1m_x, φ_1m_y, φ_1m_z];
        ψ_1_plus: [ψ_1p_x, ψ_1p_y, ψ_1p_z];
        ψ_1_minus: [ψ_1m_x, ψ_1m_y, ψ_1m_z];
```

$$\left[n_sp_x, n_sp_y, n_sp_z \right] \quad (\% \text{ o306})$$

$$\left[E_c1p_x, E_c1p_y, E_c1p_z \right] \quad (\% \text{ o307})$$

$$\left[n_sm_x, n_sm_y, n_sm_z \right] \quad (\% \text{ o308})$$

$$\left[E_c1m_x, E_c1m_y, E_c1m_z \right] \quad (\% \text{ o309})$$

$$[H_c1pp_x, H_c1pp_y, H_c1pp_z] \quad (\% \text{ o310})$$

$$[H_c1pm_x, H_c1pm_y, H_c1pm_z] \quad (\% \text{ o311})$$

$$[D_1pp_x, D_1pp_y, D_1pp_z] \quad (\% \text{ o312})$$

$$[D_1pm_x, D_1pm_y, D_1pm_z] \quad (\% \text{ o313})$$

$$[B_1p_x, B_1p_y, B_1p_z] \quad (\% \text{ o314})$$

$$[B_1m_x, B_1m_y, B_1m_z] \quad (\% \text{ o315})$$

$$[\phi_1p_x, \phi_1p_y, \phi_1p_z] \quad (\% \text{ o316})$$

$$[\phi_1m_x, \phi_1m_y, \phi_1m_z] \quad (\% \text{ o317})$$

$$[\psi_1p_x, \psi_1p_y, \psi_1p_z] \quad (\% \text{ o318})$$

$$[\psi_1m_x, \psi_1m_y, \psi_1m_z] \quad (\% \text{ o319})$$

For the current sheet

(% cmt_ch04_4537:[express(n_s_plus ~ E_c1_plus) +~ express(n_s_minus ~
i320) E_c1_minus) = -J_s_m];

$$[[-E_c1p_y n_sp_z + E_c1p_z n_sp_y - E_c1m_y n_sm_z + E_c1m_z n_sm_y, E_c1p_x n_sp_z - E_c1p_z n_sp_x + E_c1m_x n_sm_z - E_c1m_z n_sm_x]]$$

(% o320)

(% cmt_ch04_4538:[express(n_s_plus ~ H_c1_p_plus) + express(n_s_minus
i321) ~ H_c1_p_minus) = -J_s_E];

[[-H_c1pp_y n_sp_z + H_c1pp_z n_sp_y - H_c1pm_y n_sm_z + H_c1pm_z n_sm_y , H_c1pp_x n_sp_z - H_c1pp_z n_sp_x + H_c1pm_x n_sm_z - H_c1pm_z n_sm_x],
(% o321)

For the charge sheet

(% cmt_ch04_4539:[express(n_s_plus ~ D_1_p_plus) + express(n_s_minus
i322) ~ D_1_p_minus) = ρ_s_e];

[[-D_1pp_y n_sp_z + D_1pp_z n_sp_y - D_1pm_y n_sm_z + D_1pm_z n_sm_y , D_1pp_x n_sp_z - D_1pp_z n_sp_x + D_1pm_x n_sm_z - D_1pm_z n_sm_x],
(% o322)

(% cmt_ch04_4540:[express(n_s_plus ~ B_1_plus) + express(n_s_minus ~
i323) B_1_minus) = ρ_s_m];

[[-B_1p_y n_sp_z + B_1p_z n_sp_y - B_1m_y n_sm_z + B_1m_z n_sm_y , B_1p_x n_sp_z - B_1p_z n_sp_x + B_1m_x n_sm_z - B_1m_z n_sm_x],
(% o323)

For the dipole sheet

(% cmt_ch04_4541:[express(n_s_plus ~ ϕ_1_plus) + express(n_s_minus ~
i324) ϕ_1_minus) = (1/ε_0) ~ * ~ η_s_e];

[[n_sp_y ϕ_1p_z - n_sp_z ϕ_1p_y + n_sm_y ϕ_1m_z - n_sm_z ϕ_1m_y , -n_sp_x ϕ_1p_z + n_sp_z ϕ_1p_x - n_sm_x ϕ_1m_z + n_sm_z ϕ_1m_x],
(% o324)

(% cmt_ch04_4542:[express(n_s_plus ~ ψ_1_plus) + express(n_s_minus ~
i325) ψ_1_minus) = (1/ε_0) * η_s_m];

[[n_sp_y ψ_1p_z - n_sp_z ψ_1p_y + n_sm_y ψ_1m_z - n_sm_z ψ_1m_y , -n_sp_x ψ_1p_z + n_sp_z ψ_1p_x - n_sm_x ψ_1m_z + n_sm_z ψ_1m_x],
(% o325)

(% S_lk_EM:[S_lkem_x, S_lkem_y, S_lkem_z];
i327) S_lk_PM:[S_lkpm_x, S_lkpm_y, S_lkpm_z];

[S_lkem_x, S_lkem_y, S_lkem_z] (% o326)

$$[S_lkpm_x, S_lkpm_y, S_lkpm_z] \quad (\% \text{ o327})$$

```
(% cmt_ch04_4543:[
i328) 'integrate(express(div(S_lk_EM~ +~ S_lk_PM )) , S, 1, N ) +~
'integrate(q_lk_b, S,1, N)~ +~ 'integrate(r_lk_b, S, 1, N)
];
```

$$[(N-1)r_lk_b + (N-1)q_lk_b + (N-1) \left(\frac{d}{dz} (S_lkpm_z + S_lkem_z) + \frac{d}{dy} (S_lkpm_y + S_lkem_y) + \frac{d}{dx} (S_lkpm_x + S_lkem_x) \right)] \quad (\% \text{ o328})$$

```
(% cmt_ch04_4544:[S_lk_EM];
i329)
```

$$[[S_lkem_x, S_lkem_y, S_lkem_z]] \quad (\% \text{ o329})$$

```
(% cmt_ch04_4545:[
i330) S_lk_PM = ((V_bar_l_a . conjugate(J_k_a))~ +~
(conjugate(V_bar_k_a) . J_l_a))
+
((V_bar_l_m~ . conjugate(J_bar_k_m)~ ) +~ (conjugate(V_bar_k_m)~ .
J_l_m))
+
((V_bar_l_e . conjugate(J_k_e)) +~ (conjugate(V_bar_k_e)~ . J_l_e))
];
```

$$[[S_lkpm_x, S_lkpm_y, S_lkpm_z] = V_bar_l_m . J_bar_k_m + V_bar_l_e . J_k_e + V_bar_l_a . J_k_a + V_bar_k_m . J_l_m + V_bar_k_e . J_l_e + V_bar_k_a . J_l_a] \quad (\% \text{ o330})$$

The dot product below between S_bar_l and η_bar and S_bar_k should be the term by term tensor product :~(colon)

```
(% cmt_ch04_4546:[
i331) q_lk_b = 2 *~ (ω^2 *S_bar_l . η_bar . conjugate(S_bar_k)~ +~ ρ_m0
*~ U_1 . (1/τ_bar) . conjugate(U_k) )
+~
2 *~ ν_M~ *~ μ_0 *~ (ω/ω_M)^2 *~ (M_l . conjugate(M_k))
+
(1/μ_e)~ *~ ( ~ (v_1 - j_1 *~ v_0)~ . conjugate(J_k_e)~ +~
(conjugate(v_k)~ - conjugate(j_k)~ *~ v_0) . J_l_e

];
```

$$[q_{lk_b} = \frac{2(M_l.M_k)\mu_0\nu_M\omega^2}{\omega_M^2} + 2 \left(\left(S_{\text{bar}_l} \cdot \begin{pmatrix} \eta_{11} & \eta_{12} & \eta_{13} \\ \eta_{21} & \eta_{22} & \eta_{23} \\ \eta_{31} & \eta_{32} & \eta_{33} \end{pmatrix} \cdot S_{\text{bar}_k} \right) \omega^2 + \left(U_1 \cdot \frac{1}{\tau_{\text{bar}}} \cdot U_k \right) \rho_{\text{m0}} \right) + \frac{[v_0]}{\omega_M^2}$$

(% o331)

```
(% cmt_ch04_4547:[
i332) r_lk_b = -(J_bl_e . conjugate(E_ck)~ +~ J_bl_m . conjugate(H_ck_p)~ )
+
((%i*ω*ρ_bl_e)*conjugate(φ_k) +~ (%i*ω*ρ_bl_m)~ *~ conjugate(ψ_k))
];
```

$$[r_{lk_b} = \%i\rho_{bl_m}\psi_k\omega + \%i\rho_{bl_e}\phi_k\omega - J_{bl_m}.H_{ck_p} - J_{bl_e}.E_{ck}]$$

(% o332)

Apply divergence theorem (3.2.25)~to the expression for S_lk_EM in Eq (4.5.44)

```
(% E_cl:[E_cl_x, E_cl_y, E_cl_z];
i336) E_ck:[E_ck_x, E_ck_y, E_ck_z];
H_ck_p:[H_ckp_x, H_ckp_y, H_ckp_z];
H_cl_p:[H_clp_x, H_clp_y, H_clp_z];
```

$$[E_{cl_x}, E_{cl_y}, E_{cl_z}]$$

(% o333)

$$[E_{ck_x}, E_{ck_y}, E_{ck_z}]$$

(% o334)

$$[H_{ckp_x}, H_{ckp_y}, H_{ckp_z}]$$

(% o335)

$$[H_{clp_x}, H_{clp_y}, H_{clp_z}]$$

(% o336)

The three line integrals in the following expression have a minus sign $\sim(-)$ at the end of the expression inside the integral. This denotes a boundary condition where we are 'looking' toward one side of a boundary. One side is denoted as the $+\sim$ direction or side, the other is the $-$ direction or side. The first section of each line integral is 'looking' toward the $+\sim$ side of the boundary. The second section of each line integral is 'looking' toward the $-\sim$ side of the boundary. NOTE: At this point I'm not sure how to represent this in maxima, but may end up using unit vectors perpendicular to the boundary. Then multiplying the expressions inside the line integrals by the appropriate unit vector. The other approach is to use sub-expressions named with the $+\sim$ and $-$ or p and m to hold the $+\sim$ and $-$ parts of the integral expressions.


```
(%      cmt_ch04_4548_B:[
i338)  'integrate( express(div(S_lk_EM ))~ , S,1,N)~
      =
      'diff(
      'integrate(S_lk_EM . e_z, S,1,N) - 'integrate(r_kl_s, L,1,N)
      )
      ];
```

$$[(N-1) \left(\frac{d}{dz} S_{\text{lkem}_z} + \frac{d}{dy} S_{\text{lkem}_y} + \frac{d}{dx} S_{\text{lkem}_x} \right) = \text{del} ((N-1) ([S_{\text{lkem}_x}, S_{\text{lkem}_y}, S_{\text{lkem}_z}] . e_z) - (\% \text{ o338})$$

```
(%      cmt_ch04_4549:[
i339)  \tau_lk_s = -(express(J_s_e . conjugate(E_ck)) +~ express(J_s_m . conju-
      gate(H_ck_p)))
      +
      (%i*\omega*p_s_e *~ conjugate(\phi_k) +~ %i*\omega*p_s_m )
      +~
      ((\eta_s_e / \epsilon_0)~ . conjugate(%i*\omega*D_k_p)) ~ +~ ((\eta_s_m/\mu_0) .
      conjugate(%i*\omega*B_k))
      ];
```

$$[\tau_{\text{lk}_s} = \%i p_{s_e} \text{conjugate}(\phi_k) \omega + \%i p_{s_m} \omega - \%i \left(\frac{\eta_{s_m}}{\mu_0} . B_k \omega \right) - \%i \left(\frac{\eta_{s_e}}{\epsilon_0} . D_{k_p} \omega \right) - J_{s_m} . [H_{\text{ckp}_x}, H_{\text{ck}} (\% \text{ o339})$$

```
(%      S_lk_PM:[ S_lkpm_x, S_lkpm_y, S_lkpm_z];
i340)
```

$$[S_{\text{lkpm}_x}, S_{\text{lkpm}_y}, S_{\text{lkpm}_z}] \quad (\% \text{ o340})$$

```
(%      cmt_ch04_4550:[
i341)  'integrate(express(div(S_lk_PM)),S),
      'diff(~ 'integrate(S_lk_PM . e_z,S)~ , z )
      +~
      'integrate(q_lk_s,L)
      ];
```

$$[S \left(\frac{d}{dz} S_{\text{lkpm}_z} + \frac{d}{dy} S_{\text{lkpm}_y} + \frac{d}{dx} S_{\text{lkpm}_x} \right), Lq_{\text{lk}_s} + \frac{d}{dz} (S ([S_{\text{lkpm}_x}, S_{\text{lkpm}_y}, S_{\text{lkpm}_z}] . e_z))]]$$

(% o341)

```
(%
i342) cmt_ch04_4551:[
      q_lk_s = n . S_lk_PM,
      q_lk_s = - ( (n . T_bar_l_Σ)~ . conjugate(U_k) +~ conjugate(n .
      T_bar_lk_Σ) . U_l )
      +~
      λ_ex *~ ( (n . express(grad(M_l))) . conjugate(%i*ω*μ_0~ *~ M_l) +~ (n
      . conjugate(express(grad(M_k))) . (%i*ω*μ_0*M_l) ))
      -
      ((n . p_l)~ *~ conjugate(%i*ω*p_k)~ +~ (n-conjugate(p_k)) *~ (%i*ω*p_l)
      )
      ];
```

$$[q_{lk_s} = n. [S_{lkpm_x}, S_{lkpm_y}, S_{lkpm_z}], q_{lk_s} = -\%i (n - p_k) p_l \omega + \%i (n. p_l) p_k \omega + \left(\%i \left(n. \left[\frac{d}{dx} M_k, \frac{d}{dy} M \right. \right. \right.$$

(% o342)

```
(%
i343) cmt_ch04_4552:[
      'diff(P_kl(xz)~ ,z)~ +~ Q_lk(z)~ = R_lk(z)
      ];
```

$$\left[Q_{lk}(z) + \frac{d}{dz} P_{kl}(xz) = R_{lk}(z) \right]$$

(% o343)

```
(%
i344) cmt_ch04_4553:[
      P_lk= P_lk_EM +~ P_lk_PM,
      P_lk = 'integrate( (S_lk_EM~ +~ S_lk_PM)~ . e_z, S)
      ];
```

$$[P_{lk} = P_{lk_PM} + P_{lk_EM}, P_{lk} = S ([S_{lkpm_x} + S_{lk_EM}, S_{lkpm_y} + S_{lk_EM}, S_{lkpm_z} + S_{lk_EM}$$

(% o344)

```
(%
i345) cmt_ch04_4554:[
      Q_lk = Q_lk_b +~ Q_lk_s,
      Q_lk = 'integrate(q_lk_b, S)~ +~ 'integrate(q_lk_s, S)
      ];
```

$$[Q_{lk} = Q_{lk_b} + Q_{lk_s}, Q_{lk} = S q_{lk_s} + S q_{lk_b}]$$

(% o345)

```
(%
i346) cmt_ch04_4555:[
      R_lk = R_lk_EM +~ R_lk+PM,
      R_lk = 'integrate(r_lk_b, S_b)~ +~ 'integrate(r_lk_s, L_s)
      ];
```

$$[R_{lk} = R_{lk_EM} + R_{lk} + PM, R_{lk} = L_s r_{lk_s} + S_b r_{lk_b}]$$

(% o346)


```
(%      cmt_ch04_4556:[
i347)   P_lk = P_ak +~ P_bk,
        P_lk = P_ak_EM +~ P_bl_PM,
        P_lk = 'integrate((S_ak_EM +~ S_bk_PM)~ . e_z, S)
        ];
```

$$[P_{lk} = P_{bk} + P_{ak}, P_{lk} = P_{bl_PM} + P_{ak_EM}, P_{lk} = S ((S_{bk_PM} + S_{ak_EM}) . e_z)]$$

(% o347)

```
(%      cmt_ch04_4557:[
i348)   Q_lk = Q_ak +~ Q_bk,
        Q_lk = Q_ak_EM +~ Q_bk_PM,
        Q_lk = 'integrate(q_ak_b,S)~ +~ 'integrate(q_ak_s, L)
        ];
```

$$[Q_{lk} = Q_{bk} + Q_{ak}, Q_{lk} = Q_{bk_PM} + Q_{ak_EM}, Q_{lk} = Lq_{ak_s} + Sq_{ak_b}]$$

(% o348)

```
(%      cmt_ch04_4558:[
i349)   'diff(P_bk(z),z)~ +~ Q_bk(z)~ = 0
        ]
        ;
```

$$\left[\frac{d}{dz} P_{bk}(z) + Q_{bk}(z) = 0 \right]$$

(% o349)

```
(%      cmt_ch04_4559:[
i350)   'diff(P_ak(z),z)~ +~ Q_ak(z)~ =~ R_lk(z)
        ];
```

$$\left[\frac{d}{dz} P_{ak}(z) + Q_{ak}(z) = R_{lk}(z) \right]$$

(% o350)

(%
i351) cmt_ch04_4560:[

```
P_ak(z)~ = P_ak_EM +~ P_ak_PM(z),
P_ak(z)~ = 'integrate( (S_ak_EM(z)~ +~ S_ak_PM(z)) . e_z,s),
P_ak(z)~ = 'sum(A_l(z),l,-N,N)~ *~ integrate( (S_lk_EM +~ S_lk_PM)~
. e_z,S) ,
P_ak(z) = 'sum(A_l(z),l,-N,N) * integrate( (conjugate(S_lk_EM) + conju-
gate(S_lk_PM)) . e_z,S),
P_ak(z)~ = 'sum( (N_kl_EM + N_kl_PM) * A_l(z) * exp(-(conjugate(γ_k)
+ γ_l) * e_z), l,-N,N),
P_ak(z)~ = ('sum(N_kl *~ A_l(z)~ *~ exp(-γ_kl *~ z) , l,-N,N))
*~ exp(-conjugate(γ_k)~ *~ z)
];
```

[$P_{ak}(z) = P_ak_PM(z) + P_ak_EM$, $P_{ak}(z) = s((S_ak_PM(z) + S_ak_EM(z)) \cdot e_z)$, $P_{ak}(z) = (2N + 1) S$ (
(% o351)

(%
i352) cmt_ch04_4561:[

```
Q_ak(z)~ = Q_ak_b(z) +~ Q_ak_s(z) ,
Q_ak(z) = 'integrate(q_ak_b(z), S,1,N) +~ 'integrate(q_ak_s(z), L,1,N),
Q_ak(z)~ = 'sum(A_l(z)~ * ( 'integrate(q_ak_b(z), S,1,N) + 'inte-
grate(q_ak_s(z), L,1,N)) ,L,1,N),
Q_ak(z)~ = 'sum((M_kl_b +~ M_kl_s)*A_l(z)*exp(-
(conjugate(γ_k)+γ_l)*z),L,1,N),
Q_ak(z)~ = 'sum( (M_kl *~ A_l(z)*exp(-γ_l *z) ) ,L,1,N) *~ exp(-
conjugate(γ_k)*z)
];
```

[$Q_{ak}(z) = Q_ak_s(z) + Q_ak_b(z)$, $Q_{ak}(z) = (N - 1) q_ak_s(z) + (N - 1) q_ak_b(z)$, $Q_{ak}(z) = NA_l(z) ((N - 1) q$ (
(% o352)

(%
i353) cmt_ch04_4562:[

```
R_lk(z)~ = R_lk_b(z)~ +~ R_lk_s(z),
R_lk(z) = (R_lk_b(z) + R_lk_s(z)) *~ exp(-conjugate(γ_k)*z),
R_lk(z)~ = R_lk(z)~ *~ exp(-conjugate(γ_k)*z)
];
```

[$R_{lk}(z) = R_lk_s(z) + R_lk_b(z)$, $R_{lk}(z) = (R_lk_s(z) + R_lk_b(z)) \% e^{-z\gamma_k}$, $R_{lk}(z) = R_{lk}(z) \% e^{-z\gamma_k}$]
(% o353)

```
(% cmt_ch04_4563:[
i354) R_k_b(z)~ = 'integrate(
-((J_b_e . conjugate(Eck_hat)) +~ (J_b_m . conjugate(H_ck_p_hat)))~
+
(%i*~ω*ρ_b_e~ *~ conjugate(φ_k_hat)~ +~ %i*~ω*ρ_s_m *~
conjugate(ψ_k_hat))
,S )
]
;
```

$$[R_{k_b}(z) = S \left(\%i\rho_{s_m}\psi_{k_hat}\omega + \%i\rho_{b_e}\phi_{k_hat}\omega - J_{b_m} \cdot [H_{ckph_x}, H_{ckph_y}, H_{ckph_z}] - J_{b_e} \cdot E_{ckph} \right) \quad (\% \text{ o354})$$

```
(% cmt_ch04_4564:[
i355) R_k_s(z) = 'integrate(
-(J_b_e . conjugate(E_hat_ck) +~ J_b_m . conjugate(H_hat_ck_p))
+~
(%i*~ω*ρ_b_e *~ conjugate(φ_k) +~ %i*~ω*ρ_b_m *~ conjugate(ψ_k))
+
( (η_s_e / ε_0) . conjugate(%i*~ω*D_hat_k_p) + (η_s_m/μ_0) .
conjugate(%i*~ω*B_hat_k) )
,L
)
];
```

$$[R_{k_s}(z) = L \left(\%i\rho_{b_m}\psi_k\omega + \%i\rho_{b_e}\phi_k\omega - \%i \left(\frac{\eta_{s_m}}{\mu_0} \cdot B_{hat_k}\omega \right) - \%i \left(\frac{\eta_{s_e}}{\epsilon_0} \cdot D_{hat_k_p}\omega \right) - J_{b_m} \cdot H_{hat_ck} \right) \quad (\% \text{ o355})$$

```
(% cmt_ch04_4565:[
i356) 'sum( N_kl *~ 'diff(A_l.z)~ - (((conjugate(γ_k) +γ_l)~ *~ N_kl - M_kl)~
*~ A_l),L,1,N)~ * exp(-γ_l*z)~
];
```

$$\left[N \%e^{-z\gamma_l} (N_{kl} \text{ del } (A_l.z) - A_l (N_{kl} (\gamma_l + \gamma_k) - M_{kl})) \right] \quad (\% \text{ o356})$$

```
(% cmt_ch04_4566:[
i357) 'sum('sum(N_kl *~ 'diff(A_l(z),z) *~ exp(-γ_l*z)~ ,l,1,N) , k,1,N) = R_z(z)
];
```

$$\left[N^2 N_{kl} \left(\frac{d}{dz} A_l(z) \right) \%e^{-z\gamma_l} = R_z(z) \right] \quad (\% \text{ o357})$$

```
(%      cmt_ch04_4567:[
i358) 'sum('sum(N_kl * ('diff(a_l(z),z) +~ \gamma_l*a_l(z)) ,K,1,N)~ ,L,1,N)
];
```

$$\left[N^2 N_{kl} \left(a_l(z) \gamma_l + \frac{d}{dz} a_l(z) \right) \right] \quad (\% \text{ o358})$$

```
(%      E_hat_ck:[E_hck_x, E_hck_y, E_hck_z];
i362) E_hat_cl:[E_hcl_x, E_hcl_y, E_hcl_z];
      H_hat_cl_p:[H_hclp_x, H_hclp_y, H_hclp_z];
      H_hat_ck_p:[H_hckp_x, H_hckp_y, H_hckp_z];
```

$$[E_{hck_x}, E_{hck_y}, E_{hck_z}] \quad (\% \text{ o359})$$

$$[E_{hcl_x}, E_{hcl_y}, E_{hcl_z}] \quad (\% \text{ o360})$$

$$[H_{hclp_x}, H_{hclp_y}, H_{hclp_z}] \quad (\% \text{ o361})$$

$$[H_{hckp_x}, H_{hckp_y}, H_{hckp_z}] \quad (\% \text{ o362})$$

```

(% cmt_ch04_4568:[
i363) N_kl = 'integrate(
(
express(conjugate(E_hat_ck) ~ H_hat_cl_p)
+~
express(E_hat_cl~ ~ conjugate(H_hat_ck_p))
)
+
(
conjugate(phi_hat_k)~ *~ (%i*omega*D_hat_l_p)~ +~ phi_hat_l *~
conjugate(%i*omega*D_hat_k_p)
)
+
(
conjugate(psi_hat_k) *~ (%i*omega*B_hat_l)~ +~ psi_hat_l *~
conjugate(%i*omega*B_hat_k)
) . e_z
,S,1,N)
+
'integrate(
(
(
conjugate(V_bar_hat_k_a) . J_hat_ll_a)
+~
(V_bar_hat_l_a . conjugate(J_hat_k_a))
)
+
(
conjugate(V_bar_hat_k_m) . J_hat_l_m)
+
(V_bar_hat_l_m . conjugate(J_hat_k_m))
)
+
(
conjugate(V_bar_hat_k_e)~ . J_hat_l_e)~
+
(V_bar_hat_l_e . conjugate(J_hat_k_e))
)
) . e_z,
S,1,N)
];

```

$$[N_{kl} = [(N-1) ((iB_{hat_l} \psi_{hat_k} \omega - iB_{hat_k} \psi_{hat_l} \omega) . e_z - iD_{hat_k} \phi_{hat_l} \omega + iD_{hat_l} \phi_{hat_k} \omega) \\ (\% o363)$$

```
(%
i364) R_k(z)~= 'integrate(
-(J_b_e . conjugate(E_hat_ck) +~ J_b_m . conjugate(H_hat_ck_p))
+
(%i*ω*ρ_b_e *~ conjugate(φ_k) +~ ~ %i*ω*ρ_b_m *~ conjugate(ψ_k))
,S_b,1,N)
+
'integrate(
-(J_b_e . conjugate(E_hat_ck) + J_b_m . conjugate(H_hat_ck_p))
+
(%i*ω*ρ_b_e * conjugate(φ_k) + %i*ω*ρ_b_m * conjugate(ψ_k))
+~
( (η_s_e/ε_0) . conjugate(%i*ω*D_hat_k_p) +~ (η_p_m/μ_0) *~
conjugate(%i*ω*B_hat_k) )
,
L,1,N)
|;
```

$$[R_k(z) = (N-1) (\%i\rho_{-b_m}\psi_k\omega + \%i\rho_{-b_e}\phi_k\omega - \frac{\%iB_hat_k\eta_p m\omega}{\mu_0} - \%i\left(\frac{\eta_{-s_e}}{\epsilon_0} .D_hat_k p\omega\right) - J_{-b_m} . [H_hckp$$

(% 0364)

```
(% E_hat_k:[E_hk_x, E_hk_y, E_hk_z];
i368) H_hat_l_p:[H_hlp_x, H_hlp_y, H_hlp_z];
H_hat_k_p:[H_hkp_x, H_hkp_y, H_hkp_z];
H_hat_l:[H_hl_x, H_hl_y, H_hl_z];
```

$$[E_hk_x, E_hk_y, E_hk_z] \quad (\% \text{ o365})$$

$$[H_hlp_x, H_hlp_y, H_hlp_z] \quad (\% \text{ o366})$$

$$[H_hkp_x, H_hkp_y, H_hkp_z] \quad (\% \text{ o367})$$

$$[H_hl_x, H_hl_y, H_hl_z] \quad (\% \text{ o368})$$

```
(%
i369) cmt_ch04_4570:[
N_kl = 'integrate( (express(conjugate(E_hat_k) ~ H_hat_l_p)
+~ express(H_hat_l ~ conjugate(H_hat_k_p)))~ . e_z, S, 1,N)
+
'integrate(
(
((conjugate(V_hat_bar_k_a) . J_hat_l_a)~ +~ (V_bar_hat_l_a . conju-
gate(J_hat_k_a)))
+~
((conjugate(V_bar_hat_k_m)~ . J_hat_l_m)~ +~ (V_bar_hat_l_m .
conjugate(J_hat_k_m)))
+
((conjugate(V_bar_hat_k_e) . J_hat_l_e)~ +~ (V_bar_hat_l_e . conju-
gate(J_hat_k_e)))
)
~ . e_z, S,1,N)
];
```

$$[N_{kl} = (N - 1) ((V_{\text{hat_bar_}k_a} \cdot J_{\text{hat_}l_a} + V_{\text{bar_hat_}l_m} \cdot J_{\text{hat_}k_m} + V_{\text{bar_hat_}l_e} \cdot J_{\text{hat_}k_e} + V_{\text{bar_hat_}l_m} \cdot J_{\text{hat_}k_m} + V_{\text{bar_hat_}l_e} \cdot J_{\text{hat_}k_e})$$

(% o369)

```
(%
i370) cmt_ch04_4571:[
R_kl(z) = 'integrate( (J_b_e . conjugate(E_hat_k)) +~ (J_b_m .
conjugate(H_hat_k_p)), S_b, 1,N)
-
'integrate( (J_s_e . conjugate(E_hat_k)) +~ (J_s_m . conju-
gate(H_hat_k_p)) ,L_s,1,N)
-
'integrate( (J_s_eff_e . conjugate(E_hat_k)) +~ (J_s_eff_m . conju-
gate(H_hat_k)) , L_b_s,1,N)

];
```

$$[R_{kl}(z) = - (J_{s_m} \cdot [H_{\text{hkp}_x}, H_{\text{hkp}_y}, H_{\text{hkp}_z}] + J_{s_e} \cdot [E_{\text{hk}_x}, E_{\text{hk}_y}, E_{\text{hk}_z}]) (N - 1) - (J_{s_{\text{eff}_m}} \cdot H_{\text{hkp}_z})$$

(% o370)

```
(%
i371) cmt_ch04_4572:[
    (eta_s_e/epsilon_0) . conjugate(%i*omega*D_hat_k_p) + (eta_s_m/mu_0) *
    conjugate(%i*omega*B_hat_k)
    =
    phi_s*n_o ~ *conjugate(%i*omega*D_hat_k_p) + ~ psi_s * ~ n_o .
    conjugate(%i*omega*B_hat_k)
];
```

$$\left[-\frac{\%i B_{\text{hat}_k} \eta_{s_m} \omega}{\mu_0} - \%i \left(\frac{\eta_{s_e}}{\epsilon_0} \cdot D_{\text{hat}_k p} \omega \right) = -\%i D_{\text{hat}_k p} n_o \phi_s \omega - \%i (n_o \cdot B_{\text{hat}_k} \omega) \psi_s \right] \quad (\% \text{ o371})$$

4.5.3.2 Equations of mode excitation for lossless SDAM waveguides (Space Dispersive Active Media)

```
-> cmt_ch04_4573:[
'diff(A_k(z) ~ ,z) =
(1/N_k) ~ * ~ 'integrate(
-(J_b_e . conjugate(E_ck) ~ + ~ J_b_m . conjugate(H_ck_p))
+
(%i*omega*rho_b_e * ~ conjugate(phi_k) + ~ %i*omega*rho_s_m ~ * ~ conjugate(psi_k))
,S_b, 1,N)
+
(1/N_k) * 'integrate(
-(J_b_e . conjugate(E_ck) + J_b_m . conjugate(H_ck_p))
+
(%i*omega*rho_b_e * conjugate(phi_k) + %i*omega*rho_s_m * conjugate(psi_k))
((~ eta_s_e/epsilon_0) ~ . conjugate(%i*omega*D_k_p) + ~ (eta_s_m/mu_0) .
conjugate(%i*omega*B_k))

,L_s, 1,N)

];
```

$$\left[\frac{d}{dz} A_k(z) = ((N-1) (\%i \rho_{s_m} \psi_k \omega + \%i \rho_{b_e} \text{conjugate}(\phi_k) \omega - J_{b_m} \cdot [H_{\text{ckp}_x}, H_{\text{ckp}_y}, H_{\text{ckp}_z}] - J_{b_e} \cdot [H_{\text{ckp}_x}, H_{\text{ckp}_y}, H_{\text{ckp}_z}]) \right] \quad (\% \text{ o376})$$


```

-> cmt_ch04_4574:[

'diff(a_k(z),z) +~ %i*beta*a_k(z)~ =
(1/N_k) * 'integrate(
-(J_b_e . conjugate(E_ck) + J_b_m . conjugate(H_ck_p))
+
(%i*omega*rho_b_e * conjugate(phi_k) + %i*omega*rho_s_m * conjugate(psi_k))
,S_b, 1,N)
+
(1/N_k) * 'integrate(
-(J_b_e . conjugate(E_ck) + J_b_m . conjugate(H_ck_p))
+
(%i*omega*rho_b_e * conjugate(phi_k) + %i*omega*rho_s_m * conjugate(psi_k))
+
(( eta_s_e/epsilon_0) . conjugate(%i*omega*D_hat_k_p) + (eta_s_m/mu_0) .
conjugate(%i*omega*B_hat_k))

,L_s, 1,N)
];

```

$$[\%ia_k(z)\beta k + \frac{d}{dz}a_k(z) = ((N-1))/N_k + ((N-1) (\%i\rho_{s_m}\psi_k\omega + \%i\rho_{b_e}\text{conjugate}(\phi_k)\omega - J_{b_m} \cdot [H_{ckp_x} \\ (\% o396)$$


```

-> cmt_ch04_4576:[
'diff(A_k(z),z) =
(1/N_k) * 'integrate(
-(J_b_e . conjugate(E_ck) + J_b_m . conjugate(H_ck_p))
+
(%i*omega*r_b_e * conjugate(phi_k) + %i*omega*r_s_m * conjugate(psi_k))
,S_b, 1,N)
+
(1/N_k) * 'integrate(
-(J_b_e . conjugate(E_ck) + J_b_m . conjugate(H_ck_p))
+
(%i*omega*r_b_e * conjugate(phi_k) + %i*omega*r_s_m * conjugate(psi_k))
+
((eta_s_e/epsilon_0) . conjugate(%i*omega*D_k_p) + (eta_s_m/mu_0) .
conjugate(%i*omega*B_k))
,L_s, 1,N)

];

```

$$\left[\frac{d}{dz} A_k(z) = ((N-1))/N_k + ((N-1) (\%i\rho_{s_m} \psi_k \omega + \%i\rho_{b_e} \text{conjugate}(\phi_k) \omega - J_{b_m} \cdot [H_{ckp_x}, H_{ckp_y}, H_{ckp_z}]) \right. \\ \left. (\% \text{ o395}) \right]$$

```
(%
i398) cmt_ch04_4577:[
a_k = A_k(z)~*~exp(-gamma_k8z),
'diff(a_k(z),z) + gamma_k*a_k(z) =
(1/N_k)*'integrate(
-(J_b_e . conjugate(E_hat_ck) + J_b_m . conjugate(H_hat_ck_p))
+
(%i*omega*rho_b_e * conjugate(phi_hat_k) + %i*omega*rho_s_m * conjugate(psi_hat_k))
,S_b, 1,N)
+
(1/N_k)*'integrate(
-(J_s_e . conjugate(E_hat_ck) + J_s_m . conjugate(H_hat_ck_p))
+
(%i*omega*rho_b_e * conjugate(phi_hat_k) + %i*omega*rho_s_m * conjugate(psi_hat_k))
+
((eta_s_e/epsilon_0) . conjugate(%i*omega*D_hat_k_p) + (eta_s_m/mu_0) .
conjugate(%i*omega*B_hat_k))
,L_s, 1,N)
];
```

$$\left[a_k = A_k(z) e^{-\gamma_k 8z}, a_k(z) \gamma_k + \frac{d}{dz} a_k(z) = ((N-1))/N_k + ((N-1))/N_k \right] \quad (\% \text{ o398})$$

```
(%
i399) cmt_ch04_461:[
N_hat . Z(z)~ = R(z)
]
;
```

$$[N_{\text{hat}} \cdot Z(z) = R(z)] \quad (\% \text{ o399})$$

```
-> cmt_ch04_462:[
Z(k)= 'diff(a_k(z)~,z)~ +~ gamma_k * a_k(z)~,
Z(k) = 'diff(A_k(z),z) *~ exp(-gamma_k*z)
];
```

$$\left[Z(k) = a_k(z) \gamma_k + \frac{d}{dz} a_k(z), Z(k) = \left(\frac{d}{dz} A_k(z) \right) e^{-z \gamma_k} \right] \quad (\% \text{ o400})$$

```
(%      cmt_ch04_463:[
i409)  if is_active(N_kl)~ then N_kl = N_k *~ δ_kl,
        if is_reactive(N_kl) then N_kl = R_k_bar * δ_k_bar_l,
        if is_reactive(N_k) then N_k = N_k_k_bar,
        if is_reactive(Z_k) then Z_k = R_k_bar/N_k_bar,
        N_k_k_bar = conjugate(N_k_bar_k),
        conjugate(N_kbar_k) = conjugate(N_k_bar)
];
```

[if is_active(N_{kl}) then $N_{kl} = N_k \delta_{kl}$, if is_reactive(N_{kl}) then $N_{kl} = R_{k_bar} \delta_{k_bar_l}$, if is_reactive(N_k) then $N_k = N_{k_k_bar}$, if is_reactive(Z_k) then $Z_k = R_{k_bar}/N_{k_bar}$, $N_{k_k_bar} = \text{conjugate}(N_{k_bar_k})$, $\text{conjugate}(N_{kbar_k}) = \text{conjugate}(N_{k_bar})$]
 (% o409)

```
(%      cmt_ch04_464:[
i410)  conjugate(γ_k)~ +~ γ_k_bar = 0,
        if is_active(Z_k)~ then Z_k = Rk/N_k,
        if is_reactive(Z_k) then Z_k = Rk/N_k_bar
];
```

$\left[\gamma_{k_bar} + \gamma_k = 0, \text{if is_active}(Z_k) \text{ then } Z_k = \frac{Rk}{N_k}, \text{if is_reactive}(Z_k) \text{ then } Z_k = \frac{Rk}{N_{k_bar}} \right]$
 (% o410)