Parameters

$$a > 0 \tag{1}$$

$$b > 0 \tag{2}$$

$$K \ge 0$$

$$D_0 \ge 0$$
(3)

$$\tau \ge 0 \tag{5}$$

Potential

$$V(x) = \frac{a}{4}x^4 - \frac{b}{2}x^2 + \frac{b^2}{4a} = \frac{a}{4}x^2\left(x - \sqrt{\frac{2b}{a}}\right)\left(x + \sqrt{\frac{2b}{a}}\right) + \frac{b^2}{4a}$$
 (6)

Force

$$F(x) = -\frac{\partial V}{\partial x} = -ax^3 + bx \tag{7}$$

We have 2 Hypervolumes

$$\mathcal{V}_t \subset \mathbb{R}^{n_t} \tag{8}$$

$$\mathcal{V}_{\tau} \subset \mathbb{R}^{n_{\tau}} \tag{9}$$

$$\mathcal{V} = \mathcal{V}_t \times \mathcal{V}_\tau \tag{10}$$

$$n_t \ge n_\tau \tag{11}$$

$$\dot{X}_t = F(X_t, X_\tau) \,\mathrm{d}t + \sqrt{2D_0} \,\mathrm{d}B_t \tag{12}$$

$$= -aX_t^3 dt + bX_t dt + K(X_t - X_\tau) dt + \sqrt{2D_0} dB_t$$
(13)

$$F(x,x_{\tau}) = F_1(x) + F_2(x_{\tau}) = -ax^3 + (b+K)x - Kx_{\tau}$$
(14)

Fokker-Planck Equation

From the Langevin equation

$$\dot{X} = -aX^{3}(t) + bX(t) + K(X(t) - X(t - \tau)) + \sqrt{2D_{0}}\xi(t)$$
(15)

With the uncorrelated Gaussian noise ξ

$$\langle \xi(t) \rangle = 0 \tag{16}$$

$$\langle \xi(t)\xi(t')\rangle = \delta(t - t') \tag{17}$$

The infinite hierarchy

$$\partial_t \rho_1(x,t) = -\partial_x \int_{\mathcal{V}_{\tau}} \mathrm{d}x_{\tau} F(x,x_{\tau}) \rho_2(x,t;x_{\tau},t-\tau) + D_0 \partial_x^2 \rho_1(x,t) \tag{18}$$

$$\partial_t \rho_2(x, t; x_\tau, t - \tau) = -\partial_{x_\tau} \int_{\mathcal{V}} dx_{2\tau} F(x_\tau, x_{2\tau}) \rho_3(x, t; x_\tau, t - \tau; x_{2\tau}, t - 2\tau) + \dots$$
(19)

$$(20)$$

Conditional probability densities

$$\rho_1(x,t) = \langle \delta[X(t) - x] \rangle \tag{21}$$

$$\rho_2(x,t;x_\tau,t-\tau) = \langle \delta[X(t)-x]\delta[X(t-\tau)-x_\tau] \rangle \tag{22}$$

Novikov's theorem

$$\partial_t \rho_1(x,t) = -\partial_x \int dx_\tau F(x,x_\tau) \rho_2(x,t;x_\tau,t-\tau) + D_0 \partial_x^2 \rho_1(x,t)$$
(23)

$$\partial_t \rho_2(x, t; x_{\tau}, t - \tau) = -\partial_x F(x, x_{\tau}) \rho_2(x, t; x_{\tau}, t - \tau) - \partial_{x_{\tau}} \int_{\mathcal{V}_{\tau}} dx_{2\tau} F(x_{\tau}, x_{2\tau}) \rho_3(x, t; x_{\tau}, t - \tau; x_{2\tau}, t - 2\tau)$$
(24)

$$+ D_0 \left(\partial_x^2 + \partial_{x_\tau}^2 + \partial_x \partial_{x_\tau} \left(\frac{\delta X(t)}{\delta \xi(t-\tau)} \Big|_{X(t)=x}^{X(t-\tau)=x_\tau} \right) \right) \rho_2(x,t;x_\tau,t-\tau)$$
(25)

Crazy cross diffusion term for a test functional Λ of the noise process

$$\langle \Lambda[\xi]\xi(t)\rangle = \left\langle \frac{\delta\Lambda}{\delta\xi(t)} \right\rangle \tag{26}$$

Auxiliary variables

$$j = 1, 2, \dots, N \tag{27}$$

$$X(t) \to (X_0, X_1, X_2, \dots, X_N), \quad \mathbb{R} \to \mathbb{R}^{N+1}$$
 (28)

Slice the interval

$$\Delta t = \frac{\tau}{N}, \quad X_j(t) = X(t - j\Delta t) \tag{29}$$

$$\Rightarrow X_0(t) = X(t), \quad X_N(t) = X(t - \tau) \tag{30}$$

Dynamics

$$\dot{X}_j = \lim_{\Delta t \to 0} \frac{X_j(t + \Delta t) - X_j(t)}{\Delta t} \approx \frac{N}{\tau} (X_{j-1}(t) - X_j(t))$$
(31)

Langevin equation and Markovian system

$$\dot{X}_0 = F(X_0(t), X_N(t)) + \sqrt{2D_0}\xi(t) \tag{32}$$

$$\dot{X}_1 = \frac{N}{\tau} (X_0(t) - X_1(t)) \tag{33}$$

$$(34)$$

$$\dot{X}_N = \frac{N}{\pi} (X_{N-1}(t) - X_N(t)) \tag{35}$$

Projection onto dynamics of X_0

$$X_j(t)\Big|_{j>0} = \int_{-\infty}^{t} dt' \frac{N}{\tau} \exp\left(-\frac{N}{\tau}(t-t')\right) X_{j-1}(t)$$
(36)

Iterate

$$X_{j+1}(t) = \int_{-\infty}^{t} dt' \left(\frac{N}{\tau}\right)^{2} (t - t') \exp\left(-\frac{N}{\tau}(t - t')\right) X_{j-1}(t)$$
(37)

$$X_{j+2}(t) = \int_{-\infty}^{t} dt' \frac{1}{2} \left(\frac{N}{\tau}\right)^3 (t - t')^2 \exp\left(-\frac{N}{\tau}(t - t')\right) X_{j-1}(t)$$
(38)

$$X_{j+N-1}(t) = \int_{-\infty}^{t} dt' \frac{1}{(N-1)!} \left(\frac{N}{\tau}\right)^{N} (t-t')^{N-1} \exp\left(-\frac{N}{\tau}(t-t')\right) X_{j-1}(t)$$
(39)

For j = 1

$$X_N(t) = \int_{-\infty}^{t} dt' \frac{1}{(N-1)!} \left(\frac{N}{\tau}\right)^N (t-t')^{N-1} \exp\left(-\frac{N}{\tau}(t-t')\right) X_0(t)$$
(40)

Keep N not only countable, but finite. Let $j \in \{1, 2, ..., N\}$

$$\dot{X}_0 = F(X_0(t), X_N(t)) + \sqrt{2D_0}\xi(t) \tag{41}$$

$$\dot{X}_{j} = \frac{N}{\tau} (X_{j-1}(t) - X_{j}(t)) \tag{42}$$

Fokker-Planck equation for the Markovian system $\,$

$$\rho_{N+1}(x_0, \dots, x_N; t) = \langle \delta(x_0 - X_0(t)) \dots \delta(x_N - X_N(t)) \rangle$$

$$\tag{43}$$

$$\partial_t \rho_{N+1}(x_0, \dots, x_N; t) = \left(-\sum_{j=0}^N \partial_{x_j} f_j(x_j, x_{j-1}) + D_0 \partial_{x_0}^2 \right) \rho_{N+1}(x_0, \dots, x_N; t)$$
(44)

$$f_j(x_j, x_{j-1}) = \begin{cases} F(x_0, x_N), & j = 0\\ \frac{N}{\tau}(x_{j-1} - x_j), & j > 0 \end{cases}$$

$$(45)$$

Natural boundary condition

$$\lim_{x \to +\infty} \rho_{N+1}(x_0, \dots, x_N; t) = 0 \tag{46}$$

Integrating the Fokker-Planck equation over all auxiliary variables

$$\int_{\mathbb{R}} dx_1 \cdots \int_{\mathbb{R}} dx_N \, \partial_t \rho_{N+1}(x_0, \dots, x_N; t) = \int_{\mathbb{R}} dx_1 \cdots \int_{\mathbb{R}} dx_N \left(-\sum_{j=0}^N \partial_{x_j} f_j(x_j, x_{j-1}) + D_0 \partial_{x_0}^2 \right) \rho_{N+1}(x_0, \dots, x_N; t)$$

$$(47)$$

$$\int_{\mathbb{R}} dx_1 \cdots \int_{\mathbb{R}} dx_N \, \rho_{N+1}(x_0, \dots, x_N; t) = \rho_1(x_0; t)$$

$$\tag{48}$$

$$\int_{\mathbb{R}} dx_1 \cdots \int_{\mathbb{R}} dx_N \sum_{j=0}^{N} \partial_{x_j} f_j(x_j, x_{j-1}) \rho_{N+1}(x_0, \dots, x_N; t) = \partial_{x_0} \int_{\mathbb{R}} dx_N F(x_0, x_N) \rho_2(x_0, x_N; t)$$

$$(49)$$

$$\partial_t \rho_1(x_0, t) = -\partial_{x_0} \int_{\mathbb{R}} dx_N F(x_0, x_N) \rho_2(x_0, x_N; t) + D_0 \partial_{x_0}^2 \rho_1(x_0, t)$$
(50)

We take the limit $N \to \infty$ again

$$X_0(t) = X(t), \quad X_N(t) = X(t - \tau)$$
 (51)

This leads to a coordinate transformation

$$\{x_0, t\} \to \{x, t\}, \quad \{x_N, t\} \to \{x_\tau, t - \tau\}$$
 (52)

And therefore to the first Fokker-Planck equation

$$\partial_t \rho_1(x,t) = -\partial_x \int_{\mathbb{R}} dx_\tau F(x_0, x_\tau) \rho_2(x, t; x_\tau, t - \tau) + D_0 \partial_x^2 \rho_1(x, t)$$
(53)

Integrating the Fokker-Planck equation over all auxiliary variables but not x_N

$$\int_{\mathbb{D}} dx_1 \cdots \int_{\mathbb{D}} dx_{N-1} \, \partial_t \rho_{N+1}(x_0, \dots, x_N; t) = \int_{\mathbb{D}} dx_1 \cdots \int_{\mathbb{D}} dx_{N-1} \left(-\sum_{j=0}^N \partial_{x_j} f_j(x_j, x_{j-1}) + D_0 \partial_{x_0}^2 \right) \rho_{N+1}(x_0, \dots, x_N; t)$$

$$(54)$$

$$\int_{\mathbb{R}} dx_1 \cdots \int_{\mathbb{R}} dx_{N-1} \, \rho_{N+1}(x_0, \dots, x_N; t) = \rho_2(x_0, x_N; t)$$
(55)

$$\int_{\mathbb{R}} dx_1 \cdots \int_{\mathbb{R}} dx_{N-1} \sum_{j=0}^{N} \partial_{x_j} f_j(x_j, x_{j-1}) \rho_{N+1}(x_0, \dots, x_N; t) = \partial_{x_0} F(x_0, x_N) \rho_2(x_0, x_N; t) - \frac{N}{\tau} \partial_{x_N} \int_{\mathbb{R}} dx_{N-1} (x_N - x_{N-1}) \rho_3(x_0, x_{N-1}, x_N; t)$$
(56)

$$\partial_t \rho_2(x_0, x_N, t) = -\partial_{x_0} F(x_0, x_N) \rho_2(x_0, x_N; t) + \frac{N}{\tau} \partial_{x_N} \int dx_{N-1} (x_N - x_{N-1}) \rho_3(x_0, x_{N-1}, x_N; t) + D_0 \partial_{x_0}^2 \rho_2(x_0, x_N, t)$$
(57)

Taking again the limit $N \to \infty$ and interpret the auxiliary variables under the coordinate transformation

$$\left\{\frac{N}{\tau}(x_{N-1} - x_N), t\right\} \to \left\{\dot{x}, t - \tau\right\} \tag{58}$$

This forms an alternative hierarchy

$$\partial_t \rho_2(x, t; x_\tau, t - \tau) = -\partial_x F(x, x_\tau) \rho_2(x, t; x_\tau, t - \tau) + \partial_{x_\tau} \int_{\mathbb{R}} d\dot{x}_\tau \, \dot{x}_\tau \rho_3(x, t; \dot{x}_\tau, t - \tau; x_\tau, t - \tau) + D_0 \partial_x^2 \rho_2(x, t; x_\tau, t - \tau)$$

$$(59)$$

Closure via mean field approximation

$$\rho_3(x, t; \dot{x}_{\tau}, t - \tau; x_{\tau}, t - \tau) \sim \rho_2(x, t; x_{\tau}, t - \tau) \rho_1(\dot{x}_{\tau}, t - \tau)$$
(60)

We arrive at a system of 2 Fokker-Planck equations, that are delay partial integro-differential equations.

$$\partial_t \rho_1(x,t) = -\partial_x \int_{\mathbb{R}} dx_\tau F(x_0, x_\tau) \rho_2(x, t; x_\tau, t - \tau) + D_0 \partial_x^2 \rho_1(x, t)$$

$$\tag{61}$$

$$\partial_t \rho_2(x, t; x_\tau, t - \tau) = -\partial_x F(x, x_\tau) \rho_2(x, t; x_\tau, t - \tau) + \partial_{x_\tau} \int_{\mathbb{D}} d\dot{x}_\tau \, \dot{x}_\tau \rho_2(x, t; x_\tau, t - \tau) \rho_1(\dot{x}_\tau, t - \tau) + D_0 \partial_x^2 \rho_2(x, t; x_\tau, t - \tau)$$

$$\tag{62}$$

$$\frac{\partial \rho_1}{\partial t} = -\frac{\partial F_1}{\partial x} \rho_1(x, t) - F_1(x) \frac{\partial \rho_1}{\partial x} - \frac{\partial}{\partial x} \int_{\mathcal{Y}} dx_\tau F_2(x_\tau) \rho_2(x, t; x_\tau, t - \tau) + D_0 \frac{\partial^2 \rho_1}{\partial x^2}$$

$$(63)$$

$$\frac{\partial \rho_2}{\partial t} = -\frac{\partial F_1}{\partial x} \rho_2(x, t; x_\tau, t - \tau) - F_1(x) \frac{\partial \rho_2}{\partial x} - F_2(x_\tau) \frac{\partial \rho_2}{\partial x} + \frac{\partial}{\partial x_\tau} \int_{\mathcal{V}_\tau} d\dot{x}_\tau \, \dot{x}_\tau \rho_2(x, t; x_\tau, t - \tau) \rho_1(\dot{x}_\tau, t - \tau) + D_0 \frac{\partial^2 \rho_1}{\partial x^2}$$

$$(64)$$

$$\frac{\partial \rho_1}{\partial t} = \left(3ax^2 - b - K\right)\rho_1(x, t) + \left(ax^3 - xb - xK\right)\frac{\partial \rho_1}{\partial x} + K\frac{\partial}{\partial x}\int\limits_{\mathcal{V}} dx_\tau \, x_\tau \rho_2(x, t; x_\tau, t - \tau) + D_0\frac{\partial^2 \rho_1}{\partial x^2}$$

$$\tag{65}$$

$$\frac{\partial \rho_2}{\partial t} = \left(3ax^2 - b - K\right)\rho_2(x, t; x_\tau, t - \tau) + \left(ax^3 - xb - xK\right)\frac{\partial \rho_2}{\partial x} + Kx_\tau \frac{\partial \rho_2}{\partial x} + \frac{\partial \rho_2}{\partial x_\tau} \left(\int_{\mathbb{V}_\tau} d\dot{x}_\tau \, \dot{x}_\tau \rho_1(\dot{x}_\tau, t - \tau)\right) + D_0 \frac{\partial^2 \rho_2}{\partial x^2}$$
(66)