Numerical study on 2D Riemann problems using state of the art solvers

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Outline

- Introduction
- Numerical schemes
- Software and computations
- Numerical results

Introduction

Riemann problem in 2D

A Riemann problem consists of a system of conservation equation together with a piecewise constant initial condition.

$$U_{0}(x,y) = \begin{cases} u_{1} & (x,y) \in [0.5,1]^{2} \\ u_{2} & (x,y) \in [0,0.5] \times [0.5,1] \\ u_{3} & (x,y) \in [0,0.5]^{2} \\ u_{4} & (x,y) \in [0.5,1] \times [0,0.5] \end{cases}$$

Introduction

The equations

Euler equation for compressible fluids

$$U_t + F(U)_x + G(U)_y = 0$$

with

$$U = \begin{pmatrix} \rho \\ \rho u \\ \rho v \\ e \end{pmatrix} \qquad F = \begin{pmatrix} \rho u \\ \rho u^2 + p \\ \rho u v \\ u(e+p) \end{pmatrix} \qquad G = \begin{pmatrix} \rho v \\ \rho u v \\ \rho u^2 + p \\ v(e+p) \end{pmatrix}$$

Equation describing the dynamics of a compressible fluid. Conserved quantities are: ρ density, ρu momentum in x, ρv momentum in y and the total energy e.



Numerical methods

5 different numerical schemes tested:

- TVDLF
- ► HLL
- ► HLLC
- TVD-MUSCL
- ► FD

TVDLF - Total Variation Diminishing (TVD) concept

Total variation of numerical approximation of u:

$$TV(u^n) = \sum_{i=0}^{N} |u_{i+1}^n - u_i^n|$$

Scheme is total variation diminishing in time if

$$TV(u^{n+1}) \leq TV(u^n) \ \forall n$$

TVDLF

First order Lax-Friedrichs scheme:

$$U_j^{n+1} = U_j^n - \frac{\Delta t}{\Delta x} \left(F_{j+1/2} - F_{j-1/2} \right) + \frac{1}{2} \left(\Phi_{j+1/2} - \Phi_{j-1/2} \right)$$

where

$$F_{j+1/2} = \frac{F_j + F_{j+1}}{2}$$
$$\Phi_{j+1/2} = U_{j+1} - U_j$$

Scheme is TVD ⇔ CFL condition satisfied

Scheme is first order accurate

Reduce numerical diffusion (with local Courant number, maximal speed):

$$\Phi_{j+1/2} = \frac{\Delta t}{\Delta x} c_{j+1/2}^{max} \left(U_{j+1} - U_j \right)$$



TVDLF - second order spatial accuracy

Linear approximation of U and fluxes at boundary interfaces U at $x_{i+1/2}$, linear interpolation:

$$U_{j+1/2}^L = U_j^n + \frac{1}{2}\bar{\Delta U_j^n}$$

$$U_{j+1/2}^{R} = U_{j+1}^{n} - \frac{1}{2}\bar{\Delta U}_{j+1}^{n}$$

Limited slopes ΔU will be defined later

$$F_{j+1/2} = \frac{F(U_{j+1/2}^L) + F(U_{j+1/2}^R)}{2}$$

$$\Phi_{j+1/2} = rac{\Delta t}{\Delta x} c_{j+1/2}^{ extit{max}} \left(U_{j+1/2}^{ extit{R}} - U_{j+1/2}^{ extit{L}}
ight)$$



TVDLF - slope limiter

Required to ensure TVD property For example *minmod* limiter

$$\bar{\Delta U_j} = minmod(\Delta U_{j-1/2}, \Delta U_{j+1/2})$$

with

$$\Delta U_{j+1/2} = U_{j+1} - U_j$$

and

$$minmod(w_1, w_2, \dots, w_n) =$$

$$sgn(w_1)max[0, min(|w_1|, sgn(w1)w_2, \dots, sgn(w_1)w_n)]$$

The *minmod* function takes the argument with the smallest modulus when all arguments have the same signs and otherwise it is zero.

TVDLF - Temporally second order accuracy

Use Hancock's predictor step:

$$U_{j}^{n+1/2} = U_{j}^{n} - \frac{1}{2} \frac{\Delta t}{\Delta x} \left[F(U_{j}^{n} + \frac{1}{2} \bar{\Delta U}_{j}^{n}) - F(U_{j}^{n} - \frac{1}{2} \bar{\Delta U}_{j}^{n}) \right]$$

used for calculating linear extrapolations:

$$U_{j+1/2}^{L} = U_{j}^{n+1/2} + \frac{1}{2}\bar{\Delta U_{j}^{n}}$$

$$U_{j+1/2}^R = U_{j+1}^{n+1/2} - \frac{1}{2}\bar{\Delta U}_{j+1}^n$$

TVD-MUSCL

MUSCL = Monotonic Upstream Scheme for Conservation Laws

- same Hancock predictor step and upwinding as TVDLF
- upwinding is applied for characteristic variables rather than conservative variables

Characteristic variables \vec{r}^k

- linear combinations of conservative variables
- right eigenvectors of $\partial \vec{F}/\partial \vec{U}$

$$\frac{\partial F}{\partial \vec{U}}\vec{r}^k = c^k \vec{r}^k$$

- ightharpoonup eigenvalues c^k are real
- eigenvectors form a complete orthogonal basis
- ▶ left eigenvectors \vec{l}^k :

$$\vec{l}^k \cdot \vec{r}^k = \delta_{k,m}$$

TVD-MUSCL

Modify $\vec{\Phi}$:

$$\vec{\Phi} = \frac{\Delta t}{\Delta x} \sum_{k} \vec{r}^{k} |c^{k}| \vec{l}^{k} \cdot (\vec{U^{R}} - \vec{U^{L}})$$

where \vec{r}^k , c^k and \vec{l}^k are calculated for $U_{j+1/2}$

- ightarrow $ec{l}^k \cdot (ec{U^R} ec{U^L})$ determines jump in k-th characteristic variable
- multiplication by \vec{r}^k transforms result back to conservation variables

Comparison to TVDLF

- Advantage: use of eigenvalue c^k instead of largest eigenvalue $c^{max} \Rightarrow$ upwinding is accurate for each characteristic variable \Rightarrow less numerical diffusion
- ► Disadvantage: left and right eigenvectors must be calculated for each cell interface (can be very expensive)

Overview of different numerical methods

FD

Finite Differences

The flow variables are given as point-wise values U_j at locations x_j as:

$$U_j(t) = U(x_j, t)$$

Difference formulae of a given order of accuracy can be derived from Taylor expansion around the grid points

FD

Example, 1D, grid spacing Δx :

 \rightarrow a first order spatial derivative can be approximated by the centered finite difference formula:

$$\frac{U_{j+1} - U_{j-1}}{2\Delta x}$$

$$= \frac{1}{2\Delta x} \left[U_j + \Delta x \partial_x U + \frac{(\Delta x)^2}{2!} \partial_{xx} U + \dots \right]$$

$$- \frac{1}{2\Delta x} \left[U_j - \Delta x \partial_x U + \frac{(\Delta x)^2}{2!} \partial_{xx} U + \dots \right]$$

$$= \partial_x U + O((\Delta x)^2)$$

Overview of different numerical methods

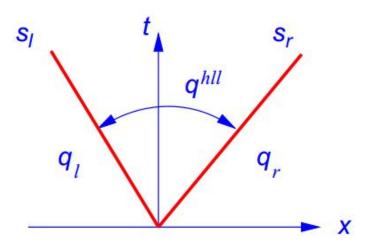
HLL

Riemann solver of Harten, Lax and van Leer System of one-dimensional conservation laws:

$$q_t + f(q)_{\times} = 0$$

with

$$q(x,0) = \begin{cases} q_I & \text{if } x < 0 \\ q_r & \text{if } x > 0 \end{cases}$$





HLL

Following approximation is proposed:

$$\widetilde{q}(x,t) = \begin{cases}
q_l & \text{if } \frac{x}{t} \leq s_l \\
q^{hll} & \text{if } s_l \leq \frac{x}{t} \leq s_r \\
q_r & \text{if } \frac{x}{t} \geq s_r
\end{cases}$$

Corresponding flux:

$$f_{i+1/2}^{hll} = \begin{cases} f_{l} & \text{if } 0 \leq s_{l} \\ \frac{s_{r}f_{l} - s_{l}f_{r} + s_{l}s_{r}(q_{r} - q_{l})}{s_{r} - s_{l}} & \text{if } s_{l} \leq 0 \leq s_{r} \\ f_{r} & \text{if } 0 \geq s_{r} \end{cases}$$

These formulas can be used in the explicit conservative formula:

$$q_i^{n+1} = q_i^n + \frac{\Delta t}{\Delta x} \left[f_{i-1/2} - f_{i+1/2} \right]$$

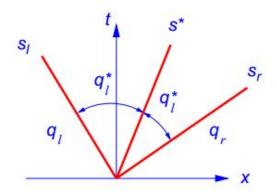
Wave speeds s_l and s_r must be estimated.



Overview of different numerical methods

HLLC

Modification of HLL scheme in which the missing contact and shear waves are restored.



Overview of different numerical methods

HLLC

$$\tilde{q}(x,t) = \begin{cases} q_{l} & \text{if } \frac{x}{t} \leq s_{l} \\ q_{l}^{*} & \text{if } s_{l} \leq \frac{x}{t} \leq s^{*} \\ q_{r}^{*} & \text{if } s^{*} \leq \frac{x}{t} \leq s_{r} \\ q_{r} & \text{if } \frac{x}{t} \geq s_{r} \end{cases}$$

$$f_{l}^{hllc} = \begin{cases} f_{l} & \text{if } 0 \leq s_{l} \\ f_{l} + s_{l}(q_{l}^{*} - q_{l}) & \text{if } s_{l} \leq 0 \leq s^{*} \\ f_{r} + s_{r}(q_{r}^{*} - q_{r}) & \text{if } s^{*} \leq 0 \leq s_{r} \\ f_{r} & \text{if } 0 \geq s_{r} \end{cases}$$

Intermediate states q_I^* and q_r^* can be determined. Wave speeds s_I , s_r and s^* need to be estimated.

MPI-AMRVAC and HPC

Maar heel even aanhalen... ledereen kent dit toch al.



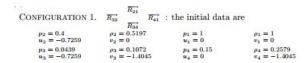
Software and computations

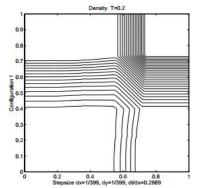
Personal working method

Talk about the scrip we wrote to set up the project and run everything.

Tell them we ran the code succesfully for small problem on the cw network on 5 computers. =¿ network was not stable or something else went wrong.

Configuration 1 - Results from paper





Configuration 1 - Own results

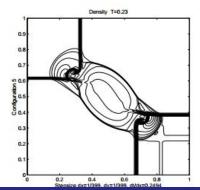
movie



2D Riemann problem

Configuration 5 - Results from paper

CONFIGURATION 5.	J_{32}^-	J_{41}^{-} : the initial data are	
$p_2 = 1$ $u_2 = -0.75$	$\rho_2 = 2$ $v_2 = 0.5$	$p_1 = 1$ $u_1 = -0.75$	$ \rho_1 = 1 \\ v_1 = -0.5 $
$p_3 = 1$ $u_3 = 0.75$	$\rho_3 = 1$ $v_3 = 0.5$	$p_4 = 1$ $u_4 = 0.75$	$ \rho_4 = 3 $ $ v_4 = -0.5 $





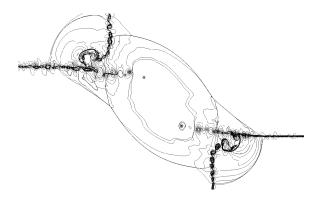
Configuration 5 - Own results

movie



Configuration 5 - Own results

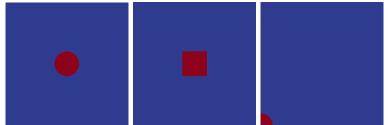
High Resolution



Setup for the explosion

We tried several parameters

- Size of the explosion centre
- Pressure difference
- Location of the explosion



Comparing two explosions

Play the movie



3d rayleigh taylor

