

SIMPLE HARMONIC MOTION

INTRODUCTION

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- 1 Simple Harmonic Motion
- 2 Energy in SHM
- 3 Other Examples
- 4 Rotating Vector picture and Complex phasors

Simple Harmonic Motion

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- Fourier theorem: *“Any periodic motion can be decomposed into harmonic (sine-form) components.”*

Simple Harmonic Motion

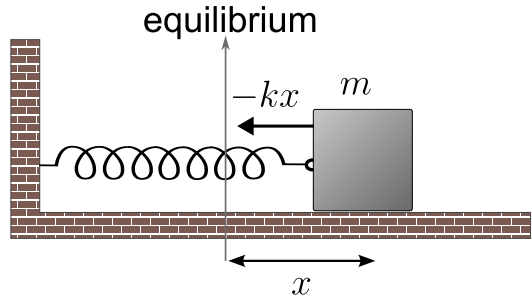
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Importance of SHM:

- Fourier theorem: *“Any periodic motion can be decomposed into harmonic (sine-form) components.”*
- **Small oscillations** about the mean of **ANY** system: nearly harmonic \implies SHM is important!

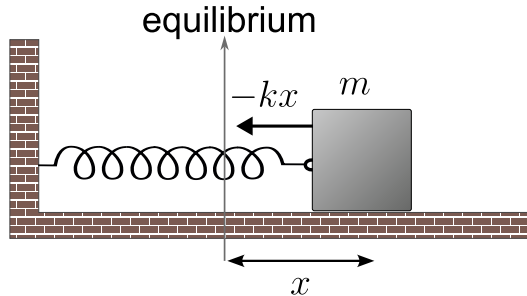
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Restoring force = $-kx$

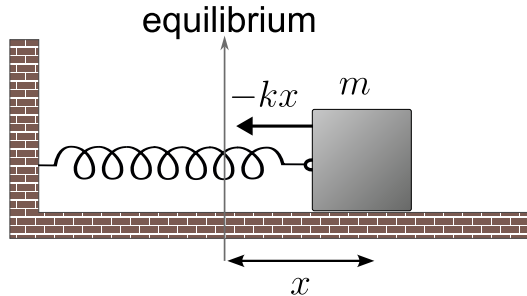


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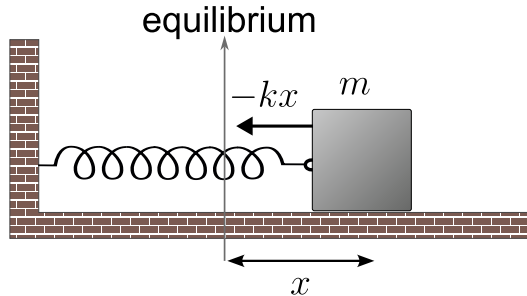


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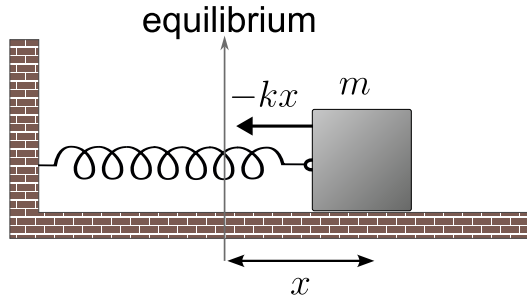
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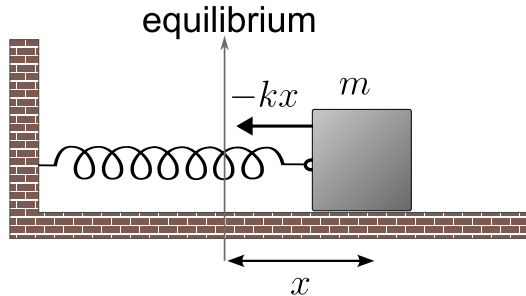
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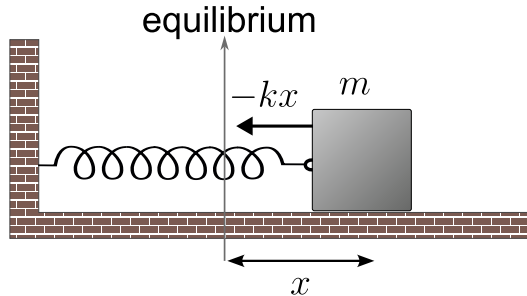
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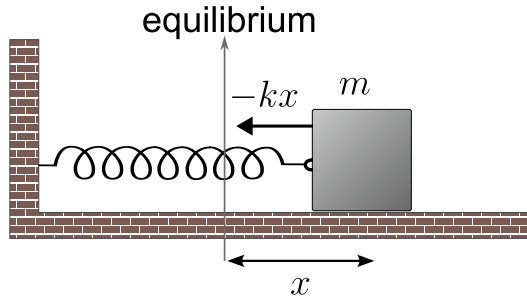
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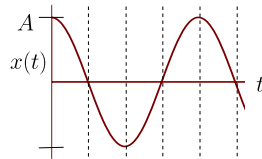


A : amplitude

ϕ : initial phase.

Characteristics of SHM

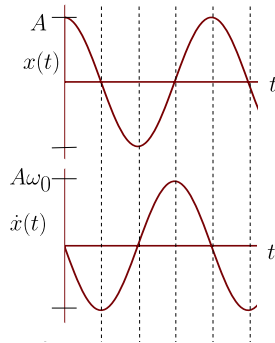
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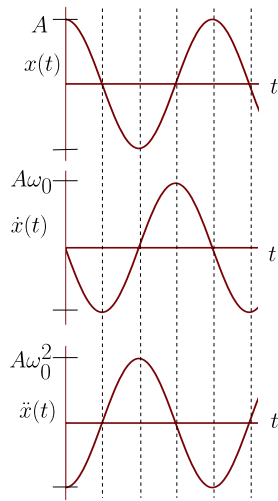


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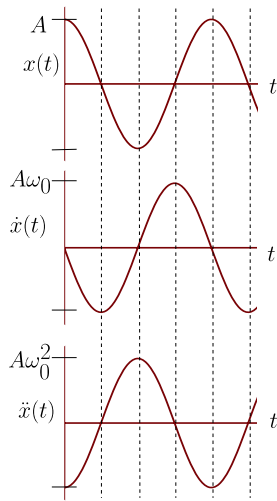
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Converse:

$m\ddot{x} = -m\omega_0^2 x$, **restoring** force.



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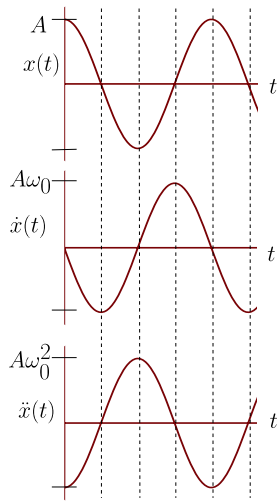
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$m\ddot{x} = -m\omega_0^2 x$, **restoring** force. $m\omega_0^2 = k$:
stiffness or spring constant.



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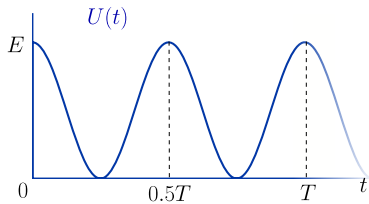
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$$\text{Phase } \tan(\phi) = \frac{v_0}{\omega_0 x_0}.$$

Energy in SHM

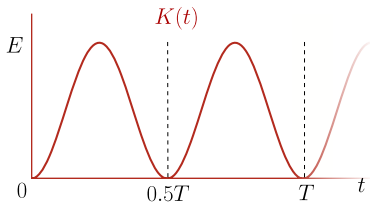
Energy in SHM



Potential Energy

$$U(t) = \frac{1}{2}kx^2$$

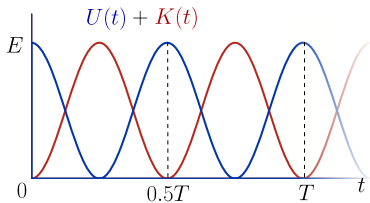
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Kinetic Energy

$$K(t) = \frac{1}{2}m\dot{x}^2$$

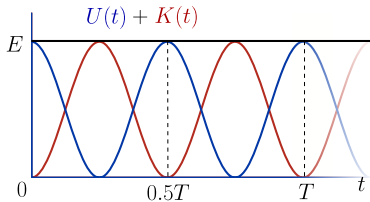
Energy in SHM



Total Energy

$$U(t) + K(t) = \frac{1}{2}kx^2 + \frac{1}{2}m\dot{x}^2 = E$$

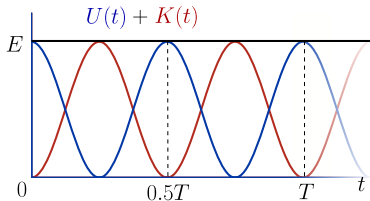
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Total Energy decided by the initial conditions:

$$\begin{aligned} U(t) + K(t) &= \frac{1}{2}kx^2 + \frac{1}{2}m\dot{x}^2 = E \\ &= \frac{1}{2}mv_0^2 + \frac{1}{2}kx_0^2 = \text{const.} \end{aligned}$$

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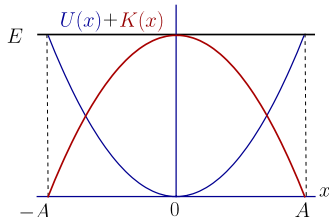
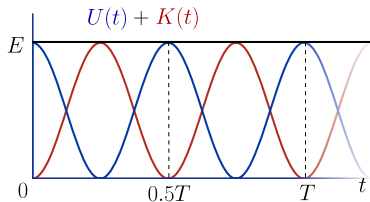


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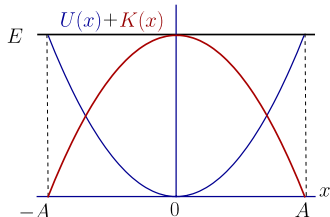
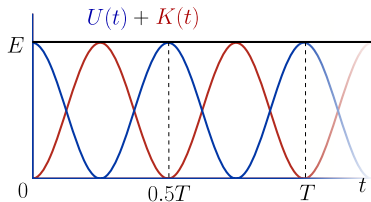
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- 1 $U_{\max} = \frac{1}{2}kA^2 = K_{\max}.$
- 2 $U_{\min} = 0$ at $x = 0$ (equilibrium)

For ANY system with a potential minimum...

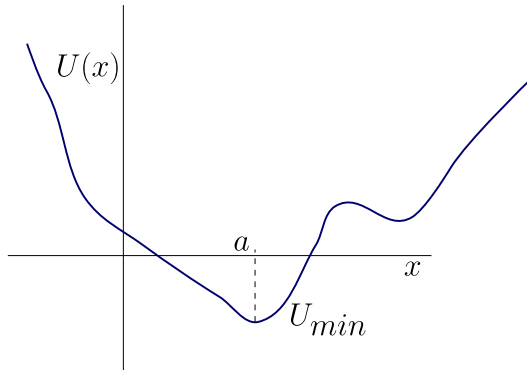
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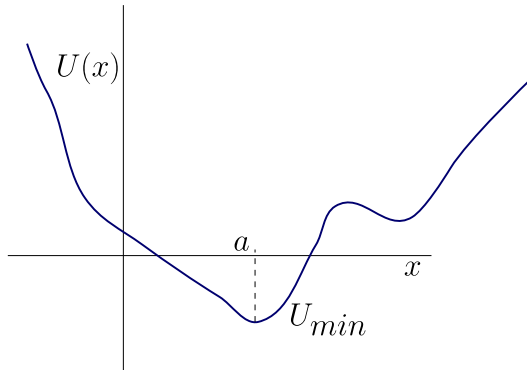
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$$U(x) \approx U(a) + xU'(a) + \frac{x^2}{2}U''(a) + \dots$$



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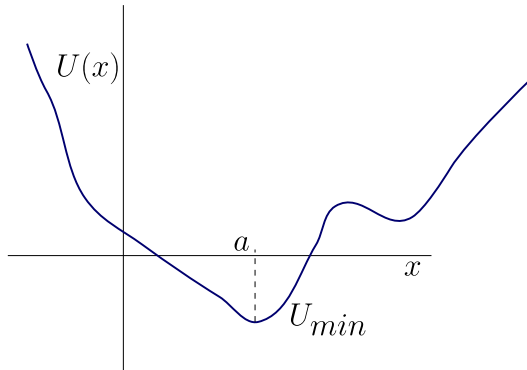
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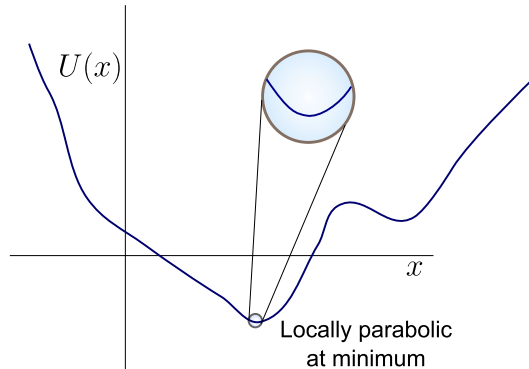
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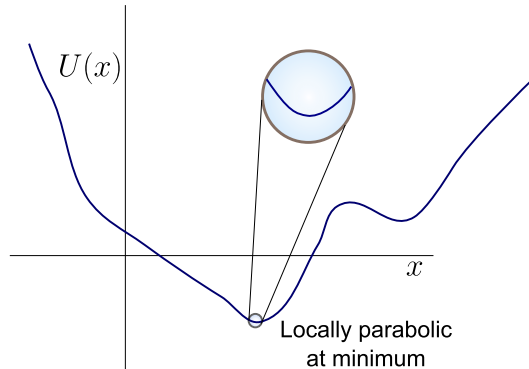
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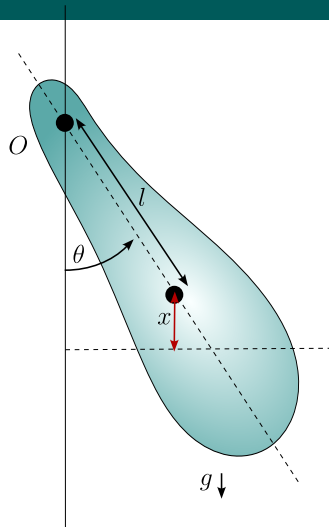
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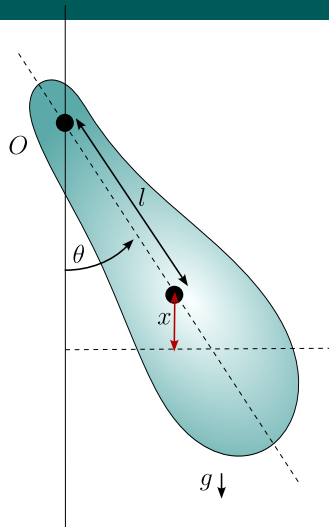
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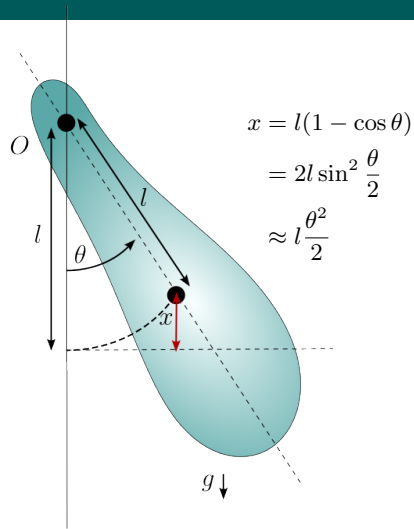
$$U(x) = mgx;$$



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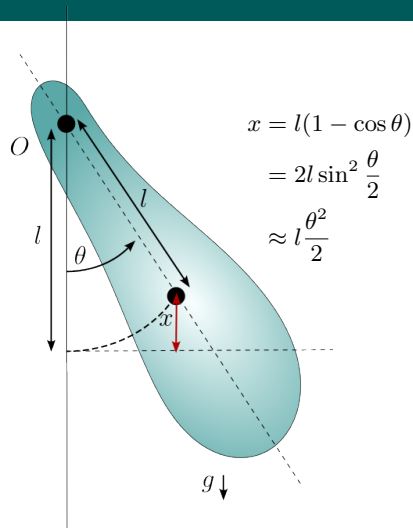
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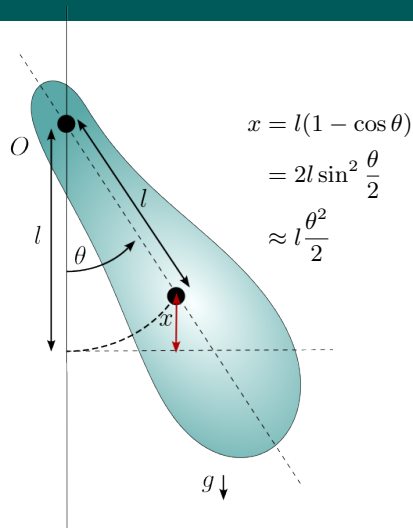
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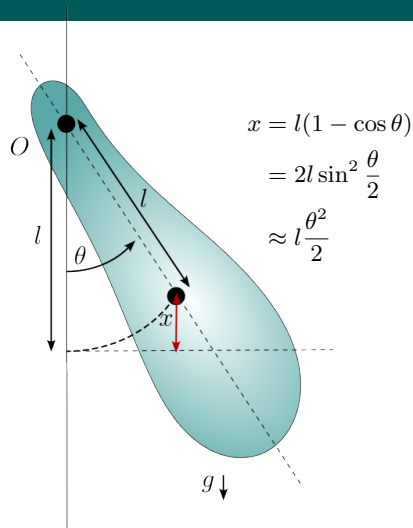
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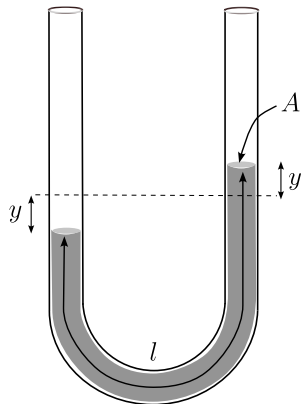
$$\therefore \omega_0^2 = \frac{mgl}{I}.$$



Non-viscous liquid in a U-tube

Displacement of level: y

Level oscillates about mean $y = 0$.

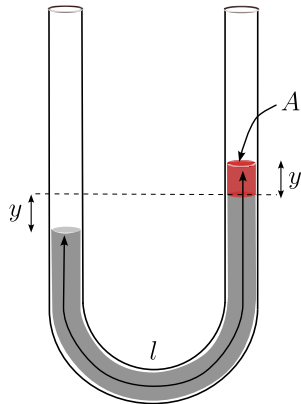


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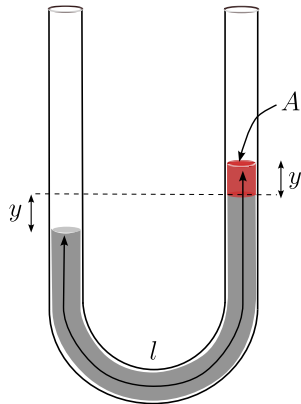
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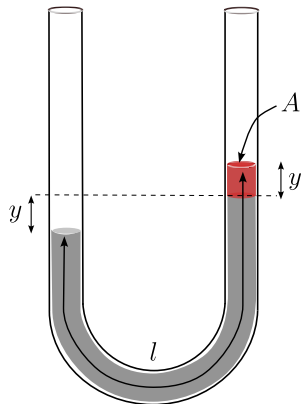
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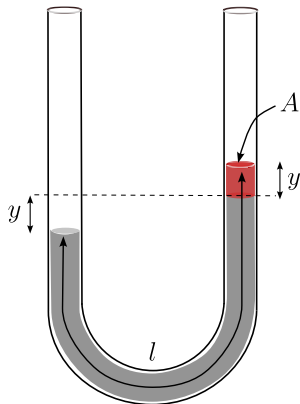
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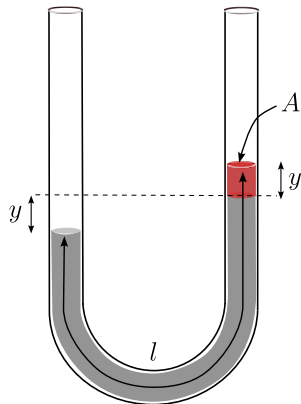
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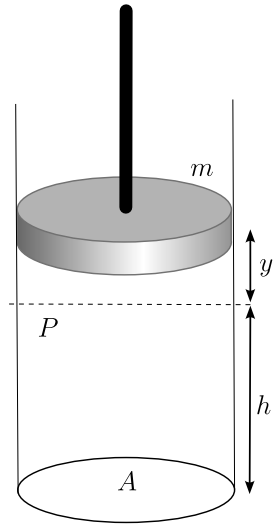
$$\therefore \omega_0^2 = \frac{2g}{l}.$$



“Air spring”

Insulated container of gas at pressure P .

Piston of mass m oscillates about $y = 0$.

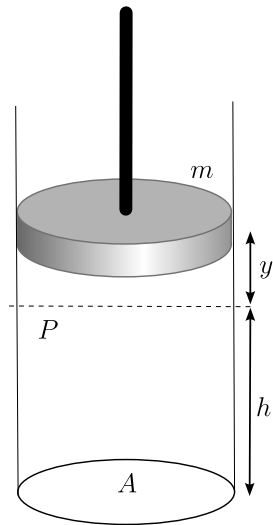


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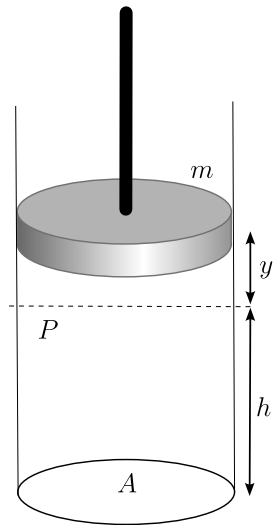
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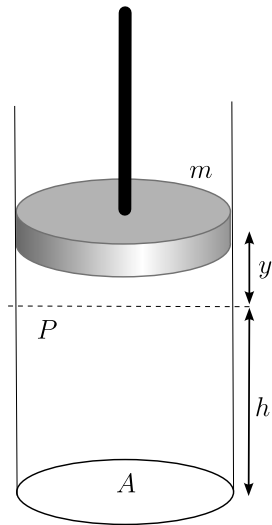
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Adiabatic change: $PV^\gamma = \text{const.}$

$$\left(\gamma = \frac{C_P}{C_V}\right)$$



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Insulated container of gas at pressure P .

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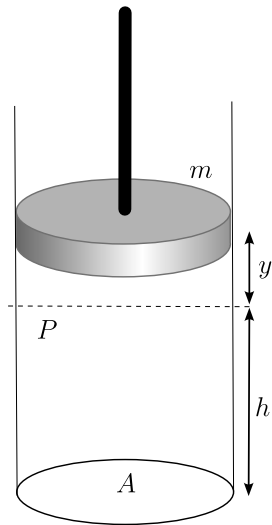
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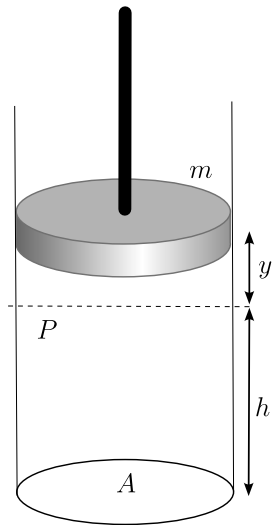
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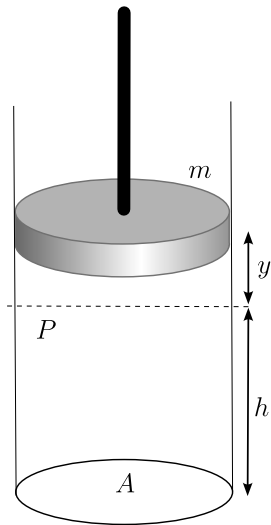
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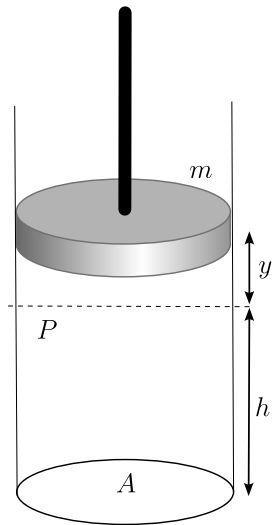
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Use: geometric calculation of superpositions of SHMs, solutions of eqns involving SHMs

SHM and Complex Phasors

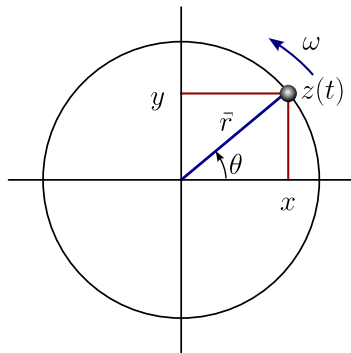
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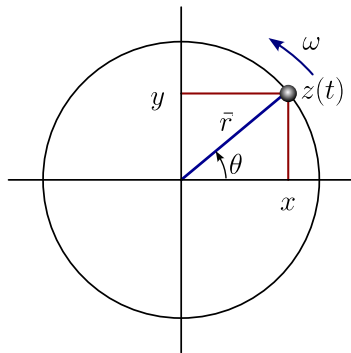
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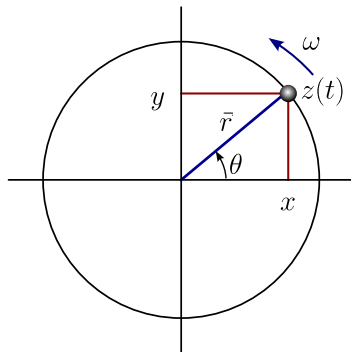


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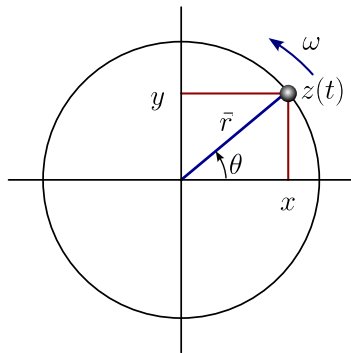
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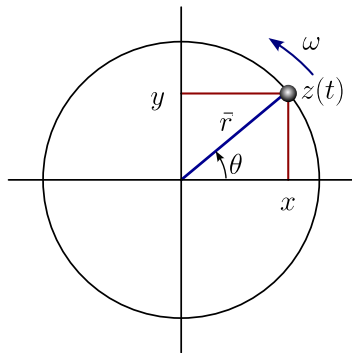
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Examples: $-1 = e^{i\pi}, \quad i = e^{i\pi/2}, \quad i^i = e^{-\pi/2}$.



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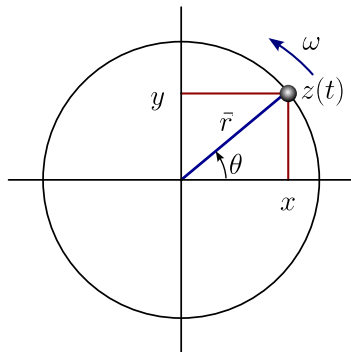
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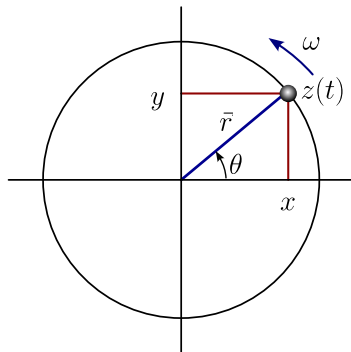
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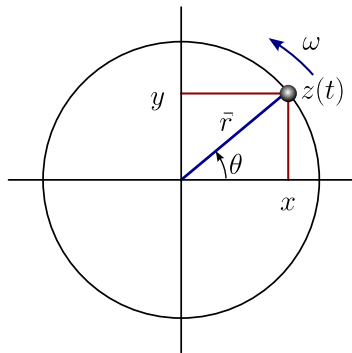
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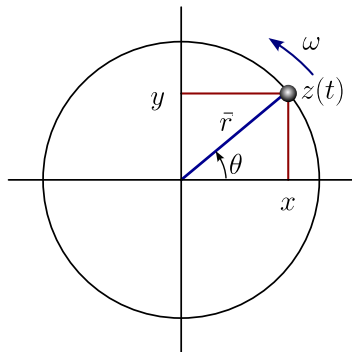
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Use: Superposition of SHMs, solution to SHM differential eqs.

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Exercise: Find C and ϕ in terms of \mathcal{A} and \mathcal{B} .