

# **Constraints**

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# Differential Constraints

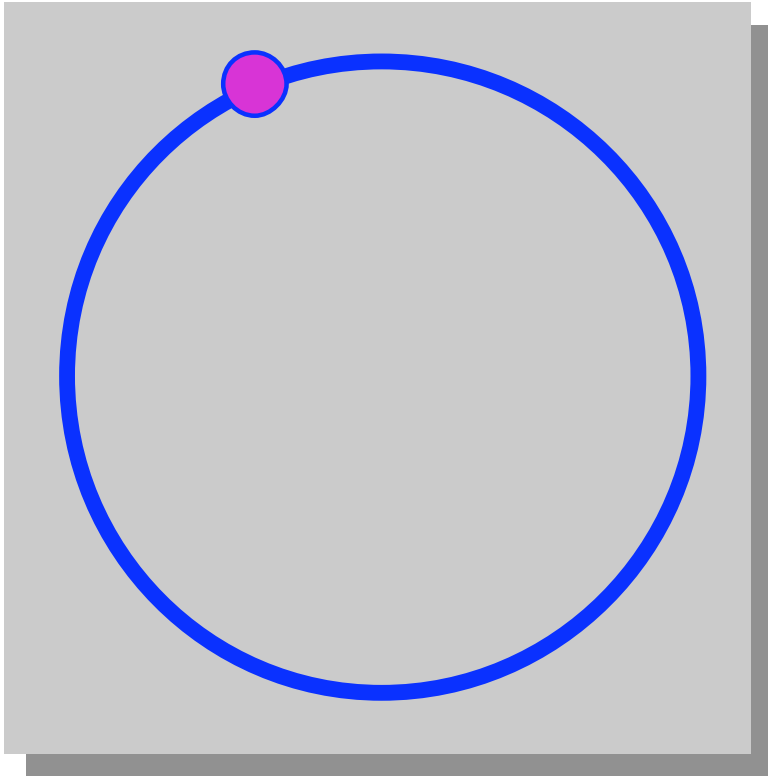
**thanks to Adrew Witkin and Zoran Popović**

# Differential Constraints

# Beyond Points and Springs

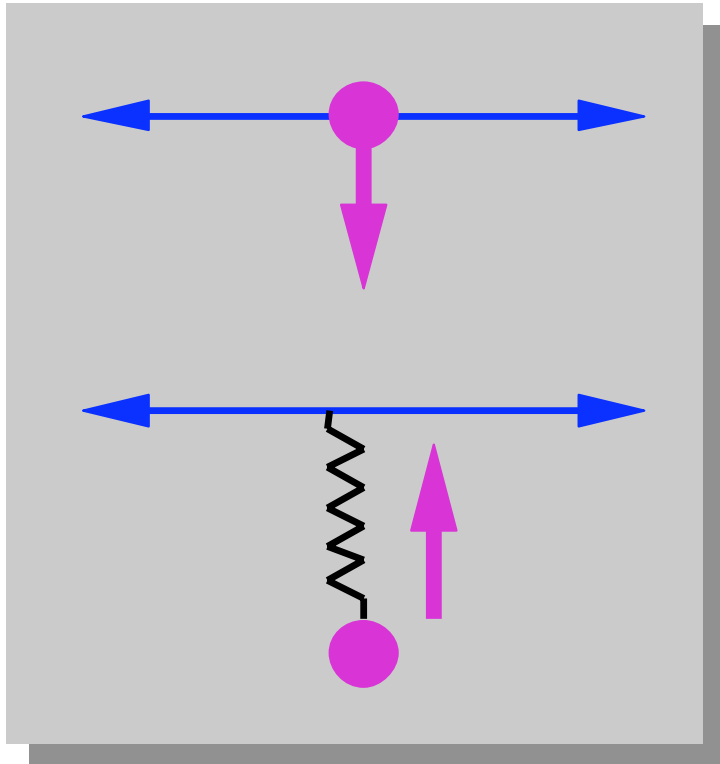
- You can make just about anything out of point masses and springs, *in principle*

# A bead on a wire



- **Desired Behavior:**
  - The bead can slide freely *along* the circle
  - It can never come off, however hard we pull
- **Question:**
  - How does the bead move under applied forces?

# Penalty Constraints

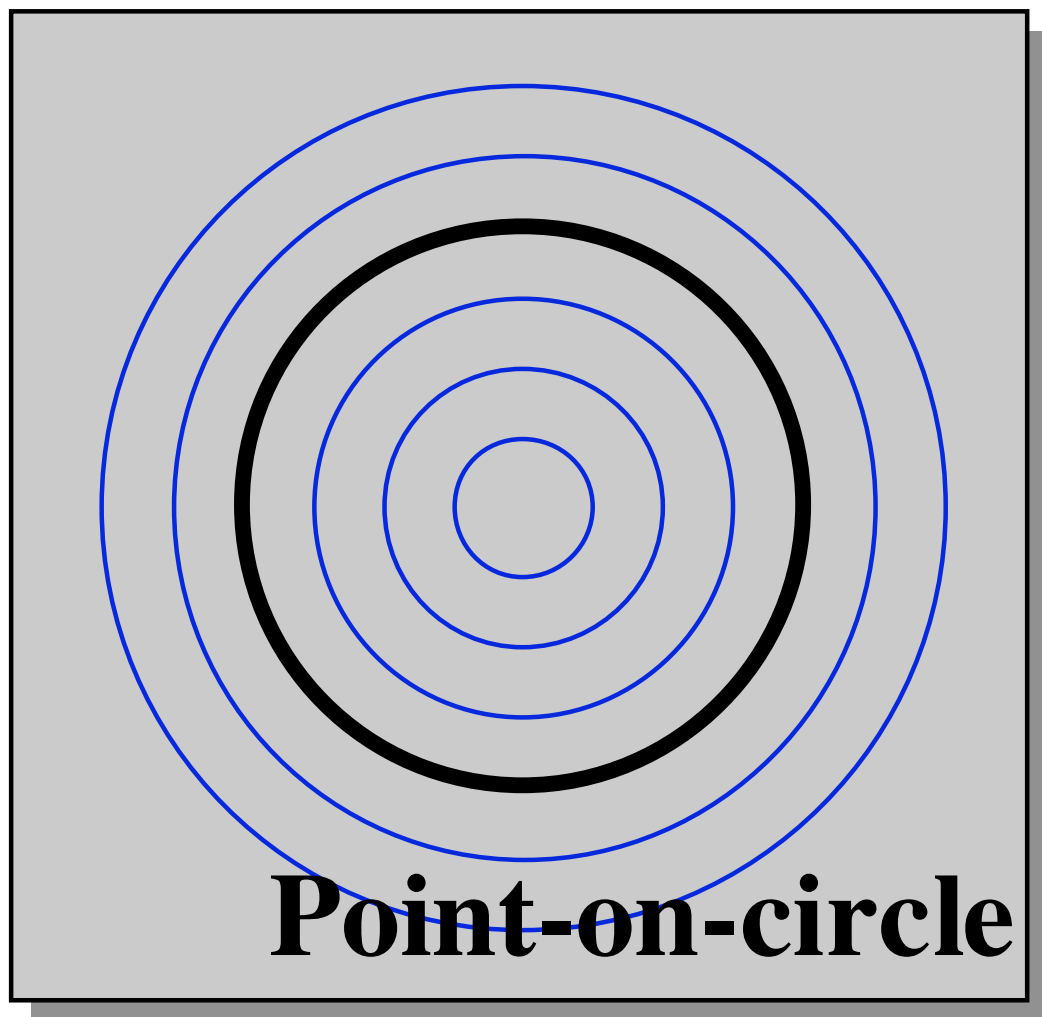


- **Why not use a spring to hold the bead on the wire?**
- **Problem:**
  - **Weak springs  $\Rightarrow$  goopy constraints**
  - **Strong springs  $\Rightarrow$  neptune express!**
- **A classic *stiff system***

## Now for the Algebra ...

- **Fortunately, there's a general recipe for calculating the constraint force**
- **First, a single constrained particle**
- **Then, generalize to constrained particle systems**

# Representing Constraints



*I. Implicit:*

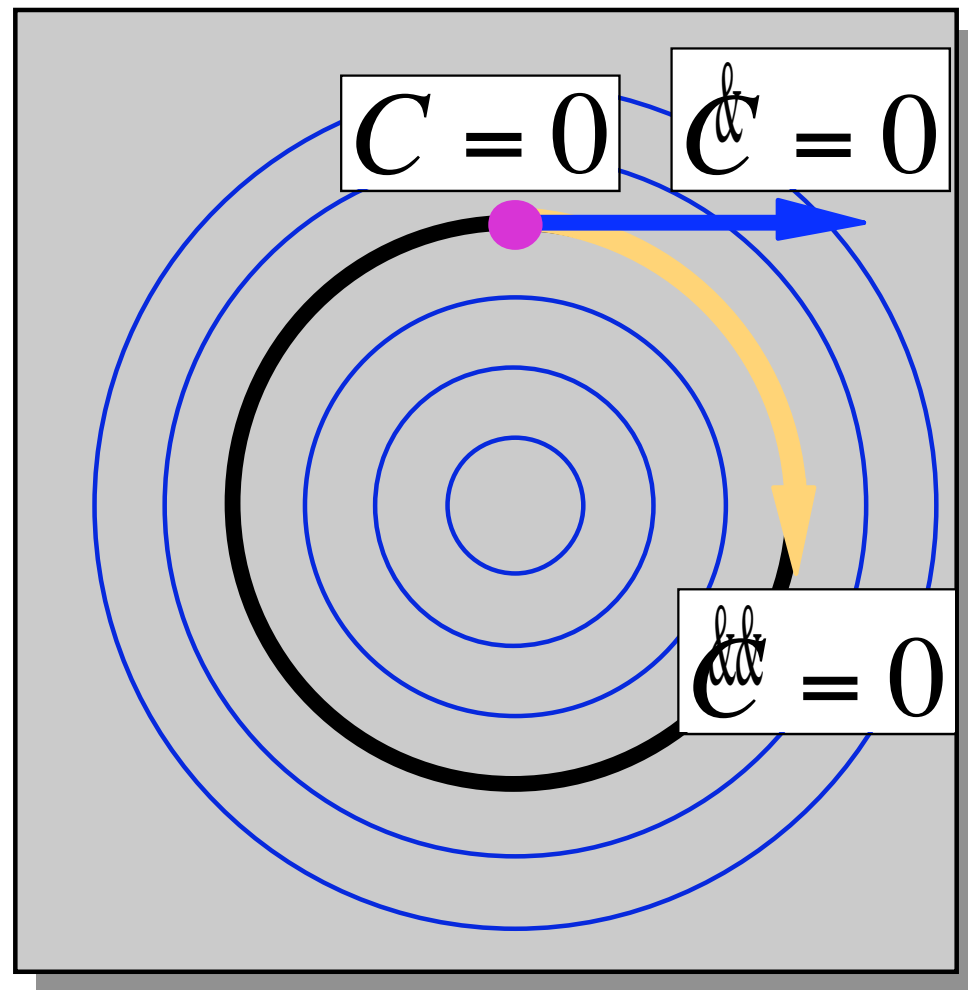
$$C(\mathbf{x}) = |\mathbf{x}| - r = 0$$

*II. Parametric:*

$$\mathbf{x} = r [\cos \theta, \sin \theta]$$



# Maintaining Constraints Differentially



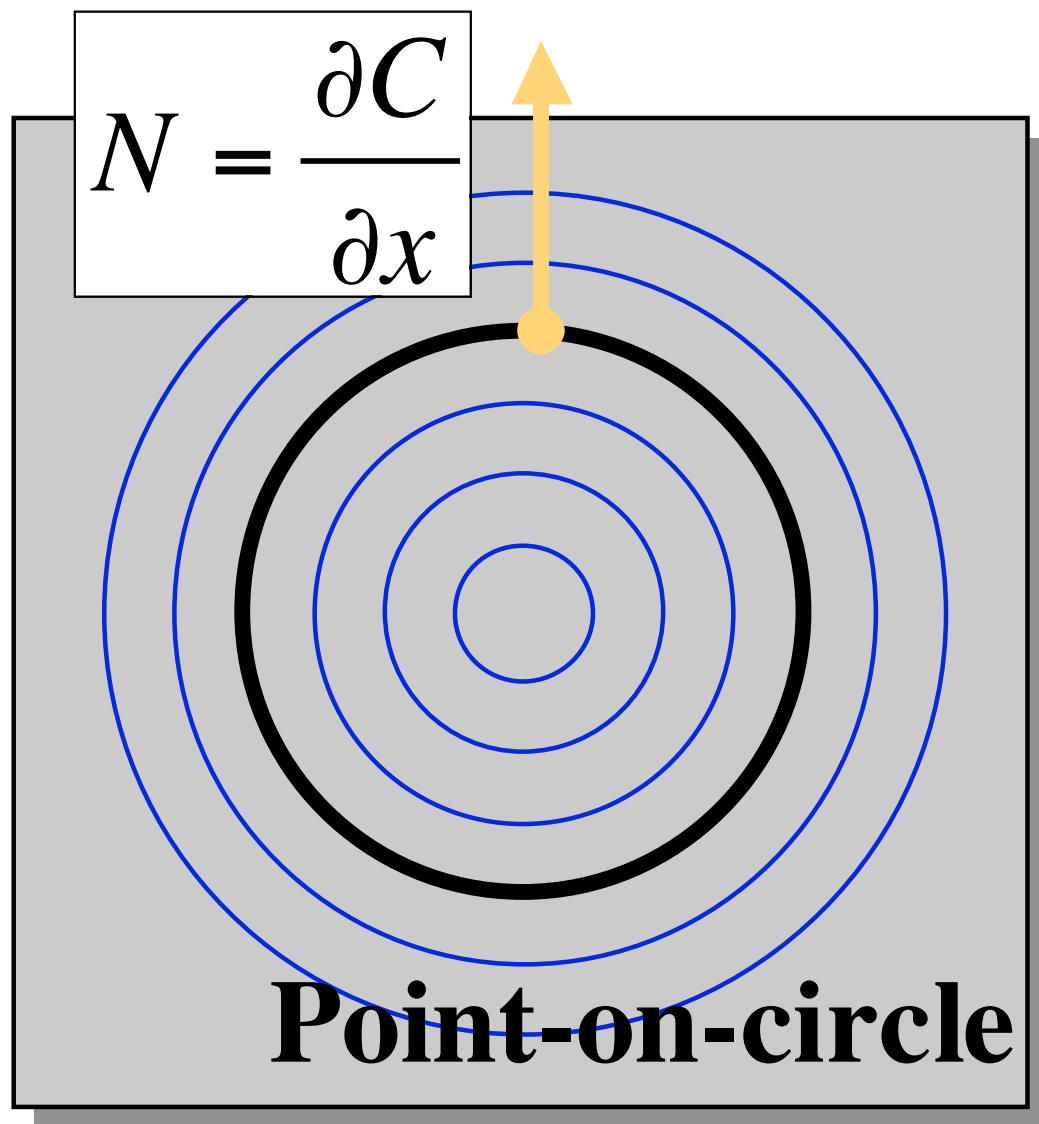
- **Start with legal position and velocity.**
- **Use constraint forces to ensure legal curvature.**

$C = 0$     legal position

$\dot{C} = 0$     legal velocity

$\ddot{C} = 0$     legal curvature

# Constraint Gradient



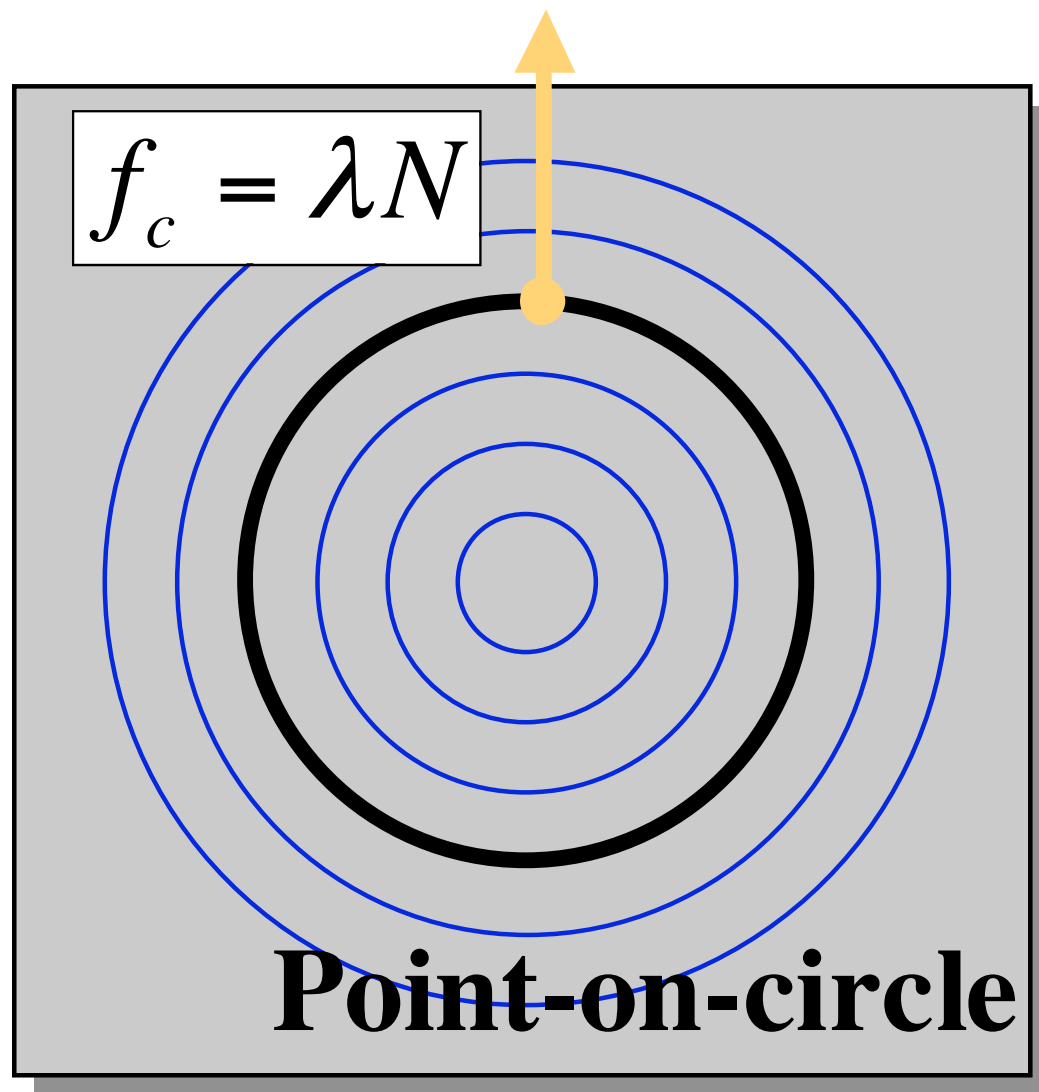
*Implicit:*

$$C(\mathbf{x}) = |\mathbf{x}| - r = 0$$

**Differentiating  $C$  gives a normal vector.**

**This is the direction our constraint force will point in.**

# Constraint Forces



**Constraint force: gradient vector times a scalar  $\lambda$**

**Just one unknown to solve for**

**Assumption: constraint is passive—no energy gain or loss**

# Constraint Force Derivation

$$C(x(t))$$

$$\dot{C} = N \cdot \dot{x}$$

$$\ddot{C} = \frac{\partial}{\partial t}(N \cdot \dot{x})$$

$$= \cancel{N} \cdot \dot{x} + N \cdot \ddot{x}$$

$$\ddot{x} = \frac{f + f_c}{m}$$

$$f_c = \lambda N$$

Set  $\ddot{C} = 0$ , solve for  $\lambda$ :

$$\lambda = -m \frac{\cancel{N} \cdot \dot{x}}{N \cdot N} - \frac{N \cdot f}{N \cdot N}$$

Constraint force is  $\lambda N$ .

$$\text{Notation: } N = \frac{\partial C}{\partial x}, \cancel{N} = \frac{\partial^2 C}{\partial x \partial t}$$

# Example: Point-on-circle

$$C = |\mathbf{x}| - r$$

$$\mathbf{N} = \frac{\partial C}{\partial \mathbf{x}} = \frac{\mathbf{x}}{|\mathbf{x}|}$$

$$\dot{\mathbf{N}} = \frac{\partial^2 C}{\partial \mathbf{x} \partial t} = \frac{1}{|\mathbf{x}|} \left[ \dot{\mathbf{x}} - \frac{\mathbf{x} \cdot \dot{\mathbf{x}}}{\mathbf{x} \cdot \mathbf{x}} \mathbf{x} \right]$$

Write down the constraint equation.

Take the derivatives.

Substitute into generic template, simplify.

$$\lambda = -m \frac{\dot{\mathbf{N}} \cdot \dot{\mathbf{x}}}{\mathbf{N} \cdot \mathbf{N}} - \frac{\mathbf{N} \cdot \mathbf{f}}{\mathbf{N} \cdot \mathbf{N}} = \left[ m \frac{(\mathbf{x} \cdot \dot{\mathbf{x}})^2}{\mathbf{x} \cdot \mathbf{x}} - m(\dot{\mathbf{x}} \cdot \dot{\mathbf{x}}) - \mathbf{x} \cdot \mathbf{f} \right] \frac{1}{|\mathbf{x}|}$$

# Tinkertoys

- **Now we know how to simulate a bead on a wire.**
- **Next: a constrained particle *system*.**
  - **E.g. constrain particle/particle distance to make rigid links.**
- **Same idea, but...**

# Compact Particle System Notation

$$\ddot{\mathbf{q}} = \mathbf{W}\mathbf{Q}$$

**q:** *3n-long state vector.*

**Q:** *3n-long force vector.*

**M:** *3n x 3n diagonal mass matrix.*

**W:** *M-inverse (element-wise reciprocal)*

$$\begin{aligned}\mathbf{q} &= [\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_n] \\ \mathbf{Q} &= [\mathbf{f}_1, \mathbf{f}_2, \dots, \mathbf{f}_n] \\ \mathbf{M} &= \begin{bmatrix} m_1 & & & & \\ & m_1 & & & \\ & & m_1 & & \\ & & & \ddots & \\ & & & & m_n \\ & & & & & m_n \\ & & & & & & m_n \end{bmatrix} \\ \mathbf{W} &= \mathbf{M}^{-1}\end{aligned}$$

$$\begin{aligned}\mathbf{C} &= [C_1, C_2, \dots, C_m] \\ \lambda &= [\lambda_1, \lambda_2, \dots, \lambda_m] \\ \mathbf{J} &= \frac{\partial \mathbf{C}}{\partial \mathbf{q}} \\ \dot{\mathbf{J}} &= \frac{\partial^2 \mathbf{C}}{\partial \mathbf{q} \partial t}\end{aligned}$$



# Solving for the Constraint Force

$$\ddot{x} = \frac{1}{m}(f + \hat{f})$$

$$C(x) = x \cdot x - 1 = 0$$

$$\dot{C}(x) = 2x \cdot \dot{x} = 0$$

$$\ddot{C}(x) = 2(x \cdot \ddot{x} + \dot{x} \cdot \dot{x})$$

$$\dot{x} \cdot \dot{x} + x \cdot \left( \frac{1}{m}(f + \hat{f}) \right) = 0$$

$$x \cdot \hat{f} = -m \dot{x} \cdot \dot{x} - x \cdot f$$

Bead on a Wire

$$\ddot{q} = W(Q + \hat{Q})$$

$$C(q) = 0$$

$$\dot{C} = \frac{\partial C}{\partial q} \dot{q} = J \dot{q}$$

$$\ddot{C} = \dot{J} \dot{q} + J \ddot{q}$$

$$\dot{J} \dot{q} + J W(Q + \hat{Q}) = 0$$

$$J W \hat{Q} = -\dot{J} \dot{q} - J W Q$$

General Case



# Force Must be a Linear Combination of Constraint Gradients

$$T = \frac{1}{2} M \dot{x} \cdot \dot{x}$$

$$\dot{T} = M \dot{x} \cdot \ddot{x} = 0$$



$$\dot{x} \cdot \hat{f} = 0$$

$$\hat{f} = \lambda \vec{1}$$

$$T = \frac{1}{2} \dot{q}^T M \dot{q}$$

$$\dot{T} = \dot{q}^T M \ddot{q} = 0$$

$$= \dot{q}^T M W \hat{Q}$$

$$= \dot{q}^T \hat{Q} = 0$$

$$\Rightarrow \hat{Q} = J^T \lambda$$

Bead on a Wire

General Case



## Final Solution for the $\lambda$ Multipliers

$$\lambda \dot{x} \cdot \dot{x} = -m \dot{x} \cdot \ddot{x} - x \cdot f$$
$$\lambda = \frac{-m \dot{x} \cdot \ddot{x} - x \cdot f}{\dot{x} \cdot \dot{x}}$$

**Bead on a Wire**

$$\lambda = (J^T W J)^{-1} (-J^T x - J^T W Q)$$

**General Case**



# Particle System Constraint Equations

Matrix equation for  $\lambda$

$$[\mathbf{J}\mathbf{W}\mathbf{J}^T]\lambda = -\dot{\mathbf{J}}\dot{\mathbf{q}} - [\mathbf{J}\mathbf{W}]\mathbf{Q}$$

Constrained Acceleration

$$\ddot{\mathbf{q}} = \mathbf{W}[\mathbf{Q} + \mathbf{J}^T\lambda]$$

Derivation: just like bead-on-wire.

More Notation

$$\mathbf{C} = [C_1, C_2, \dots, C_m]$$

$$\lambda = [\lambda_1, \lambda_2, \dots, \lambda_m]$$

$$\mathbf{J} = \frac{\partial \mathbf{C}}{\partial \mathbf{q}}$$

$$\dot{\mathbf{J}} = \frac{\partial^2 \mathbf{C}}{\partial \mathbf{q} \partial t}$$

# Drift and Feedback

- In principle, clamping  $\ddot{C}$  at zero is enough
- Two problems:
  - Constraints might not be met initially
  - Numerical errors can accumulate
- A feedback term handles both problems:

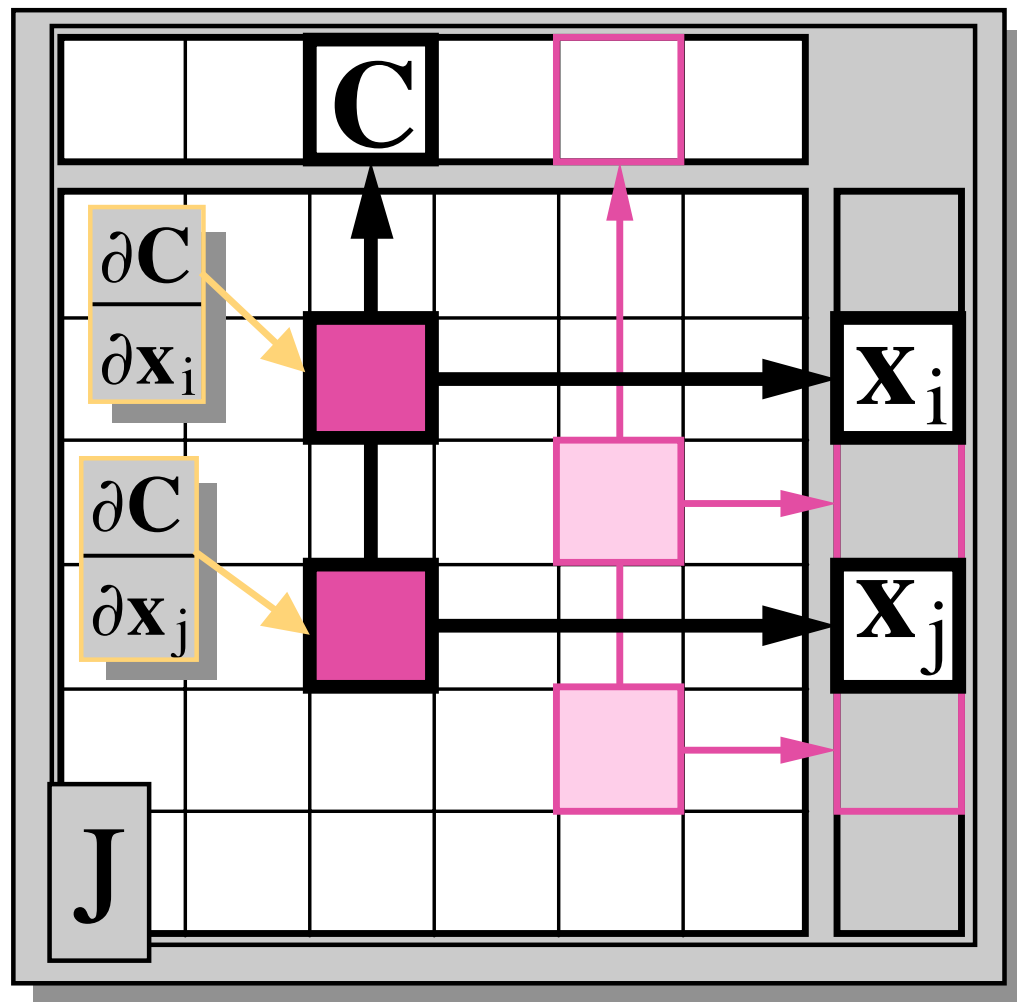
$$\ddot{C} = -\alpha C - \beta \dot{C}, \text{ instead of } \ddot{C} = 0$$

$\alpha$  and  $\beta$  are magic constants.

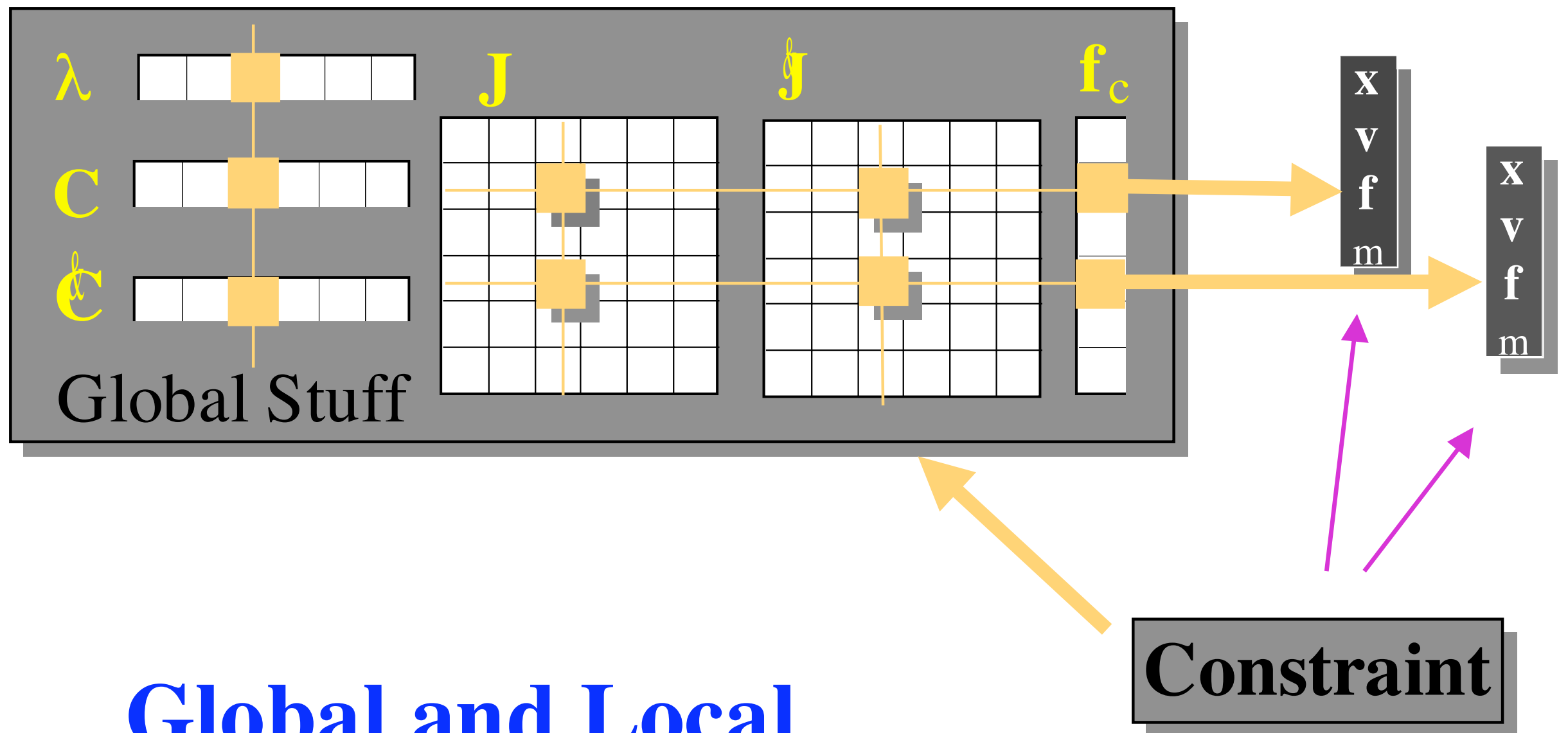
# How do you implement all this?

- **We have a global matrix equation.**
- **We want to build models on the fly, just like masses and springs.**
- **Approach:**
  - **Each constraint adds its own piece to the equation.**

# Matrix Block Structure

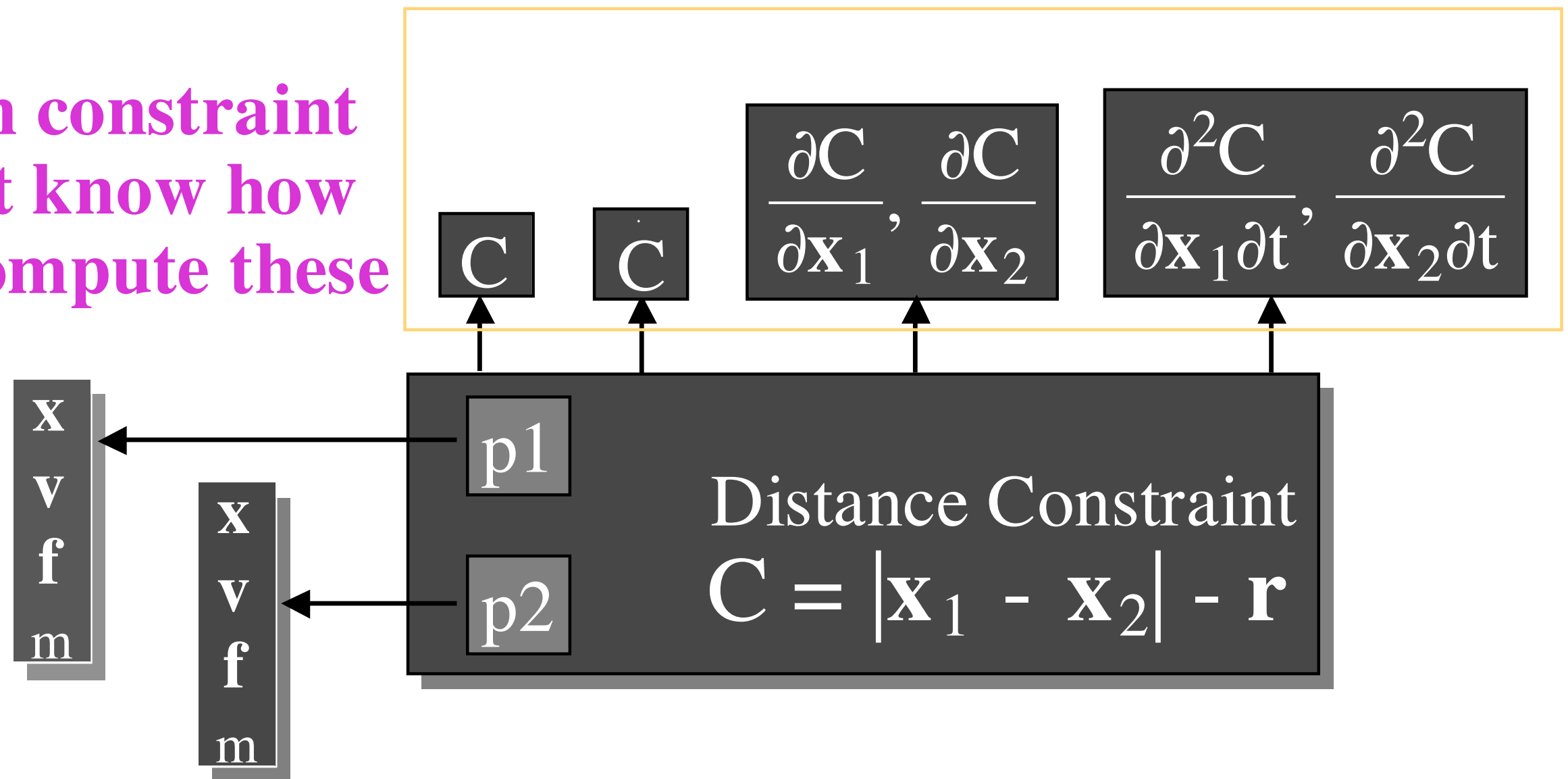


- Each constraint contributes one or more *blocks* to the matrix.
- Sparsity: many empty blocks.
- Modularity: let each constraint compute its own blocks.
- Constraint and particle indices determine block locations.



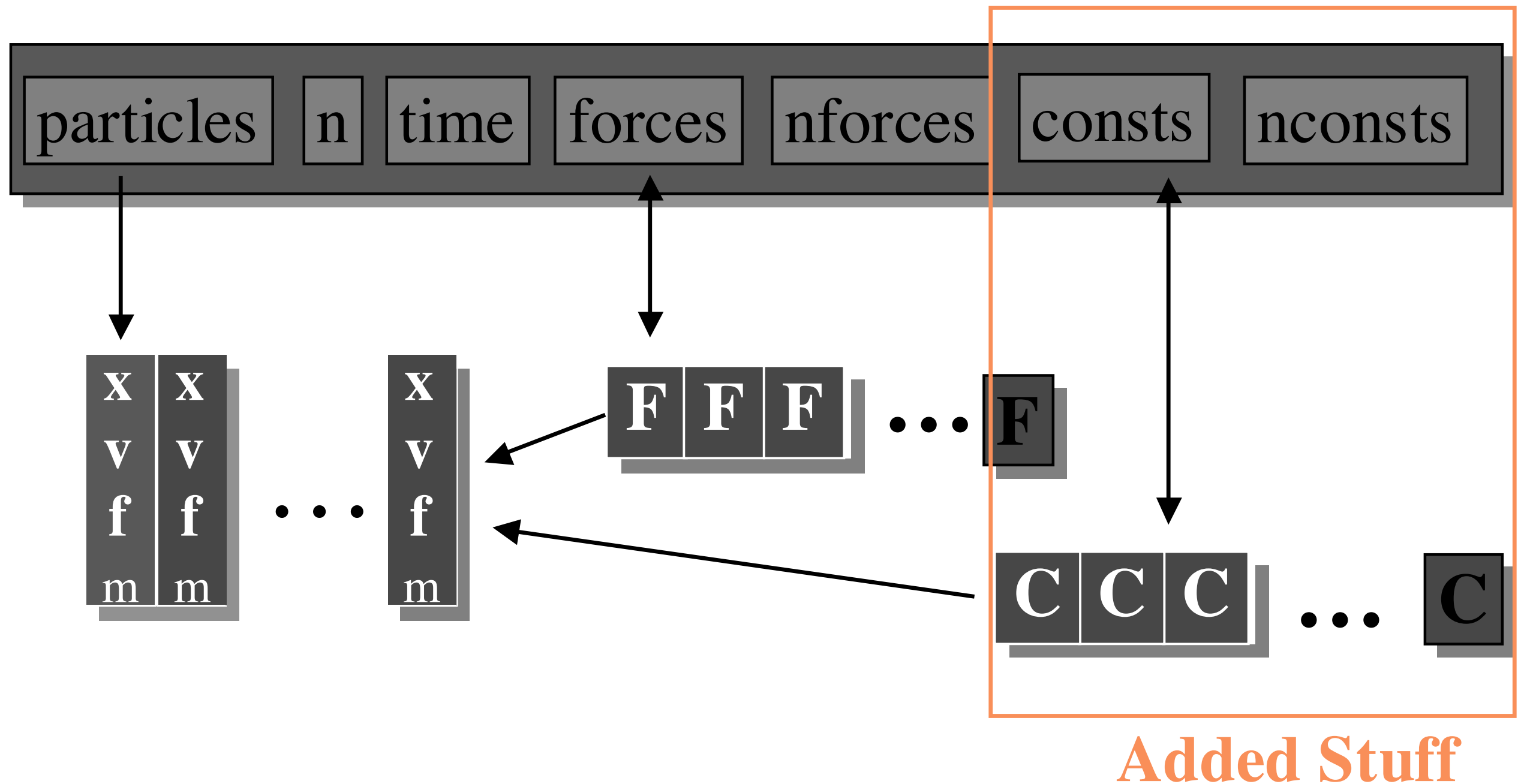
# Constraint Structure

Each constraint must know how to compute these

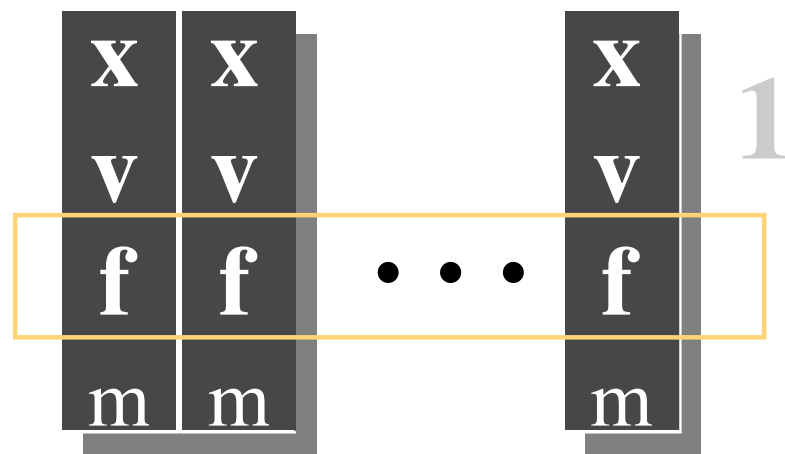




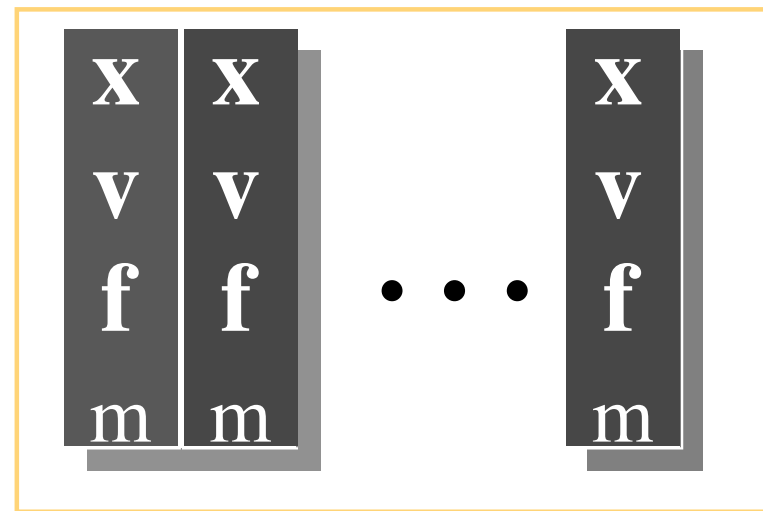
# Constrained Particle Systems



# Modified Deriv Eval Loop



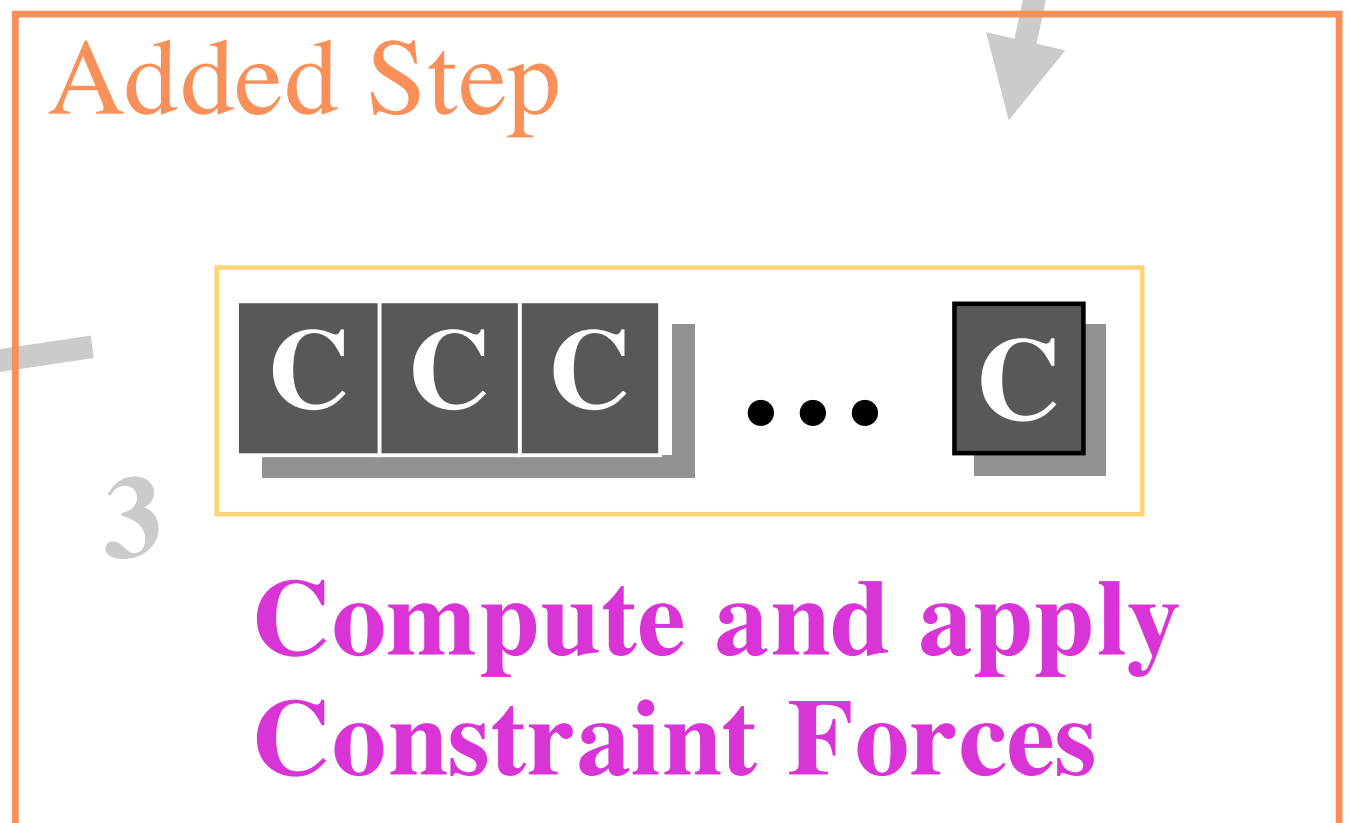
Clear Force  
Accumulators



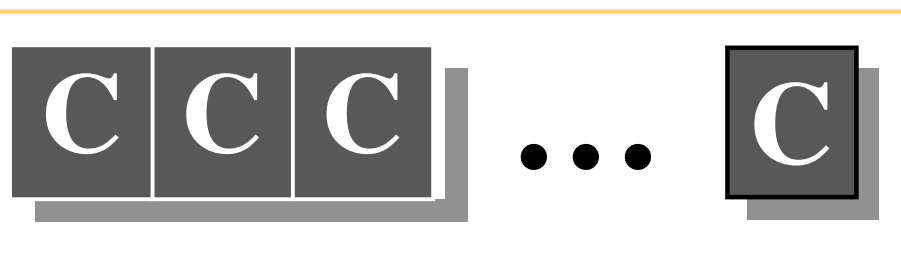
Return to solver



Apply forces



Added Step



Compute and apply  
Constraint Forces

# Constraint Force Eval

- **After computing ordinary forces:**
  - **Loop over constraints, assemble global matrices and vectors.**
  - **Call matrix solver to get  $\lambda$ , multiply by  $J^T$  to get constraint force.**
  - **Add constraint force to particle force accumulators.**

# Impress your Friends

- The requirement that constraints not add or remove energy is called the *Principle of Virtual Work*.
- The  $\lambda$ 's are called *Lagrange Multipliers*.
- The derivative matrix, J, is called the *Jacobian Matrix*.