

Theory Meets Experiments 2025:

Direct Detection across Dark Matter mass ranges

Wave Dark Matter Experiments

Chiara P. Salemi
(UC Berkeley and LBNL)

Me + disclaimers

My research:

- ABRACADABRA → DMRadio
- BREAD
- Qubit-based sensors for axions, low-mass particle DM, CEvNS



UC Berkeley



Disclaimers:

- Content is subject to my bias
- Few references included – please ask me if you would like more references on any topic
- Several topics/images reproduced from others'. Thank you to many, including: Karl van Bibber, Julia Vogel, Bianca Giaccone, Axel Lindner, Lindley Winslow, Ben Safdi, Kevin Zhou

Expectations

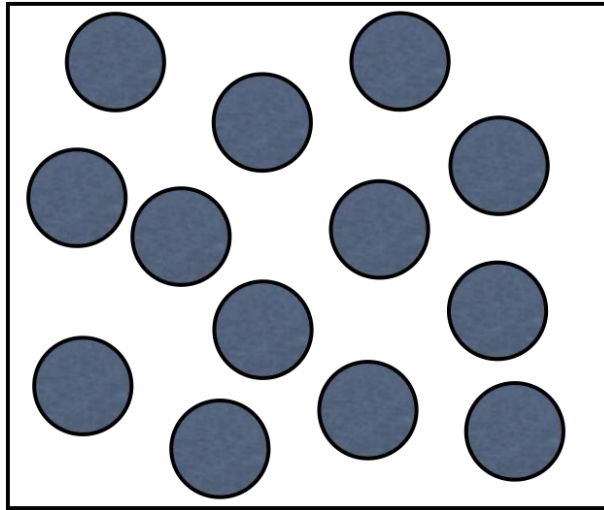
- Please ask lots of questions!
- I'm happy to have the discussion be guided by people's interests (within reason)
 - If there is a topic that sparks a lot of interest, we can add it to the last lecture to go more in-depth

Lecture 1: axion-photon experiments overview

- Wave dark matter detection
- Axion interactions with photons
- Sources of axions and their implications for experiments

Wave vs particle dark matter

- **A definition for wave DM:**
 - de Broglie wavelength $>$ average particle separation
 - Equivalently, occupation number > 1



Particle dark matter



Wave dark matter

Wave vs particle dark matter

- **A definition for wave DM:**
 - de Broglie wavelength $>$ average particle separation
 - Equivalently, occupation number > 1

de Broglie Wavelength - $\lambda_{dB} \approx \frac{2\pi}{mv}$

Occupancy Number - $N \approx \frac{\rho_{DM}}{m} \lambda_{dB}^3$

Wave vs particle dark matter

- **A definition for wave DM:**

- de Broglie wavelength $>$ average particle separation
- Equivalently, occupation number > 1

de Broglie Wavelength - $\lambda_{dB} \approx \frac{2\pi}{mv}$ **Occupancy Number** - $N \approx \frac{\rho_{DM}}{m} \lambda_{dB}^3$

- **Axion** ($m \sim 10^{-9}$ eV): $\lambda_{dB} \sim 10^4$ km with $N \sim 10^{44}$
- **WIMP** ($m \sim 100$ GeV): $\lambda_{dB} \sim 10^{-16}$ km with $N \sim 10^{-36}$

where $\rho_{DM} = 0.4 \text{ GeV/cm}^3$

Wave vs particle dark matter

- **A definition for wave DM:**

- de Broglie wavelength $>$ average particle separation
- Equivalently, occupation number > 1

de Broglie Wavelength - $\lambda_{dB} \approx \frac{2\pi}{mv}$

Occupancy Number - $N \approx \frac{\rho_{DM}}{m} \lambda_{dB}^3$

- Axion ($m \sim 10^{-9}$ eV): $\lambda_{dB} \sim 10^4$ km with $N \sim 10^{44}$

- WIMP ($m \sim 100$ GeV): $\lambda_{dB} \sim 10^{-16}$ km with

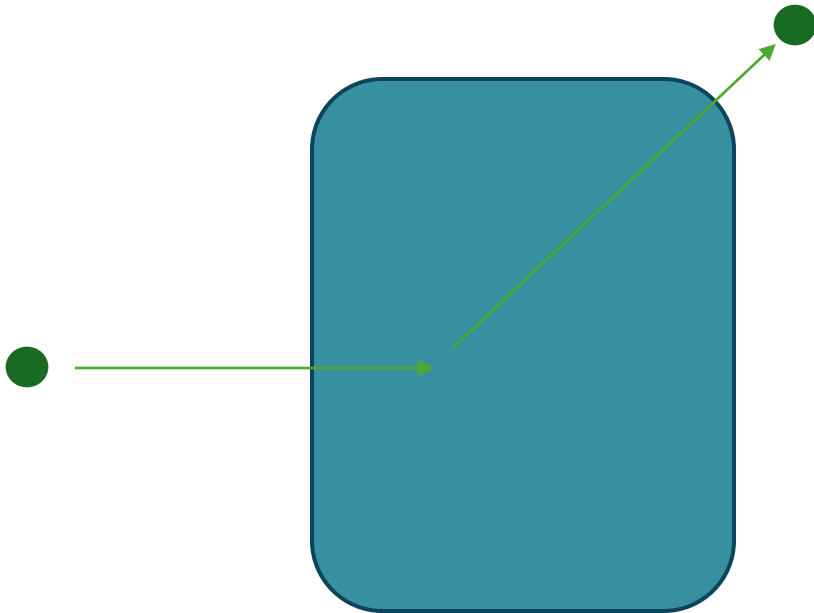
where $\rho_{DM} = 0.4 \text{ GeV/cm}^3$

What does this mean for how we build our DM detectors? Compare to detectors from last week!

What does this mean for detection?

- Coherent and very high occupation number: can describe dark matter in terms of classical waves
- In a detector, we should think of the collective interactions of the field of DM instead of the individual particle interactions

Not a normal scattering experiment



Not a normal scattering experiment



Not a normal scattering experiment

- Too light to use standard particle detectors




Not a normal scattering experiment

- Too light to use standard particle detectors
- But very numerous!



Wave DM candidates

- QCD axions 
- Axion-like particles (ALPs)
- Dark photons

Wave DM candidates

- QCD axions ★
- Axion-like particles (ALPs)
- Dark photons

} same way to detect, just different coupling strengths



QCD axions: coupling dependent on m_a

ALPs: coupling independent of m_a

Wave DM candidates

- QCD axions ★
 - Axion-like particles (ALPs)
 - Dark photons
- } same way to detect, just different coupling strengths
- } many axion detectors also have dark photon sensitivity, but not all!

Axion interactions

$$\mathcal{L} = \underbrace{-\frac{1}{4}g_{a\gamma\gamma}aF_{\mu\nu}\tilde{F}^{\mu\nu}}_{\text{coupling to photons}} - \underbrace{\frac{i}{2}g_d a \bar{N} \sigma_{\mu\nu} \gamma_5 N F_{\mu\nu} + g_{aNN}(\partial_\mu) \bar{N} \gamma^\mu \gamma_5 N}_{\text{coupling to nucleons}} + \underbrace{g_{aee}(\partial_\mu) \bar{e} \gamma^\mu \gamma_5 e}_{\text{coupling to electrons}}$$

All interactions are feeble! ($g \propto 1/f_a$)

Axion interactions

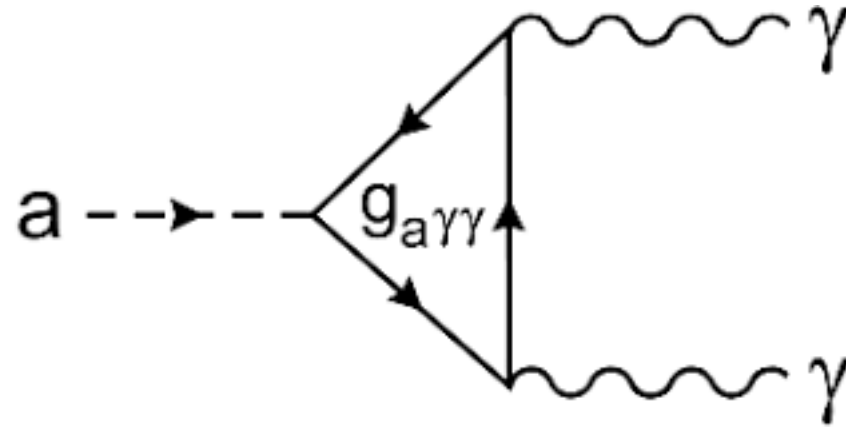
$$\mathcal{L} = \underbrace{-\frac{1}{4}g_{a\gamma\gamma}aF_{\mu\nu}\tilde{F}^{\mu\nu}}_{\text{coupling to photons}} - \underbrace{\frac{i}{2}g_d a \bar{N}\sigma_{\mu\nu}\gamma_5 N F_{\mu\nu} + g_{aNN}(\partial_\mu)\bar{N}\gamma^\mu\gamma_5 N}_{\text{coupling to nucleons}} + \underbrace{g_{aee}(\partial_\mu)\bar{e}\gamma^\mu\gamma_5 e}_{\text{coupling to electrons}}$$

most experiments

All interactions are feeble! ($g \propto 1/f_a$)

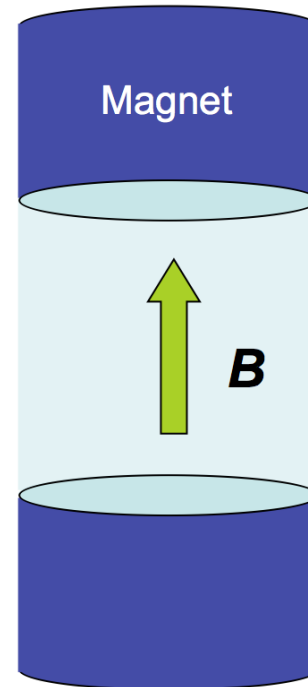
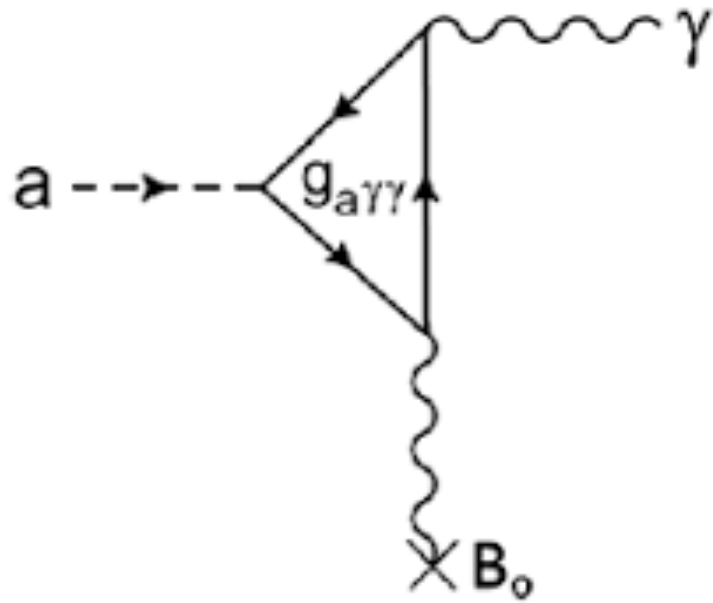
Primakoff effect

Axion couples to 2 photons through loop

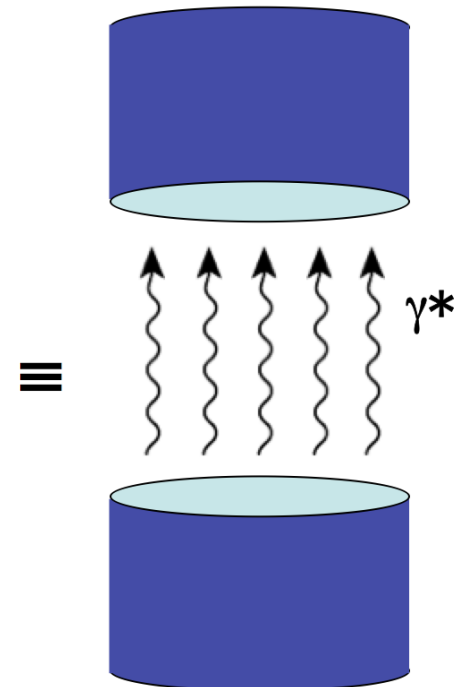


Primakoff effect

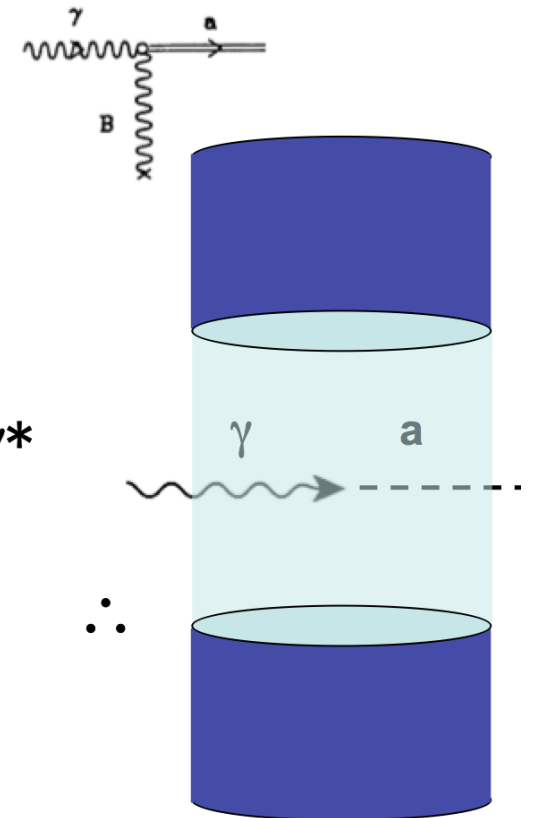
Often drawn with one virtual photon from B field



Classical EM field



Sea of virtual photons



Primakoff Effect

Modified E&M

Modified QED Lagrangian

$$\mathcal{L}_{QED} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}g_{a\gamma\gamma}aF_{\mu\nu}\tilde{F}^{\mu\nu}$$

Modified E&M

Modified QED Lagrangian

$$\mathcal{L}_{QED} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - g_{a\gamma\gamma}a\mathbf{E} \cdot \mathbf{B}$$

Modified E&M

Modified QED Lagrangian

$$\mathcal{L}_{QED} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - g_{a\gamma\gamma}a\mathbf{E} \cdot \mathbf{B}$$

Modified Source-Free Maxwell Equations

$$\nabla \cdot \mathbf{E} = -g_{a\gamma\gamma}\mathbf{B} \cdot \nabla a$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} - g_{a\gamma\gamma}\left(\mathbf{E} \times \nabla a - \frac{\partial a}{\partial t}\mathbf{B}\right)$$

Modified E&M

Modified QED Lagrangian

$$\mathcal{L}_{QED} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - g_{a\gamma\gamma}a\mathbf{E} \cdot \mathbf{B}$$

Modified Source-Free Maxwell Equations

$$\nabla \cdot \mathbf{E} = -g_{a\gamma\gamma}\mathbf{B} \cdot \nabla a$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} - g_{a\gamma\gamma}\left(\mathbf{E} \times \nabla a - \frac{\partial a}{\partial t}\mathbf{B}\right)$$

$$g_{a\gamma\gamma} \propto \frac{1}{f_a} \ll 1$$

Modified E&M

Modified QED Lagrangian

$$\mathcal{L}_{QED} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - g_{a\gamma\gamma}a\mathbf{E} \cdot \mathbf{B}$$

Modified Source-Free Maxwell Equations

$$\nabla \cdot \mathbf{E} = -g_{a\gamma\gamma}\mathbf{B} \cdot \nabla a$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} - g_{a\gamma\gamma}\left(\mathbf{E} \times \nabla a - \frac{\partial a}{\partial t}\mathbf{B}\right)$$

$$g_{a\gamma\gamma} \propto \frac{1}{f_a} \ll 1$$

Exploiting Ampère's Law

$$\nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} - g_{a\gamma\gamma} (\cancel{\mathbf{E} \times \nabla a} - \frac{\partial a}{\partial t} \mathbf{B})$$

Exploiting Ampère's Law

$$\nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} - g_{a\gamma\gamma} (\cancel{\mathbf{E} \times \nabla a} - \frac{\partial a}{\partial t} \mathbf{B})$$

$$\mathbf{J}_{eff} = g_{a\gamma\gamma} \frac{\partial a}{\partial t} \mathbf{B}$$

- Time-varying axion field acts as an effective current (or displacement current) in the presence of a magnetic field

Exploiting Ampère's Law

$$\nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} - g_{a\gamma\gamma} (\mathbf{E} \times \nabla a - \frac{\partial a}{\partial t} \mathbf{B})$$

$$\begin{aligned} \mathbf{J}_{eff} &= g_{a\gamma\gamma} \frac{\partial a}{\partial t} \mathbf{B} \\ &= g_{a\gamma\gamma} \sqrt{2\rho_{DM}} \cos(m_a t) \mathbf{B} \end{aligned}$$

$$a(t) = \frac{\sqrt{2\rho_{DM}}}{m_a} \sin(m_a t)$$

- Time-varying axion field acts as an effective current (or displacement current) in the presence of a magnetic field
- Frequency is the mass of the axion

Axion-photon experiment ingredients

1. Strong magnetic field* (axion-to-photon converter)

$$\mathbf{J}_{\text{eff}} = g_{a\gamma\gamma} \sqrt{2\rho_{DM}} \cos(m_a t) \mathbf{B}$$

*can use DC field or RF field

Axion-photon experiment ingredients

1. Strong magnetic field (axion-to-photon converter)
2. Coupling system (couple photons/EM signal into detector)

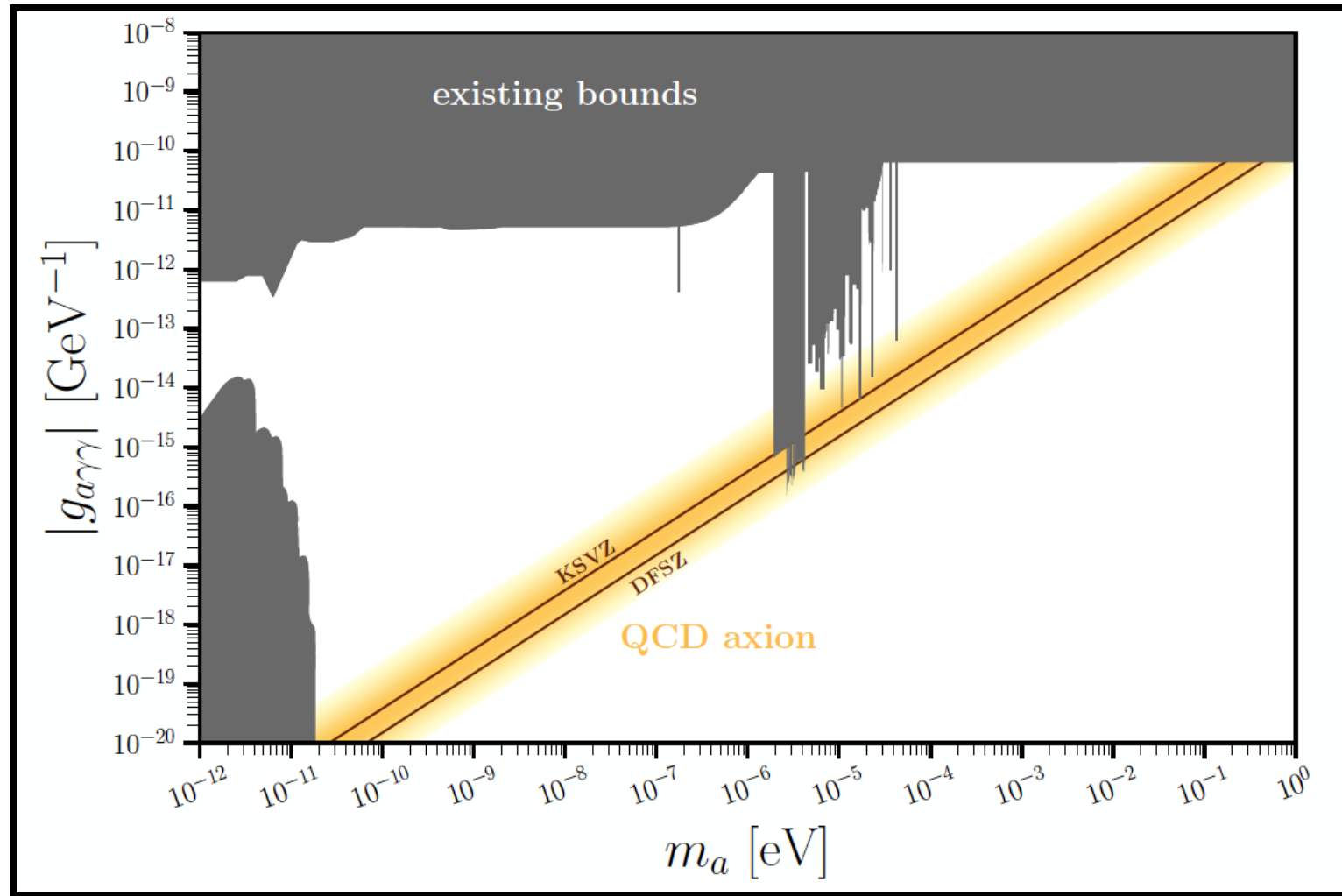
$$\mathbf{J}_{\text{eff}} = g_{a\gamma\gamma} \sqrt{2\rho_{DM}} \cos(m_a t) \mathbf{B}$$

Axion-photon experiment ingredients

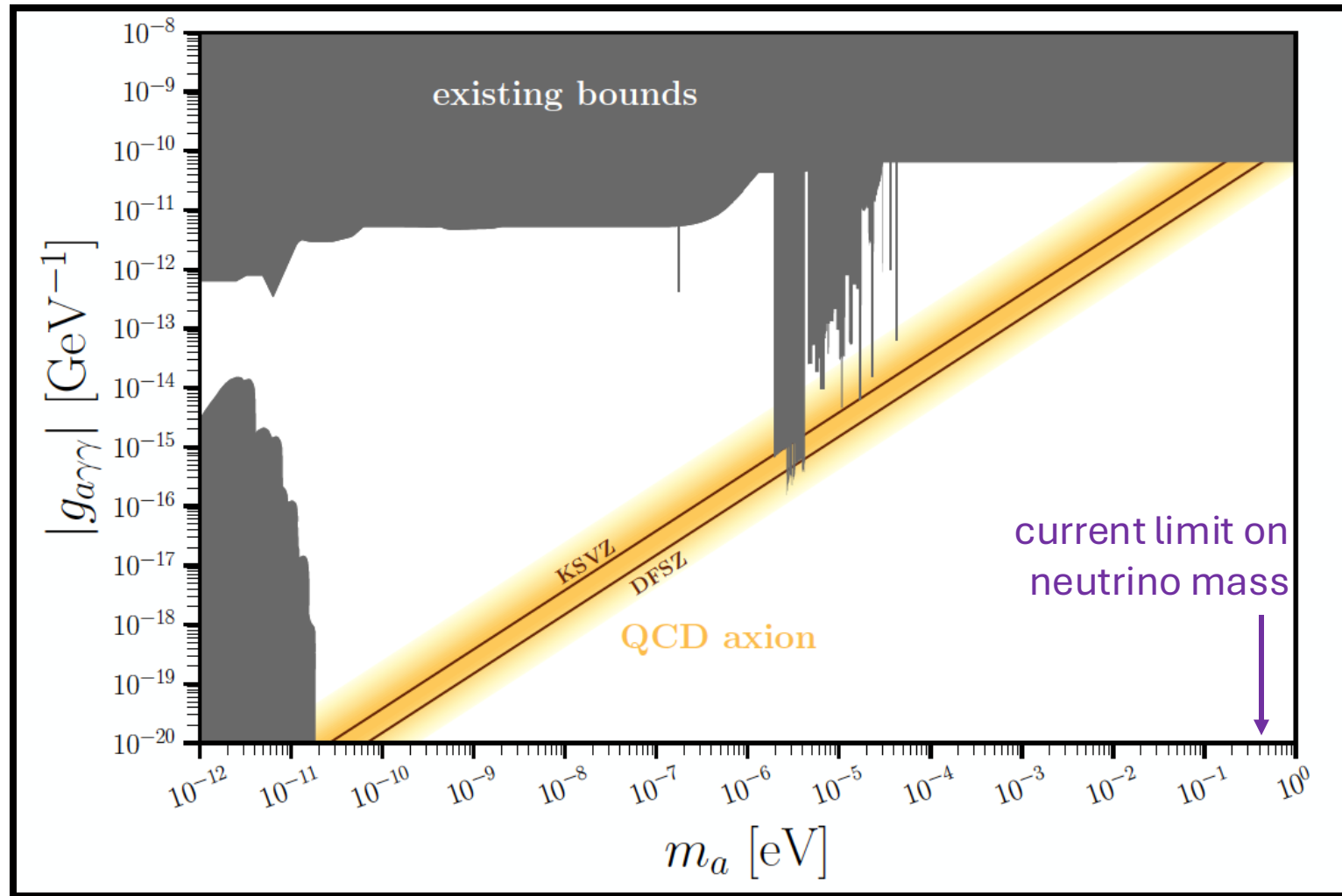
1. **Strong magnetic field** (axion-to-photon converter)
2. **Coupling system** (couple photons/EM signal into detector)
3. **Very sensitive detector** (amplify a tiny signal)

$$\mathbf{J}_{\text{eff}} = g_{a\gamma\gamma} \sqrt{2\rho_{DM}} \cos(m_a t) \mathbf{B}$$

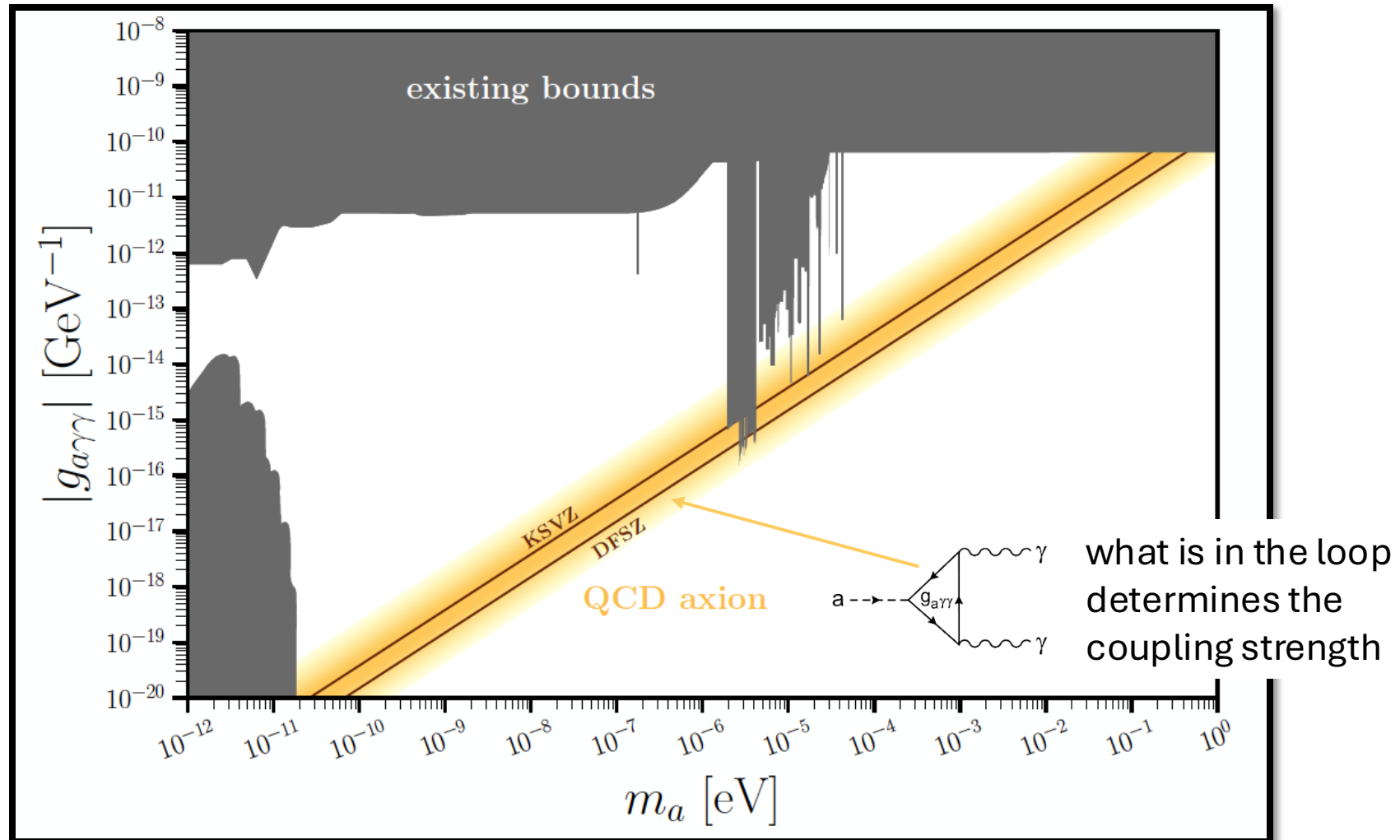
Parameter space to search



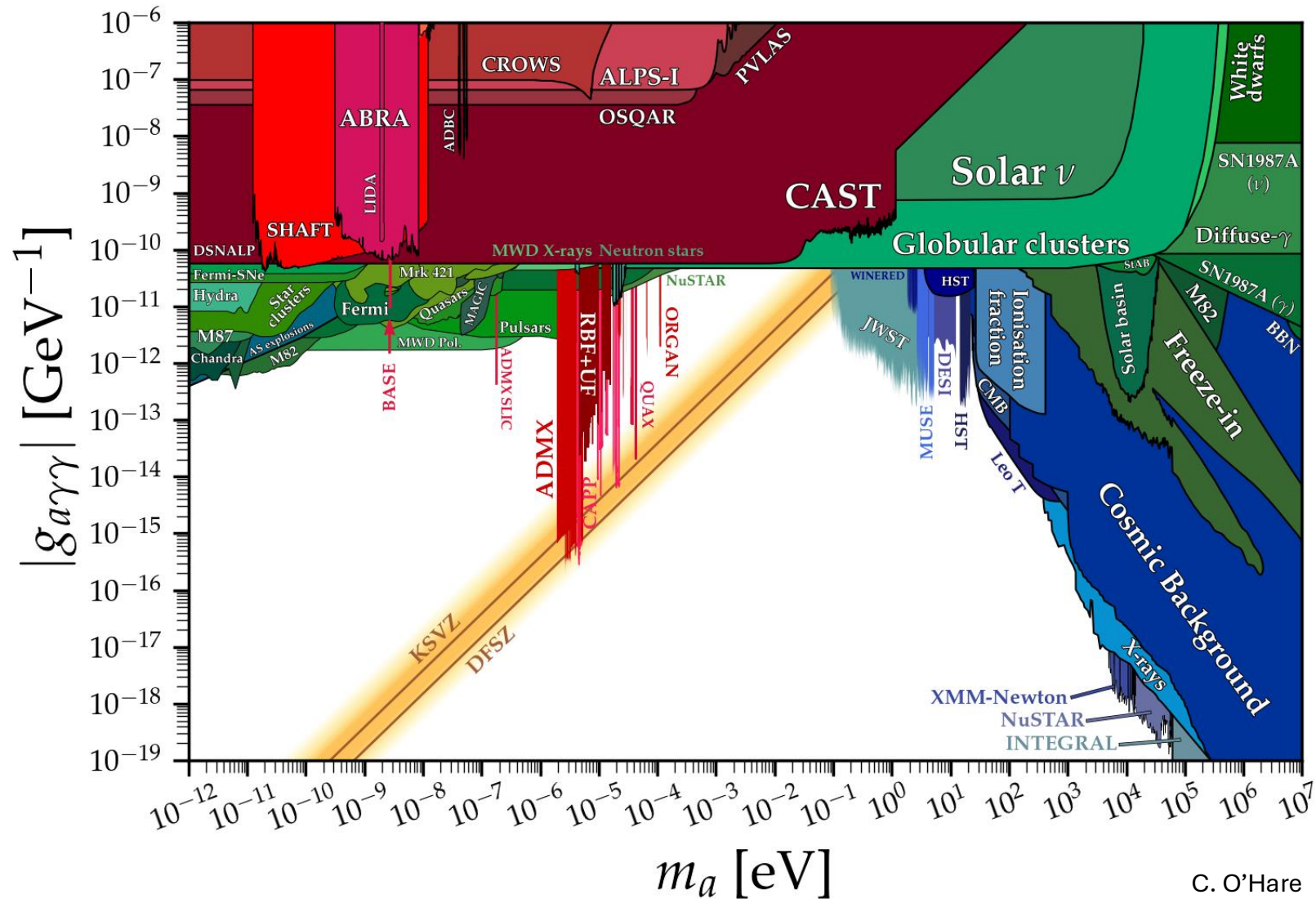
Parameter space to search



Parameter space to search

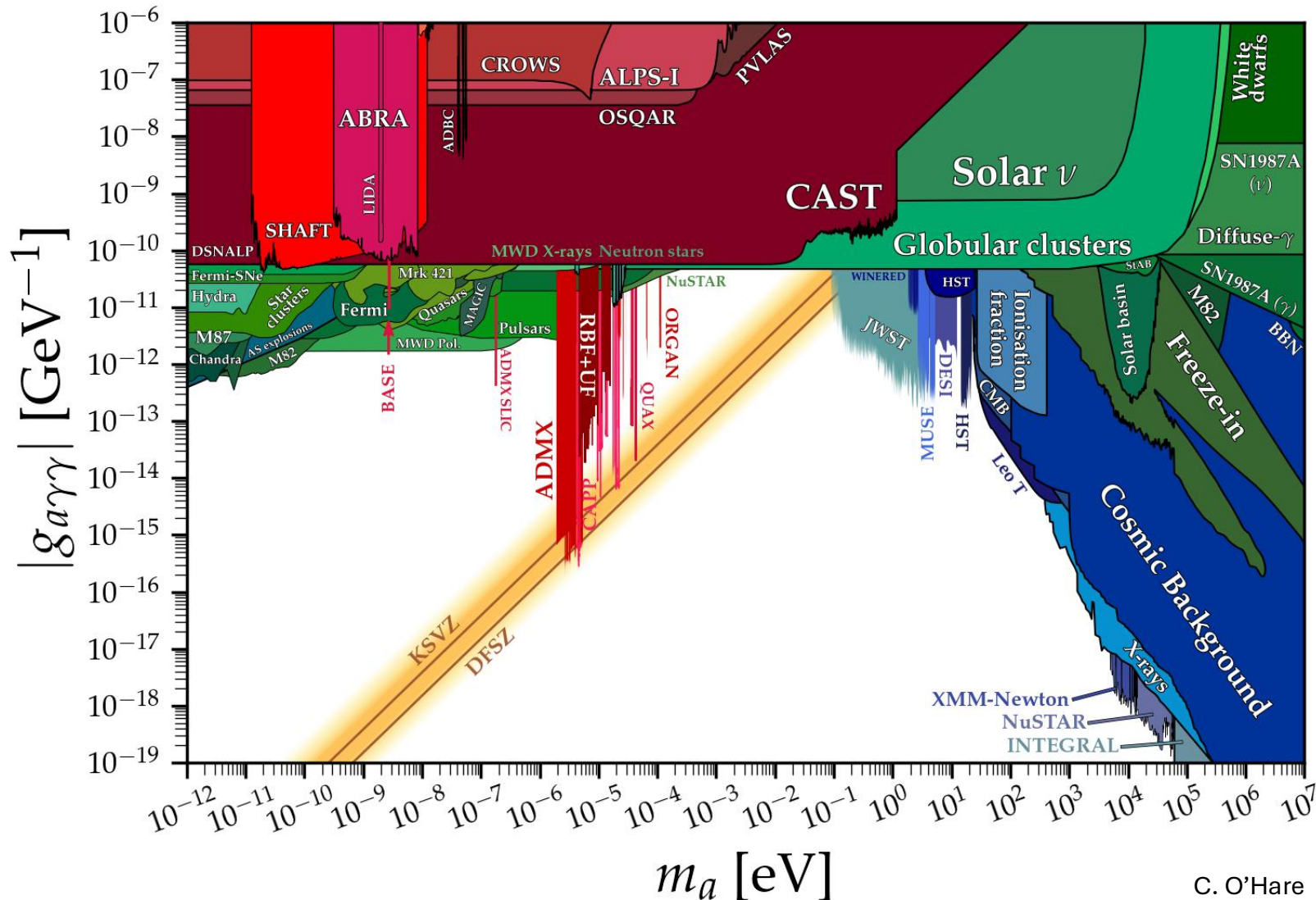


Existing bounds



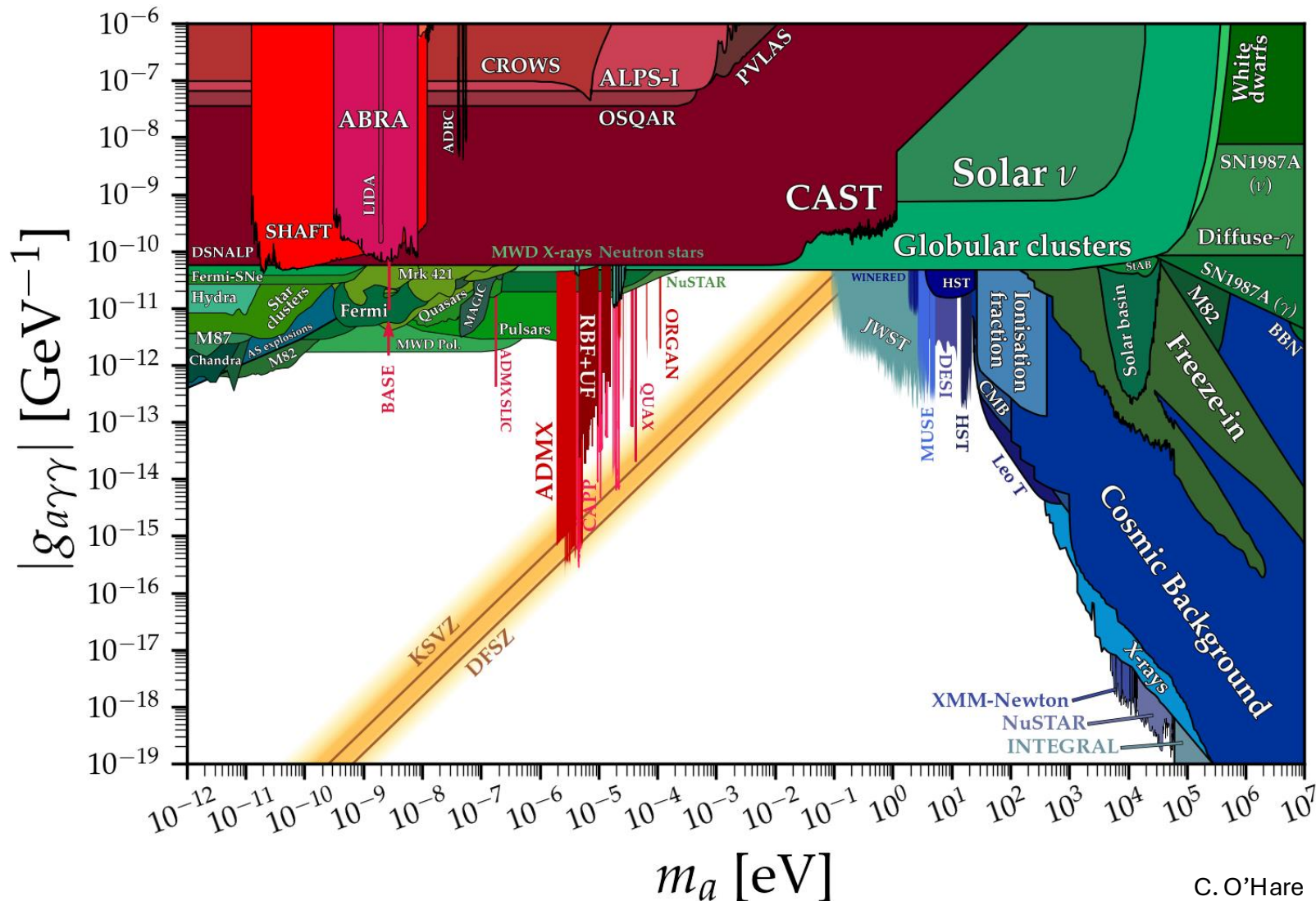
Existing bounds

- * Many different sources of axions (not all dark matter)
- * Many different experimental techniques
- must recognize different assumptions and systematics



Existing bounds

- * Many different **sources of axions** (not all dark matter)
- * Many different experimental techniques
- must recognize different assumptions and systematics



Sources of axions: astrophysical “indirect” searches

- Astrophysics provides much larger E and B fields than those achievable on earth
- Entire cosmos is the laboratory
- But... astrophysics is *hard*: lots of systematics
- Many targets and methods

Sources of axions: astrophysical “indirect” searches

- Astrophysics provides much larger E and B fields than those achievable on earth
- Entire cosmos is the laboratory
- But... astrophysics is *hard*: lots of systematics
- Many targets and methods

a few examples...
(many details under the hood)

Stellar cooling

- **Basic principle:**

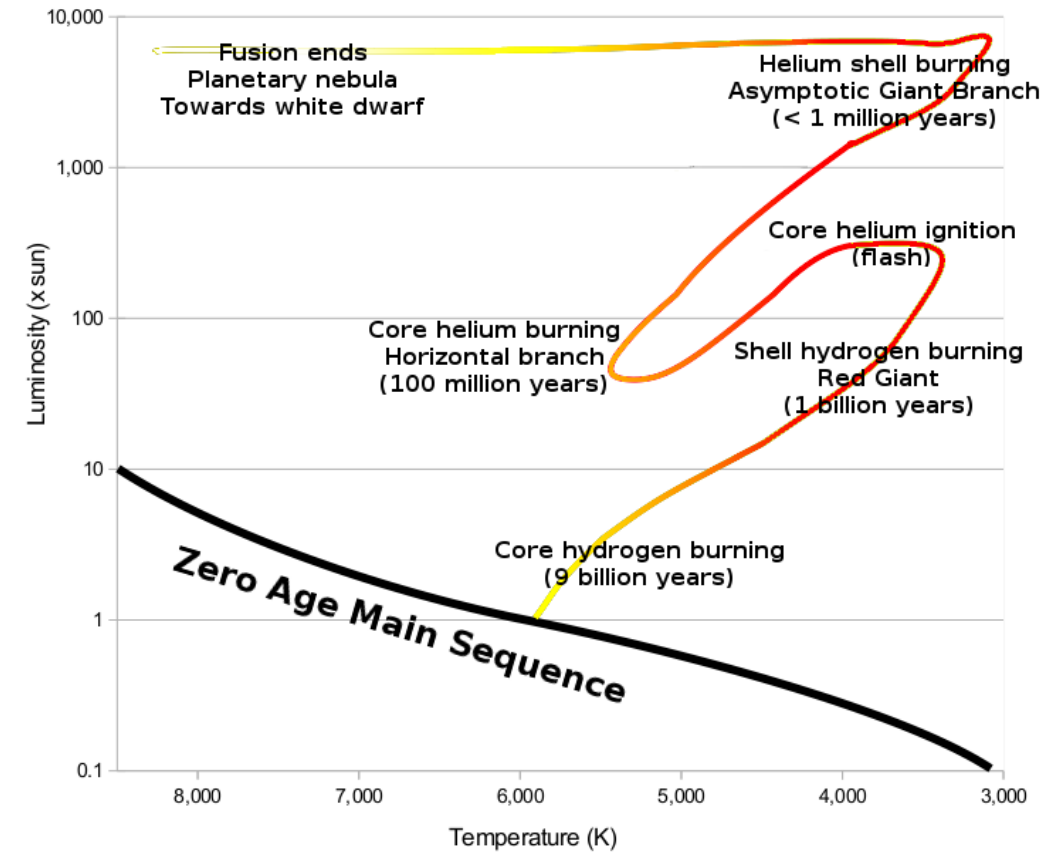
- Axions should be produced inside stars
- Once created, they will stream out, removing energy from stellar core
- If they take away too much energy, the process would affect stellar evolution

Stellar cooling in globular clusters

- Globular clusters
 - O(millions) stars
 - Some of oldest objects in milky way
 - Only relatively low-mass stars remain

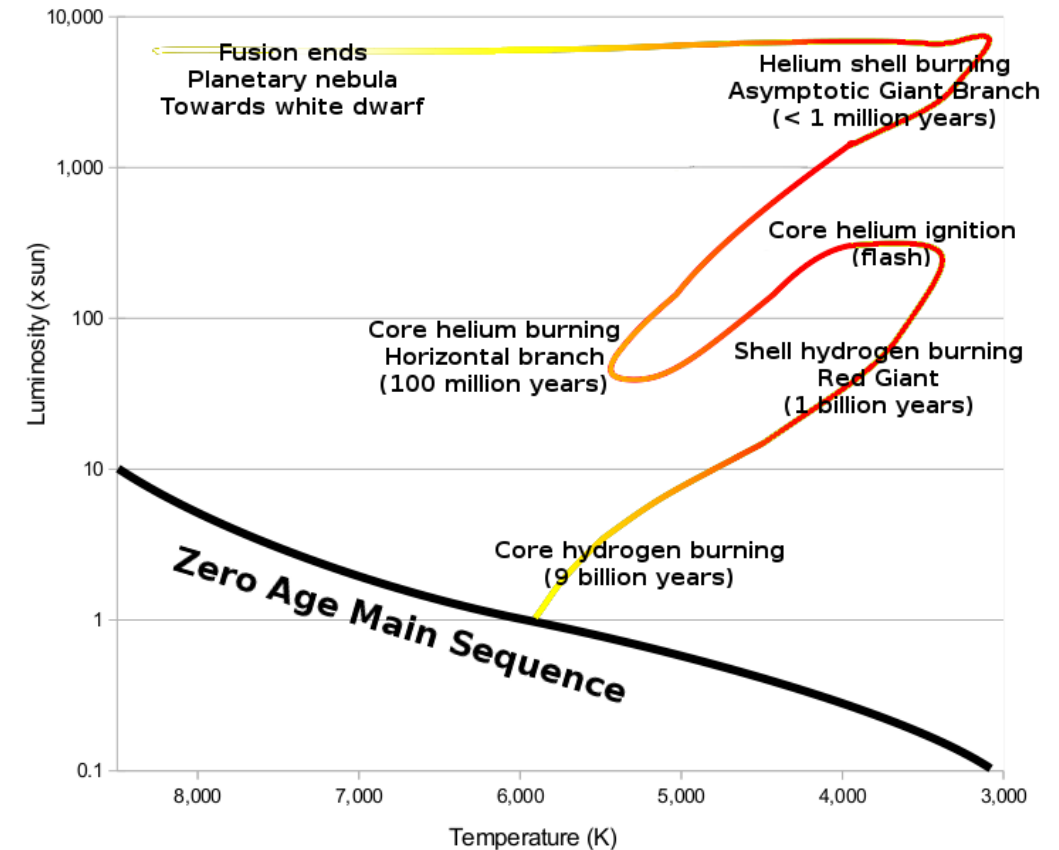
Stellar cooling in globular clusters

- Globular clusters
 - O(millions) stars
 - Some of oldest objects in milky way
 - Only relatively low-mass stars remain
- Stellar evolution proceeds in phases
 - Main sequence: core H burning
 - Red giant branch: H burning shell
 - Horizontal branch: core He burning





Stellar cooling in globular clusters

- Globular clusters
 - O(millions) stars
 - Some of oldest objects in milky way
 - Only relatively low-mass stars remain
- Stellar evolution proceeds in phases
 - Main sequence: core H burning
 - Red giant branch: H burning shell
 - Horizontal branch: core He burning
- Number of stars in each phase tells you how long stars stay in each phase



Stellar cooling in globular clusters

- Globular clusters
 - O(millions) stars
 - Some of oldest objects in milky way
 - Only relatively low-mass stars remain
- Stellar evolution proceeds in phases
 - Main sequence: core H burning
 - Red giant branch: H burning shell  minimal loss to axions via photon coupling
 - Horizontal branch: core He burning  significant loss to axions via photon coupling
- Number of stars in each phase tells you how long stars stay in each phase

Axion conversion in plasma

- Photons propagate through dense plasma inside the stars
- EM waves cannot propagate in plasma below the plasma frequency

$$\omega_p \approx \sqrt{\frac{4\pi\alpha_{\text{EM}}n_e}{m_e}} \approx 1 \times 10^{-12} \text{ eV} \sqrt{\frac{n_e}{10^{-3}/\text{cm}^3}}$$

- Conversion probability after traveling distance L in Coulomb field E

$$p_{\gamma \rightarrow a} \sim g_{a\gamma\gamma}^2 E^2 L^2$$

Back to globular clusters

- **RGB:** plasma frequency $>$ temperature (no photon propagation)
- **HB:** plasma frequency $<$ temperature (photon propagation)
- Only get axion conversion if we have propagating photons!
- Calculate ratio:

$$R = \frac{n_{\text{HB}}}{n_{\text{RGB}}}$$

Look for gamma ray bursts from supernovae

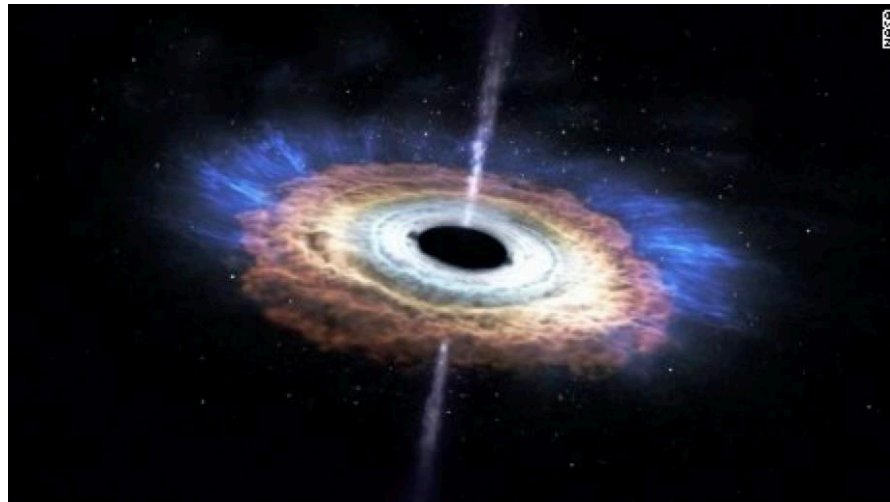
- **SN1987a:** Type II (core collapse) supernova 50kpc from Earth
- Hot photons propagating in core-collapse B-fields create axions
- Axion reconvert to gamma rays in galactic B-fields
- No gamma ray burst found in Gamma-Ray Spectrometer (Solar Maximum Mission)
- Limits:

$$g_{a\gamma} \lesssim 5.3 \times 10^{-12} \text{ GeV}^{-1}, \text{ for } m_a \lesssim 4.4 \times 10^{-10} \text{ eV},$$



Totally different method: black hole superradiance

- Axions with Compton wavelength \sim black hole size can extract angular momentum from the black hole, spinning it down
- If a rapidly spinning black hole exists with a given mass, that strongly constrains axions with the corresponding wavelength

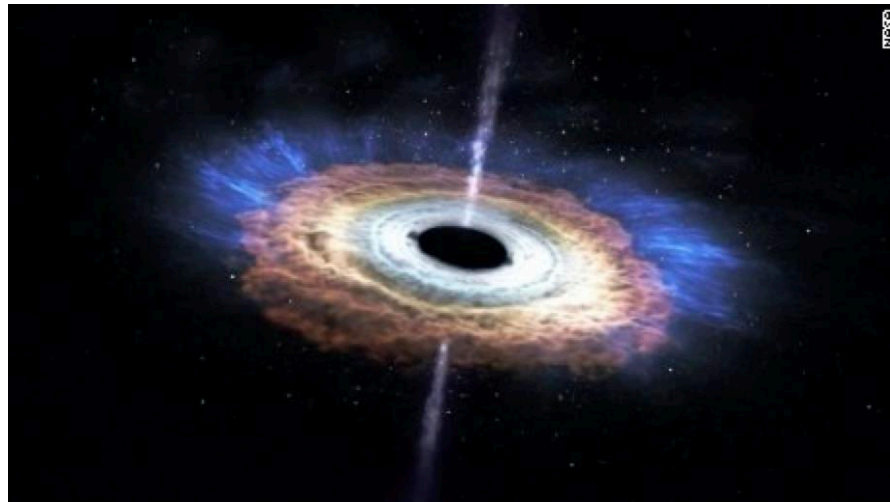


Note: exponential process--
axions don't have to be a large
fraction of the DM!

Totally different method: black hole superradiance

- Axions with Compton wavelength \sim black hole size can extract angular momentum from the black hole, spinning it down
- If a rapidly spinning black hole exists with a given mass, that strongly constrains axions with the corresponding wavelength!

Note: can also look for gravitational wave signatures from “gravitational atom”



Note: exponential process-- axions don't have to be a large fraction of the DM!

Small group activity!

link to github
generating this plot:



- Pick one of the blue or green exclusions from Ciaran's plot
- Lightning lit review: figure out the basic concept and discuss as a small group
- At the end, elect someone to briefly summarize the concept to everyone

