

UNIVERSITY OF CALIFORNIA, SAN DIEGO

**A Search for New Physics producing Jets, Large MT₂, and Disappearing Tracks in 13 TeV
Proton-Proton Collision at CERN's Large Hadron Collider**

A dissertation submitted in partial satisfaction of the
requirements for the degree
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in

Physics

by

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PUBLICATIONS

The CMS collaboration, Sirunyan, A.M., et al., *Searches for physics beyond the standard model with the M_{T2} variable in hadronic final states with and without disappearing tracks in proton-proton collisions at $\sqrt{s} = 13$ TeV*, Eur. Phys. J C (2019)

ABSTRACT OF THE DISSERTATION

A Search for New Physics producing Jets, Large MT₂, and Disappearing Tracks in 13 TeV Proton-Proton Collision at CERN’s Large Hadron Collider

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This work presents two searches for new physics characterized by pair-production of strongly interacting particles, each decaying to hadronic jets and a particle that is not detectable. The searches use the full 13 TeV proton-proton collision dataset produced by CERN’s Large Hadron Collider and recorded by the CMS detector from 2016 to 2018, with total integrated luminosity 137 fb^{-1} . The presence of particles interacting too weakly to be detected is inferred using imbalance in the transverse momentum of the collision products, and sensitivity to pair-production is enhanced by requiring large values of the kinematic variable M_{T2} in events with

at least two jets. The first search is inclusive, binning events using the total hadronic transverse energy, the total number of jets, the number of jets reconstructed as originating from a bottom quark, and either the value of M_{T2} in multijet events, or the transverse momentum of the jet in monojet events. The second search extends the first, by requiring the presence of a disappearing track in the event, and adds binning in the length and transverse momentum of the disappearing track. Both searches are sensitive to a variety of extensions to the Standard Model that include dark matter candidates. Of greatest interest, the results set constraints on pair production of squarks and gluinos as predicted by R-parity conserving supersymmetric extensions of the Standard Model, in which the lightest supersymmetric particle is a neutralino. The first search is sensitive to any decay chain terminating in Standard Model hadrons plus the neutralino, while the second specifically targets, with greatly enhanced sensitivity, decay chains containing an intermediate long-lived chargino. These constraints are the most stringent yet produced by any experiment.

Chapter 1

Introduction

Elementary particle physics aims to describe nature at its most fundamental level. This thesis presents analyses of data obtained from high energy proton collisions, which search for evidence for or against new theoretical models of elementary particle physics. This section introduces the theoretical context of these searches, describing first the theoretical model presently serving as the null hypothesis of elementary particle physics, then a few issues identified with this model, and finally some proposed solutions for these issues.

1.1 The Standard Model of Particle Physics

Modern elementary particle physics finds itself in a peculiar state. Physicists have constructed a theoretical model, now called the Standard Model of Particle Physics (the Standard Model; SM), that is the most quantitatively accurate scientific model of any kind, consistent with experiment in almost every laboratory test, even when experimental uncertainty is smaller than a part per ten billion [55, 12], with its most famously accurate prediction displayed in Table 1.1.

It is important to take this time to reflect on the Standard Model's astonishing accuracy, because the Standard Model is known with certainty to be imperfect, for a few reasons. A subset are discussed in Section 1.2. As a result, the Standard Model is sometimes treated with disdain,

	μ_e
Experiment	2.00231930436 (56)
Theory	2.002319304363 (15)

Table 1.1: A comparison of the leading experimental measurement of the electron’s magnetic moment [55], and the Standard Model’s theoretical prediction [12], in terms of the classical prediction, the Bohr magneton. Famously, this is the most accurate verification of theory by experiment in all of science. The numbers in parentheses are the uncertainties, on the same order as the last digits printed.

belying its quantitative success as a model of nature. This attitude is born from frustration. Particle physicists have a theory *known* to be incomplete, only an approximation, and yet so accurate an approximation that no evidence in favor of any specific proposed extension has ever been found. Modern particle physicists hope to find experimental clues by testing the Standard Model using any available means, including, as in this work, by comparing its predictions for the outcomes of high energy proton collisions with experimental data.

1.1.1 A Brief Description

The Standard Model is based in quantum field theory (QFT), which provides the general toolkit used to make calculations, and identifies the observed elementary particles as excitations of these underlying fields. Since every elementary particle corresponds directly to an underlying field and vice-versa, for example “the electron particle” and “the electron field” are colloquially treated as synonyms.

The SM itself, like every model of elementary particle physics based in QFT, consists of a list of particles present in nature and a description of their interactions with each other. The particles can be grouped in several ways, but by far the most significant is to group bosons and fermions.

The bosons of the Standard Model are the photon (typically indicated by γ), the W^+ and W^- , the Z , the eight gluons, and the Higgs. Of these, the Higgs has no intrinsic angular momentum (“spin-0”) and the others have intrinsic angular momentum of exactly \hbar (“spin-1”).

Boson	Mass (GeV) [73]	Interaction
Photon (γ)	0	Electromagnetism
Gluon (g)	0	Strong (QCD)
W^\pm	80.379 ± 0.012	Weak (charged current)
Z	91.1876 ± 0.0021	Weak (neutral current)
Higgs (h)	125.18 ± 0.16	Weak (symmetry breaking and mass)

Table 1.2: The Standard Model bosons, their masses, and associated interactions. Note that there are eight gluons, but they are not experimentally distinct.

One of the major insights obtained from QFT is that the existence of spin-1 particles requires the existence of fermions that interact with them, and vice versa. A fermion that interacts with a boson is said to be “charged under” the boson. Moreover, pure fermion interactions are forbidden, so that fermion-fermion interactions require a “mediating” boson under which both fermions are charged, while boson-boson interactions are perfectly acceptable. For this reason, bosons are said to mediate fundamental interactions. The photon mediates electromagnetism, and the gluons collectively mediate the strong interaction, as described by Quantum Chromodynamics (QCD). For historical reasons, the distinct interactions mediated by the W^+ and W^- , the Z, and the Higgs are all collectively called the weak interaction. The bosons of the Standard Model are summarized in Table 1.2.

The interactions have very different character due to differences in the mediating bosons. The photon and electromagnetism are most familiar. Electromagnetism is relatively simple, as a result of having only a single-component charge and a single mediating boson, and is relatively easily studied due to its long range, as a consequence of the masslessness of the photon. The weak interaction, by contrast, is less well-known in part due to its very short range, a consequence of the large masses of the mediating bosons. In addition to the $1/r^2$ behavior familiar from electromagnetism, forces are also suppressed by an exponential term, $e^{-\frac{mc^2}{\hbar c} r}$, where m is the mass of the mediating boson and $\hbar c \approx 200$ GeV-pm. This term vanishes for massless bosons like the photon, but since the weak bosons have masses on the order of 100 GeV, the weak interaction becomes negligible after only a few picometers despite having an intrinsic interaction

strength comparable to electromagnetism. Thus, the weak interaction has an impact only on nuclear and elementary particle physics. The strong interaction is also short ranged, with effects becoming significant only on the order of femtometers, but for an entirely different reason. As its name suggests, it is also intrinsically stronger than the other interactions. It will be discussed in Section 1.1.3.

The fermions of the Standard Model are all spin- $\frac{1}{2}$, and come in two major groups. The leptons are those fermions which do not participate in the strong interaction, and the quarks are those that do. The leptons can be further subdivided into the electrically charged leptons, like the electron, and the electrically neutral neutrinos, which having neither strong nor electromagnetic interactions can interact only weakly, and are ghost-like particles as a result. The quarks can be further subdivided into the up-type quarks, which are positively charged, and the down-type quarks, which are negatively charged. There are three of each quark, one for each of the charges of the strong interaction, but like the eight gluons, they are not experimentally distinct.

As is evident in Table 1.3, the fermions of the Standard Model occur in triplets with identical properties aside from mass: the up-type quarks, the down-type quarks, the charged leptons, and the neutrinos are all roughly three versions of the same particle with different masses. The fermions are therefore said to be arranged in three similar generations, with generation 1 having the least mass and generation 3 the greatest.

If each generation were truly independent of the others, then each would be independently stable, as there would be no way to cross from, say, generation 3 to generation 1. However, it is an experimental fact that quarks can decay across generations via their interactions with the W boson, albeit much more slowly than within the same generation. This is well-accommodated in the Standard Model [21, 49], in which there is no reason in general for the charge eigenstates with respect to any given boson to be equal to those of any other. Famously, the Higgs boson is responsible for the masses of the Standard Model's elementary particles, with each particle acquiring a mass proportional to the strength of its interaction with the Higgs. The Higgs

Fermion	Symbol	Mass (GeV) [73]	Electric charge
Up quark	u	$0.0022^{+0.0005}_{-0.0004}$	$+\frac{2}{3}$
Charm quark	c	$1.275^{+0.025}_{-0.035}$	$+\frac{2}{3}$
Top quark	t	173.0 ± 0.4	$+\frac{2}{3}$
Down quark*	d	$0.0047^{+0.0005}_{-0.0003}$	$-\frac{1}{3}$
Strange quark	s	$0.095^{+0.009}_{-0.003}$	$-\frac{1}{3}$
Bottom quark	b	$4.18^{+0.04}_{-0.03}$	$-\frac{1}{3}$
Electron	e	$0.0005109989461(31)$	-1
Muon	μ	$0.1056583745(24)$	-1
Tau	τ	$1.77686(12)$	-1
Neutrinos	ν_e, ν_μ, ν_τ	0^\dagger	0

Table 1.3: A table of the Standard Model fermions, first the quarks then the leptons. The vastly different experimental character of quarks due to their confinement by the strong interactions make their masses much more difficult to measure than those of the charged leptons. Note that every Standard Model fermion also has an antiparticle, with identical properties except for opposite charge. (*) The charge and mass eigenstates of (by convention) down-type quarks are different, so that for example “the mass of the down quark” does not exist. Otherwise, decays across generations would be impossible. In practice, they are so nearly equal that the mass of the down quark is equated with that of the lightest mass eigenstate, etc. (†) The Standard Model predicts that neutrinos are massless but they experimentally have nonzero masses, albeit one million times smaller than the electron’s (see Section 1.2).

interaction eigenstates, being the mass eigenstates, are privileged as the only charge eigenstates that are also eigenstates of the Hamiltonian, and so are the only states that can be more than transient, the states of so-called “real” particles. The hybrid nature of the real states with respect to W interactions produces effective generation-crossing interactions for quarks described by the Cabibbo-Kobayashi-Maskawa (CKM) matrix, with values measured experimentally to be approximately [73],

$$\begin{bmatrix} |V_{ud}| & |V_{us}| & |V_{ub}| \\ |V_{cd}| & |V_{cs}| & |V_{cb}| \\ |V_{td}| & |V_{ts}| & |V_{tb}| \end{bmatrix} \approx \begin{bmatrix} 0.974 & 0.224 & 0.004 \\ 0.218 & 0.997 & 0.042 \\ 0.008 & 0.039 & 1.019 \end{bmatrix} \quad (1.1)$$

where V_{xy} indicates the effective coupling between quarks x and y . If the eigenstates were exactly equal, the CKM matrix would be diagonal, with $V_{xx} = 1$ and $V_{xy} = 0$ ($x \neq y$), prohibiting trans-generation decays. As can be seen in Equation 1.1, the CKM matrix is *nearly* diagonal, meaning trans-generation decays are slow but not impossible. This causes the third generation bottom quark to have a longer lifetime than the second generation charm quark despite having a larger mass, because a bottom quark must cross a generation to decay (the top quark is more massive) using $|V_{ub}| \approx 0.004$ or $|V_{cb}| \approx 0.042$, while a charm quark does not (the strange quark, also of the second generation, is less massive) and may decay using $|V_{cs}| \approx 0.997$. In fact, the bottom quark lifetime is long enough that bottom hadron decay lengths are macroscopic, on the order of millimeters. This is experimentally relevant, as discussed in Section 2.2.2. Similar physics is possible for leptons since neutrinos are now known to be massive, but is not incorporated in the Standard Model, in which neutrinos are massless (see Section 1.2).

Further description of the Standard Model’s interactions is more conveniently done using Feynman Diagrams, in the next section.

1.1.2 Feynman Diagrams and Perturbative Expansions

The interactions of the Standard Model are expressed as a Lagrangian density of intimidating complexity. Applying the Euler-Lagrange equations to this Lagrangian produces nonlinear differential equations that have no known exact solutions. However, it is possible to extract an approximate solution as a perturbative expansion in powers of the coupling, a parameter that indicates the intrinsic strength of an interaction. This procedure at first seems barely feasible, as it is no trivial task to find all the contributions to the leading order term, then the next to leading order term, and so on by simple inspection of the Lagrangian. Fortunately, physicist Richard Feynman was able to devise a diagrammatic method for expressing the terms of the perturbative expansion that is far more intuitive.

First, one may inspect each term of the Lagrangian to assemble so-called vertices, the building blocks of diagrams. Each field that occurs in a term represents one line, and all the lines of a term intersect at a central point to form the vertex. For example, the Standard Model Lagrangian contains a term in which the electron field appears twice and the photon field appears once, from which one obtains the vertex shown in Figure 1.1. Each vertex represents one power of the coupling. The diagrams with the fewest vertices, then, are the lowest order diagrams in the perturbative expansion.

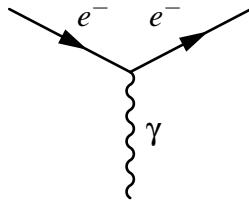


Figure 1.1: The Standard Model Lagrangian contains a term in which the electron field appears twice, and the photon field once, which corresponds to this Feynman diagram vertex. This vertex, and analogous vertices in which another electrically charged fermion replaces the electron, are the vertices of electromagnetism. Traditionally, straight lines with arrows represent fermions, and a wavy line represents a photon (or a W or Z). The arrows on fermion lines mark whether a fermion is matter or antimatter, depending on whether the arrow points generally in the same (matter) or opposite (antimatter) direction as the flow of time (left to right).

Typically, time runs left to right in Feynman diagrams. The vertex in Figure 1.1, then, depicts an incoming electron absorbing or emitting a photon, and continuing along. All Feynman vertices can be freely rotated, effectively changing the flow of time. For example, one can rotate Figure 1.1 to produce Figure 1.2. Figure 1.2 depicts the annihilation of an electron and its antiparticle, the positron, into a photon. While this process is allowed by the Standard Model as shown, it is not allowed kinematically, as it is impossible to conserve energy and momentum with only a single massless particle in the final state. To produce the leading order diagram for electron-positron annihilation that is allowed by kinematics, two vertices must be *connected* as shown in Figure 1.3, by attaching two identical external lines to form an internal line.

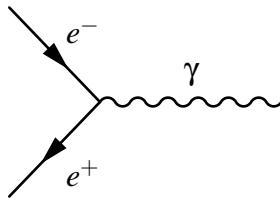


Figure 1.2: This diagram is a rotation of Figure 1.1. Instead of an incoming electron absorbing or emitting a photon, this diagram represents an electron and its antiparticle, the positron, annihilating into a photon. While permitted as an interaction by the Standard Model, this process is not kinematically allowed, making Figure 1.3 the leading order diagram for electron-positron annihilation.

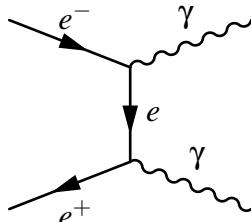


Figure 1.3: This diagram connects vertices like those of Figures 1.1 and 1.2 to produce the leading order kinematically allowed diagram for electron and positron annihilation. Whether the internal line is an electron or positron is ambiguous, due to the relativity of simultaneity. In some reference frames, the positron emits a photon first, then annihilates with the electron, and so the internal line is a positron. In others, the electron emits a photon first, then annihilates with the positron, and the internal line is an electron.

All of the diagrams shown thus far have been “tree-level,” with no internal loops. These tend to be the leading order diagrams, but diagrams with internal loops also contribute to the

amplitude, usually at sub-leading order. Figure 1.4 is one such diagram, contributing to the next-to-leading order term of the electron-positron annihilation amplitude.

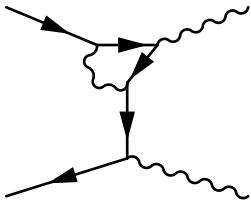


Figure 1.4: This diagram is one of the next-to-leading order contributions to the electron-positron annihilation amplitude, containing an internal loop.

In all of the diagrams we have reviewed thus far, the electron could be replaced with any electrically charged fermion, such as a muon or up quark. Collectively, all of these diagrams and others like them, constructed by combining variants of Figure 1.1, constitute electromagnetism.

The strong interaction as described by Quantum Chromodynamics (QCD) has a similar foundational diagram, shown in Figure 1.5 (upper), in which the photon is switched out for a gluon, and the fermion must be a quark. Unlike electromagnetism, the strong interaction also contains interactions between the mediating bosons, the gluons, alone. These vertices are partly responsible for the dramatically different physics of the strong interaction compared to electromagnetism. See Section 1.1.3 for more details.

Finally, some of the vertices of the weak interaction are shown in Figure 1.6.

In order to calculate rates and distributions of elementary particle physics processes, one assembles diagrams using the minimal number of vertices, then the next to minimal number, and so forth up to the desired precision, then follows an algorithm to convert each diagram to its equivalent mathematical expression. The result of evaluating this expression is the quantum mechanical amplitude for the process. This tedious procedure is now performed almost entirely by computers, using software such as the MadGraph generator [10].

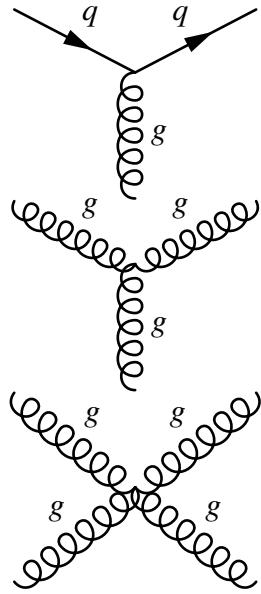


Figure 1.5: The Standard Model Lagrangian contains terms in which appear (upper) one of the quark fields twice, and a gluon field once, (middle) 3 separate gluon fields, and (lower) 4 gluon fields. These are the vertices of QCD. Unlike the lines of other bosons, gluon lines are traditionally drawn as springs.

1.1.3 Quantum Chromodynamics, Hadronization, and Jets

Quantum Chromodynamics is the theory of the strong interactions incorporated into the Standard Model, describing the interactions of quarks and gluons. States composed of quarks that are neutral with respect to the strong interactions (“colorless,” to avoid confusion with electromagnetic neutrality) are called hadrons, analogous to atoms of electromagnetism. There are two distinct ways to create a colorless state, due to the unique 3-component structure of the QCD charge. The first is to have a quark and antiquark with +1 and -1 units, respectively, of the same component of the QCD charge. These mutually cancel like a proton and electron in electromagnetism, forming a hadron called a meson. The second is to have 3 quarks, each with 1 unit of a different component of the QCD charge, or 3 antiquarks each with -1 unit of a different component. This kind of hadron is called a baryon, and there is no electromagnetic analogue to this kind of neutral object. The only stable hadron is the proton, the lightest baryon, excepting neutrons bound to protons by the QCD equivalent of dipole forces in atomic nuclei.

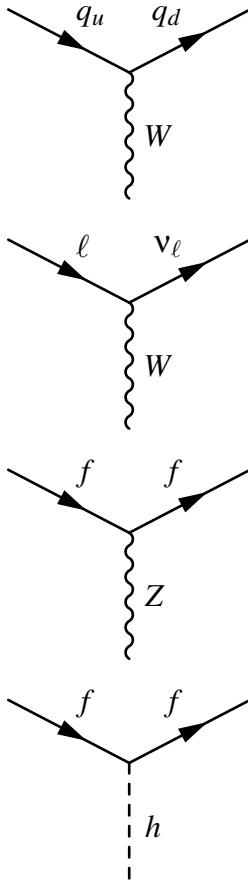


Figure 1.6: The Standard Model Lagrangian contains terms in which appear (upper) an up-type quark field and down-type quark field once each, and the W field once, (second) a similar diagram with a charged lepton and its associated neutrino replacing the quarks, (third) a fermion field appears twice, and the Z field once, and (lower) a non-neutrino fermion field appears twice, and the Higgs once. These are the vertices of the weak interaction that include both fermions and bosons. The weak interaction also includes many boson-only interactions similar to Figure 1.5 (lower), but these are omitted for brevity. The Z diagram (third) is very similar to the interaction of the photon with fermions, with the important addition that the fermion here can also be an electrically-neutral neutrino. The lines of spin-0 bosons are traditionally drawn as dotted lines, as in the lower diagram.

The formation of hadrons is called hadronization, and the hadronization process after a high energy collision tends to produce objects called jets that are important in experiments.

To understand hadronization and jets, it is necessary to understand a feature of the strong interaction intimately related to why it is short ranged despite having a massless mediator, as was mentioned in the previous sections. Unlike the familiar electromagnetism, in which the potential

energy associated with two charges decreases as the separation of the charges increases, the QCD potential energy increases without bound as two quarks are separated, roughly linearly proportional to distance [50], as shown in Figure 1.7. Eventually, the potential energy stored in the quark system exceeds the rest energy of a new hadron. At this threshold, it becomes energetically preferable to convert some of the stored energy into new hadrons, creating a new quark-antiquark pair in order to reset the quark separation distance, than to allow further separation of the quarks. As a result, any attempt to separate two quarks cannot ultimately succeed, and all quarks are bound inside colorless hadrons. As gluons *also* carry color charge (that is, gluons interact with themselves as shown in Figure 1.5), they are similarly confined, and the reach of the strong interaction is limited to only a few femtometers.

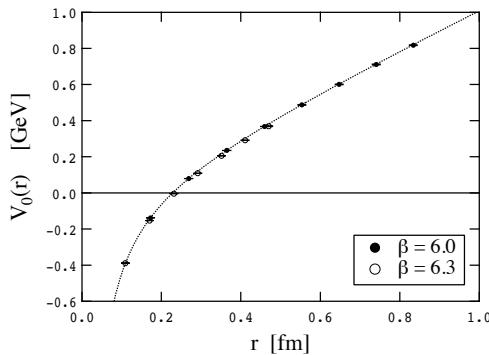


Figure 1.7: The potential energy associated with two static quarks predicted by QCD as a function of distance, calculated numerically using the lattice technique. The potential increases roughly linearly with distance once a quark is no longer inside a hadron. The strong coupling used for the numerical calculation is denoted by β . Taken from [50].

During a high energy collision involving a hadron like a proton, quarks and gluons are customarily ejected, and eventually reach a great enough distance to trigger the formation of new hadrons. But, the ejection is oftentimes so violent that even this new hadron fragments into

new daughter hadrons, then these fragment, and so on. This fragmentation process produces a spray of hadrons, hadron decay products, and hadrons formed from QCD radiation analogous to bremsstrahlung, all traveling in approximately the same direction as the original ejected quark or gluon, which are collectively called a hadron jet.

Unfortunately, while the production of jets is understood qualitatively, it is difficult to predict quantitatively due to yet another peculiar feature of QCD. The perturbative expansion described in the previous section is performed in powers of the interaction's coupling, with one factor of the coupling per vertex. In order for a truncated perturbation series to be a good approximation, the coupling must be significantly less than 1. Otherwise, diagrams with more vertices, representing higher order terms in the expansion, are in general more important than diagrams with fewer vertices, and no finite truncation can be accurate. Although the QCD coupling is less than 1 at high energy, making the perturbative approach still viable for very high energy collisions, the coupling explodes at low energies, which earns the strong interaction its name. Even at intermediate energies where QCD is technically perturbative, the number of diagrams that must be evaluated for an accurate result is still impractically large.

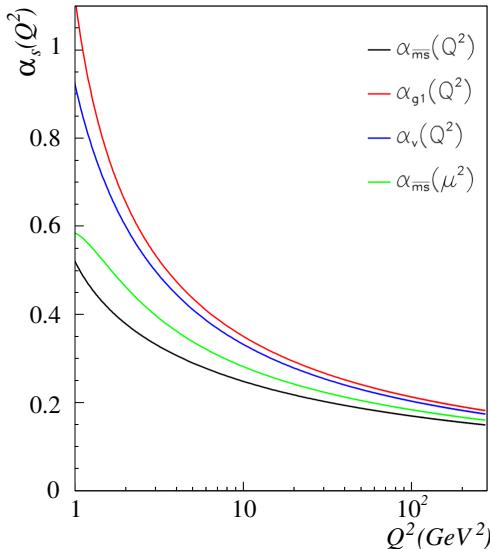


Figure 1.8: The QCD coupling explodes at low energy, making QCD non-perturbative for any interaction energy below around 1 GeV. The different lines are obtained for different methods of performing the calculation, specifically different renormalization schemes. Taken from [37].

Non-perturbative approaches to QCD exist, including the lattice numerical method used to produce the plot shown in Figure 1.7, but they are extremely computationally expensive for even the most simple calculations. It is not feasible to perform non-perturbative calculations of low energy QCD physics from first principles in an environment as complex as high energy hadron collisions. For the hadronization and fragmentation process, heuristic approximations tweaked to match data are used instead, typically the Lund String Model as implemented in the Pythia software package (see Section 2.4). Additionally, relatively low energy quarks and gluons can be emitted during the collision itself, called Initial State Radiation (ISR). A diagram of $u\bar{u} \rightarrow Z \rightarrow e^+e^-$ that includes a single ISR gluon is shown in Figure 1.9. The most reliable way to predict ISR is to measure it in a very similar physics process, then convert this measurement to an expected rate in the physics process of interest. For example, ISR in $Z \rightarrow e^+e^-$ events is essentially identical to ISR in $Z \rightarrow \nu\nu$ events.

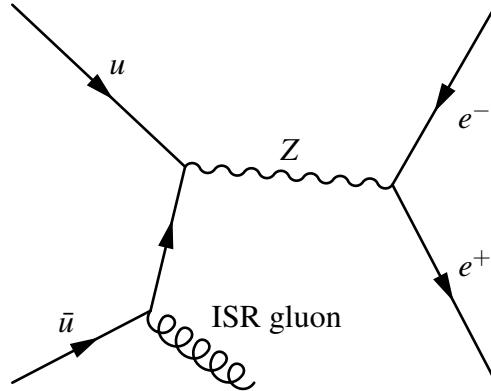


Figure 1.9: Initial State Radiation emitted as part of the hard collision, in this case $u\bar{u} \rightarrow Z \rightarrow e^+e^-$, can be at relatively low energy. This gluon will hadronize and add a jet to the event. QCD ISR is difficult to predict accurately due to the intractability of QCD at low energy.

In summary, QCD's intractability at low energy combined with its raw strength make precise predictions of the outcome of hadron collisions difficult to obtain, despite the Standard Model's overall success in describing the interactions of elementary particles. The Standard Model prediction cannot always be obtained directly from first principles with satisfactory accuracy.

1.1.4 Parton Distribution Functions

Being composite particles bound by QCD, the internal structure of hadrons is rich and difficult to understand quantitatively, due to the impossibility of using perturbation theory at the low energies that prevail in hadrons. Of greatest interest is the internal structure of the proton, the only stable hadron.

The proton is often said to be a bound state of two up quarks and one down quark. This is true enough—the *net* composition of a proton is two up quarks and one down quark—but is an oversimplification. The quarks in a proton constantly emit and absorb gluons, which can split to quark-antiquark pairs, which can in turn annihilate back to gluons. Viewed at a small enough length scale, the proton is not a simple object, and not even a bound state of 3 quarks, but rather a complex cloud of quarks and gluons with various energies. This is problematic for any attempt to predict the outcome of proton collisions, because one cannot be sure which of these subcomponents, called partons, will be the component of the proton that actually experiences a collision. Due to the intractability of QCD at low energy, the composition of a proton is known only from a fit to data, and called a Parton Distribution Function (PDF). PDFs are typically expressed as the probability to find a parton of a given species carrying a given fraction of the proton’s energy, as shown in Figure 1.10. These PDFs display a few sensible features described in the figure’s caption. A large number of proton collisions collectively sample all the possible initial partons, with momentum distributions described by the PDFs. Thus, a dataset composed of proton collisions naturally probes a wide range of collision energies with a variety of colliding particles, making proton collisions well-suited to searches for new particles of unknown properties.

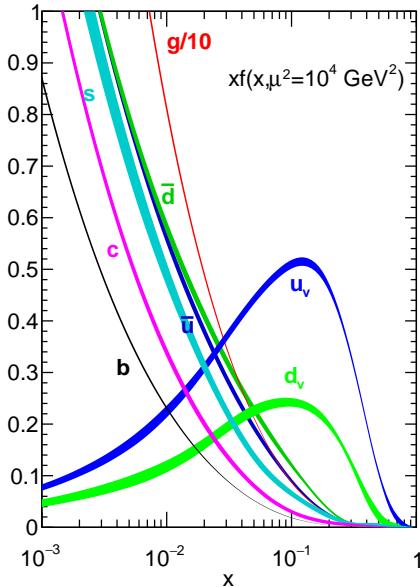


Figure 1.10: The plot shows the measured parton distribution function for a proton at an interaction energy of 100 GeV. The horizontal axis x is the fraction of the proton’s total momentum, and the vertical axis is the probability distribution. The finite thickness of the curves indicates the uncertainty. As expected for a proton, uud , the probability to collide with an up quark is around twice the probability to collide with a down quark at any energy. The v subscript on these two curves indicate that these are the “valence” quarks, the only quarks that are present as more than transient products of gluon splitting. It is much less likely to find an up or down antiquark, and roughly equally likely to find either, because these antiquarks are produced only by gluons splitting to quark-antiquark pairs. Since they are also produced only transiently by gluon splitting, the probability to find a heavier quark, namely s , c , or b , is exactly equal to the probability to find a heavy antiquark. Since the heavy quark and antiquark curves overlap, only the quark curves are represented. Quarks aside, notice that the gluon curve (red) has been suppressed by a factor of 10 to fit on the plot. To a decent approximation, high energy proton colliders in practice collide gluons. Taken from [16].

1.2 Some Problems with the Standard Model

The Standard Model is extremely successful, as epitomized by the result in Table 1.1, but not entirely satisfactory. In fact, its success in laboratory experiments is a source of frustration for elementary particle physicists, as there are precious few hints as to how its deficiencies can be repaired. This section introduces a few of the known problems with the Standard Model.

1.2.1 Gravity, and the Standard Model as an Effective Field Theory

In Section 1.1, there is never any mention of gravity as an interaction between particles. This is because the Standard Model entirely omits gravity from its description of nature, which means the Standard Model cannot possibly be the final theory. Instead, the Standard Model must be only an *effective theory*, a low energy approximation of nature that must be replaced by a more fundamental theory at some cutoff energy. In the case of the Standard Model, this cutoff cannot be any larger than the energy at which gravity becomes relevant in the interactions of elementary particles, called the Planck scale.

Unlike the fields of the Standard Model, whose couplings are dimensionless, the gravitational coupling G has dimensions of inverse energy squared. An effective field theory of the weak interactions developed by Enrico Fermi, which models the weak interactions of fermions without mediating bosons, has a similar coupling, $G_F \approx \frac{1}{(100 \text{ GeV})^2}$. The energy scale embedded in the coupling is a message from nature, hinting that this effective theory must be replaced by something more fundamental in order to model physics at the scale of around 100 GeV or more. Nowadays, we recognize this energy as the mass scale of the weak interaction's mediating bosons. At interaction energies on the order of 100 GeV or larger, there is enough energy to produce the boson itself as a resonant excitation of the field, a real particle, and the underlying theory is revealed. Making an analogous inference from the gravitational coupling, one obtains the maximum possible cutoff energy for the Standard Model, the Planck energy of around 10^{19} GeV, at which gravity must be incorporated into the quantum theory.

Moreover, the Standard Model becomes self-inconsistent at very high energies, even ignoring the gravity problem [43, 62]. The couplings of both electromagnetism and of the Higgs to itself increase with increasing energy, until eventually diverging at very high energy scales, a theoretical consistency issue called quantum triviality. As a result, the Standard Model's formulation of electromagnetism cannot be used for any energy greater than around 10^{34} GeV [43]. Although this energy is far greater than the Planck energy, at which the Standard Model must fail

anyway due to its exclusion of gravity, it is a completely independent indication that the Standard Model must be only a low energy limit of some more fundamental theory that supersedes it at high energy, one that does not exhibit these kinds of problems. Even if one supposes that gravity's exclusion from the Standard Model is not actually a problem, that attempting to treat gravity like the other interactions is misguided, the Standard Model still cannot be the final theory.

The Higgs self-interaction similarly diverges above the Planck scale, although had the Higgs mass turned out to be only around 200 GeV rather than the measured 125 GeV, the problem would arise below the Planck scale, and a Higgs mass as small as around 600 GeV would have made the triviality scale only a few TeV, accessible to current experiments [62].

Thus, there exist known finite energy scales at which the Standard Model *must* fail, due both to its exclusion of gravity and to internal mathematical consistency issues, and the actual cutoff energy may be much smaller than these upper limits. Every experiment probing an unexplored energy scale has a chance to be the first to observe a hint as to what this more fundamental theory looks like.

1.2.2 The Hierarchy Problem

As described in the previous section, the energy scales of the weak interaction and of gravity differ by approximately 17 orders of magnitude. If the Standard Model emerges from some more fundamental theory at the Planck energy, then what effect has pushed the weak scale down 17 orders of magnitude? The most troublesome expression of this apparently inexplicable hierarchy of energy scales centers on the mass of the Higgs boson. In the Standard Model, the Higgs mass m_h is given approximately by [15, 35],

$$m_h^2 \approx 2\mu^2 + (\Delta m_h)^2 \quad (1.2)$$

where μ is a free parameter called the bare mass of the Higgs that would most naturally be 0, and Δm_h is an adjustment that the Higgs acquires via its interactions with other fields in the Standard Model,

$$\Delta m_h^2 \approx \frac{3}{4\pi^2} (\lambda^2 - \lambda_t^2 + \dots) \Lambda^2 \quad (1.3)$$

where λ is the Higgs' coupling to itself, λ_t is its coupling to the top quark, and Λ is the energy at which the Standard Model is replaced by a more fundamental theory that is assumed not to contribute any further corrections to the Higgs mass, at most the Planck Energy. The couplings of the Higgs to all the fermions of the Standard Model appear in the correction with the same form as the top's λ_t , but the coupling of the top quark is by far the most important, since at ~ 1 it is by far the largest, as can be seen from the mass column in Table 1.3. The Higgs' self-coupling is somewhat smaller [35]. Therefore, $\Delta m_h^2 \sim -\Lambda^2$. If one takes $\Lambda \approx 10^{19}$ GeV, the Planck Energy, then Equation 1.2 is deeply implausible after also inserting the experimental value $m_h \approx 125$ GeV,

$$(125)^2 \approx 2\mu^2 - (10^{19})^2. \quad (1.4)$$

The Higgs mass of 125 GeV emerges from almost perfect cancellation of two terms with values around 10^{19} GeV. A model that produces an output of a very different order of magnitude from the generating input is said to be finely tuned. The Standard Model, epitomized by the Higgs mass, is *egregiously* finely tuned. This is the Hierarchy Problem.

Of course, the Standard Model's cutoff need not be the Planck Energy. If Λ is only 1000 GeV, then the Hierarchy Problem as stated disappears, and one can expect to discover new physics in elementary particle interactions at this experimentally accessible scale. Still, merely not being the Standard Model is insufficient; the replacement theory also needs to introduce a mechanism to stabilize the Higgs mass, lest it make its own contributions or allow the Standard Model fields to continue pushing it to large values. By far the most popular proposal for the kind of model that may supersede the Standard Model at the TeV scale is called Supersymmetry,

discussed in Section 1.3.

1.2.3 Massive Neutrinos

In the Standard Model, all fundamental particles acquire mass through their interactions with the Higgs, and the neutrinos are strictly massless. However, recent measurements of neutrinos have established that they undergo flavor oscillation, whether produced by cosmic rays in the atmosphere [39], or by nuclear processes in reactors on Earth [11] or in the Sun’s core [8]. This means that, just as for quarks, the neutrino mass eigenstates are not equal to the interaction eigenstates, and more importantly that the neutrinos’ mass eigenstates are not equal *to each other* [38]. This is only possible, of course, if at least 2 of the 3 neutrinos are massive. Although difficult particles to measure, experiments have managed to map out the neutrino mixing matrix analogous to the CKM matrix for quarks, the Pontecorvo-Maki-Nakagawa-Sakata matrix [73],

$$U_{PMNS} \approx \begin{bmatrix} 0.800 \leftrightarrow 0.844 & 0.515 \leftrightarrow 0.581 & 0.139 \leftrightarrow 0.155 \\ 0.229 \leftrightarrow 0.516 & 0.438 \leftrightarrow 0.699 & 0.614 \leftrightarrow 0.790 \\ 0.249 \leftrightarrow 0.528 & 0.462 \leftrightarrow 0.715 & 0.595 \leftrightarrow 0.776 \end{bmatrix} \quad (1.5)$$

where each element of the matrix shows the range consistent with experiment. Mysteriously, U_{PMNS} is not approximately diagonal, unlike the CKM matrix. Nonzero neutrino masses are entirely inconsistent with the Standard Model and, to date, are the only observables measured in laboratories that the Standard Model has failed to predict accurately.

1.2.4 Astrophysical Evidence of Dark Matter

Over the last century, astrophysical observations have assembled conclusive evidence that most of the Universe’s gravitating mass has at most a neutrino-like interaction cross section.

Some of the oldest evidence was obtained from galactic rotation curves. In general, stars

far from the cores of galaxies have much greater orbital velocity than would be inferred from the visible mass to their interior, implying either that general relativity is not a good model of gravity on large distance scales, or that most of the gravitating mass in galaxies is weakly interacting and arranged as an extended halo [58]. Recent observations have concluded with good precision that this is true also in the Milky Way [54].

Observation of the Cosmic Microwave Background (CMB) has allowed for measurement of the Universe’s energy density, and the forms taken by that energy density. The best fit is a Universe composed of about 6 times as much “dark” matter as visible matter, consistent with the ratio that would explain galactic rotation curves [31]. No modification of gravity has ever been found that can satisfactorily explain both the rotation curve observations, and explain these energy density measurements.

Observations of the Universe at more recent times have found a few cases of colliding galaxy clusters, most famously the Bullet Cluster, that offer further evidence for the existence of dark matter. In galaxy cluster collisions, the luminous matter can collide and produce a very hot plasma due to its relatively strong interactions, while dark matter continues on unaffected. The luminous matter can be accounted for directly via the emissions of the hot plasma, while the location of the mass, visible or not, can be measured by observing the gravitational lensing of the light of more distant galaxies by the foreground colliding clusters [27]. The result of this analysis for the Bullet Cluster is shown in Figure 1.11. The visible matter, revealed by its x-ray emissions, is caught in the collision zone at the center, while the majority of the gravitating matter has continued on unaffected, and is detectable only through its powerful gravitational lensing on either side of the luminous plasma. The inferred ratio of dark matter and visible matter is consistent with the amount that would explain both rotation curves and the energy densities measured using the CMB. Again, no modification to gravity has been found that can simultaneously explain galactic rotation curves and the energy density measurements from the CMB, and also the gravitational lensing observed in the Bullet Cluster and similar collisions.

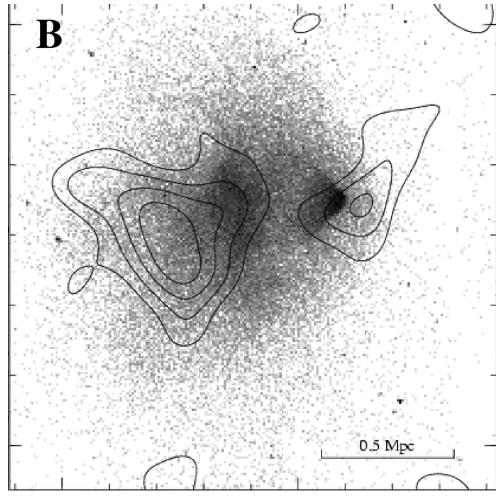


Figure 1.11: The distribution of luminous matter emitting x-rays (black dots) and the location of mass inferred through gravitational lensing of background sources (contours) are offset. The plasma emitting x-rays has been slowed and shocked as a result of the astronomically recent collision between the galaxy clusters, and so has been left behind by the dark matter, which constitutes most of the mass and was unaffected by the collision due to its negligible interaction cross section. Taken from [27].

Finally, astronomers have recently identified a few small galaxies that appear to contain little or no dark matter [75]. While this is an extraordinarily rare occurrence and effectively impossible in a large galaxy like the Milky Way, small galaxies are subject to larger statistical fluctuations, and the Universe has a rather large sample size of galaxies. It is difficult to imagine a modification of gravity that could, in most observations, cause there to appear to be around 6 times as much invisible mass in the Universe as visible mass, but in some small galaxies appear as if there is no modification at all.

If dark matter is not composed of a weakly interacting particle beyond the Standard Model, the only remaining possibility is that it is composed of larger non-radiating masses like rogue planets or black holes produced in the high-density primordial Universe and surviving to the present day. This possibility is taken seriously, and can be investigated by searching for small

microlensing events, where the light of a distant star is briefly gravitationally lensed when one of these bodies passes between an Earth-based telescope and the star. As there is around 6 times as much dark mass as visible in the galaxy, such an event should be relatively common. No excess of such events is observed in microlensing and complementary searches, so that if all dark matter has exactly the same mass, it cannot have any mass in the range $10^{-7}M_{\text{Sun}} < M_{DM} < 10^5M_{\text{Sun}}$ [42]. This leaves open the possibility that dark matter could be composed of black holes of small mass, with Schwarzschild radius around 0.1 mm or less, or that it could be composed of bodies with a range of masses. However, both possibilities are disfavored. It is suspected that small primordial black holes would generically be captured by stars, eventually migrate to the core as the star dies, and destroy the compact remnant via accretion [24], but no such events are observed. As for the latter possibility, the data places such strong constraints on the acceptable mass distribution that it is necessary to construct one carefully by hand to avoid the limits [42]. Particle dark matter, despite its own history of non-observation, remains the possibility most easily reconciled with the data.

There is especially strong hope for detecting particle dark matter at the TeV scale, due to a suggestive coincidence called the Weakly Interacting Massive Particle (WIMP) Miracle. The characteristic energy scale of the weak interaction is ~ 100 GeV. A hypothetical particle with mass in the suitable range, on the order of 1 to 10,000 GeV, interacting via the Standard Model's weak interaction or a similar interaction, would have been produced in the early Universe and survived to the present day at the proper density to account for the observed cold dark matter [44]. Furthermore, it would be an astonishing coincidence if the dark matter particle is only distantly related to the Standard Model, perhaps only by gravitational interactions, and yet somehow by dumb luck ended up with a relic mass density on the same order of magnitude as that of the Standard Model particles. Viewed in light of indications that new physics ought to be found at the TeV scale from the Hierarchy Problem and other issues with the Standard Model, the case for particle dark matter at the TeV scale is compelling. The chief argument against WIMPs is

that they have not yet been observed in experiments despite concentrated efforts, yet much of the parameter space remains untouched. It is a high priority of current experiments to extend sensitivity into these regions.

1.3 Supersymmetry

Among all the proposals for extending the Standard Model and repairing many of its issues at the TeV scale, supersymmetric extensions are by far the most popular and intensively researched. Supersymmetry proposes that for every Standard Model fermion (boson), there exists a *superpartner* boson (fermion) with nearly identical properties aside from its spin. The superpartners of spin- $\frac{1}{2}$ particles are spin-0, and the superpartners of spin-1 and spin-0 particles are spin- $\frac{1}{2}$. For example, the superpartner of the electron is a spin-0 particle called the scalar electron, usually contracted to “selectron.” The fermionic superpartners of bosons instead receive the “-ino” suffix, for example “higgsino.” If supersymmetry is realized at the TeV scale, it would resolve many of the Standard Model’s problems, and would be experimentally testable using TeV-scale particle colliders. This section introduces some of the theoretical reasons for the popularity of supersymmetry, and how it may be targeted experimentally.

1.3.1 Theoretical Appeal

A full discussion of the theoretical appeal of supersymmetry is beyond the scope of this thesis, so this section briefly discusses only a pair of the most compelling points that are most closely connected to why supersymmetry at the TeV scale is of great interest. First, supersymmetry provides a well-motivated dark matter candidate, a concrete realization of the WIMP Miracle. Second, it resolves the Hierarchy Problem and is almost certainly the only solution that possibly can; if supersymmetry is not realized at the TeV scale, it likely means that the Standard Model’s fine tuning has somehow been misinterpreted and is not in fact a problem.

Supersymmetry's dark matter candidate arises as a side effect of a solution to an unrelated problem. Naive implementations of supersymmetry make possible the decay $p \rightarrow \pi^0 e^+$ via mediation by a squark (the pion is the lightest meson and the only hadron lighter than the proton). The proton would be highly unstable for any reasonable squark mass, in conflict with the lower limit on this decay mode of some 10^{32} years, and the any-mode limit of more than 10^{29} years [73]. However, it is possible to have a stable proton within supersymmetric models, by proposing that there exists a conserved multiplicative quantum number called R-parity and assigning superpartners odd R-parity and Standard Model particles even R-parity. This saves the proton, but as a side effect also makes the lightest supersymmetric particle (LSP) stable, since decays to final states including other superpartners are kinematically impossible and decays to final states containing only Standard Model particles cannot conserve R-parity [73].

At first, LSP stability appears to be a problem as severe as proton decay. If the LSP is stable, then there should be a large relic population persisting from the primordial plasma, indeed this relic population may represent the majority of the mass of the Universe, but no such population has ever been observed. However, astrophysical observations have assembled a compelling case that a majority of the mass of the Universe is not of the Standard Model, making it possible that stable LSPs *have* been observed, as dark matter, via their gravitational interactions. Of course, this position is only tenable if one of the superpartners has the properties necessary to be a dark matter candidate. The superpartners of neutrinos, the sneutrinos, are an obvious first guess, but their interaction cross sections are too large in generic implementations of supersymmetry, so that a sneutrino LSP requires a special purpose-built model to be viable [74]. Fortunately, the neutral superpartners of the weak bosons, the neutralinos, could have even smaller interaction cross sections than sneutrinos, and are ideal dark matter candidates. Moreover, as part of supersymmetry's solution to the Hierarchy Problem, the neutralinos in general and especially the higgsino are constrained to be relatively light [56, 15, 14], making a neutralino LSP plausible on more than phenomenological grounds.

The supersymmetric solution to the Hierarchy Problem is founded on a simple observation. In Equation 1.3, the Higgs' coupling to itself, a boson, and to the top quark, a fermion, appear with opposite signs. This is general: the contributions of bosons and of fermions to the Higgs mass have opposite signs. Among the Standard Model fields, this is not very helpful, since the top quark simply dominates the other fields. But, if every boson had a fermion counterpart with exactly the same properties and vice versa, all contributions to Equation 1.3 would exactly cancel. Of course, experiments have long excluded the existence of, say, a selectron of the same mass and interactions as the electron, so supersymmetry must be spontaneously broken if it exists, with the superpartners pushed to a higher mass scale. A variety of well-motivated mechanisms for this have been proposed [59, 56]. Whatever the mechanism, the scale of supersymmetry breaking would lead directly to the observed Higgs mass, implying that superpartners and especially the higgsino should begin to appear at energies not far beyond the observed Higgs mass [56, 59, 15, 14].

1.3.2 Experimental Signatures

Since every supersymmetric particle must decay to the nearly undetectable LSP, the primary signature of superpartner production in colliders is large missing energy and momentum, carried away by the undetectable LSP. Additionally, conservation of R-parity requires that superpartners be produced in pairs, so that these events are characterized by *two* heavy invisible particles. As experimental constraints on the visible superpartner masses increase, the likelihood that the decays are extremely energetic also increases. Two particles with large masses of perhaps 2 TeV would be produced only rarely in state of the art accelerators, but their events would be extraordinarily high energy, unless the LSP mass is also large and absorbs most of the decay energy.

By definition, the LSP is the most energetically accessible superpartner. Even so, the preferred supersymmetry search strategy is to use the invisibility of the LSP to identify decay chains of more massive but more strongly interacting superpartners, as the LSP itself must have a

very low production rate due to its weak interaction cross section, if it exists. At hadron colliders, the best candidate superpartners are those participating in the strong interactions, squarks and gluinos, due to the relatively large production cross sections for strongly interacting particles (see Figure 1.12). Unfortunately, while the mass of the LSP is restricted to a relatively low value if supersymmetry is to solve the Hierarchy Problem, the masses of other superpartners are not as strongly constrained. Attempts to observe the LSP in decays of more easily produced superpartners require that these intermediate states have accessible masses, which is not necessarily the case for any experiment to date [15, 14]. It may be necessary to search for direct LSP production if the other superpartners are too massive, which will require a vast dataset due to the low production rate.

As the present datasets are still too small for highly sensitive direct LSP searches, most current analyses are two-dimensional, targeting a relatively easily pair-produced heavy superpartner and the LSP ultimately produced at the end of its decay chain.

1.3.3 Simplified Models

Due to the large number of particles and free parameters in even the simplest supersymmetric extensions of the Standard Model, it is not possible to make exact calculations that hold for every possible realization of supersymmetry, and every supersymmetry simulation is to some extent dependent on the exact choices of free parameters. To make calculations tractable, the community has widely adopted use of simplified models [9]. Simplified models are effective field theories that assume all superpartners except those of interest are at an inaccessible mass scale. Effectively, instead of extending the Standard Model with every superpartner simultaneously, only a few particles, perhaps only the gluino and the LSP, are added. This makes it possible to calculate pair-production cross sections as a function of mass, allows for the decay channel(s) to be set without worrying about the model’s self-consistency, and so on without excessive computational effort. The gluino pair-production cross section as a function of mass in such a simplified model

is shown in Figure 1.12. Results expressed in these models are straightforwardly reinterpreted as

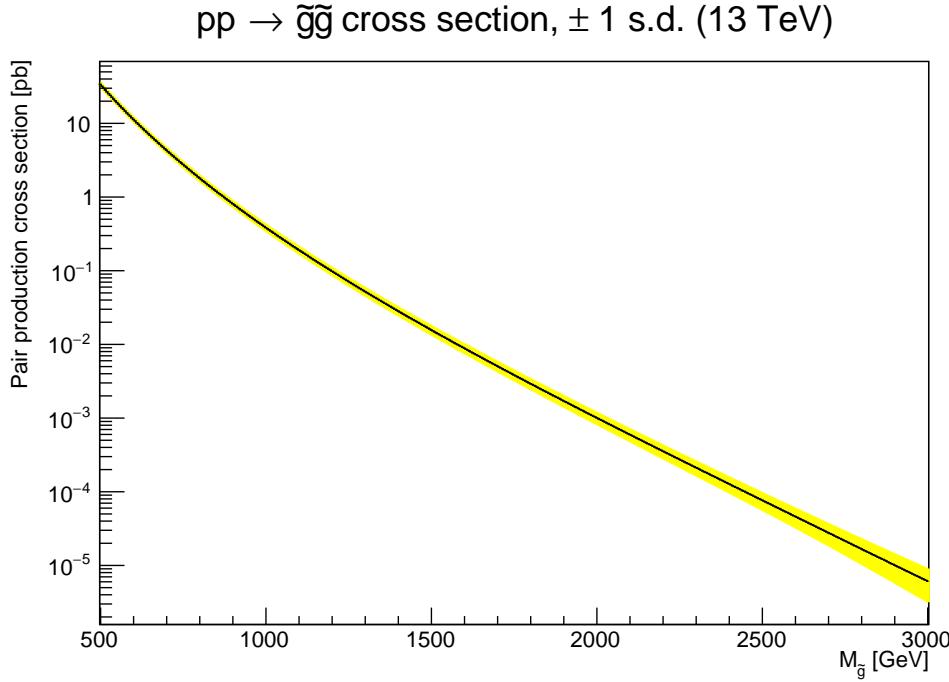


Figure 1.12: In simplified models, it is possible to calculate superpartner production cross sections from theory. Here, the theoretical gluino pair production cross section in 13 TeV proton-proton collisions is shown in black, with the one standard deviation uncertainty shown as a yellow band. The cross section drops rapidly with increasing mass. Based on cross section values used in [30], calculated in [20]. Compare Figure 5 (upper) from [20], of which this is a simplified reproduction.

general constraints on more realistic models with more complicated phenomenologies.

1.4 Other Models

While supersymmetry is generally of greatest interest and so drives the design of many searches for new physics, it shares its primary signatures with some other hypothetical extensions of the Standard Model. Two are worth special mention, since they are considered explicitly by the analysis discussed in Section 3.1.

Some models propose the existence of “leptoquarks,” strongly-interacting bosons that have vertices in which one fermion is a lepton and one is a quark, with one example shown in

Figure 1.13. These models have recently attracted some interest due to their ability to explain some minor anomalies in a few rare meson decays [61, 4, 5]. As leptoquarks can be pair-produced and decay to a neutrino and a quark, they produce a very similar experimental signature to a squark decaying to a quark and LSP, in the limit where the LSP is nearly massless. Therefore, searches for squarks decaying to LSPs can easily be reinterpreted as searches for leptoquarks.

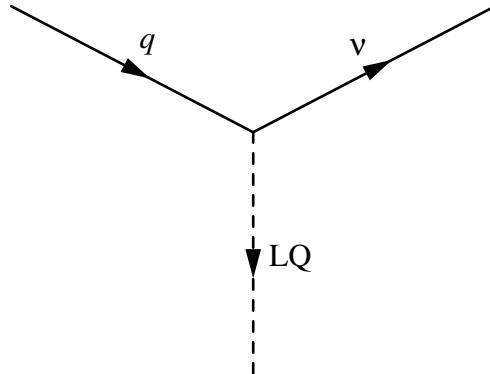


Figure 1.13: Leptoquarks are defined by vertices that connect quarks and leptons. In this example, the lepton is a neutrino.

A second model, this time motivated by a small excess in events with large missing energy and relatively low energy, proposes a generic strongly-interacting scalar ϕ mediating Standard Model interactions with a fermionic dark matter candidate ψ [13] as shown in Figure 3.2, with parameters tuned to match the minor observed anomaly. This model need not involve pair-production and could in principle have a smaller mass scale than supersymmetry, and so can produce events with lower energy and lower numbers of hadronic jets than would be typical for supersymmetric models. While not nearly as well-motivated theoretically as supersymmetry, it provides a framework for analyzing anomalies in relatively low activity kinematic regions that superpartner pair-production events cannot easily populate.

In general, while supersymmetry drives the design of many searches for new physics, the signatures of supersymmetry are in practice very general, so that searches nominally targeting supersymmetry are in practice sensitive to a wide variety of models of new physics. Supersymmetric models are used as a benchmark means of communicating results that is familiar throughout

the elementary particle physics community. Ultimately, the focus of experiments is on providing evidence to support or refute any potential explanations of observations that conflict with the Standard Model prediction.

Chapter 2

The Large Hadron Collider and the CMS Detector

The searches for new physics discussed in Chapter 3 analyze data recorded by the Compact Muon Solenoid (CMS) experiment. The CMS detector straddles beams generated by the Large Hadron Collider (LHC), a particle collider located at CERN in Geneva, Switzerland. CMS directly measures the products of collisions between the protons in these beams, digitizes the measurements, and records a small subset of these events. The data is then processed by reconstruction algorithms and shared across the worldwide CMS computing grid. This chapter describes the LHC and CMS.

2.1 The Large Hadron Collider

A detailed description of the LHC is available in Reference [19]. This section provides a summary of the essential information.

2.1.1 A Brief Description

Although it sometimes collides atomic nuclei, typically Pb-208, the LHC is primarily a proton-proton collider, and that is the operational mode considered in this work. The LHC is the highest energy and highest luminosity proton-proton collider ever constructed, thus far achieving collisions with center of momentum frame energy as large as 13 TeV, produced by two counter-rotating beams of 6.5 TeV protons, at sustained rates on the order of one billion collisions per second.

Producing such high energy collisions at such high intensity requires that the LHC be a large ring, 27 km in circumference. The circular shape allows multiple opportunities for the protons in each beam to collide, rather than a single all-or-nothing pass, and additionally allows the beams to be circulated for as long as necessary to accelerate them up to the intended energy, at a rate of 0.5 MeV per revolution. The size makes it possible for the protons to be confined to a circular path, even at very high energies. The beams are constrained to travel in the ring using powerful magnetic fields, and the maximum feasible magnetic field strength is the primary limit on the maximum achievable beam energy. The LHC beams require an 8.36 T field to stay on track, provided by superconducting niobium-titanium electromagnets carrying a current of 11,080 A. Superconductors stable at greater magnetic field strengths are rare and expensive, and the magnets must be superconducting to carry such enormous currents. Larger rings reduce the beams' curvature and allow for greater energies with the same magnetic field strength.

Protons are chosen for the beams rather than electrons, the only other obvious potential choice, for two reasons. First, protons are composite particles, so their total energies are distributed among their many subcomponents, referred to as partons in this context (see Section 1.1.4). Collisions between protons with energy 6.5 TeV each are in fact collisions between a parton in each proton with some fraction of the total energy. This naturally scans across the accessible energy range, providing sensitivity to any new state with mass on the TeV scale. Were the beams composed of elementary particles like electrons, the collider would need to scan the beam energy

to provide sensitivity to any possible mass. Therefore, electron colliders are poorly suited for new discoveries compared to proton colliders, and better suited for studying a state with known mass in greater detail. Second, and more importantly, the mass of the proton is much greater than that of the electron, reducing energy losses to synchrotron radiation. Because synchrotron losses are relatively small, circular colliders of protons are limited chiefly by the magnetic field strength, as discussed above. Electron colliders, on the other hand, are limited by synchrotron losses. Illustrating this, the Large Electron Positron (LEP) collider previously occupied the same tunnel that now houses the LHC, and managed a maximum collision energy of only 209 GeV [51].

The protons are not collided one-by-one, as this would be impossibly tedious. Instead, they are grouped in bunches of over 100 billion, spaced by 25 ns to make it possible for the detectors to distinguish the products of one bunch crossing from those of the preceding and following crossings. During each bunch crossing, a few dozen proton pairs (see Figure 2.13) collide simultaneously, of which either one or none are interesting and the rest are undesirable noise called pileup, discussed in Section 2.3.1. These bunch crossings occur at a few points around the LHC tunnel occupied by detectors. Among them are the CMS detector, described in Section 2.2, and its sister detector ATLAS, which are both general-purpose detectors designed to study any and all phenomena at the TeV scale, and operate in parallel to provide independent checks of any measurements and discoveries.

2.1.2 Luminosity Delivered

The sensitivity of searches performed in CMS scales straightforwardly with the total number of collisions provided by the LHC. Thus far, the LHC has provided CMS with nearly 200 fb^{-1} across all operating periods, with the rate accelerating as the machine is pushed to the limits of its design, as shown in Figure 2.1. The searches discussed in the next chapter use only the 137 fb^{-1} recorded by CMS during the 13 TeV runs from 2016 to 2018. As the total inelastic proton-proton collision cross section is approximately 69.2 mb at 13 TeV, the total number of

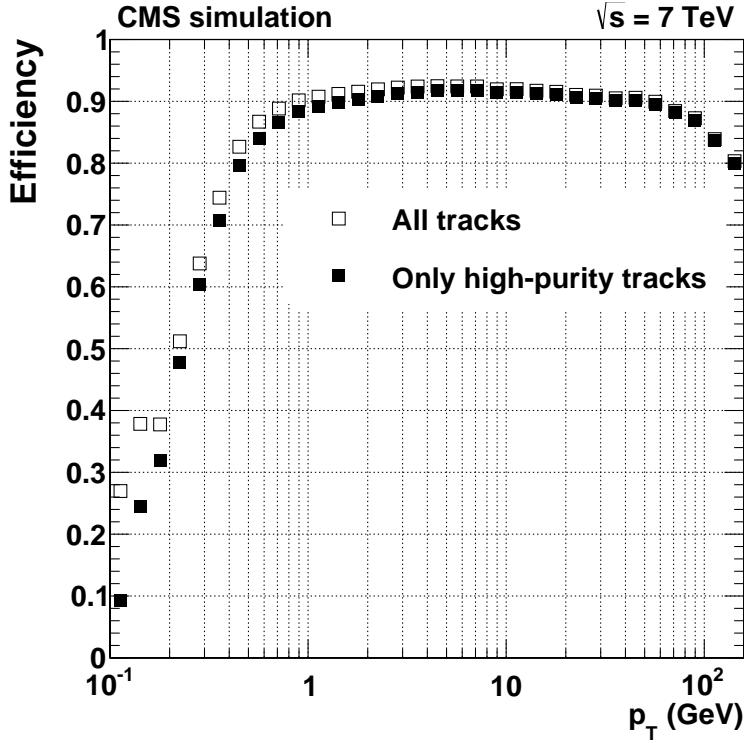


Figure 2.1: The integrated luminosity delivered to CMS over time by the LHC, in blue, and the portion CMS managed to record, in yellow. As can be seen, CMS is not 100% efficient at recording collisions. It is important to note that this inefficiency is a straightforward failure to observe the events, *not* the purposeful rejection of events by the trigger system discussed in Section 2.3.1. Taken from [1].

collisions observed by CMS from 2016 to 2018 is nearly 10^{16} . While this vast sample size makes very rare genuine physics processes potentially observable, it also presents a challenge in that very rare and difficult to model detector failures may pollute the dataset. It is prudent to adopt a conservative approach to any searches for rare new physics.

2.2 The CMS Detector

While the LHC provides the collisions, it is the Compact Muon Solenoid (CMS) detector that ultimately measures the products and reconstructs each event. A detailed description of all of the CMS hardware and software is available in Reference [18], and a description of the general

event reconstruction technique used at CMS, called Particle Flow, is available in Reference [68]. This section serves as a general overview of how CMS works, sufficient to understand the searches for new physics presented later in this work.

The particles stable enough to be detected directly by CMS come in 5 distinct varieties, namely electrons, muons, photons, charged hadrons (largely protons, pions, and charged kaons), and neutral hadrons (largely neutrons and neutral kaons). Figure 2.2 depicts a generic detection of each kind of particle. All other particles are too unstable to be measured directly, except the neutrinos. Neutrinos are, uniquely among Standard Model particles, both stable and invisible to CMS. Their presence in the products of a collision is inferred by measuring all of the visible content of an event, and noting that some energy is missing, as discussed in Section 2.2.6. Any stable weakly interacting particles beyond the Standard Model produced in a collision at CMS would present a similar signature.

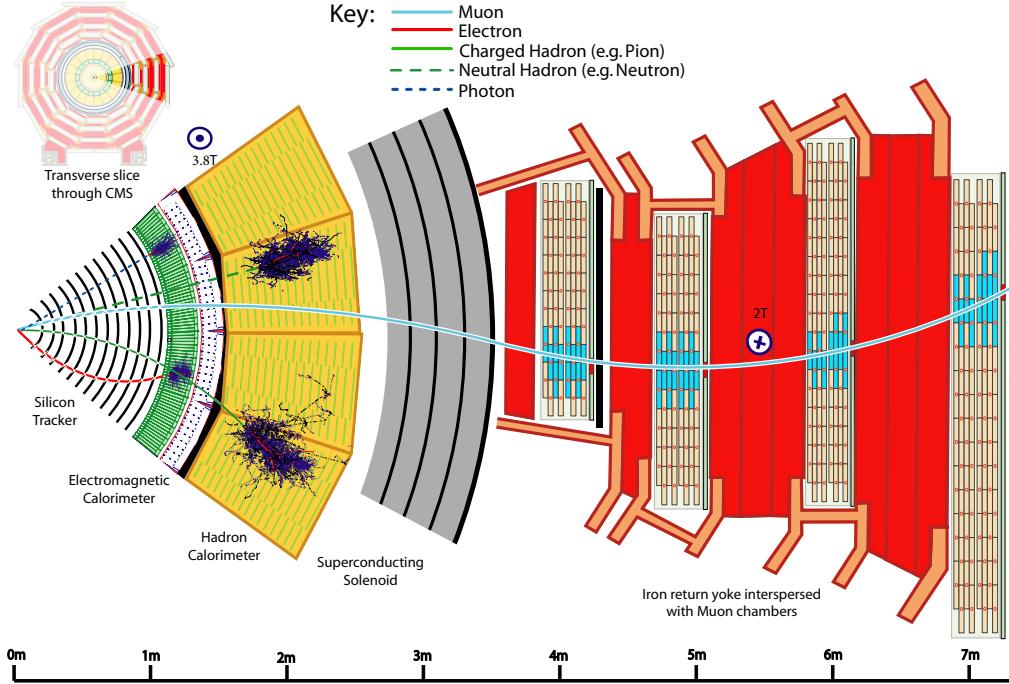


Figure 2.2: The design of CMS causes particles of different types to leave distinct experimental signatures, allowing them to be distinguished. The rest of this section references and explains this image. Taken from [68].

The CMS detector is a cylinder composed of a central barrel section with two flat endcaps, each of which contain layers of different detection systems designed to provide the information necessary to distinguish the various particles that can be generated by collisions, and measure their positions and energies. Starting from the beam pipe, they are the tracker, the electromagnetic calorimeter, the hadronic calorimeter, and the muon system. The tracker and the calorimeters are inside a solenoid that produces a powerful magnetic field, while the muon system is subject only to the external, returning field of the solenoid. Altogether, the detector measures about 7.5 meters in radius and 21 meters in length. As suggested by the cylindrical detector shape, the experiment uses a cylindrical coordinate system, in which the beam axis is assigned to the z axis, and the angle in the transverse plane, ϕ , is measured starting from a line connecting CMS to the center of the LHC ring. The trajectories of particles are usually described using this angle ϕ alongside the pseudorapidity η ,

$$\eta = -\log \left[\tan \left(\frac{\theta}{2} \right) \right] \quad (2.1)$$

where θ is the angle made with the z axis [18]. This usage is motivated by the fact that the average flux of particles produced by collisions in a window of fixed $\Delta\eta = \eta_2 - \eta_1$ is roughly constant for any η_1 and η_2 .

This section describes each of the subsystems, as well as the objects that are constructed using their measurements.

2.2.1 Magnet

As its name suggests, a major feature of the CMS detector is a superconducting solenoid, the largest ever built, which provides a constant 3.8 T magnetic field in the tracker and calorimeters. This magnetic field curves the trajectories of charged particles, allowing for a determination of their charges and assisting in the measurement of their masses and energies, contributing to highly accurate identification of particles.

CMS is compact in part to make it possible to fit the bulk of the detector inside this single solenoid. This maximizes the strength and consistency of the magnetic field, which in turn maximizes the precision with which particle positions, and the curvatures of their trajectories, can be measured.

In Figure 2.2, one can see the path of an electron (orange) curving sharply, in a direction indicating that it is negatively charged. Similarly, the path of a charged hadron (solid green) curves in the opposite direction, indicating a positive charge. A positively charged muon (blue) curves first one way, then the reverse once the field flips after exiting the solenoid. Neutral particles (dotted lines) do not curve at all.

2.2.2 Tracker

The silicon tracker is the innermost detector system and is primarily responsible for measuring the positions of charged particles. In Figure 2.2, the trajectories of neutral particles are represented with dotted lines, to indicate that they do not actually leave those tracks in the detector. The tracker is critically important for the disappearing tracks search covered in Section 3.2, and so is described here in more detail than the other detector systems.

As shown in Figure 2.3, the tracker is composed of several discrete layers. When a charged particle passes through a layer, it usually causes a measurable current to flow in the silicon, and is said to have produced a “hit.” As shown in Figure 2.4, the probability for a genuine particle to leave a hit is near but not quite unity, so a missing hit where one would be expected is credible, but uncharacteristic of genuine tracks. Each event produces an intimidating array of these discrete hits throughout the tracker, which must then be connected by the track reconstruction algorithm into particle tracks, each consisting of at least three distinct hits. CMS uses an iterative procedure described in Reference [33], in which clean, high energy, isolated tracks are reconstructed first so that their hits may be removed from the pool, followed by the next most obvious tracks, and so on, dramatically simplifying the computational complexity of each step. The early iterations

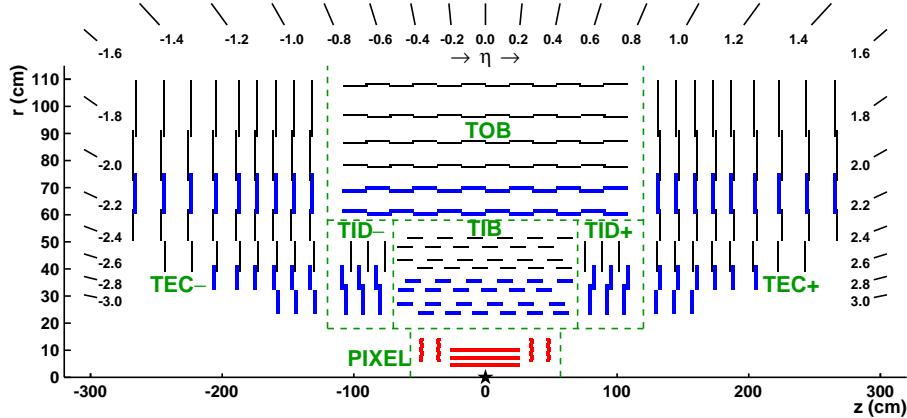


Figure 2.3: A slice of the CMS tracker in the $r - z$ plane, with the inner pixel tracker (red) in the 3-layer configuration used until 2016. In 2017 and 2018, the pixel tracker was upgraded with a 4th layer. The tracker is symmetric about the beam axis ($r = 0$); one should imagine rotating this image around the horizontal axis to form a cylinder, to visualize the tracker’s 3-dimensional form. The labels in green text denote tracker subsystems, the Tracker End Cap, the Tracker Inner Disk, the Pixel Tracker, the Tracker Inner Barrel, and Tracker Outer Barrel. The interior subsystems have more miniaturized detector elements and, therefore, better resolution. Taken from [33].

that produce the tracks judged to be of highest quality find only those tracks with transverse momenta at least 0.3 GeV, and no iteration can produce a track with transverse momentum below 0.1 GeV, so that particles of such low energy are effectively invisible to CMS. Note that, due to the strong magnetic field present in the tracker, charged particles with energies lower than around 0.1 GeV travel in such tight spirals that reconstruction becomes unfeasible. Although this is normally insignificant, since these particles are not typically important to an event’s overall characterization, this fact allows for the existence of disappearing tracks, in which a high energy charged particle decays to a trackless neutral particle and a charged particle of such low energy that it is not reconstructed.

The performance is outstanding, especially for tracks passing the sophisticated “high purity” selection applied by the vast majority of CMS analyses. Figure 2.5 shows the fake rate of this procedure, less than 5% for tracks in the intermediate energy range, and further suppressible with additional cleaning selections like those of the disappearing tracks search

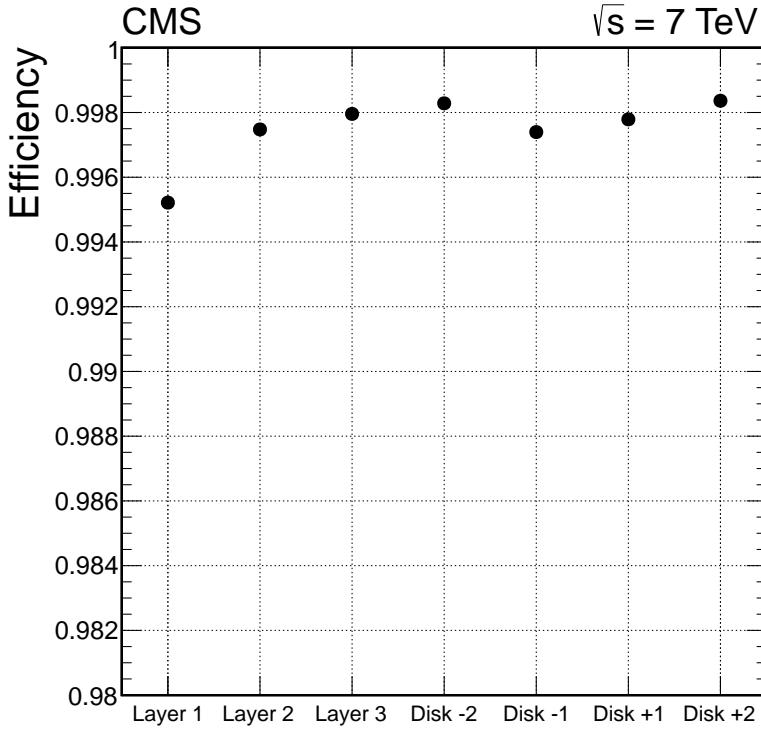


Figure 2.4: The efficiency for a genuine charged particle to produce an expected hit in the various components of the CMS pixel tracker, in the configuration used in 2016 and earlier. Fake tracks may have many missing expected hits, but more than one or perhaps two is extremely rare in tracks produced by genuine particles. The probability for a neutral particle to produce a hit is small, and the probability to leave enough to produce a track is negligible. Taken from [33].

covered in Section 3.2. Figure 2.6 shows the efficiency to reconstruct tracks successfully for the three most common varieties of charged particles measured by CMS, namely muons (left), electrons (center), and pions (right), approaching unity in all cases for higher energy tracks in the barrel region of the tracker.

The tracker detection elements are smaller and consequently provide better localization in interior layers of the tracker, with the pixel detector components only approximately $100\text{ }\mu\text{m}$ across. As a general guideline, no element in the tracker is permitted to have an expected probability of a hit greater than 1% per bunch crossing, to maintain a low probability of tracks intersecting. This structure produces outstanding sub-millimeter resolution of a track's origin, shown in Figure 2.7, allowing tracks originating from collisions only microns apart to be separated,

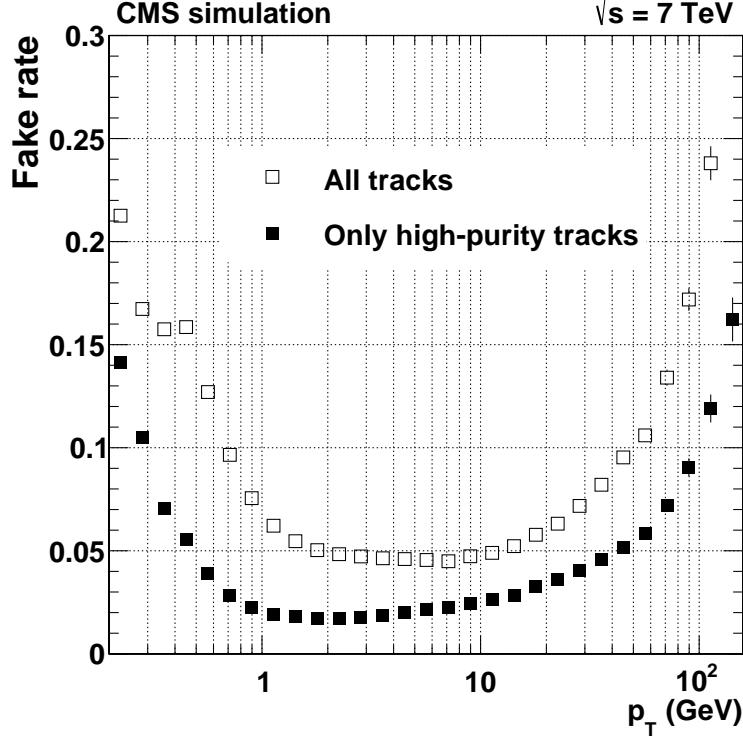


Figure 2.5: The CMS track fake rate as a function of p_T , in the configuration used in 2016 and earlier. The sophisticated high purity selection aims to suppress the fake rate without badly affecting efficiency and is applied in the vast majority of CMS analyses. Taken from [33].

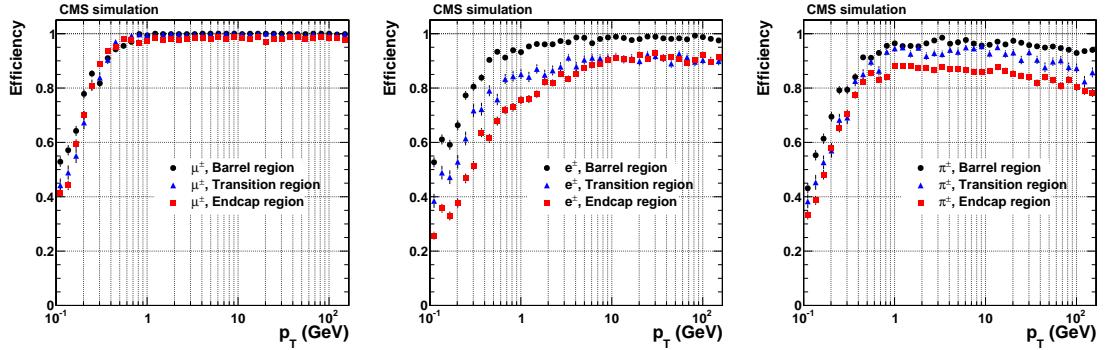


Figure 2.6: The CMS track reconstruction efficiency for (left) muons, (center) electrons, and (right) pions, split by detector region as a function of p_T , in the configuration used in 2016 and earlier. Muons are by far the cleanest objects at CMS as they are massive enough to avoid strong perturbations by tracker material, and lack nuclear interactions. Electrons are lighter and so are vulnerable to tracker interactions, experiencing strong bremsstrahlung losses, while pions can undergo nuclear interactions. Taken from [33].

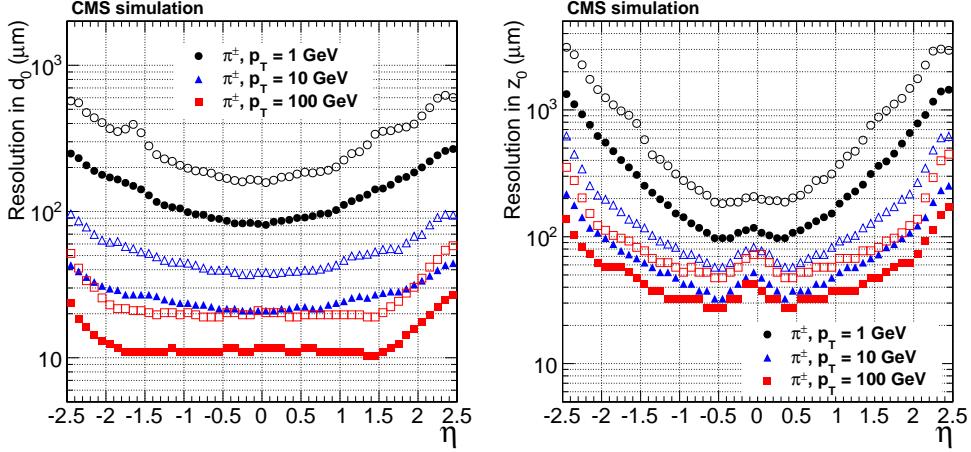


Figure 2.7: The CMS tracker has micron-level impact parameter resolution, for both (left) the transverse plane and (right) the z direction. Solid markers indicate the half-width of the 68% confidence interval, while open markers indicate 90% confidence. With this resolution, a track reconstructed with an impact parameter of more than a few hundred microns is likely a fake, from a pileup vertex, or from a secondary vertex produced by a decaying particle like a bottom hadron. Taken from [33].

and tracks originating from decays of particles with decay lengths on the order of a millimeter to be identified. The former is important for rejection of pileup, and the latter for identification of jets originating from bottom quark hadrons.

The CMS pixel tracker was upgraded after 2016 for the 2017 and 2018 runs, increasing the number of layers from three to four. Comparisons between the track reconstruction efficiency and fake rate in 2016 and 2017 are shown in Figure 2.8. The new pixel tracker significantly improves performance, for all tracks.

Vertices and b-Tagging

Almost every track ultimately originates from a single point in a proton-proton collision. Tracks reconstructed in the tracker can be extrapolated back to their origins to form vertices, locations at which multiple tracks intersect and the inferred location of the initial hard interaction between partons. All the content associated with a single vertex constitutes a single event, subject to analysis in an attempt to determine the nature of the physics that occurs in the hard interaction.

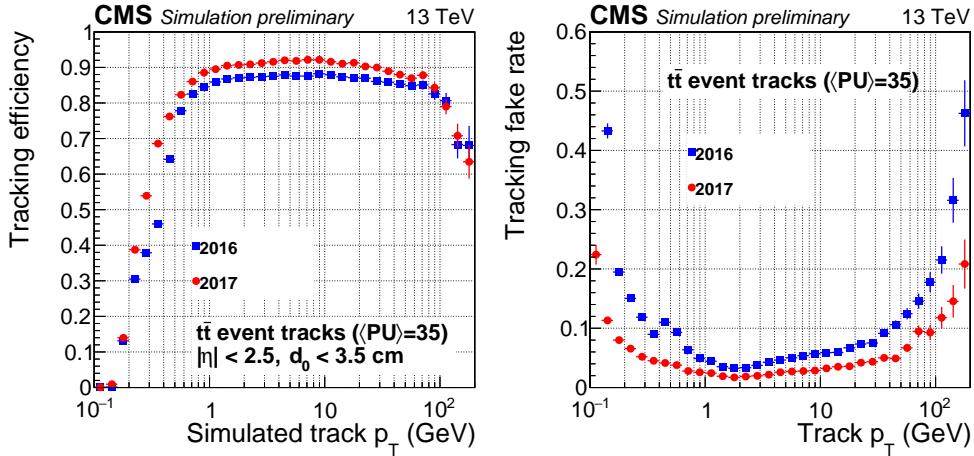


Figure 2.8: A comparison of the performance before (blue squares) and after (red circles) the pixel tracker upgrade following the 2016 data taking period. The upgrade increased the efficiency (left) and reduced the fake rate (right) for tracks of all momenta. Taken from [28].

Some tracks do not extrapolate to a proton-proton primary vertex back but instead a point a few millimeters away. While a few millimeters is not even enough to exit the beam pipe, such a decay length is nevertheless well within the tracker’s resolution. These features are called secondary vertices, and are a distinctive feature of jets associated with hadrons containing bottom quarks. These hadrons, uniquely, have decay lengths on the order of millimeters. Charmed hadrons tend to decay somewhat more quickly, while hadrons containing only up, down, and strange quarks tend to live long enough to reach the calorimeters. These secondary vertices, along with some other indicators including the muon content of a jet, the number of tracks, and the mass, are leveraged by sophisticated algorithms that can identify jets originating from bottom quarks produced in the hard process with high efficiency [29]. Such jets are said to be “b-tagged.” The fake rate for a given efficiency achieved by various b-tagging algorithms is shown in Figure 2.9, with the older algorithms listed first and the newer, machine learning-based algorithms last. The analyses discussed in Section 3 apply the DeepCSV algorithm (purple) at the 1% light quark fake rate working point, which achieves a successful b-tag rate of nearly 70% and a charm quark mistag rate of just over 10%.

Although the ability to identify a jet’s source flavor is intrinsically useful, the identification

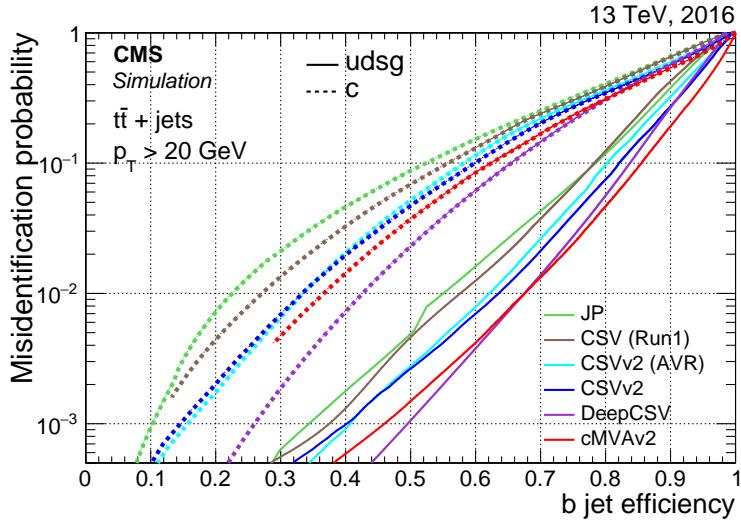


Figure 2.9: The fake rate of various b-tagging algorithms (different colored lines) are plotted against the successful tag rate for genuine b-jets. Although charm quarks are distinguishable from bottom quarks in a variety of ways, including shorter lifetimes, they are still similar enough to cause a meaningful fake rate. These curves are plotted as dotted lines. The lighter quark fake rate is much lower, shown as solid lines. The algorithm of most interest is DeepCSV, a machine learning algorithm used in the analyses discussed in Section 3. Taken from [29].

of bottom quarks is of special interest because they tend appear in the majority of Higgs and top quark decays, and for their relative rarity in generic QCD multijet events compared to lighter quark flavors, which means that requiring many b-tags greatly suppresses background. Supersymmetry searches consider the case of bottom and top superpartners separately due in large part to the number of b-jets expected in the final state, and the consequently different, smaller backgrounds that dominate these models.

As a general rule, the tracker intends to interact with particles only to the extent absolutely necessary to measure their positions. Strong interactions between the tracker and the particles it is measuring modify the particles' tracks and affect their energies, and are therefore undesirable. Unfortunately, electrons are so strongly affected by bremsstrahlung that the tracker is quite thick in terms of radiation lengths, as shown in Figure 2.10 (left). Therefore, electrons at CMS are experimentally observed more as clouds of bremsstrahlung photons by the time they reach the electromagnetic calorimeter at CMS, and must be reconstructed with a special tracking procedure,

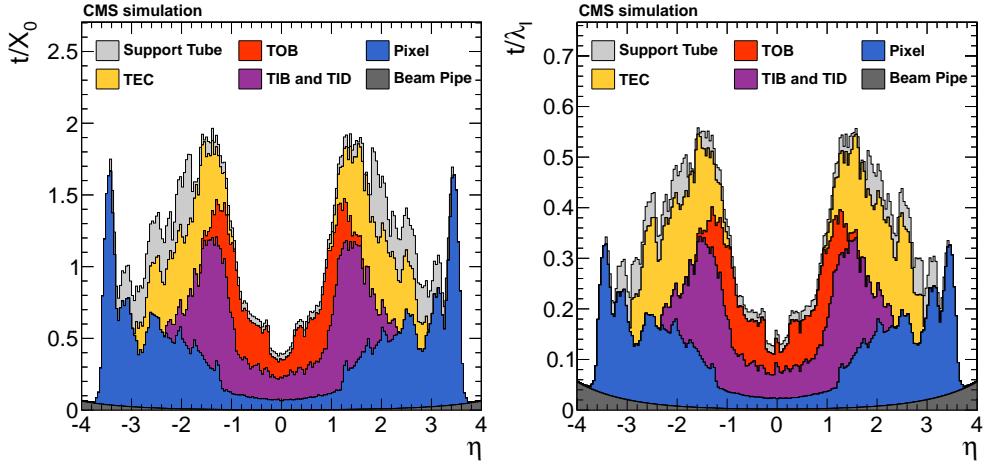


Figure 2.10: The CMS tracker material budget in (left) radiation lengths (ie electromagnetic) and (right) interaction lengths (ie nuclear), in the configuration used in 2016 and earlier. While the CMS tracker is not intended to shower particles, only to measure positions as gently as possible, some amount of interaction between tracker material and collision products is inevitable. Taken from [33].

the Gaussian-sum filter described in [6]. In terms of interaction lengths, the CMS tracker is much thinner, as shown in Figure 2.10 (right). The vast majority of hadrons therefore reach the calorimeters without having undergone a nuclear interaction. Figure 2.11 shows that over 80% of high energy charged pions reach the calorimeters intact. Even so, 80% is not 100%. Pion tracks occasionally terminate in a shower induced by a collision with a nucleus inside the tracker. In rare cases, these showers can contain *zero* visible particles of any kind. These disappearing pions are an important background source for the disappearing tracks search in Section 3.2.

2.2.3 Electromagnetic Calorimeter

The primary purpose of the electromagnetic calorimeter (ECAL) is to observe the energies of photons, which do not leave tracks. It is also important in the observation of electrons, which produce a jet of photons due to bremsstrahlung in the tracker. Other electromagnetically interacting particles will also leave small ECAL deposits, but much less due to their lesser tendency to produce photons. In Figure 2.2, the electron (red) and photon (dotted blue) produce

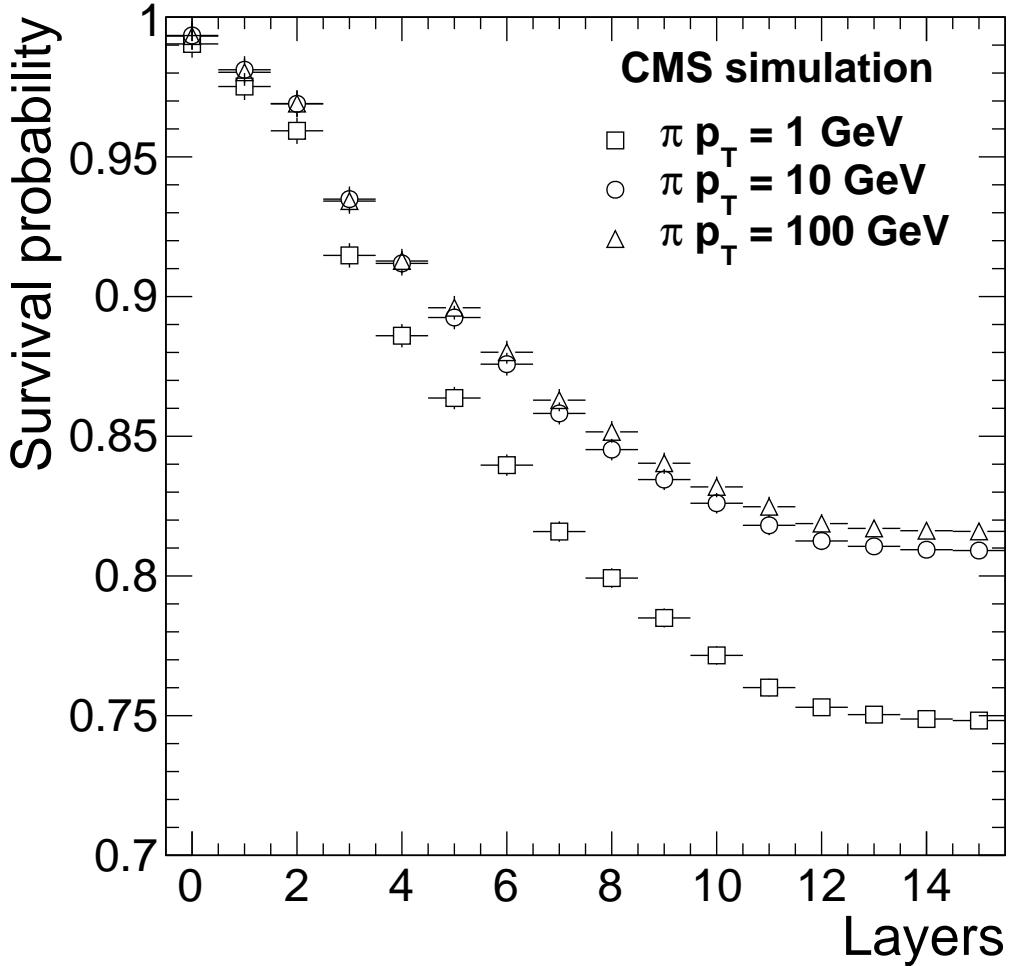


Figure 2.11: The probability for pions to undergo nuclear interactions in the CMS tracker by layer, in the configuration used in 2016 and earlier. In rare cases, the pion may shower in such a way that no charged daughters have high enough energy to be reconstructed, causing its track to disappear. Taken from [33].

strong showers in the ECAL (green), while the other particles leave negligible deposits.

The design of the CMS ECAL prioritizes good resolution of photon energies and positions over all else, due to the experimental importance of the Higgs boson's decay to a pair of isolated photons. To accomplish this, the ECAL is constructed from tens of thousands of scintillating lead tungstate crystals, which have several desirable properties.

First, these crystals have a Moliere radius of only 2.2 cm [18]. Aside from providing generally good position resolution, this small Moliere radius most importantly allows a prompt

isolated photon, perhaps from a Higgs decay, to be distinguished from two closely overlapping photons produced by the decay of a neutral pion, because the two separate showers produced by the photon pair are clearly separated in the ECAL cells. These two nearby showers may merge into a single large shower indistinguishable from that of a single photon in a material with a larger Moliere radius. Similarly, photons produced inside jets can be distinguished from truly isolated photons, since the showers of other particles in the jets can be discriminated from the photon's own shower. Although the size of the shower is restricted, a typical one will still extend over a few adjacent crystals as shown in Figure 2.2, so that a 5-by-5 cluster of crystals, or potentially a smeared supercluster for electrons that have radiated over a larger area, is the essential object of ECAL reconstruction. Studying the distribution of energy measurements across the crystals in a cluster, as described in [26], allows for better fake rejection, and energy resolution of 1% [25].

Second, lead tungstate also has an extremely short radiation length, only 0.89 cm compared to the roughly 2 radiation lengths of the entire 1 m radius tracker, so it contains electrons very efficiently [18].

Third, scintillation in lead tungstate happens quickly, with 80% of the light produced within 25 ns [18]. This speed is critical for coping with the 25 ns bunch spacing of the LHC's proton beams, so that the ECAL deposits produced by adjacent bunch crossings can be distinguished. Due to this rapid shower production and careful calibration of the electronics, the ECAL achieves a timing resolution of roughly 0.2 ns for particles with momenta on the order of tens of GeV [36].

Finally, lead tungstate is resistant to radiation damage, so that the ECAL can continuously operate in the high flux environment of LHC proton collisions. Unfortunately, it is not immune to radiation damage, and so the performance of the ECAL degrades over time. By the end of data taking in 2018, roughly 20% of ECAL cells no longer consistently detected showers of photons and electrons. This leads to a background for the disappearing tracks search in Section 3.2 consisting of electrons that transform to photons inside the tracker, no longer leaving a track, but are not later detected as photons in the ECAL. Fortunately, 80% of the ECAL is still performing

well despite years of high radiation flux, so this background can be almost entirely rejected by mapping and vetoing the poorly performing cells, using events in which a Z boson decayed to a pair of electrons, exactly one of which is lost due to this effect, as described in Section 3.2.

2.2.4 Hadronic Calorimeter

Charged hadrons are too massive to shower significantly in the ECAL, and neutral hadrons are invisible to both the ECAL and the tracker. However, both can be made to shower by capitalizing on their powerful nuclear interactions, and the hadronic calorimeter (HCAL) accomplishes this, as depicted in Figure 2.2 for a neutral hadron in dotted green and a charged hadron in solid green. The performance of the HCAL is very important to analyses selecting events containing jets and missing energy, as HCAL measurement errors lead directly to errors in jet energies, which in turn cause errors in the inferred missing energy in the event.

The essential plan of the CMS HCAL design is to induce hadron showers by putting a high density of atomic nuclei in their paths in the form of 5 cm brass plates, then to measure the products of these showers using plastic scintillators, and from this reconstruct the energies of the original hadrons. To ensure that as few hadrons as feasible escape detection, whether from the primary interaction or from secondary showers, there are multiple layers of this shower-inducing metal plate and scintillator combination. Nevertheless, some hadrons do manage to “punch through” the HCAL, especially if they have extremely high energy, and even successful measurements are only accurate to within roughly a factor of 2 [34], unless the hadron has very high energy. On the whole, the CMS HCAL performance is the weakest of any detector system. However, the energy of a typical jet is not measured only with the HCAL, and errors made on different hadrons within a jet tend to cancel, so that CMS achieves a typical jet energy resolution of approximately 10%, for the jets of most interest [18]. Still, jets are occasionally badly mismeasured, most often due to HCAL measurement errors, and it is important for analyses selecting events with large missing energy to filter out these mismeasurement events, see Section 3.1.3.

Accuracy of energy measurements aside, the ability to detect neutral hadrons consistently in the HCAL is important to the disappearing tracks search discussed in Section 3.2. A charged hadron can decay to a neutral hadron inside the tracker, and thus allow known Standard Model physics to produce an apparent disappearing track. Unlike a true disappearing track, however, the neutral hadron is detectable in the HCAL, revealing what happened in the tracker and rejecting this Standard Model background. As with the ECAL, poorly performing regions of the HCAL can be vetoed to ensure this background is wholly rejected.

2.2.5 Muon System

The muon system is the outermost detector of CMS, and the only one outside the solenoid. Muons, like charged hadrons, are too heavy to produce significant showers in the ECAL, and unlike hadrons, do not have nuclear interactions allowing them to be trapped in the HCAL. Instead, they penetrate through the entire detector, leaving a track in the tracker, largely disappearing inside the calorimeters, and then reappearing in the muon system. Unlike other particles, then, the energies of muons must be measured largely by analyzing the curvature of their tracks. As one of the design goals of the Compact Muon Solenoid experiment is ultra-precise measurements of muons, the muon system is large partly in order to provide a long look at the muon's track for momentum measurement purposes, and partly to provide a positive identification of a particle as a muon, rather than a hadron that managed to escape the HCAL, as only muons can penetrate through so many layers of detector material. As installing a second tracker would be far too expensive and excessive, the muon system instead uses only four widely spaced layers of detectors operating based on ionization of gas by passing muons, which while generally worse than the tracker system in terms of space and time resolution, are much more economical for covering such a large volume [18]. The design of CMS has proven successful. Muons are by far the cleanest, most well-measured particles at CMS, with momentum resolution better than 1% and high reconstruction efficiency, as shown for tracker muon tracks in Figure 2.6.

While muons are the only charged Standard Model particle that can consistently penetrate to the muon system, the possibility exists that heavy long-lived charged particles beyond the Standard Model may be produced in LHC collisions and leave hits in the muon system. The disappearing tracks search in Section 3.2 targets models that could potentially produce such a signature. However, that search considers only tracks that disappear inside the tracker, and in fact explicitly vetoes muon-like tracks, so it is not sensitive to this potential signature. Adding sensitivity to this powerful signature could be an interesting extension in the future.

2.2.6 Missing Transverse Energy

Once muons are measured in the muon system, all Standard Model particles have been accounted for, except neutrinos. Unfortunately, the neutrino interaction length, even at LHC energies, is on the order of billions of meters at Earthly densities, so that the probability of a typical LHC neutrino experiencing even a single interaction anywhere inside the CMS detector is negligible [40]. However, the presence of neutrinos can be inferred by measuring all other particles inside the event, and observing that the transverse momentum does not balance, indicating either the presence of at least one undetected particle, or a detector mismeasurement. At a hadron collider like the LHC, only the transverse momentum can be used for this procedure, because the exact center of momentum frame of the primary collision is unknown along the beam axis. The actual colliding particles are not the protons themselves, but rather subcomponents of the proton called partons. The partons are approximately at rest with respect to the detector in the transverse plane, so that their transverse momenta are known to sum nearly to zero, but have unknown momenta along the beam axis, drawn from a probability distribution as described in Section 1.1.4. Unbalanced momentum along this axis may indicate the presence of an invisible particle, but more often it indicates simply that the initial momentum of the collision along the beam axis was nonzero. The observable is not the missing energy, then, but the missing transverse energy, written as MET, \cancel{E}_T , or E_T^{miss} . This observable is tentatively identified with the unobservable

transverse momentum of a neutrino or other invisible particle or particles.

Although the invisible momentum is only available in the transverse plane, it is still a very useful quantity. For instance, consider the leptonic decay of a W boson, $W^\pm \rightarrow \ell^\pm v$. The W boson has mass approximately 80 GeV with width approximately 2 GeV [73]. Even allowing for detector resolution smearing of around 10% in addition to the intrinsic width of the W resonance, a lepton and neutrino system originating from an on-shell W decay can never have mass greater than approximately 100 GeV. It follows directly that transverse mass M_T , the mass calculated using energy and momentum in the transverse plane, also cannot exceed 100 GeV, since $E_T \leq E$ and $p_T \leq p$. The analyses discussed in Chapter 3 leverage this fact to determine whether a lepton or disappearing track may potentially have originated in a leptonic W decay.

Of course, while the missing transverse energy is sometimes the undetected energy of a single neutrino or unknown invisible particle, it can also be a consequence of mismeasurement. In fact, even events with relatively large amounts of missing energy at CMS are most often produced by detector mismeasurement (see Figure 3.4, in which “Multijet” events in yellow have \cancel{E}_T due to detector mismeasurement), even though such large errors are rare, because neutrinos are rarer still. Searches for invisible particles naturally wish to eliminate these events with fake \cancel{E}_T using cleaning selections. The \cancel{E}_T can also potentially be the vector sum of the transverse momenta of *multiple* undetected particles. One of the most powerful observables that can be used to reject events with fake \cancel{E}_T and indicate than an event may have multiple invisible particles is M_{T2} [52], a generalization of M_T .

M_{T2}

Although the presence of invisible particles can be inferred via the observation of a nonzero \vec{E}_T^{miss} , it is not trivial to distinguish events in which \vec{E}_T^{miss} corresponds to the transverse momentum of a single invisible particle, and events in which it is the vector sum of the momenta of multiple invisible particles. It can also be difficult to distinguish the presence of genuine

invisible particles from cases of detector mismeasurement. While no conclusions about invisible content can ever be drawn with perfect certainty for any single event, the situation is not hopeless. There exist tools that can be used to make quantitative inferences concerning an event's invisible content, including the M_{T2} observable. Specifically, M_{T2} is designed to identify events consistent with the presence of two invisible particles in the final state, produced in symmetric decay chains of pair-produced heavy particles. This also tends to reject events in which \cancel{E}_T was produced by detector mismeasurement, as these events tend to have no symmetry whatsoever.

The algorithm begins by dividing the visible portion of the event into two hemispheres. To do this, an event must have at least two visible objects, so in an event with only hadron jets, like the majority of the events at CMS, M_{T2} can be calculated only for multijet events. The method used to assign each visible object to a hemisphere varies. One popular method, and the one used in [30], begins by identifying the pair of objects with largest system mass, M_{12} ,

$$M_{12} = \sqrt{(E_1 + E_2)^2 - (\vec{p}_1 + \vec{p}_2)^2}, \quad (2.2)$$

and assigning each as the seed of a separate hemisphere. In an event containing two identical pair-produced particles undergoing similar decay chains, the pair of objects with largest mass is unlikely to originate from the same particle, as the mass of the full system is strictly greater than the mass of either particle.

Then, each object is associated to the seed of lesser Lund distance, D_L , which is defined for seed i as [69, 70],

$$D_L = (E_i - p_i \cos \theta) \frac{E_i}{(E_i + E)^2} \quad (2.3)$$

where E_i is the seed's energy, p_i is the seed's momentum, θ is the angle between the momenta of the seed and of the object under consideration, and E is that object's energy. Considering the

massless limit of D_L provides some intuition,

$$D_L = \frac{1 - \cos \theta}{(1 + \frac{E}{E_i})^2}. \quad (2.4)$$

The Lund distance is small if the seed and object point in nearly the same direction so that $\cos \theta \sim 1$, and if E_i is small. Thus, selecting based on minimal Lund distance prefers to assign objects to the less energetic seed when possible to keep the hemispheres energetically balanced, so long as the object and the seed point in similar directions. Once all objects have been assigned to a seed, two new seeds are defined as the sum of all the four-momenta in each hemisphere, and the assignment procedure repeats using these seeds. This continues until no object changes hemisphere after an iteration. Ultimately, the algorithm assigns objects to the correct hemisphere with efficiency ranging from approximately 70–85%, with higher efficiencies obtained in events with fewer objects and therefore reduced combinatorics [32].

The next step of the calculation of M_{T2} takes these finalized hemispheres as inputs. Under the hypothesis that the observed missing energy is the vector sum of the unobserved momenta of two invisible particles, the algorithm considers every possible vector decomposition,

$$\vec{E}_T^{\text{miss}} = \vec{v}_1 + \vec{v}_2 \quad (2.5)$$

and calculates M_T for hemisphere i with respect to \vec{v}_i . The maximum of these two values of M_T is taken as the *candidate* M_{T2} for this decomposition. This choice is motivated by a desire to infer the true mass of the decaying pair-produced particle, again, under the hypothesis that there are two pair-produced particles undergoing symmetric partially invisible decays. For the correct \vec{E}_T^{miss} decomposition and hemisphere assignments, the two calculated transverse masses are correct, and both are no larger than the true mass of the pair-produced particles. The larger, of course, is nearer to that true mass.

The only remaining step is to select which of the candidate values of M_{T2} , equivalently

which \vec{E}_T^{miss} decomposition, is correct. The algorithm selects the smallest candidate M_{T2} value found, the “minimum maximum,” so that the final M_{T2} is

$$M_{T2} = \min_{\vec{E}_T^{\text{miss}} = \vec{v}_1 + \vec{v}_2} [\max(M_T^1, M_T^2)] \quad (2.6)$$

where hemispheres 1 and 2 are assembled as described above. Since the correct missing energy decomposition is among those considered and manifestly cannot produce a value of M_T greater than the true mass of the parent particle when hemisphere assignments are correct, selecting the “minimum maximum” guarantees that M_{T2} can be no larger than the mass of the parent particle. Just as M_T cannot exceed the mass, M_{T2} is unlikely to exceed the true M_T .

Of course, it is possible for M_{T2} to be small even in genuine signal events for which it is designed to be large, and it is possible for especially pathological mismeasurement events to have large M_{T2} despite M_{T2} being a powerful rejector of these backgrounds (see Section 3.1.3). No observable is so magical that it can accept 100% of signal and reject 100% of background. Still, M_{T2} is one of the most powerful discriminants available for analyses targeting pair-produced particles decaying semi-invisibly, as can be seen for instance in Figure 3.4. Indeed, the selection of $M_{T2} > 200$ GeV applied by the analyses discussed in Section 3 is the single most important part of their baseline selection, dramatically reducing especially the detector mismeasurement background.

2.3 CMS Event Reconstruction

After observing the signatures of particles in the detector, CMS must reconstruct what happened in the event. A fully detailed description is well beyond the scope of this thesis and indeed, well beyond the scope of any single document. This section will describe only some of the most essential pieces that have not already been covered by previous sections.

2.3.1 Pileup and the Trigger System

Previous sections have discussed some of the challenges faced by the detector in reconstructing a single collision. However, the large event rate at CMS produces its own set of challenges. First, the sheer number of bunch crossings per second produces too high a data rate to save the measurements of all but a small subset of collisions [18]. Second, the collision of protons in bunches rather than singly to maintain a collision rate produces multiple simultaneous collisions in each bunch crossing, called pileup, which must be disentangled.

The Trigger System

The LHC provides CMS with roughly 1 billion collisions per second spread across 40 million bunch crossings per second, one per 25 ns. This enormous event rate is necessary to provide sensitivity to very rare physics. For instance, the expected number of gluino pair-production events observed at CMS is approximately 1 across the entire 2016, 2017, and 2018 datasets, for a gluino mass of 3 TeV. Unfortunately, as each bunch crossing constitutes roughly 1.5 MB of data, it is feasible to write only approximately 100 events per second to disk, for 150 MB/s [18]. All but one bunch crossing in 400,000 must be rejected, a feat accomplished by a two-tier trigger system [47].

The Level-1 trigger system is implemented in hardware and considers gross features of the event that are rapidly computable and qualify it for further consideration. It rejects the vast majority of bunch crossings, reducing the rate from 40 MHz to 100 kHz. The rejected events are almost all mundane events containing only a few low energy hadron jets. Some of the signatures that may qualify an event for further consideration include measurements in the muon system consistent with an isolated muon of a minimum p_T , or preliminary signs that the event may have large missing energy. The Level-1 trigger system is allotted only 3.2 μ s to make a decision, of which less than 1 μ s is available to perform a limited reconstruction, and so is subject to errors that can cause it to incorrectly reject some borderline events, so that the trigger is not a perfect

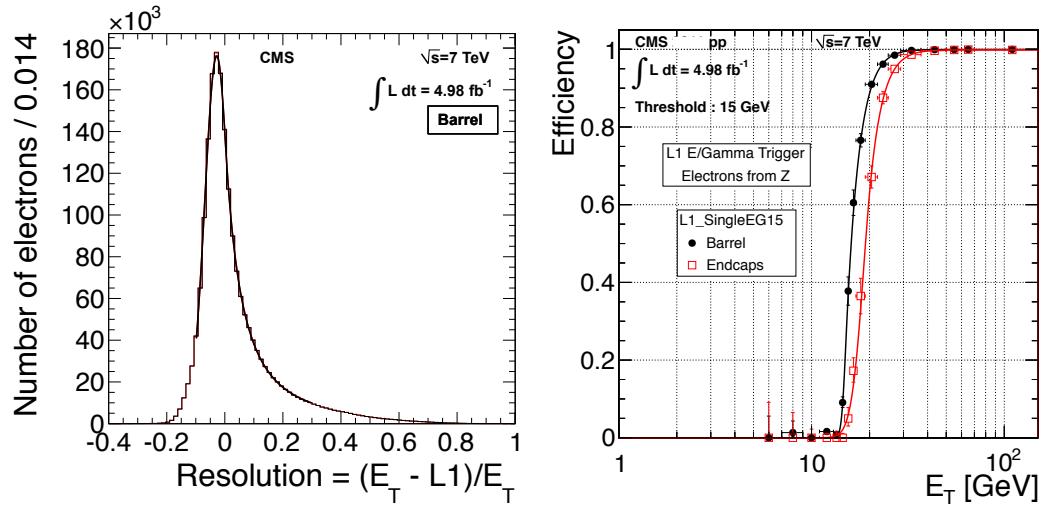


Figure 2.12: (Left) The electron energy resolution of the Level-1 CMS trigger obtained in 7 TeV collisions, in the barrel. Positive values correspond to incorrectly small measurements of the electron energy, and may cause candidate electrons to fall short of the minimum energy requirement, leading to a trigger inefficiency. (Right) The Level-1 electron trigger turn-on curve as a function of offline reconstructed electron energy, with a nominal threshold of 15 GeV. The black markers are for electrons in the barrel, corresponding to the electron population whose energy resolution is shown in the plot at left, and red is for electrons in the endcaps. To avoid exposing itself to trigger turn-on systematic uncertainties, an analysis may be forced to select only events possessing an electron with energy greater than approximately 30 GeV, the beginning of the efficiency plateau, despite the trigger nominally accepting events containing an electron with energy as low as 15 GeV. Taken from [47].

step function. For instance, the electron energy resolution of the Level-1 reconstruction is shown in Figure 2.12 (left), along with the resulting trigger turn-on curve (right). To avoid the need to simulate this trigger turn on, analyses typically select only events sufficiently distinctive that they are accepted by the trigger with 100% efficiency.

Events passing the Level-1 trigger are subjected to a more careful analysis by the High Level Trigger (HLT), implemented in software on standard servers. With 13,000 machines working in parallel, the HLT is allotted 90 ms per event, enough time for reconstruction of the portion of the event identified by the Level-1 trigger as interesting at near offline quality. One event in 1000 considered by the HLT is sufficiently spectacular to qualify to be written to disk, and these constitute the CMS dataset available for final offline analysis.

Pileup Correction

In addition to the high bunch crossing rate, the reconstruction also must cope with the large number of simultaneous collisions in each bunch crossing, called pileup. The mean number of pileup collisions was 32 in 2018 [1], and a significant number of bunch crossings exceeded 60 pileup collisions, as can be seen in Figure 2.13. Given that only one bunch crossing in 400,000 contains a collision interesting enough to be written to disk, the number of interesting collisions per bunch crossing is at most 1, and the others must be subtracted.

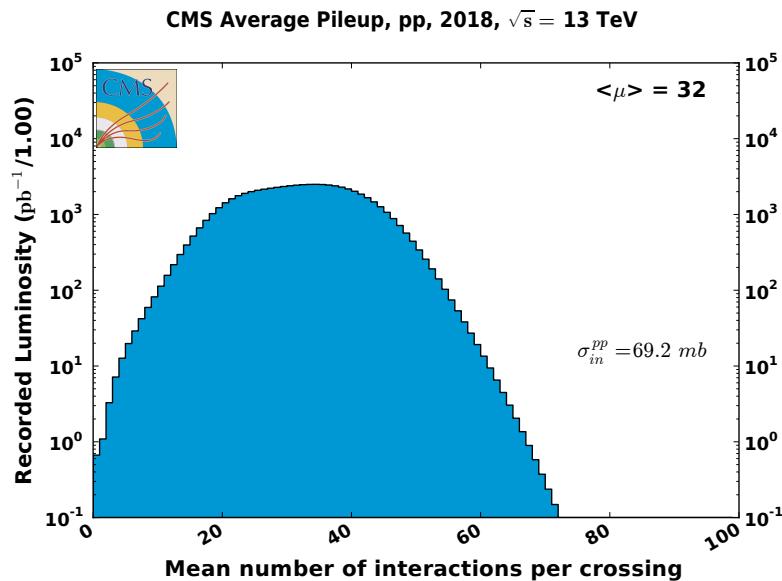


Figure 2.13: The distribution in 2018 of the number of simultaneous collisions per bunch crossing, so-called pileup, provided by the LHC and observed by CMS. The average bunch crossing produced 32 simultaneous collisions, with an inferred total inelastic collision cross section of 69.2 mb. Taken from [1].

Barring errors in track reconstruction, charged pileup can be subtracted perfectly [22]. Charged particles produce tracks, which can be associated to vertices. Tracks associated to vertices other than the primary vertex of interest can be subtracted. This is not possible for the neutral content of an event, which is known only from calorimeter deposits.

There exist two popular heuristic methods for subtracting neutral pileup. The first method, called Delta-Beta, proposes that the neutral energy from a pileup vertex is equal everywhere to

exactly one half of the charged pileup, because particles with charges +1, -1, or 0 tend to be produced with roughly equal probability and energy. This is true on average, but of course is only approximately true in any given instance. The second method, called Effective Area, benefits from larger pileup and has gained popularity recently. This method divides the detector into many small regions and calculates the median energy density in these regions, ρ . By taking the median, the calculation naturally mitigates the impact of outlier regions of the detector with large contributions from the primary vertex, which tends to be much more energetic than any of the pileup vertices. This median energy deposit is taken to be representative of the pileup contribution in every region of the detector, true by construction within statistical fluctuations, and subtracted everywhere to leave only the energy believed to originate from the primary vertex. This algorithm is only as accurate as the statistical fluctuations of pileup energy are small, and so benefits directly from increased pileup vertex counts. At present, both algorithms have similar performance and see wide use. The search in Section 3.1 uses Effective Area, while the extension to this search in Section 3.2 uses the Delta-Beta method to correct the isolation of disappearing tracks similarly.

Pileup also makes reconstruction more difficult in a more general sense by overwhelming reconstruction algorithms. Of special importance to the disappearing tracks analysis in Section 3.2 is its effect on track fake rates. High pileup vertex counts cause proportionally increased hit counts in the tracker. With more hits to consider, the track reconstruction algorithm is more prone to make errors, and connect hits to form tracks that are not all associated with the same genuine particle. For instance, the algorithm may output a fake track using three hits produced by particles from three different pileup vertices that happen to fall in a line extrapolating to the primary vertex. The tracker is also overwhelmed at the hardware level, with higher particle flux causing higher operating temperatures, and therefore more noise, and overflowing readout buffers during periods where the beam luminosity, and therefore the typical pileup, is especially high [76], as shown in Figure 2.14. Events containing disappearing tracks, especially very short tracks, disproportionately originate from events at the extreme high end of the pileup distribution,

indicating that pileup and, more generally, high beam luminosity directly causes background for this search by inducing track reconstruction errors.

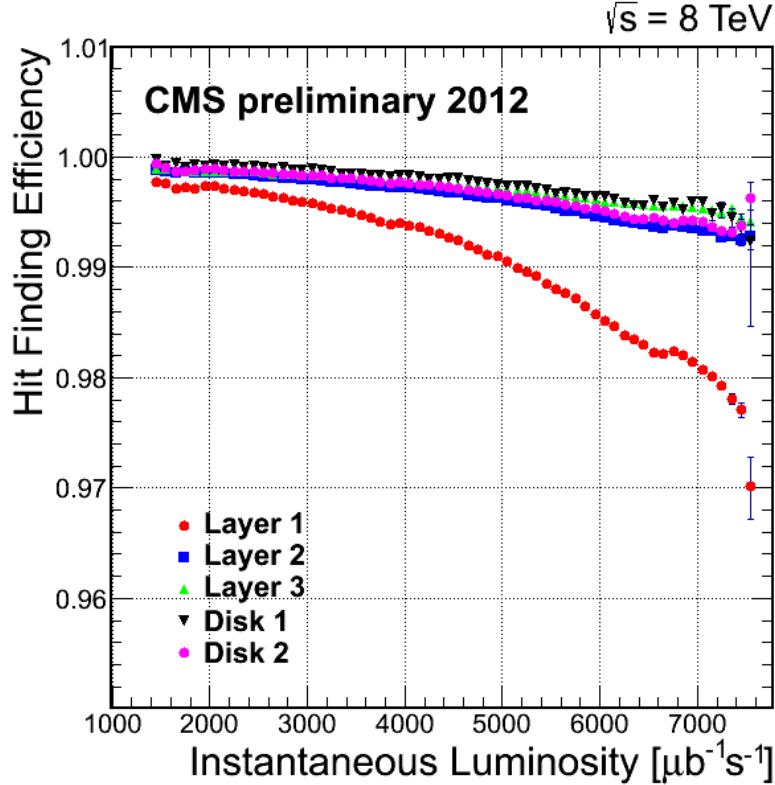


Figure 2.14: The efficiency for a charged particle to leave a hit drops at high instantaneous beam luminosity, especially in the first layer of the tracker (red). Periods of high luminosity tend to produce events with high pileup. Taken from [76].

2.3.2 Particle Flow

The general purpose offline reconstruction algorithm used in most cases at CMS is called Particle Flow [68]. Particle Flow (PF) attempts to reconstruct every individual particle in an event by integrating information from all detector components, producing so-called PF candidates. Rather than, say, obtain an energy measurement entirely from the calorimeters and ignore the energy measurement that can be made using the curvature of a particle's track, PF uses the tracker and the calorimeters as checks on each other, and intelligently assigns a final energy based on

its beliefs about the relative trustworthiness of each measurement in each case. This extraction of all available information from the event rather than the traditional segregated inputs from different detector components is in part responsible for CMS’s outstanding reconstruction. These achievements include a high b-tagging efficiency, which is maximized by considering all possible clues simultaneously using a machine learning approach, rather than for instance using only the secondary vertex and track multiplicity information provided by the tracker [68], and efficient identification of electrons, as shown in Figure 2.15.

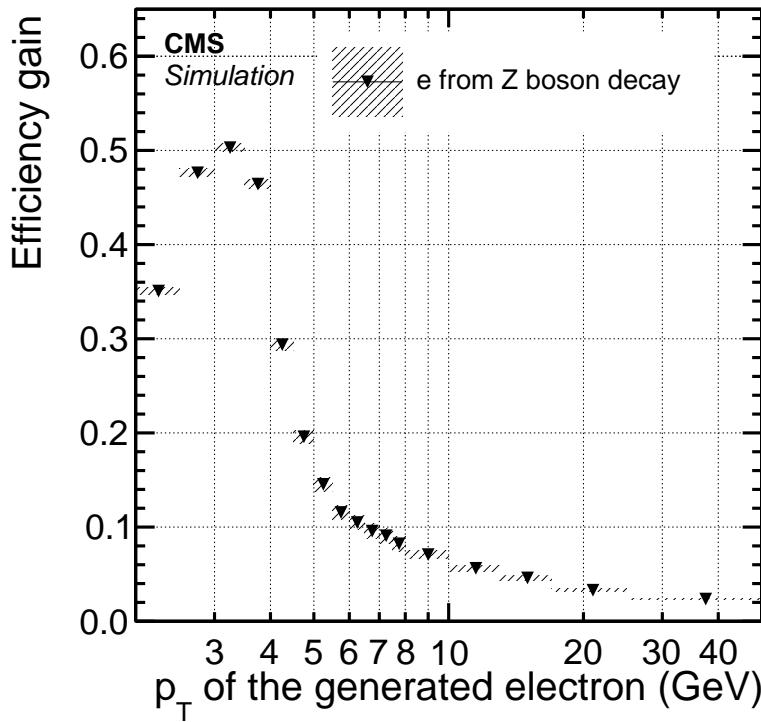


Figure 2.15: The identification efficiency of electrons produced by the decay of a Z boson is dramatically improved by using information from both the tracker and the ECAL, rather than the ECAL alone as has been the standard approach in other detectors at hadron colliders. The absolute efficiency gain, $\text{Eff}_{\text{ECAL+Tracker}} - \text{Eff}_{\text{ECAL}}$, is plotted. Taken from [68].

This approach can allow PF to reject fakes or badly mismeasured particles. For instance, PF rejects tracks that do not extrapolate to significant calorimeter deposits, indicating that they are neither electrons nor hadrons, and also do not have a corresponding track in the muon chamber. In most cases, PF is correct to reject these tracks, which are almost assuredly fakes. In rare cases, the

track is of such high quality that it passes the “high purity” selection described in Section 2.2.2, indicating that the tracker is highly unlikely to have made an error. These are called “Lost Tracks,” and are the foundational object for the track selection used in the Section 3.2 disappearing tracks analysis. The fact that PF rejected these tracks despite a strong tracker measurement flags them as worthy of further study, serving as a useful clue for searches for new physics.

2.3.3 Jet Clustering

Jets produced by hadron fragmentation are constructed from PF candidates using the Anti- k_t algorithm [23]. A major goal of such an algorithm is including all of the content associated with a hadron’s fragmentation in a single jet, and nothing else. Additionally, it is desirable for jets output by an algorithm not to be affected by very low energy particles that may be simple manifestations of noise, so that the same event measured by multiple detectors with different properties would produce nearly identical jets in every case. Older jet clustering algorithms simply used a fixed cone size, and included in jets anything that fell inside a cone centered on a few of the larger energy deposits in the calorimeters. Modern algorithms start at the particle level and associate particles to jets using distance measures similar to the Lund distance (Equation 2.4) used in M_{T2} . The Anti- k_t algorithm uses the measure of distance between particles i and j defined as

$$d_{ij} = \min \left(\frac{1}{k_{ti}^2}, \frac{1}{k_{tj}^2} \right) \frac{\Delta R_{ij}^2}{R^2} \quad (2.7)$$

where $\Delta R_{ij}^2 = (\eta_i - \eta_j)^2 + (\phi_i - \phi_j)^2$, R is a tunable jet cone size parameter, and k_{ti} is the p_T of particle i . The algorithm gets its name from the appearance of k_t^2 in the denominator rather than the numerator of the minimum as in its predecessor, the k_t algorithm. The most important advantage of this algorithm relative to its predecessor is the suppression of the effect of very low energy particles on jet composition and shape by using the minimum of $1/k_t$ rather than of k_t . This causes low energy particles to cluster not simply with whatever the nearest particles are, but

rather with the nearest *high energy* particle, since d_{ij} is always large for two low energy particles, while it is suppressed by the smallness of the high energy particle’s $1/k_t^2$ in any pairing with a low energy particle. Nearby hadrons producing partially overlapping jets can thereby be distinguished, by clustering the lower energy particles around nearby but separate high energy cores, instead of forming one fat jet from the agglomeration of all particles falling inside the cone, and jets can pick up low energy particles that may have been radiated a moderate distance away from the core of the source jet, instead of these outliers clustering amongst themselves to form spurious jets. Thus, jets clustered by the modern Anti- k_t algorithm closely correspond to the source hadrons, making the N_{jet} observable more useful than ever before, and also assisting flavor tagging.

2.4 Simulation

The Standard Model and hypothetical models of new physics predict what happens at the primary interaction. These parton-level predictions need to be converted into final state particles, and then a prediction of the observable signatures produced by these particles in the detector. To accomplish this, CMS employs a multi-tier Monte Carlo simulation process.

2.4.1 Objectives

Simulation is used both to understand sources of background and to predict the appearance and yield of signals. Although simulation is subject to uncertainties that cause analyses to prefer using data-driven background predictions, it still captures the gross features of all but the most pathological background sources. For understanding the expected appearance of signals, it is the only option available.

2.4.2 Limitations and Challenges

Simulation of collider physics faces four primary challenges.

First and foremost is a finite computation budget. All simulations must be performed numerically, using a Monte Carlo approach. While the physical collider can produce 1 billion event per second and thereby produce events with extremely rare properties, such an event rate is orders of magnitude out of reach of software. This is not necessarily a problem for rare signal models, as a computer can be made to simulate only the signal model of interest, and ignore the much more common background physics, thereby producing simulated events even faster than the collider can, sometimes much faster for the rarest signals. However, much more common backgrounds are typically difficult to simulate in the phase space of greatest interest. For instance, QCD events with large M_{T2} are extremely rare as discussed in Section 3.1.3, but QCD events comprise nearly all of the 1 billion collisions per second observed at CMS, such that even the high M_{T2} event rate is non-negligible. Software simply cannot produce a statistically robust sample of these events using the traditional Monte Carlo sampling approach. Section 3.1.5 discusses a successful workaround for this problem.

The other three problems affect three separate stages of the simulation pipeline.

First is an uncertainty arising from theory due to the perturbative approximation used to extract predictions from the Standard Model at parton level. Although the Standard Model is capable of producing astonishingly accurate predictions, famously including over 10 digits of accuracy in its prediction of the electron’s magnetic moment [55], calculating predictions to such precision is a laborious undertaking that cannot practically be managed for every conceivable process [12]. In practice, most theoretical calculations are truncated at only a few orders, or even only to leading order, and so are not as precise. At CMS, most of these parton-level calculations are performed numerically by the MadGraph generator [10]. This issue is mitigated in part by using data-driven predictions for background, but this is not possible for signal, contributing to significant theoretical uncertainties in signal predictions. Additionally, the Standard Model’s theoretical prediction is dependent on experimental inputs, which have an unavoidable measurement uncertainty.

Second is an uncertainty arising due to QCD’s non-perturbativity at low energy, as described in Section 1.1.3. At energy scales around 1 GeV and below, the QCD coupling exceeds 1, so that the typical perturbative approach is useless for low energy QCD [37]. As the LHC is a hadron collider, QCD dominates the underlying structure of every event, and although the initial collision may be at a high enough energy scale to allow for perturbative calculation of QCD effects, the energy eventually becomes spread among enough particles to reach the non-perturbative regime. Thus, the conversion of the output of a parton-level generator like MadGraph to the final state particles, chiefly the fragmentation of hadrons to produce jets, can only be performed using heuristics rather than from first principles. CMS simulation typically uses the Pythia program to perform this second step, which implements the Lund string framework of hadronization [71]. Moreover, the internal structure of the proton itself is dictated by non-perturbative QCD, creating the need for empirical PDFs as discussed in Section 1.1.4.

Finally, the interaction of final state particles with the detector depends on the exact state of the detector at every moment. The detector’s state is dynamic, and fine details like which components momentarily overheated or the impact of induced radioactivity in detector components, can only be represented approximately at best. CMS employs two simulation methods for particle interactions with the detector. The first is a detailed simulation of the interaction of simulated particles with as detailed a model of the detector as knowledge permits, implemented in the GEANT program [7] and called full-simulation. Naturally, this is the gold standard, and is applied to simulation campaigns that are either of broad interest to the entire collaboration, or require and can afford the additional computational complexity. The second approximates typical detector measurements using empirical smearing functions to represent, for example, the detector’s ECAL resolution for electrons. This is called fast simulation. Fast simulation does not accurately capture the tails of distributions, having limited knowledge of rare events, but accurately predicts the bulk. This makes fast simulation suitable for producing large samples of signal events, but unsuitable for background. Signal events *define* the parameter

space of interest, and so the bulk of signal events populate the desired region. Typically, only background events deep into the tails of distributions manage to appear signal-like, and so sources of background must be modeled with full simulation.

In the end, the CMS simulation pipeline cannot be perfect, but is largely robust. Figure 2.16 shows a comparison between simulation and data of the M_{T2} distribution in events containing exactly one low M_T electron or muon and either (left) 0 or (right) 1 b-tagged jet, and additionally with requirements on the minimum H_T and \cancel{E}_T . These are small corners of phase space, combining contributions from multiple distinct physics processes, and yet still the simulation's agreement with data is remarkable, overcoming all of the challenges described above.

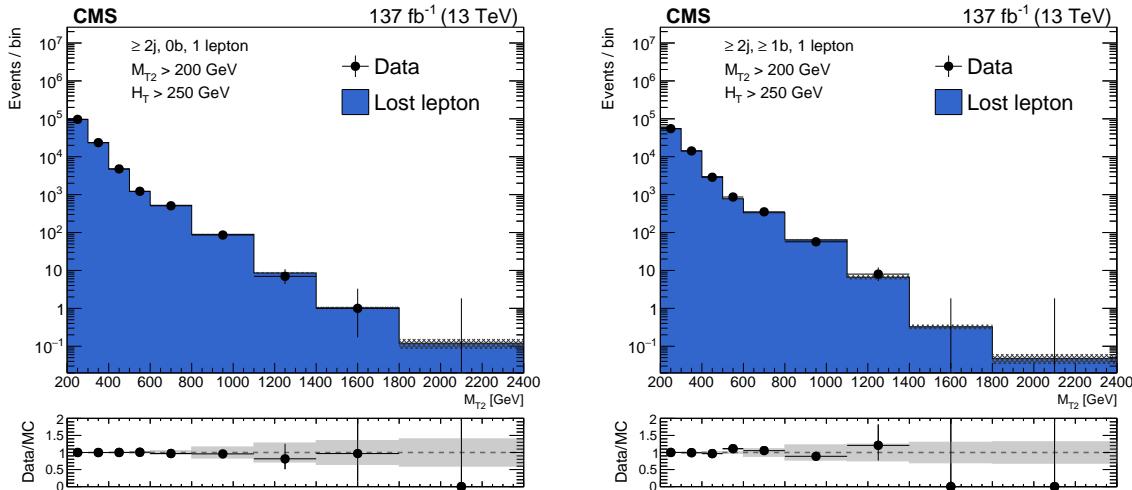


Figure 2.16: Modeling of M_{T2} in simulation is checked against data with 0 (left) and 1 (right) b-tagged jet, in events satisfying $H_T > 250$ GeV, $\cancel{E}_T > 250$ GeV if $H_T < 1200$ GeV and $\cancel{E}_T > 30$ GeV otherwise, containing one lepton with $M_T < 100$ GeV, consistent with a $W^\pm \rightarrow \ell^\pm v$ decay. The simulation combines all contributing Standard Model processes. Even though the simulation is being asked to predict a narrow corner of phase space, and M_{T2} is a complex observable depending in detail on all components of the detector, the agreement is remarkable. Taken from [30].

Chapter 3

The Full 13 TeV M_{T2} Analysis

This section presents two searches for new physics in all-hadronic final states with substantial \cancel{E}_T as inferred through large M_{T2} in 13 TeV proton-proton collisions recorded by the CMS detector [30].

The \cancel{E}_T selection targets invisible particles, motivated by the evidence for dark matter discussed in Section 1.2.4. If dark matter is at least partially composed of a particle or particles associated with physics at the weak scale, then it may be produced in LHC collisions and, while not detectable itself, be inferred via an excess of events with imbalanced transverse momentum.

Considering only all-hadronic events with large M_{T2} is part of a divide and conquer strategy employed by the CMS collaboration. CMS analyses also search for dark matter in events containing leptons, e.g. [66], but there is no way to know whether dark matter will be found in either or both. By splitting the searches, each can optimize for its unique event characteristics. Likewise, the collaboration also considers all-hadronic events that do not necessarily have large M_{T2} [67], since while M_{T2} in many cases provides enhanced sensitivity, it can reduce sensitivity to some models, for instance in cases where the decay energy is very small, or those in which invisible particles are not produced in pairs.

The first search is inclusive, and is a continuation of previous analyses based on smaller

datasets, most recently using the 35.9 fb^{-1} dataset recorded in 2016 [63]. Accordingly, it is referred to as the classic search.

The second search is a new extension of the first, requiring additionally the presence of a disappearing track in a selected event. Disappearing tracks have been targeted as an observable of interest by both CMS [45, 64] and ATLAS [3, 2] previously, but with lower energy or smaller datasets, and different methods.

Both searches set the strongest constraints to date on a variety of hypothetical extensions to the Standard Model possessing pair-produced dark matter candidates, most notably R-parity conserving supersymmetry.

3.1 Classic Search

3.1.1 General Description

The classic search’s core, defining selections are for events with large missing transverse energy \cancel{E}_T as inferred via M_{T2} , large total hadronic transverse energy, called H_T , and no leptons.

The primary motivation for the analysis is either finding or setting constraints on particle dark matter. The experimental signature of dark matter is its undetectability; its presence can only be inferred through large imbalance of transverse energy. The M_{T2} analysis further focuses on the case of pair-produced dark matter using its eponymous observable. For reasons discussed in Section 1.3, R-parity conserving supersymmetry is a model of special interest to the theoretical community, and this model’s dark matter candidate would always appear in pairs. The M_{T2} variable tends to increase sensitivity to the pair-production scenario as described in Section 2.2.6.

The selection of events with large H_T is motivated by the CMS trigger as discussed in Section 2.3.1, which is in turn motivated by the inferred large energy scale of new physics. The decays of heavy new particles to relatively light Standard Model particles will result in, typically, very energetic events. It is possible that the dark matter candidate is itself very heavy, nearly as

massive as the new particle that decays to it. In this scenario, while the events are indeed very energetic, much of the energy is lost to the invisible portion of the event, and the visible, hadronic energy can be relatively small. Accordingly, sensitivity in the compressed scenario is weaker, as many of the new physics events are too low energy to pass the kinematic selections. The M_{T2} selection is especially inefficient, in this scenario. However, sensitivity is not negligible, and the search is sensitive to any mass splitting in many cases.

The last selection, a veto of leptons, is chiefly motivated by CMS’s divide and conquer strategy. Final states including leptons are considered in other searches. However, the veto does allow for partial suppression of the Standard Model neutrino background, since neutrinos are often produced alongside a charged lepton in decays of the W boson, as discussed in Section 3.1.3.

Ideally, any observation of an event with large \cancel{E}_T and M_{T2} , large H_T , and no leptons, would constitute a discovery. Unfortunately, this is not the case, as the Standard Model is capable of producing events with all of these signatures. Neutrinos and simple mismeasurement of the event can produce large missing energy signatures, events with large H_T are uncommon in the Standard Model but hardly impossible, and the vast majority of events in a proton-proton collider have no leptons. Some number of events will have all of these properties without any new physics occurring, and most of the analysis’ efforts are spent estimating how often these background processes occur, as discussed in Section 3.1.3.

3.1.2 Signals

The classic search generically targets any new physics that produces high H_T events with jets and missing energy from undetected particles. Diagrams for several candidate models are shown in Figures 3.1, 3.2, and 3.3. The diagrams in Figure 3.1 of gluino and squark pair production as predicted by supersymmetric extensions of the Standard Model are of greatest interest, and while the analysis is sensitive to a wide variety of similar hypothetical models,

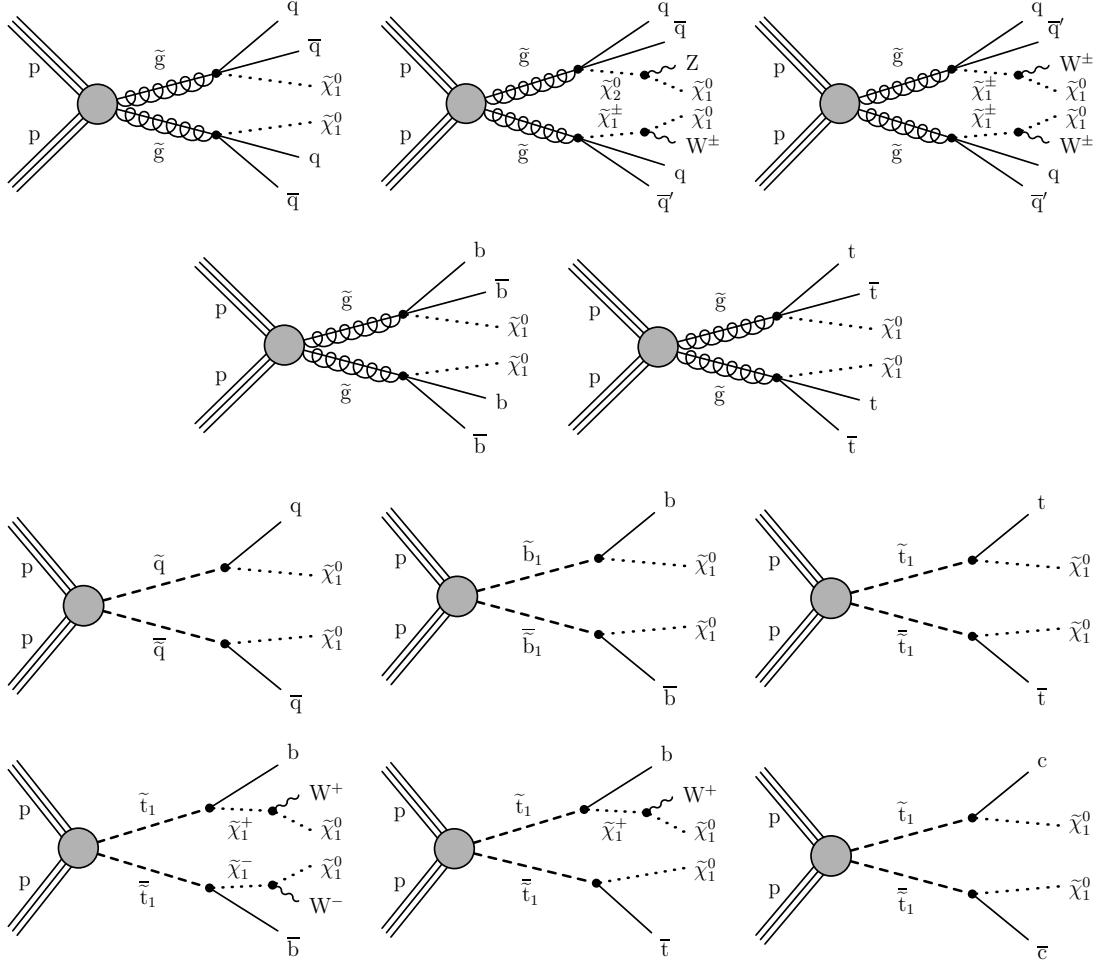


Figure 3.1: Diagrams of gluino (upper five) and squark (lower six) pair production, as predicted by supersymmetric extensions of the Standard Model. The classic search considers five potential gluino decay chains. At upper left, the gluinos decay to light flavor quarks (up, down, strange, or charm), and the lightest neutralino, $\tilde{\chi}_1^0$. At upper center, the gluinos decay to light flavor quarks, but rather than decaying directly to $\tilde{\chi}_1^0$, the gluinos decay either to the second neutralino, $\tilde{\chi}_2^0$, which subsequently decays to a Z boson and $\tilde{\chi}_1^0$, or to the lightest chargino, $\tilde{\chi}_1^\pm$, which subsequently decays to the W boson and $\tilde{\chi}_1^0$, with equal probability. At upper right, both gluinos undergo the $\tilde{\chi}_1^\pm$ to W decay chain. On the left of the second row, both gluinos decay to bottom quark pairs and $\tilde{\chi}_1^0$. On the right of the second row, both gluinos decay to top quark pairs and $\tilde{\chi}_1^0$. Additionally, the classic search considers six potential modes of squark production and decay. On the left of the third row, light flavor squarks decay to light flavor quarks and $\tilde{\chi}_1^0$. In the center of the third row, bottom squarks decay to bottom quarks and $\tilde{\chi}_1^0$. On the right of the third row, top squarks decay to top quarks and $\tilde{\chi}_1^0$. At lower left, top quarks decay to bottom squarks and $\tilde{\chi}_1^\pm$, which subsequently decay to the W boson and $\tilde{\chi}_1^0$. At lower center, top squarks may undergo either the $\tilde{\chi}_1^\pm$ decay chain, or a direct decay to a bottom squark and $\tilde{\chi}_1^0$, with equal probability. At lower right, each top squark decays to a charm squark and $\tilde{\chi}_1^0$. Taken from [30].

it is optimized for these. As discussed in Section 1.3.3, these models simplify the enormous parameter space of supersymmetric extensions by assuming that all superpartners except those in the process are so massive that they can be neglected entirely. The upper five diagrams all depict pair production of gluinos, \tilde{g} , the supersymmetric partner of the gluon. As each gluino decays to two quarks, each of which will typically produce a jet, gluino pair production events tend to have many jets. Many different decay chains are possible, and the analysis considers five representative benchmarks.

In the first benchmark, the gluinos decay to light flavor quarks (up, down, strange, or charm), and the lightest neutralino, $\tilde{\chi}_1^0$.

In the second benchmark, the gluinos decay to light flavor quarks, but rather than decaying directly to $\tilde{\chi}_1^0$, the gluinos can decay either to the second neutralino, $\tilde{\chi}_2^0$, which subsequently decays to a Z boson and $\tilde{\chi}_1^0$, or to the lightest chargino, $\tilde{\chi}^\pm$, which subsequently decays to the W boson and $\tilde{\chi}_1^0$. Each of these decays can occur with equal probability. Both the W and Z will themselves decay, usually to a pair of quarks, which will in turn produce jets. So, relative to the first signal model, the second tends to have greater jet multiplicity. The Z can also decay to neutrinos, and the W can decay to a neutrino and a lepton that is not reconstructed. In these scenarios, this benchmark trades some jets for an enhanced missing energy signature. If the W or Z decays leptonically and a lepton is successfully reconstructed, events from these signals may fail the lepton veto and instead end up in control regions, biasing the background prediction. The procedure used to handle such signal contamination of control regions is described in Section 3.1.6.

The third benchmark, in which both gluinos undergo the $\tilde{\chi}^\pm$ to W decay chain, is similar.

In the fourth benchmark, both gluinos decay to the bottom quark and $\tilde{\chi}_1^0$. As described in Section 2.2.2, it is possible to identify jets that originated from a bottom quark; such a jet is said to be “b-tagged.” The classic search bins in the number of b-tagged jets in order to enhance sensitivity to signals of this type.

In the fifth and final gluino pair-production benchmark, both gluinos decay to top quarks and $\tilde{\chi}_1^0$. The top quark decays with probability near unity to a bottom quark and a W boson. The extra W bosons, compared to the direct bottom decay model, can either add jets or leptons and neutrinos, as previously discussed. This signal tends to produce the most remarkable events of any signal model considered, with very large H_T , N_{jet} , and $N_{b\text{-tag}}$, but also loses many events to the lepton veto as any of the four W bosons is liable to produce a lepton.

The lower six diagrams of Figure 3.1 all depict pair production of squarks, the supersymmetric partners of quarks. Squark decays directly produce only one quark each, compared to two for gluino decays, so squark pair-production events tend to have fewer jets than gluino pair-production events.

In the first squark benchmark diagram, on the left of the third row, a pair of light flavor squarks (up, down, strange, or charm) is produced and each decays to a light flavor squark and $\tilde{\chi}_1^0$.

In the second benchmark, bottom squarks are produced and decay to bottom quarks and $\tilde{\chi}_1^0$. These events tend to have b-tagged jets, but fewer than in gluino decays to bottom squarks.

In the third benchmark, top squarks are pair produced and decay to top quarks and $\tilde{\chi}_1^0$. The relationship between this process and the previous one is similar to the relationship between the fifth and fourth gluino benchmarks, respectively.

The fourth benchmark is very similar to the third. Instead of the top squark decaying directly to a top quark, which then decays to a bottom quark and W boson, the squark decays to a bottom quark directly and $\tilde{\chi}^\pm$, which subsequently produces the W. While the final state particles are identical, their kinematics can be very different depending on the distribution of masses realized in nature. For instance, if the top squark and $\tilde{\chi}^\pm$ mass splitting is very small, the bottom quark jet in this benchmark may have such low p_T in a typical event that it is difficult to reconstruct.

In the fifth benchmark, each top squark may undergo the $\tilde{\chi}^\pm$ decay chain or decay directly to a bottom squark and $\tilde{\chi}_1^0$, with equal probability, mixing the two previous models.

In the last benchmark, each top squark decays to a charm quark and $\tilde{\chi}_1^0$. This decay chain could dominate when the top squark and $\tilde{\chi}_1^0$ mass splitting is too small to allow decay to on-shell top quarks, and the mass of $\tilde{\chi}^\pm$ is larger than that of the top squark.

In each of these models, the mass of the squark or gluino and the mass of $\tilde{\chi}_1^0$ are free parameters. Large squark and gluino masses cause low production rates, as the production cross sections drop rapidly with increasing mass, as discussed in Section 1.3.3 and shown in Figure 1.12. The mass of $\tilde{\chi}_1^0$ does not affect the production rate, but it does affect the character of the events. When the mass splitting between the gluino or squark and $\tilde{\chi}_1^0$ is small, only a small portion of the event's energy ends up in the visible decay products, and most is lost to the rest energy of $\tilde{\chi}_1^0$. These events have low H_T and \cancel{E}_T , low N_{jet} , and when applicable, low $N_{\text{b-tag}}$, and more closely resemble background. As the mass splitting increases, more energy shifts to the visible portion of the event, making events less background-like and increasing sensitivity. The analysis considers a grid of potential mass points for each model, and this sensitivity pattern is evident in the curves shown in Section 3.1.8.

The classic search also considers non-supersymmetric models.

The first is referred to as the mono- ϕ model. In this model, a colored boson much like a squark is produced singly, rather than pair-produced as mandated in squark models to conserve R-parity, and decays to a quark and invisible fermion, as shown in Figure 3.2. This model has an especially low number of jets and low M_{T2} , and so would appear in more background-like bins than most supersymmetric models. In fact, it was originally proposed [13] to explain a potential excess in bins of this sort in the previous edition of the classic search, published based on 2016 data [63], and so there is some interest in whether such an excess persists in the larger dataset. While the analysis is not optimized for this kind of model, it still has some residual sensitivity to any model characterized by jets and missing energy in the final state, including mono- ϕ .

The second non-supersymmetric model, and final model considered explicitly by the classic search, is a leptoquark extension of the Standard Model. As discussed in Section 1.4,

while squarks can decay to a quark and $\tilde{\chi}_1^0$, leptoquarks can decay to a quark and a neutrino. A neutrino is experimentally effectively indistinguishable from a low mass $\tilde{\chi}_1^0$, such that leptoquarks and squarks produce nearly identical final states, as first established in a reinterpretation of the 2016 edition of the classic search [65]. Thus, it is a relatively simple exercise to reinterpret squark results in the low mass $\tilde{\chi}_1^0$ limit as leptoquark results. A few leptoquark production diagrams are shown in Figure 3.3.

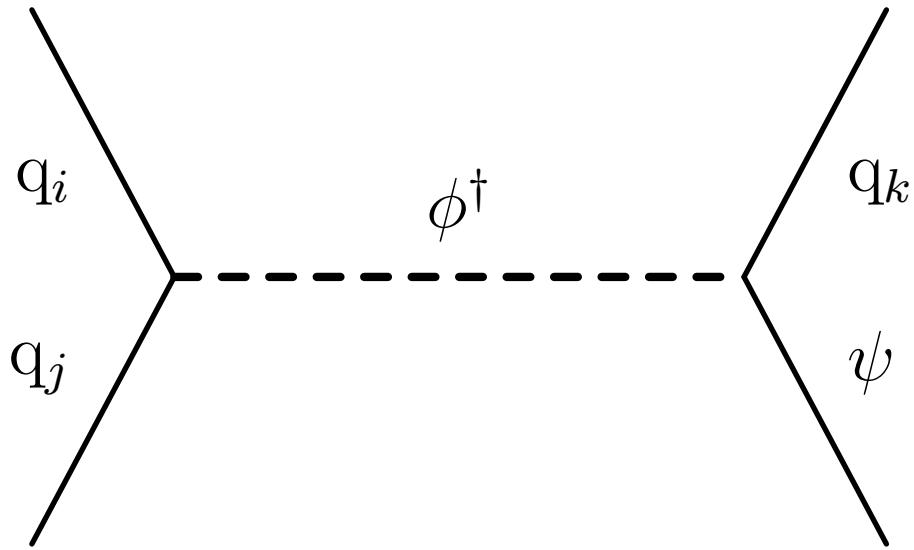


Figure 3.2: Diagram for the mono- ϕ model, in which a colored scalar ϕ is resonantly produced, and decays to an invisible massive Dirac fermion ψ and an SM quark. Note that ϕ is not pair-produced, in contrast to otherwise-similar squarks. Taken from [30].

By necessity and unlike backgrounds, the expected event yields for signal models are obtained directly from simulation. For this reason, while most uncertainties are shared with background, signal estimates are exposed to a few extra uncertainties, and some of the shared uncertainties increase in magnitude.

First, while background is normalized to control region counts and so automatically scaled to the correct integrated luminosity, signal counts must be scaled by an independent measurement of the integrated luminosity. This measurement is only precise to within a few per cent, leading

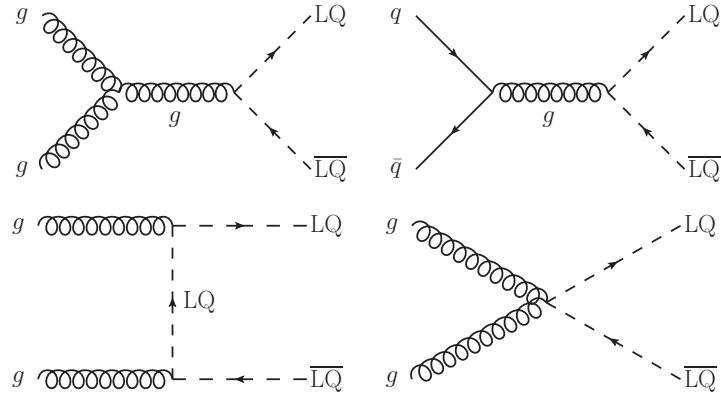


Figure 3.3: Diagrams of leptoquark pair production. Each leptoquark decays to a neutrino and a quark. Taken from [30].

Table 3.1: Systematic uncertainties in the signal yields for the simplified models of BSM physics. The large statistical uncertainties in the simulated signal sample come from a small number of bins with low acceptance, which are typically not among the most sensitive bins contributing to a given model benchmark point. Taken from [30].

Source	Range [%]
Integrated luminosity	2.3–2.5
Limited size of MC samples	1–100
b -tagging efficiency, heavy flavors	0–40
b -tagging efficiency, light flavors	0–20
Lepton efficiency	0–20
Jet energy scale	5
Fast simulation \cancel{E}_T modeling	0–5
ISR modeling	0–30
μ_R and μ_F	5

to a few per cent uncertainty on the expected event count.

Second, additional jets obtained from ISR must be modeled in simulation for signal, while for background, the control regions provide a precise measurement of the relevant ISR jet production. In bins with many jets, where ISR jet production rates are less well-known and simultaneously more important, this uncertainty is as large as 30%.

Finally, signal models must be simulated in the fast-simulation framework discussed in Section 2.4.2 to make the enormous array of signal scenarios considered computationally feasible. A systematic is assessed to cover potential errors made by this approximation.

The remaining uncertainties are described in the next section, in the context of background estimates.

3.1.3 Backgrounds

All of the targeted signals are characterized by all-hadronic events with large missing energy, but observing such an event does not constitute discovery due to the existence of backgrounds that can produce the same basic signature. These backgrounds can be broadly divided into the detector mismeasurement background, in which apparent missing energy is generated not by genuine undetected particles but by an error in reconstruction, and the neutrino background, in which the missing energy is produced by genuine neutrinos as predicted in the Standard Model. The neutrino background can be subdivided into neutrinos originating from $W^\pm \rightarrow \ell^\pm v$, in which the presence of a charged lepton allows the neutrino to be rejected with good efficiency, and those originating from $Z \rightarrow vv$, in which the final state is entirely invisible and cannot be efficiently vetoed.

Mismeasurement

The most problematic background is that caused by detector mismeasurement. Nearly every mismeasured event is a QCD multijet event, because nearly every event at a proton-proton collider is a QCD multijet event and the probability of mismeasurement is roughly flat across events. Accordingly, the mismeasurement background is also referred to as the QCD multijet background. While the detector makes mistakes only very rarely, QCD events are so relatively common that this background is still dominant in the raw dataset, before any cleaning selections.

Estimating the QCD background is challenging because it requires highly detailed knowledge of the detector’s idiosyncrasies. Rather than risk falsely discovering a signal or failing to identify one that is present due to misprediction of this background, the analysis adopts selections designed to suppress it, until the background is sufficiently minor that large relative error in its

prediction is acceptable.

The first and most powerful of these selections uses the M_{T2} variable described in Section 2.2.6, namely $M_{T2} > 200$ GeV, where the value is chosen to achieve the desired suppression of the mismeasurement background. Although the M_{T2} selection is somewhat expensive in the sense that it eliminates a significant fraction of signal, especially for signals with a small mass splitting and signals like mono- ϕ that are not pair-produced, it is crucial for suppressing the mismeasurement background. At $H_T > 1500$ GeV, the mismeasurement background extends unacceptably beyond $M_{T2} \sim 200$ GeV, and the M_{T2} selection is tightened to $M_{T2} > 400$ GeV.

The second selection uses the observable

$$\Delta\phi_{min}^{1,2,3,4} = \text{Min} \left(\left| \phi_{\cancel{E}_T} - \phi_i \right| \right)$$

where ϕ_i indicates the ϕ coordinate of the i th p_T jet and $\phi_{\cancel{E}_T}$ is the ϕ coordinate of the missing energy vector. Stated simply, $\Delta\phi_{min}^{1,2,3,4}$ is the smallest angular separation in the transverse plane of the missing energy vector and any of the four highest p_T jets. Close overlap between a jet and the missing energy vector indicates a high probability that the jet was badly mismeasured, and that this mismeasurement is the source of \cancel{E}_T in the event. The selection applied is $\Delta\phi_{min}^{1,2,3,4} > 0.3$, where the value is chosen to achieve strong background rejection without too great a loss of signal efficiency. Only the four highest energy jets are used because the probability of *some* jet overlapping the missing energy vector approaches unity as the number of jets increases, so the selection would nearly always veto high N_{jet} events, and the high N_{jet} bins are sufficiently background-depleted that aggressive background rejection is not as necessary. The effect of this selection is depicted in Figure 3.5, after an M_{T2} selection of only 100 GeV.

The last selection rejects events in which a suspiciously large fraction of the missing energy comes from very soft objects. The minimum p_T for selected jets is 30 GeV. The missing energy vector from only selected jets is denoted $\vec{\cancel{H}}_T$. The missing energy vector used in the

analysis, \vec{E}_T , uses all PF candidates, including those outside jets or in jets with $p_T < 30$ GeV. If these two quantities are very different, it means that a large portion of the \vec{E}_T in the event was generated by these low p_T objects, a sign that something may have gone wrong in reconstruction. Specifically, the selection is $|\vec{H}_T - \vec{E}_T| / |\vec{E}_T| < 0.5$.

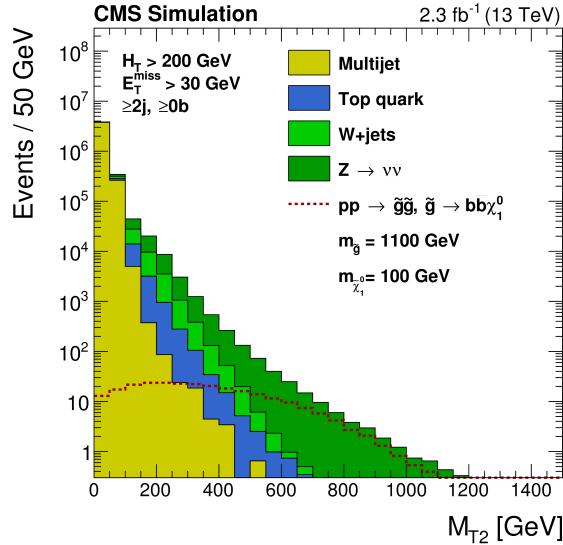


Figure 3.4: The distributions in M_{T2} of the QCD background (filled yellow) and the neutrino backgrounds (filled blue and green) are stacked and overlaid with an example signal point (gluino pair production and decay to bottom quarks, in red). Even with other mismeasurement-suppression selections applied, the mismeasurement background still dominates without $M_{T2} > 200$ GeV. Taken from [46].

The residual mismeasurement background is estimated using a procedure called Rebalance and Smear that was newly implemented for this edition of the classic analysis, and is described in Section 3.1.5.

Lost Lepton

The lepton veto rejects most events containing neutrinos originating from the decay of a W boson, $W^\pm \rightarrow \ell^\pm \nu$. However, the lepton is not always successfully reconstructed, usually because it is a τ that decays hadronically and is mistaken for a meson, and this residual so-called lost lepton background must be estimated.

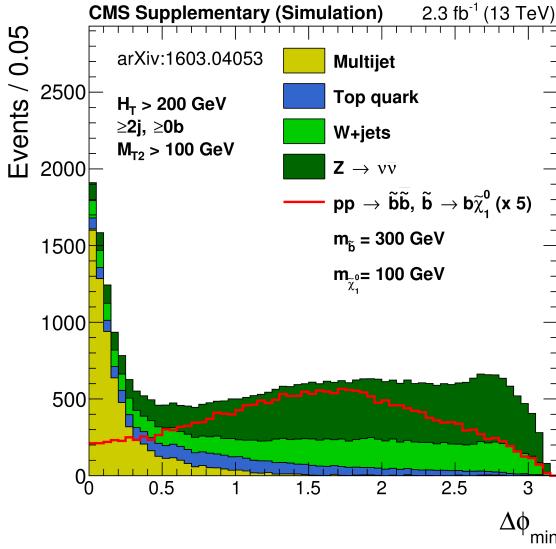


Figure 3.5: The distributions in $\Delta\phi_{min}^{1,2,3,4}$ of the QCD background (filled yellow) and the neutrino backgrounds (filled blue and green) are stacked and overlaid with an example signal point scaled up by a factor of 5 (bottom squark pair production, in red), with the M_{T2} selection relaxed to 100 GeV. Note that most QCD lies below the cut value of 0.3. Taken from [46] supplementary materials.

A first attempt might be to simulate events containing W bosons, scale the simulation to the desired luminosity, and count the events in which the lepton is not reconstructed. This procedure would have reasonable accuracy, but producing accurate simulations is not trivial, and the result would be subject to an array of systematic errors. Schematically, if N_{LL}^{Data} is the actual number of lost lepton events in data and N_{LL}^{MC} is the number predicted by Monte Carlo simulation, the simulation will mispredict by some factor ϵ such that $N_{LL}^{\text{Data}} = \epsilon N_{LL}^{\text{MC}}$.

It is possible to do better than ϵ with data driven techniques. In addition to the lost lepton events, one can also ask the simulation for its prediction of the number of W events in which the lepton is not lost, single lepton events, N_{SL}^{MC} , where the lepton is required to be an electron or muon since τ reconstruction is much more difficult. Every part of this simulation is exactly identical to the lost lepton simulation, subject to almost the same errors, with the major exception being the predicted lepton reconstruction efficiency. Call this new error factor δ so that $N_{SL}^{\text{Data}} = \delta N_{SL}^{\text{MC}}$. Then $\epsilon/\delta = \gamma$ is the portion of the misprediction due to the simulation's imperfect knowledge

of the lepton reconstruction efficiency and a few other more minor uncorrelated effects, a small portion of the total. The ratio $R_{\text{MC}}^{0\ell/1\ell} = N_{LL}^{\text{MC}}/N_{SL}^{\text{MC}}$ is subject only to this relatively small error γ , since the fully correlated errors cancel. The prediction of N_{LL}^{Data} follows directly,

$$N_{LL}^{\text{Data};\text{Est}} = R_{\text{MC}}^{0\ell/1\ell} N_{SL}^{\text{Data}}. \quad (3.1)$$

The input N_{SL}^{Data} is measured in a control region populated by single lepton events observed in data, in a kinematic region identical to the corresponding lost lepton signal region. The remaining systematic uncertainty is only about 15% in most signal regions.

Table 3.2: Summary of systematic uncertainties in the lost-lepton background prediction, together with their typical size ranges across the search bins. (\dagger) In every topological region, bins along the M_{T2} axis with insufficient statistics to allow a fully data-driven estimate are assigned an M_{T2} shape from simulation, normalized to the integral of the impacted bins in the data control region. Taken from [30].

Source	Range [%]
Limited size of data control samples	5–100
Limited size of MC samples	0–50
e/μ efficiency	0–10
τ efficiency	0–3
b -tagging efficiency	0–3
Jet energy scale	0–5
M_T (lepton, \vec{E}_T) selection efficiency	0–3
M_{T2} shape uncertainty †	0–40
Renormalization and factorization scale variation	0–5
$t\bar{t}b\bar{b}/t\bar{t}jj$ weight	0–25

Other uncertainties affecting the lost lepton background prediction are summarized in Table 3.2, together with their typical size ranges across the search bins. The first two uncertainties due to limited data control region and Monte Carlo simulation sample sizes are purely statistical. The remainder are systematic.

The efficiency to reconstruct leptons is, as mentioned, the primary uncertainty of the simulation with respect to data, that does not drop out after taking a ratio of lost- and single-lepton

events.

The impact of the τ reconstruction efficiency is smaller, as it is not used in the single lepton control region, however uncertainty on the rate at which the τ veto fails to reject events contributes to uncertainty in the expected yield.

The expected transfer factor from the single lepton to lost lepton regions varies slightly with respect to the binning variables, H_T , N_{jet} , $N_{b-\text{tag}}$, and M_{T2} , and so is recalculated in simulation for each bin. Since the transfer factor varies, they may be incorrectly estimated if simulation does not accurately reproduce the b -tagging efficiency and jet energy assignments in real data.

To restrict the single lepton control region to leptons consistent with originating from a leptonic W decay, the lepton and \cancel{E}_T system is required to satisfy $M_T(\text{lepton}, \cancel{E}_T) < 100 \text{ GeV}$, since an on-shell W can never decay to a lepton-neutrino system more massive than this, and $M_T \leq M$.

The renormalization and factorization scales are generic issues with theoretical calculations, discussed in Section 2.4.2.

Finally, the $t\bar{t}bb/\bar{t}\bar{t}jj$ weight is a special systematic assessed to cover an ad hoc procedure that corrects a particularly severe data-simulation discrepancy affecting the relative rates of $b\bar{b}$ ISR and light flavor ISR in simulated $t\bar{t}$ events. This simulation error is not significant for the majority of CMS analyses, however it is important for analyses like this search that include bins with large $N_{b-\text{tag}}$, for which $t\bar{t}bb$ is the only meaningful background.

In closing, it should be noted that while the lost lepton background tends to be subdominant relative to the Invisible Z background discussed in the next section, it is the largest background in certain high N_{jet} , high $N_{b-\text{tag}}$, high H_T bins since Z events do not populate these bins efficiently, while $t\bar{t}$ pair-production events do, and all $t\bar{t}$ events contain two W bosons that may decay leptonically.

Invisible Z

The $Z \rightarrow vv$ background is predicted in a similar fashion to the lost lepton background, using a control region populated with $Z \rightarrow \ell^+\ell^-$ events. Again, the leptons are restricted to pairs of electrons and pairs of muons, since τ reconstruction is much more difficult. Figure 3.6 (right) shows the similarity in the M_{T2} distributions of simulated $Z \rightarrow vv$ events and observed $Z \rightarrow \ell^+\ell^-$ events in which one pretends that the leptons are invisible, demonstrating both that these events are kinematically very similar as expected, and that the simulation is accurate. The ratio $R_{MC}^{vv/\ell^+\ell^-}$ can be extracted from Monte Carlo simulation, and is dominated by only the lepton reconstruction efficiency uncertainty. In the Standard Model, this ratio is almost exactly 3, but it is significantly larger experimentally because it is possible for one of the leptons in $Z \rightarrow \ell^+\ell^-$ not to be well-reconstructed, causing the affected event not to be counted. The value $N_{\ell^+\ell^-}^{\text{Data}}$, unlike N_{SL}^{Data} , is not trivial to extract, since a non-negligible fraction of double lepton events come from sources other than a single Z boson, almost always one each from a pair of W bosons. $N_{\ell^+\ell^-}^{MC}$ and N_{vv}^{Data} are exclusively $Z \rightarrow \ell^+\ell^-$ and $Z \rightarrow vv$ events, respectively, and for maximal cancellation of systematic uncertainties, it is desirable that $N_{\ell^+\ell^-}^{\text{Data}}$ be purified to the greatest extent possible.

Fortunately, there is an experimental handle on the contamination. When a Z boson decays to a lepton pair, the flavor is always identical, either two electrons or two muons. When a pair of W bosons each decay to a lepton-neutrino pair, their choices of lepton are uncorrelated. Half the time, the flavors will be identical as for a Z event, and half the time, one W will decay to a muon and the other to an electron. The first case is the undesired impurity in $N_{\ell^+\ell^-}^{\text{Data}}$, and events of the second type are used to populated a different-flavor control region used to predict the impurity, N_{DF}^{Data} . It is nearly sufficient simply to subtract N_{DF}^{Data} from $N_{\ell^+\ell^-}^{\text{Data}}$ since the different flavor and same flavor W events occur at the same rate. However, it is possible that the detector is slightly more or less efficient at reconstructing events where both leptons are the same flavor than events in which they are different flavors, so the different flavor count must be scaled slightly to compensate, by a factor $R^{SF/DF}$. $R^{SF/DF}$ is measured in data in $\ell^+\ell^-$ events that are kinematically

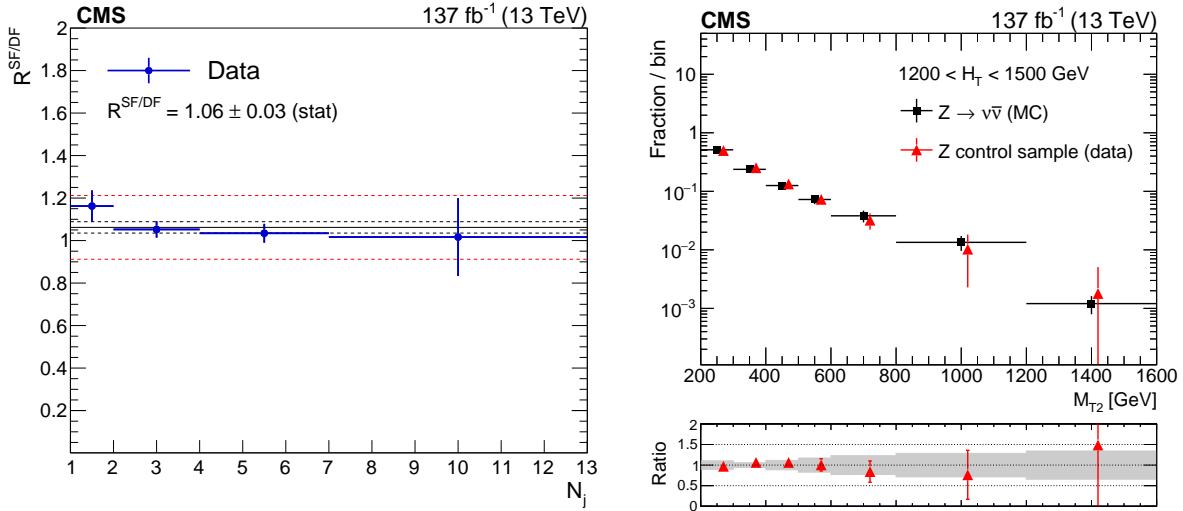


Figure 3.6: (Left) The ratio of the number of same flavor to different flavor lepton pairs in data, $R_{SF/DF}$, which is a component of the $Z \rightarrow vv$ background estimate from $Z \rightarrow \ell^+\ell^-$ events. (Right) A comparison of the M_{T2} shape between data $Z \rightarrow \ell^+\ell^-$ events and simulated $Z \rightarrow vv$ events. The two processes should be kinematically identical, and the similarity of the distributions indicates that the Monte Carlo simulation models the processes well. Taken from [30].

inconsistent with originating from a Z decay, chiefly due to a requirement that the invariant mass of the lepton pair be at least 20 GeV away from the Z mass. One finds that $R^{SF/OF} \approx 1.06 \pm 0.15$, as shown in Figure 3.6 (left). The final prediction is

$$N_{vv}^{Data;Est} = R_{MC}^{vv/\ell^+\ell^-} (N_{\ell^+\ell^-}^{Data} - R^{SF/OF} N_{DF}^{Data}). \quad (3.2)$$

Being irreducible, the $Z \rightarrow vv$ plus jets background is dominant in the vast majority of bins.

The uncertainties in the $Z \rightarrow vv$ background prediction are summarized in Table 3.3 together with their typical size ranges across the search bins. Only the uncertainty in $R^{SF/DF}$ is unique to $Z \rightarrow vv$, while the others are shared with the lost-lepton estimate.

Table 3.3: Summary of systematic uncertainties in the $Z \rightarrow vv$ background prediction, together with their typical size ranges across the search bins. (\dagger) In every topological region, bins along the M_{T2} axis with insufficient statistics to allow a fully data-driven estimate are assigned an M_{T2} shape from simulation, normalized to the integral of the impacted bins in the data control region. Take from [30].

Source	Range [%]
Limited size of data control samples	5–100
Limited size of MC samples	0–50
Lepton efficiency	0–5
Jet energy scale	0–5
Uncertainty in $R^{\text{SF/DF}}$	0–5
M_{T2} shape uncertainty †	0–40

Observable	Selection	Notes
M_{T2}	$> 200 \text{ GeV}$	Only for multijet events. Increased to $M_{T2} > 400 \text{ GeV}$ for $H_T > 1500 \text{ GeV}$ to maintain QCD suppression.
p_T^{jet1}	$> 250 \text{ GeV}$	Only for monojet events.
H_T	$> 250 \text{ GeV}$	Motivated by available triggers. Background events at lower H_T are too common for these events to be always recorded.
E_T	$> 250 \text{ GeV}$	Relaxed to $E_T > 30 \text{ GeV}$ for $H_T > 1200 \text{ GeV}$. Motivated by available triggers.
$\Delta\phi_{min}^{1,2,3,4}$	> 0.3	Auxiliary mismeasurement suppression.
$ \vec{H}_T - \vec{E}_T / E_T$	< 0.5	Auxiliary mismeasurement suppression.
N_{lep}	$= 0$	The lepton veto; rejects the majority of $W^\pm \rightarrow \ell^\pm v$ background.

Table 3.4: A summary of the baseline event selection for the classic M_{T2} search. Events are required to have large H_T , no leptons, and significant missing energy unlikely to be the product of detector mismeasurement or a single undetected particle.

3.1.4 Baseline Selection

The properties of these signals and backgrounds motivate the baseline event selection summarized in Table 3.4. The M_{T2} selection primarily suppresses the mismeasurement background, by many orders of magnitude. The $\Delta\phi_{min}^{1,2,3,4}$ and H_T selections also help to suppress this background further, until it is smaller than the genuine \cancel{E}_T backgrounds. The H_T and \cancel{E}_T selections are chosen primarily so that all of the selected events pass the online trigger. Backgrounds are so common at $H_T < 250$ GeV, and $\cancel{E}_T < 250$ GeV for $H_T < 1200$ GeV, that the experiment is unable to record all of the events observed, making this part of the parameter space a poor region to search for a rare signal in any case. The lepton veto rejects most $W^\pm \rightarrow \ell^\pm v$ events, so that only events in which the lepton is not reconstructed remain in the signal region, as described in Section 3.1.3, and serves to narrow the analysis’ focus to the all-hadronic final state as part of CMS’s larger research program.

In addition to this baseline selection, the analysis bins in M_{T2} , H_T , N_{jet} , and N_{b-tag} to enhance signal sensitivity as described in Section 3.1.5.

3.1.5 Upgrades in 2019

The classic search has been performed before at 13 TeV, once in 2015 [46] using a 2.3 fb^{-1} dataset, and again in 2016 [63] using 35.9 fb^{-1} . The update in 2019 uses the full dataset from 2016, 2017, and 2018, totaling 137 fb^{-1} . As the analysis is dominated by statistical uncertainties in its most sensitive bins, the increased statistical power is the primary improvement in the update.

The update includes two other major upgrades.

The first improves the estimate of the mismeasurement background using a technique called Rebalance and Smear.

The second leverages the increased statistics to expand the signal region binning, better targeting signal models with more extreme jet and b-tagged jet multiplicities, and M_{T2} .

Rebalance and Smear

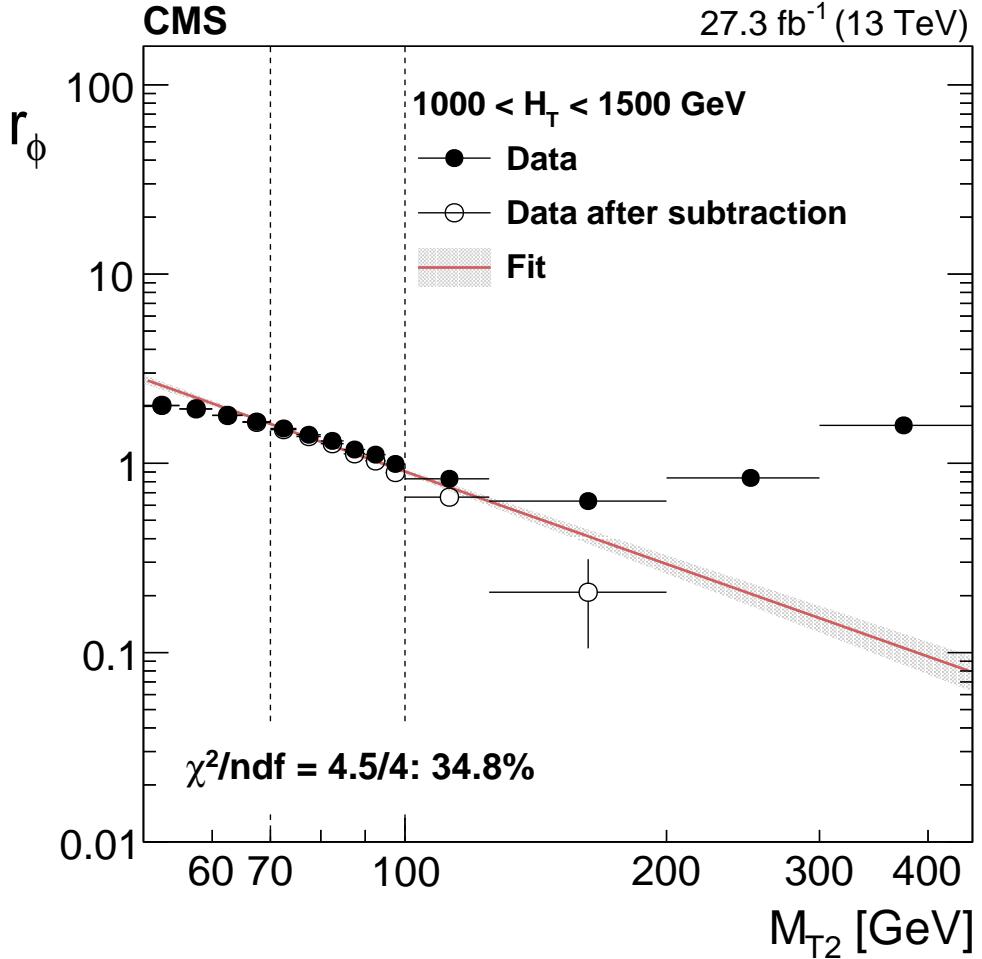


Figure 3.7: The fit of $r_\phi = N_{\Delta\phi_{min}^{1,2,3,4} > 0.3} / N_{\Delta\phi_{min}^{1,2,3,4} < 0.3}$ as a function of M_{T2} obtained in the 2016 classic search, for the $1000 < H_T < 1500$ GeV H_T band. The fit is performed in events in the M_{T2} band $70 < M_{T2} < 100$ GeV and extrapolated to the $M_{T2} > 200$ GeV signal region. Black points represent raw data, while white points represent data after contribution from genuine E_T backgrounds is subtracted. Taken from [63].

In older versions of the classic search [46, 63], the mismeasurement background estimate used the $\Delta\phi_{min}^{1,2,3,4}$ observable and extrapolated across M_{T2} . Events at low M_{T2} and at low $\Delta\phi_{min}^{1,2,3,4}$ are both dominated by QCD mismeasurement. The suppression effect of M_{T2} is so strong that even events at low M_{T2} and *high* $\Delta\phi_{min}^{1,2,3,4}$ are QCD dominated. As essentially every low M_{T2} event is a QCD event, low M_{T2} events can be used to measure the ratio r_ϕ of QCD mismeasurement events at

high and low $\Delta\phi_{min}^{1,2,3,4}$. High M_{T2} events with low $\Delta\phi_{min}^{1,2,3,4}$ can then serve as a control region for estimating the QCD mismeasurement background, $N_{\Delta\phi_{min}^{1,2,3,4}>0.3} = r_\phi N_{\Delta\phi_{min}^{1,2,3,4}<0.3}$. Unfortunately, r_ϕ decreases with increasing M_{T2} , so that instead the dependence must be fit to a power law at low M_{T2} and extrapolated to high M_{T2} , as shown in Figure 3.7. This procedure has obvious potential for statistical errors in the fit, systematic error in extrapolating the fit to high M_{T2} , and potential non-multijet contamination of the low $\Delta\phi_{min}^{1,2,3,4}$ control region at high M_{T2} , producing total relative error at least 40% and as large as 180%. Although the impact of these errors on the overall analysis is controlled by suppressing the mismeasurement background as described in Section 3.1.3, it is desirable to replace this procedure with a more robust one.

Rebalance and Smear (R&S) achieves the desired improvement. R&S begins with a sample of multijet events from data with H_T on the order of hundreds of GeV and at least two jets with $p_T > 10$ GeV. This is as nearly unbiased a sample of QCD events as is allowed by available triggers. These events will generically have some small but nonzero \cancel{E}_T due to imperfect measurement of the jet energies, dictated by the detector's resolution, jets from pileup, very low energy objects in the event, and potentially a very small contribution from genuine neutrinos produced in the hadron decay chains inside the jets. The p_T values assigned to the jets are then adjusted, finding the most likely assignment of p_T values subject to the hypothesis that the true \cancel{E}_T is very nearly zero, and that jet mismeasurements of a given size occur with probability given by jet response templates. As shown in Figure 3.8, large mismeasurements are rare. Thus, the Rebalancing step tries to get the \cancel{E}_T as close to zero as feasible without making more improbable adjustments to the jets than necessary. The output of the Rebalancing step is a large sample of real QCD events with maximally accurate jet p_T assignments.

Then, each of these events goes through a Smearing step. This step randomly assigns a new p_T to each jet in the event according to the same jet response templates. The vast majority of the time, the new event looks as unremarkable as the original event before Rebalancing. Rarely, the output event passes the baseline selection and represents a potential mismeasurement event

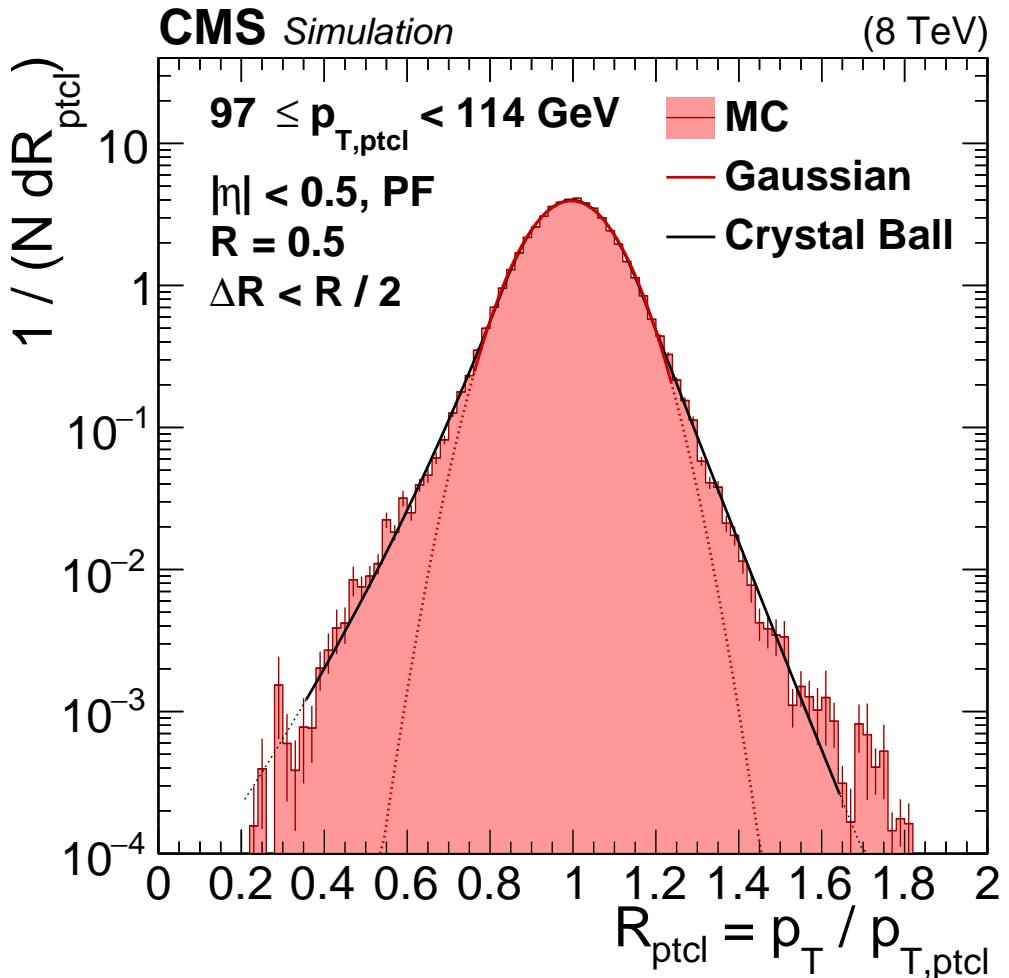


Figure 3.8: $p_{T,\text{ptcl}}$ is the “particle level” p_T , the total energy that the jet’s particle components actually had, while p_T is the energy measured by the detector. The horizontal axis is the ratio of these two quantities, and the vertical axis the probability that a jet’s measured energy will differ from its true energy by that ratio. The response curve shown here applies to jets with particle level p_T in the band $97 < p_T < 114 \text{ GeV}$, and measured in the CMS barrel. Large mismeasurements of jet energy are much less probable than small mismeasurements. The core is well-described by a Gaussian, but the tails are highly non-Gaussian and best described by a Crystal-Ball function. Taken from [48].

that might lie in the data signal region. The Smearing can be repeated as many times for each event as is desired, subject to computing limitations. The number of Smeared events falling into each signal region, from this sample of known equivalent integrated luminosity, can then be converted into the expected number of events that would fall into the signal region in data, due to jet mismeasurement.

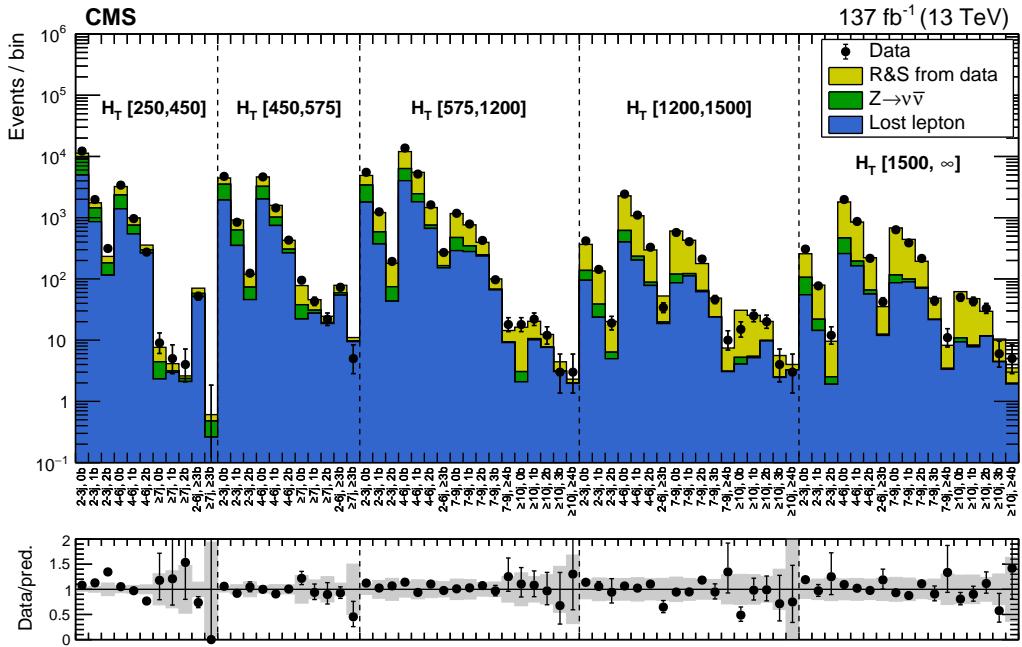


Figure 3.9: The Rebalance and Smear jet mismeasurement background estimate is validated in the $\Delta\phi_{min}^{1,2,3,4} < 0.3$ control region. The multijet mismeasurement event count predicted by Rebalance and Smear is shown in yellow, and contributes most of the event counts in this control region. The combined background estimate (filled histogram) is consistent with the data observation (black data points). Taken from [30].

Of course, many of the output events will have $\Delta\phi_{min}^{1,2,3,4} < 0.3$, falling outside of the signal region into the QCD-dominated low $\Delta\phi_{min}^{1,2,3,4}$ control region. This allows for a validation in data of the R&S procedure, shown in Figure 3.9. The total background estimate is consistent with data in the $\Delta\phi_{min}^{1,2,3,4} < 0.3$ control region across all of the analysis bins, with the R&S estimated counts contributing most of the predicted events.

R&S achieves a significant improvement over the old r_ϕ based system, effectively eliminating statistical error, and improving the worst case relative error on the mismeasurement estimate from 180% to less than 50%. Systematic uncertainties are summarized in Table 3.5 together with their typical size ranges across the search bins. The sources of these uncertainties are described either in the text above, or are shared with the uncertainties affecting the genuine neutrino backgrounds.

Table 3.5: Summary of systematic uncertainties in the multijet background prediction, together with their typical size ranges across the search bins. (\dagger) The parameter σ_T^{soft} is the width of the assumed underlying irreducible \cancel{E}_T distribution for a multijet event, used as an input to the Rebalancing step, and so-called because such \cancel{E}_T is dominated by low energy objects. The final output is largely independent of the precise choice for this parameter and the shape of the underlying distribution, but a systematic is assessed to cover any observed variations. Taken from [30].

Source	Range [%]
Jet energy resolution	10–20
Tails of jet response in templates	17–25
σ_T^{soft} modeling †	1–25
N_{jet} modeling	1–19
$N_{\text{b-tag}}$ modeling	1–16

Expanded Binning

Although the baseline selection defines the class of events in which a signal of interest may be found, most signals will produce a significant number of events in only a subset of the phase space. For instance, a signal producing 4 top quarks in the final state (see Figure 3.1 second row, right) will almost always produce events with $N_{\text{jet}} \geq 7$, $N_{\text{b-tag}} > 0$, and $H_T \sim 1 \text{ TeV}$. If the entire set of selected events is considered together, the entire background is combined, potentially hiding a signal. If instead the phase space is divided into many separate regions, the background in each region is smaller, so that the background count is as small as possible in the small subset of regions that any given signal actually populates. This sensitivity enhancement motivates very fine binning of the signal region. The classic search uses M_{T2} , H_T , N_{jet} , and $N_{\text{b-tag}}$ as binning variables.

In the 2016 version of the classic search [63], the most extreme N_{jet} bin was $N_{\text{jet}} \geq 7$, and the most extreme $N_{\text{b-tag}}$ bin was $N_{\text{b-tag}} \geq 3$. This limitation was imposed by limited statistics. The background estimate for each bin is performed separately, and the observed counts are subject to Poisson statistical fluctuations. Binning too finely causes any potential sensitivity gains from better isolating signal to be lost to greater uncertainty in the expected background. The analysis binning was updated for the latest edition of the classic search [30], anticipating the improved

statistical power. The full set of new bins, along with the predicted background counts and observed event counts in data, is available in Appendix ??.

The new binning extends the old binning in three ways. First, new $N_{\text{jet}} \geq 10$ bins were added to the $H_T > 1200$ GeV regions. This allows for enhanced sensitivity to signals with extremely high jet multiplicity, such as the 4 top signal previously mentioned. Similarly, new $N_{b-\text{tag}} \geq 4$ bins were added to these same regions, targeting the same signal. Sensitivity to some mass points of this signal model roughly doubled due to the addition of these bins, which have negligible background but appreciable signal counts. Finally, M_{T2} binning was generally made narrower and the last bin moved to larger M_{T2} values, for all signal regions, until the expected background in the last bin was on the order of 1 event. In all, there are 282 classic search bins, enhancing sensitivity to a broad array of potential signal models.

3.1.6 Signal Contamination

The data driven estimate procedures for the neutrino backgrounds use control regions defined using leptons. Signals capable of producing leptons can contaminate these control regions, increasing the observed control region counts above the actual Standard Model production rate. This leads directly to an overprediction of background. For instance, consider stop pair production followed by decay to tops, as shown in Figure 3.1 (row 3, right). Both tops will decay to a W and a bottom quark. If one of the W bosons decays leptonically and the lepton is reconstructed, the event will likely enter the single lepton control region. The lost lepton background, described in Section 3.1.3, will be overpredicted,

$$N_{LL}^{\text{Data};\text{Est}} = R_{MC}^{0\ell/1\ell} N_{SL}^{\text{Data}} = R_{MC}^{0\ell/1\ell} \left(N_{SL}^{\text{Data};\text{SM}} + N_{SL}^{\text{Data};\text{Sig}} \right). \quad (3.3)$$

The background overprediction is $\Delta N = R_{MC}^{0\ell/1\ell} N_{SL}^{\text{Data};\text{Sig}}$. For analysis purposes, the overprediction is modeled in simulation and treated as a reduction of the expected signal counts in every affected

bin.

$$N_{SR}^{\text{Sig;Adjusted}} = N_{SR}^{\text{Sig;Raw}} - \Delta N \quad (3.4)$$

This adjustment has the nice property that all terms are linear in the signal strength, so that it does not need to be recalculated for every signal strength considered when performing statistical analysis of the results. The same fraction of the signal is lost at every signal strength.

As a result of this loss of sensitivity due to signal contamination, the classic analysis is less effective, all things equal, when used to search for signals that sometimes produce leptons than the naive expectation based on the leptonic versus hadronic branching ratios. This is an inevitable consequence of performing an all-hadronic search.

3.1.7 Results

The full set of results for every classic search signal region, including every background prediction and the observed count, are available in Appendix ???. The full set of results integrating over the M_{T2} binning are displayed in Figure 3.10 (upper), and the results including the M_{T2} binning for the $575 < H_T < 1200$ GeV bins, the largest set, are shown in Figure 3.10 (lower). The observed counts are consistent with the background-only hypothesis, and the results are used to set exclusion limits at 95% confidence level on the signals discussed in Section 3.1.2.

3.1.8 Limits

The CL_S statistical analysis procedure applied in this analysis is described in [57]. It begins with maximum likelihood fits to the background-only and signal-plus-background hypotheses, for each signal model, considering each mass point separately. The likelihoods are products of Poisson probabilities for each signal region bin, with log-normal constraint terms for each systematic affecting the predicted counts. Correlation between uncertainties affecting different bins are fully accounted for. As stated in the previous section, the background-only hypothesis is consistent

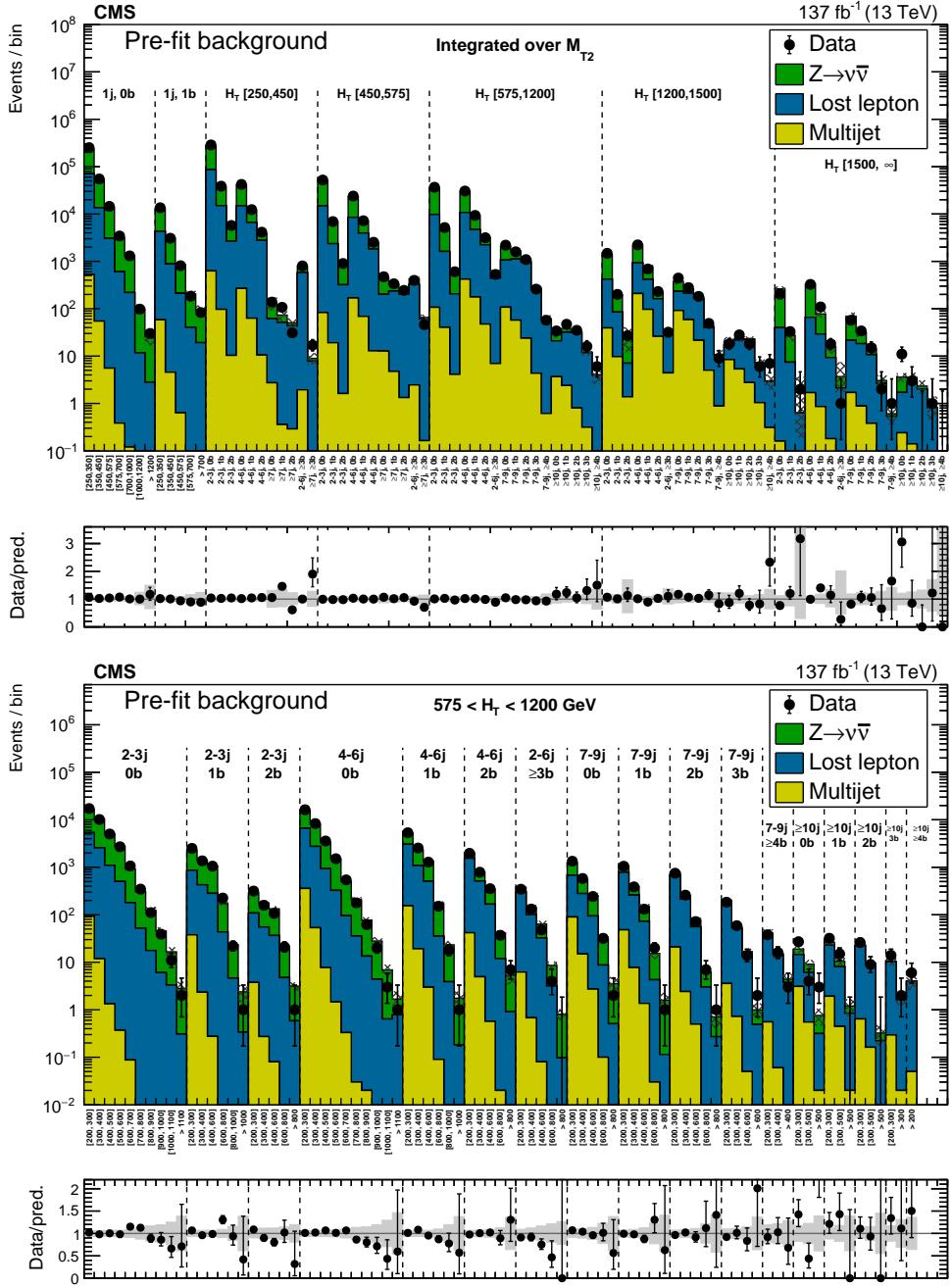


Figure 3.10: Comparison of predicted background and observed data events in the classic search (upper) integrated over M_{T2} and (lower) for each of the medium H_T search regions. Taken from [30].

with the observation. So, the parameter of interest for each signal model is the maximum signal cross section that can be excluded with 95% confidence level. That is, the maximum possible

production rate that the signal model could have, without requiring a combination of background and signal fluctuations more improbable than 1 in 20 to be consistent with the data. If the production cross section excluded at 95% CL is less than the theoretical cross section for a given signal model, that signal model is said to be excluded at 95% CL.

The simplified supersymmetric extensions to the Standard Model shown in Figure 3.1 have only two free parameters, the mass of the pair-produced superpartner, and the mass of the lightest supersymmetric particle, the dark matter candidate $\tilde{\chi}_1^0$. Plotting the gluino or squark mass on the horizontal axis and the mass of $\tilde{\chi}_1^0$ on the vertical axis, then marking the mass points excluded at exactly 95% CL, produces the exclusion curves shown in Figures 3.11–3.14. Points to the lower left of these curves are excluded, while points above and to the right are not, as they require fluctuations no more improbable than 1 in 20 to be consistent with present data.

The CMS dataset was collected only once, and this dataset is subject to fluctuations. The multiple exclusion curves shown on each plot, some in red and some in black, provide an indication of how unusual the limits generated by this dataset were. Suppose that the background’s expected event count in a given region, summing across this entire dataset, is B . It is clearly possible that the actual number of background events that occur in the dataset in this bin could be around $B + \sqrt{B}$, or $B - \sqrt{B}$. In the first case, the analysis will draw unexpectedly weak limits on signals that populate this bin, as the additional background appears consistent with signal. In the second case, the analysis will draw unexpectedly strong limits on signals that populate this bin, as there are not even enough events to cover the expected background count.

The curves shown in red are the median, 1 standard deviation, and 2 standard deviation expected exclusion curves, in the background-only hypothesis. Together, these indicate the typical range of exclusion curves across many imaginary CMS datasets in which there is no signal to find. The curves in black are those that were observed in the CMS dataset as actually recorded. When the black observed curves extend out beyond the red expected curves, the analysis likely benefited from a downward fluctuation of background. In cases where the observed curves swing

inward compared to the expected curves, the analysis may have experienced an analogous unlucky upward fluctuation of background that mimicked a small signal, or alternatively there may be a genuine small signal lurking in the data that is resisting exclusion!

The general shape of the curves is set by a combination of the falling cross section with increasing gluino or squark mass shown in Figure 1.12, and a loss of signal efficiency and signal versus background discriminatory power when the mass splitting between the gluino or squark and $\tilde{\chi}_1^0$ is small. On the lower right hand side of the plots lie signals that produce spectacular, energetic events, but at a very low rate due to the large mass of the gluino and squark, and accordingly small pair-production cross section. Moving up along the $\tilde{\chi}_1^0$ mass axis towards the upper right, the events remain spectacular, so the exclusion curves remain roughly vertical, limited only by the production cross section. Eventually, the mass splitting becomes small enough that the loss of signal efficiency and signal versus background discriminatory power becomes significant, and the exclusion curve turns to the left, towards lower mass squarks or gluinos and higher production rates. Decreasing the gluino or squark mass at fixed $\tilde{\chi}_1^0$ mass reduces the mass splitting, so the curve generally must also drop down somewhat along the $\tilde{\chi}_1^0$ mass axis to maintain a reasonable splitting. Eventually, the curve intercepts the $M_{\tilde{\chi}_1^0} = M_{\tilde{g}}$ or $M_{\tilde{\chi}_1^0} = M_{\tilde{q}}$ line. At this mass and below, all points are excluded even in the limit of zero mass splitting.

Figure 3.11 shows the exclusion curves for gluino pair production and decay to light quarks (upper), light quarks and the Z boson (lower left), and light quarks and the W boson (lower right). Figure 3.12 shows the exclusion curves for gluino pair production and decay to bottom (left) and top (right) quarks.

Figure 3.13 shows the exclusion curves for pair production of light-flavor squarks (upper left), bottom squarks (upper right), and top squarks in which the top squark decays to a top quark (lower). The light-flavor figure contains two curves, one which assumes that there is only a single low mass light flavor squark, and another that assumes that there are eight light flavor squarks of (approximately) degenerate mass, which implies a production cross section eight times larger.

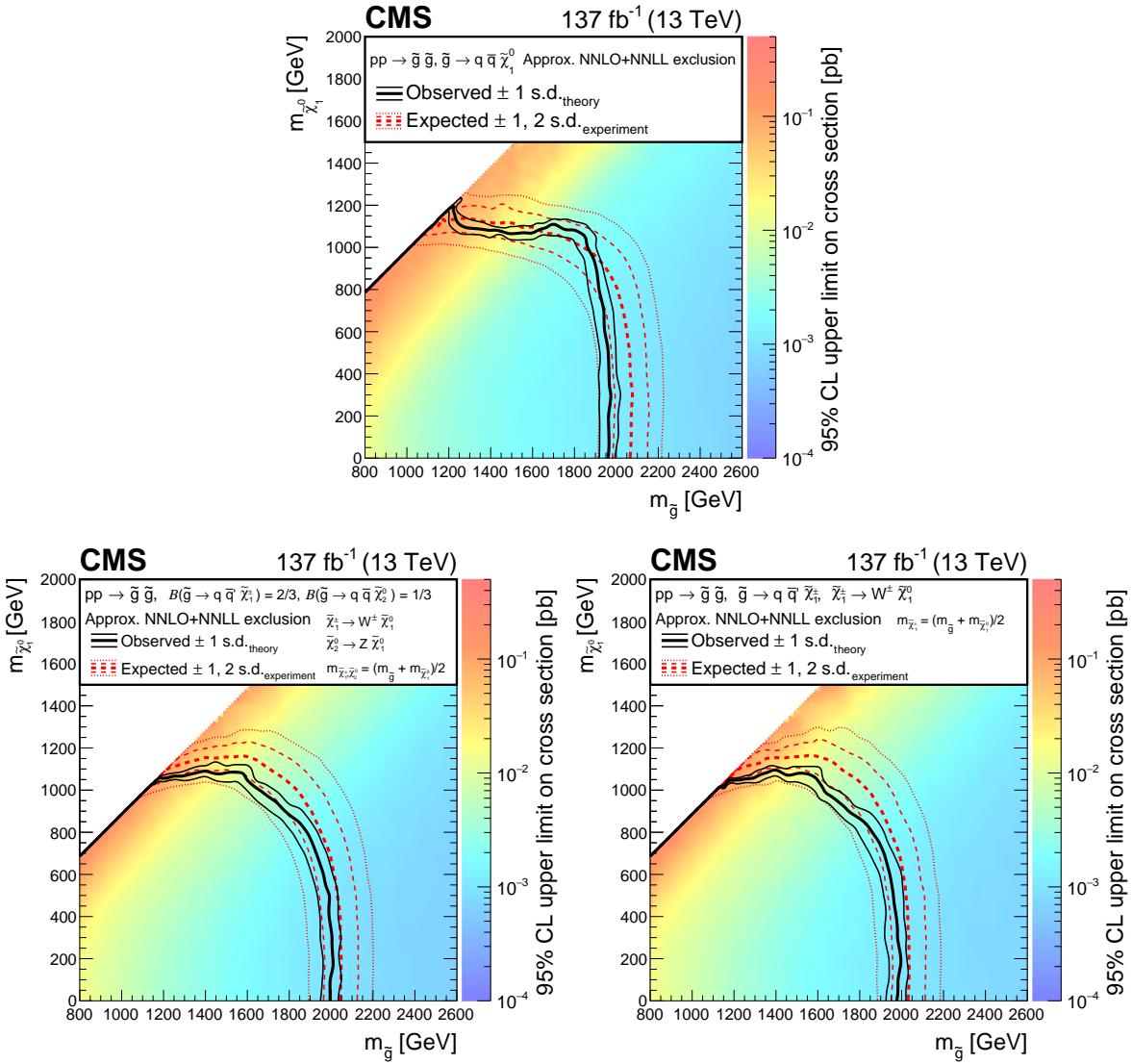


Figure 3.11: Exclusion limits at 95% CL_S for gluino pair production and decay to a pair of light quark jets and (upper) $\tilde{\chi}_1^0$, (lower left) a democratic split between $\tilde{\chi}_1^0, \tilde{\chi}_2^0$, which then decays to a Z boson and $\tilde{\chi}_1^0$, and $\tilde{\chi}^\pm$, which decays to a W boson and $\tilde{\chi}_1^0$, and (lower right) $\tilde{\chi}^\pm$ then W and $\tilde{\chi}_1^0$ with 100% branching fraction. Taken from [30].

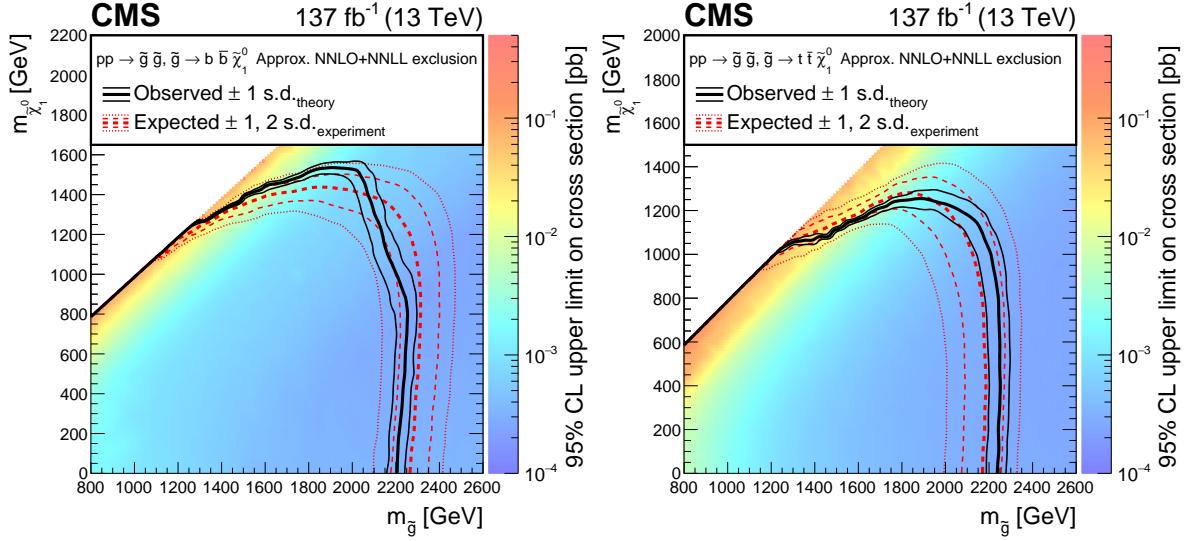


Figure 3.12: Exclusion limits at 95% CL_S for gluino pair production and decay to (left) bottom quarks and (right) top quarks. Taken from [30].

The other top squark decay modes, in which the top decays to (upper left) a bottom quark and $\tilde{\chi}^\pm$, which subsequently decays to a W boson and $\tilde{\chi}_1^0$, (lower) a charm and $\tilde{\chi}_1^0$, and (upper right) either $\tilde{\chi}^\pm$ and a bottom quark or $\tilde{\chi}_1^0$ and a top quark, are shown in Figure 3.14. The charm decay channel is only shown for signal models with small mass splittings, as the top squark would strictly prefer to decay to a top quark and $\tilde{\chi}_1^0$ than to a charm quark and $\tilde{\chi}_1^0$ if kinematically allowed.

Figure 3.15 shows limits placed on the mono- ϕ model. Here, the horizontal axis is the mass of the singly-produced scalar ϕ , and the vertical axis the mass of the invisible fermion ψ . A star indicates the mass point proposed by the original authors in [13] as most phenomenologically interesting, which is not yet excluded. It is worth emphasizing that the background model is nevertheless consistent with data; this signal is simply very difficult to exclude with the M_{T2} analysis methodology. To save computing resources, the mono- ϕ model is only simulated in the phenomenologically interesting subset of the mass plane, similarly to the top squark to charm model in Figure 3.14 (lower), hence the large white space beneath the considered range of masses.

Figure 3.16 shows limits for leptoquarks decaying to (upper left) a light flavor quark and neutrino, (upper right) a bottom quark and neutrino, and (lower) a top quark and neutrino. As

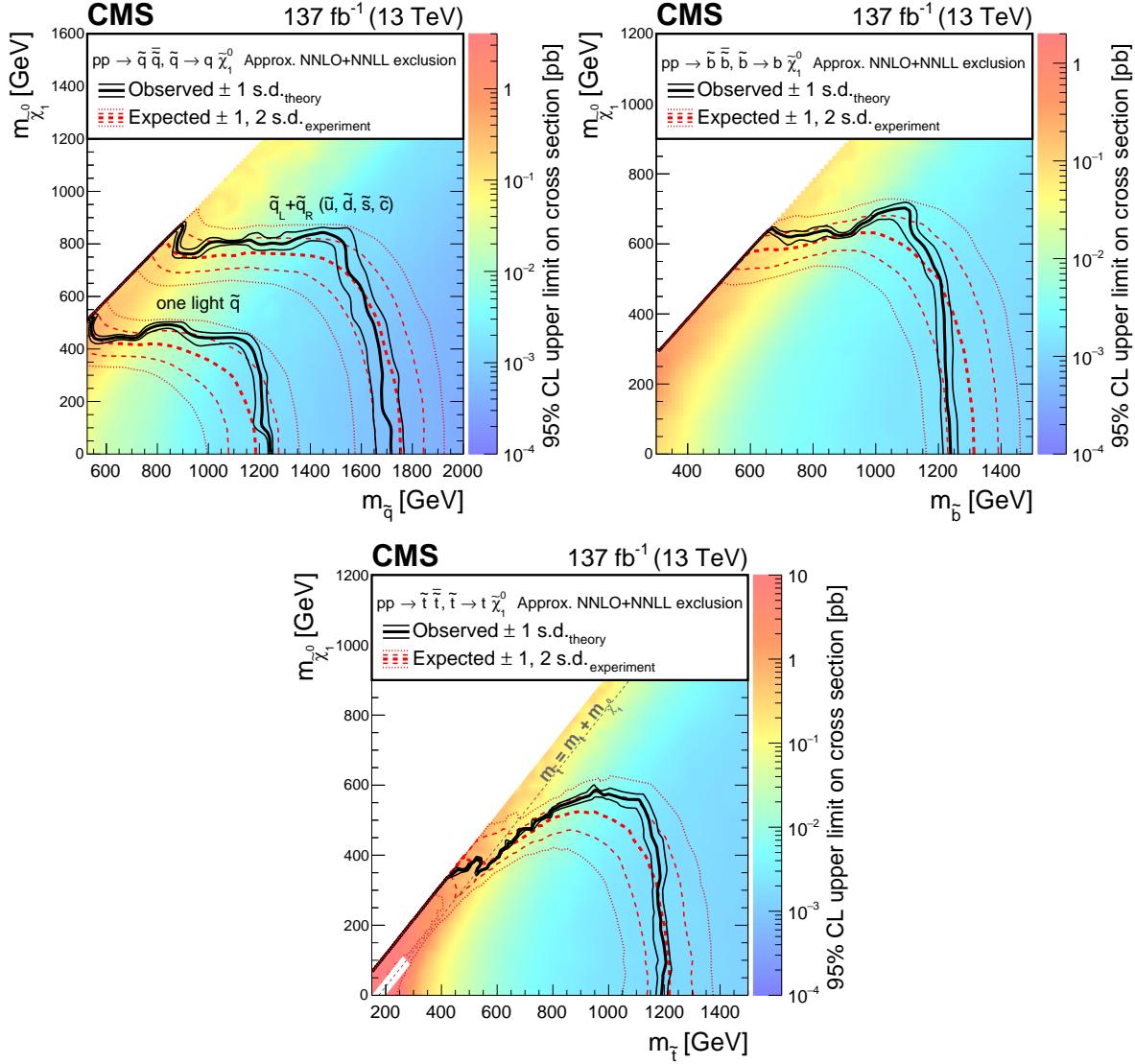


Figure 3.13: Exclusion limit at 95% CL_S for (upper left) light-flavor squark pair production, (upper right) bottom squark pair production, and (lower) top squark pair production. Taken from [30].

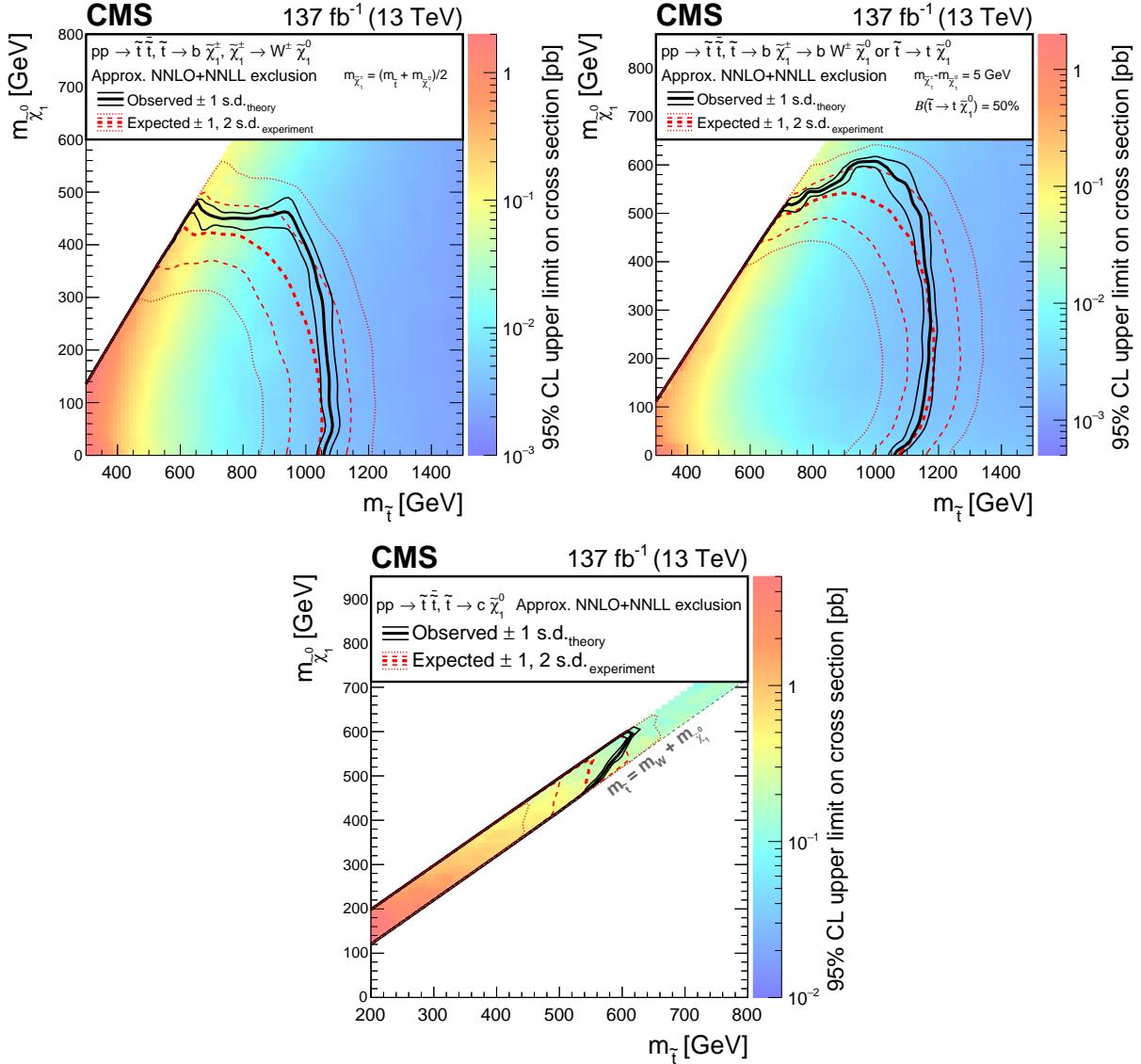


Figure 3.14: Exclusion limits at 95% CLs for top squark pair production and decay to (upper left) a bottom quark and $\tilde{\chi}_1^\pm$, which subsequently decays to a W boson and $\tilde{\chi}_1^0$, (upper right) either a bottom quark and $\tilde{\chi}_1^\pm$ or top quark and $\tilde{\chi}_1^0$, or (lower) a charm quark and $\tilde{\chi}_1^0$. Taken from [30].

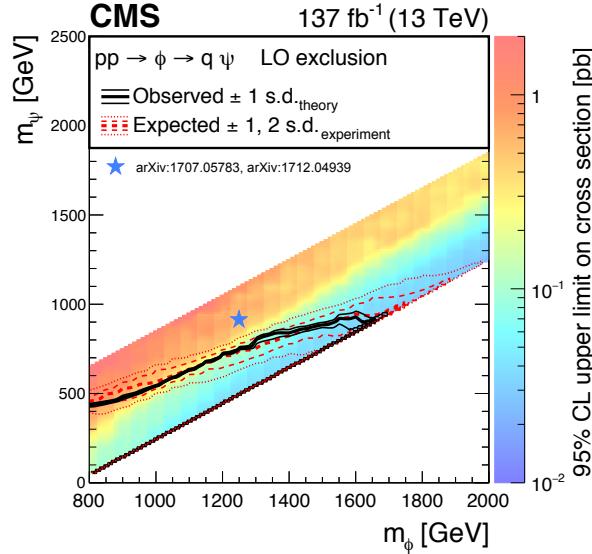


Figure 3.15: Exclusion limit at 95% CL_S for the mono- ϕ model. Only the portion of the mass plane of phenomenological interest was simulated. The star indicates the original authors' proposed best fit mass point, which remains unexcluded. Taken from [30].

the mass of the neutrinos, in contrast to the mass of $\tilde{\chi}_1^0$, are known to be approximately zero, the leptoquark limits are one dimensional, in the leptoquark masses.

The limits produced by this edition of the classic search improve upon the limits set by the previous edition [63] by hundreds of GeV, and in most cases are the strongest constraints on their respective signal models yet produced by any experiment.

3.1.9 Future of the Classic M_{T2} Search

The current limits produced by the classic M_{T2} search are impressive, and are unlikely to improve much in the near future. The pair production cross section for squarks and gluinos drops rapidly with mass, as shown in Figure 1.12. Generically, a factor of 10 improvement in sensitivity is necessary to push the exclusion limits outward by around 500 GeV along the horizontal axis. At this mature statistical stage, a factor of 10 improvement in sensitivity requires a factor of 100 increase in the integrated luminosity, unfeasible in the near future. Improvement along the $\tilde{\chi}_1^0$ axis is similarly difficult due to large backgrounds and low signal efficiency for models with small

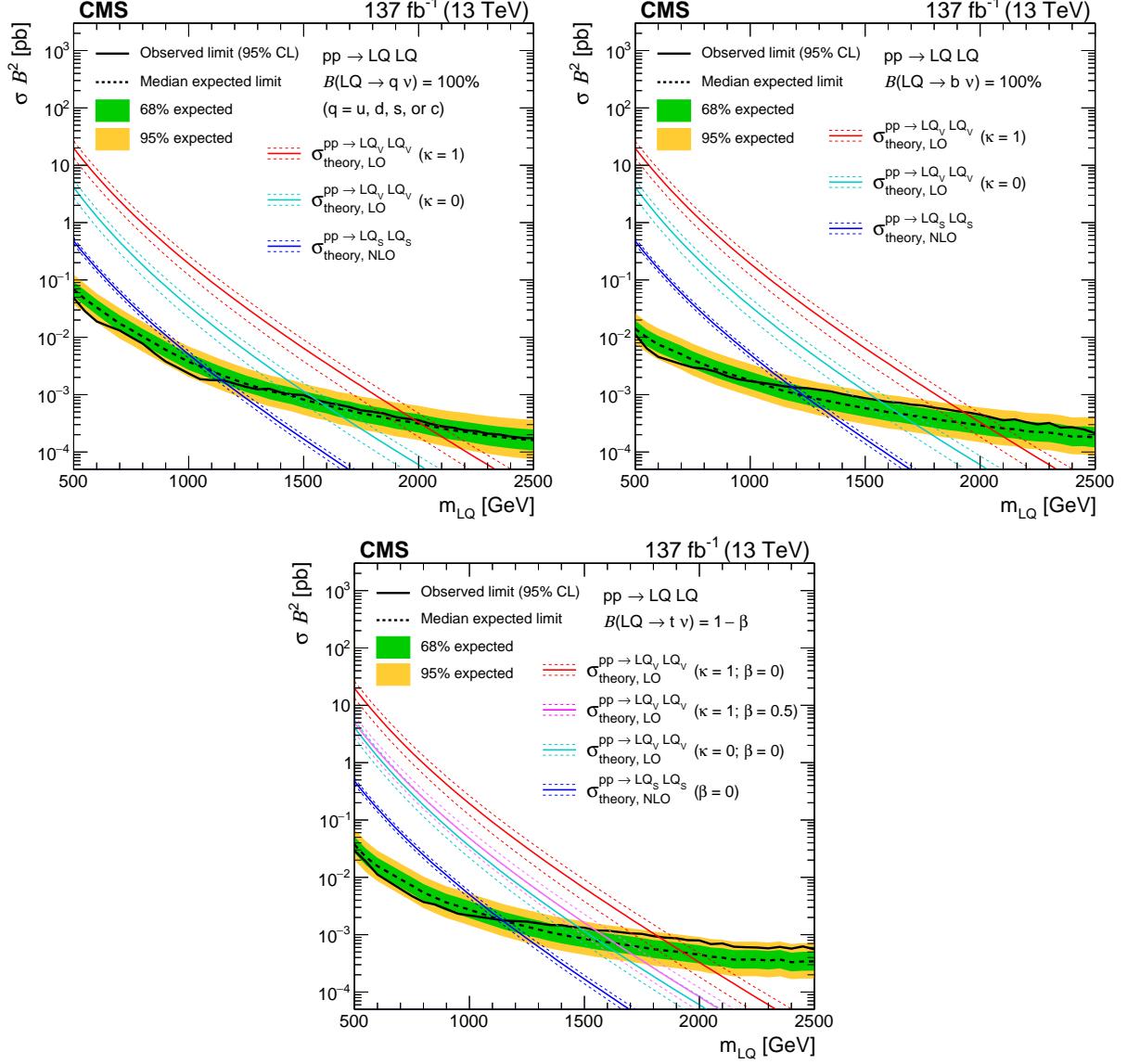


Figure 3.16: Upper limits at 95% CL_S on the leptoquark production cross sections as a function of leptoquark mass. Unlike the limits on supersymmetric models, in which the mass of $\tilde{\chi}_1^0$ is a free parameter, the limits on leptoquarks are 1-dimensional in the leptoquark mass since the neutrino masses are known to be approximately zero. Taken from [30].

mass splittings. Therefore, attention in the near future will turn to other new techniques, one of which is discussed in the next section.

3.2 Disappearing Tracks Search

3.2.1 General Description

As searches like that discussed in the previous section approach the practical limits of their sensitivities at the LHC, interest has grown in targeting plausible signals with peculiar features that more standard searches do not exploit. Among these features are those produced by relatively long-lived particles (LLPs) that do not decay promptly at the collision point, but instead travel a macroscopic distance into the detector before decaying. An overview of some possibilities in the context of supersymmetry is provided in [53].

One of these signatures is a disappearing track, produced when a charged particle, such as $\tilde{\chi}^\pm$, travels into the tracker, then decays to an invisible particle, such as $\tilde{\chi}_1^0$, with a mass so nearly equal to that of the decaying charged particle that the visible products of the decay are too low energy to be reconstructed. The supersymmetric parameter space is vast and there is no way to know which if any version of supersymmetry is realized in nature, but models including LLPs that would produce disappearing tracks at CMS have been argued to be especially plausible on theoretical grounds [41, 56, 17]. The M_{T2} analysis is already well-optimized for all-hadronic supersymmetry, so it is natural to search for all-hadronic decays of these models by extending it with a disappearing track search. This extension begins with exactly the same set of selected events as the classic search, and then further selects the subset of events possessing a disappearing track.

3.2.2 Challenges

The disappearing tracks search must confront a pair of significant challenges.

First, the background (see Section 3.2.4) entirely consists of events in which reconstruction failed, analogous to the classic search’s detector mismeasurement background that was discussed in Section 3.1.3, and is difficult to understand and estimate for many of the same reasons. Simulation cannot be trusted to account for all of the detector details that lead to extremely rare reconstruction errors. While the classic search can mitigate the impact of this issue with selections that make the mismeasurement background subdominant compared to the better understood neutrino backgrounds, this is not possible for the disappearing track search since no such genuine background exists. Moreover, since the background is dependent on details of the detector, it is strongly affected when these details change. Most prominently, the CMS pixel detector was entirely replaced between the 2016 and 2017 data taking periods [60], and this profoundly affected CMS track reconstruction, to the extent that the disappearing track search must treat the 2016 data entirely separately from the 2017 and 2018 data. The detector’s state also evolves more subtly over the data taking period, due in large part to inevitable radiation damage, and the analysis must track, study, and account for the effect of this evolution on the occurrence of disappearing tracks.

Second, disappearing tracks are extremely rare. This is of course one of the motivations for the disappearing tracks extension—they are a powerful tool to discriminate between background and likely signal events—but it also makes them difficult to study.

Nevertheless, the disappearing tracks search produced a data-driven background estimation procedure, described in Section 3.2.5, that was successfully validated in data, and achieved the best sensitivity to disappearing tracks in supersymmetric models to date.

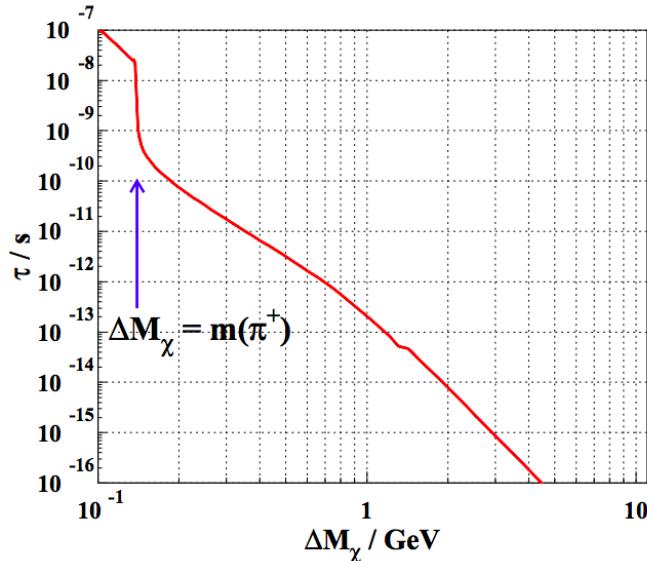


Figure 3.17: The rest frame lifetime of $\tilde{\chi}^\pm$ as a function of the $\tilde{\chi}^\pm$ - $\tilde{\chi}_1^0$ mass splitting. The lifetime is macroscopic for mass splittings less than 1 GeV in a broad class of supersymmetric models, a regime that is realized when the superpartner mass scale is much larger than the Standard Model mass scale. This mass splitting is so small that the charged daughter of $\tilde{\chi}^\pm$ decay is too low energy for track reconstruction. Thus, $\tilde{\chi}^\pm$ lives long enough to produce a short track in the CMS tracker, then disappears when it decays to $\tilde{\chi}_1^0$ and a lost charged daughter. For mass splittings less than $M_\pi \approx 140$ MeV, indicated on the plot, the lifetime of $\tilde{\chi}^\pm$ is so long that it typically does not decay inside the tracker. As a result, the sweet spot for the disappearing tracks search is mass splittings on the order of hundreds of MeV. The discontinuity at a mass splitting of 1.4 GeV is not physical. Taken from [17].

3.2.3 Signal Models

For the class of supersymmetric models discussed in [56] & [17], the relative mass splitting of $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_1^0$ is proportional to $(M_W/\mu)^4$, where μ is a parameter that roughly represents the energy scale of supersymmetry. Non-observation of superpartners thus far at the LHC suggests that μ is probably large, on the order of TeV, implying $(M_W/\mu)^4 < 10^{-4}$. Such a tiny mass splitting leads to a remarkably long $\tilde{\chi}_1^\pm$ lifetime, shown in Figure 3.17.

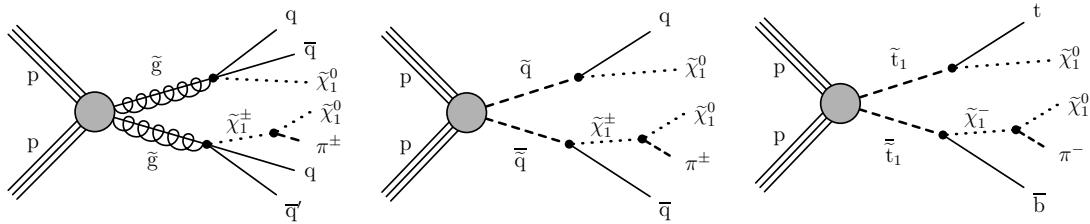


Figure 3.18: Diagrams for (left) gluino, (center) light-flavor squark, and (right) top squark pair production, in which the gluinos and squarks can decay via a long-lived $\tilde{\chi}^\pm$. In this analysis, the gluino is taken to decay with branching fraction $\frac{1}{3}$ each to the $\tilde{\chi}_1^0$, $\tilde{\chi}_1^-$, and $\tilde{\chi}_1^+$. Squarks decay with branching fraction $\frac{1}{2}$ each to the $\tilde{\chi}_1^0$ and the $\tilde{\chi}_1^\pm$ allowed by charge conservation. The $\tilde{\chi}^\pm$ mass is greater than the $\tilde{\chi}_1^0$ mass by hundreds of MeV, so that the charged product of the $\tilde{\chi}^\pm$ decay is too soft to be detected. Taken from [30].

Motivated by these models, the disappearing track search considers signals that include an intermediate $\tilde{\chi}^\pm$ in the superpartner decay chain, similar to the models in Figure 3.1 (upper right, lower left, and lower center) considered by the inclusive search. The $\tilde{\chi}^\pm$ is taken to be long-lived, only hundreds of MeV more massive than $\tilde{\chi}_1^0$. Specifically, the rest frame $\tilde{\chi}^\pm$ lifetime, τ_0 , is varied from $c\tau_0 = 1$ cm to $c\tau_0 = 2000$ cm. This range of lifetimes is selected because the CMS tracker extends from $r \approx 10$ cm to $r \approx 100$ cm, and a significant fraction of $\tilde{\chi}^\pm$ decays must occur inside the tracker for the disappearing track search to have meaningful sensitivity. If the decay occurs outside the tracker, the track does not disappear. The full set of models considered is shown in Figure 3.18.

In the model at left, gluinos are pair-produced and decay with equal probability either to $\tilde{\chi}_1^+, \tilde{\chi}_1^-,$ or $\tilde{\chi}_1^0$, and light-flavor quarks.

In the central model, light flavor squarks are pair-produced and decay with equal probability either to the allowed charge of $\tilde{\chi}^\pm$, or to $\tilde{\chi}_1^0$, and a light-flavor quark.

In the model at right, top squarks are pair-produced and decay with equal probability either to the allowed charge of $\tilde{\chi}^\pm$, or to $\tilde{\chi}_1^0$, and a bottom quark.

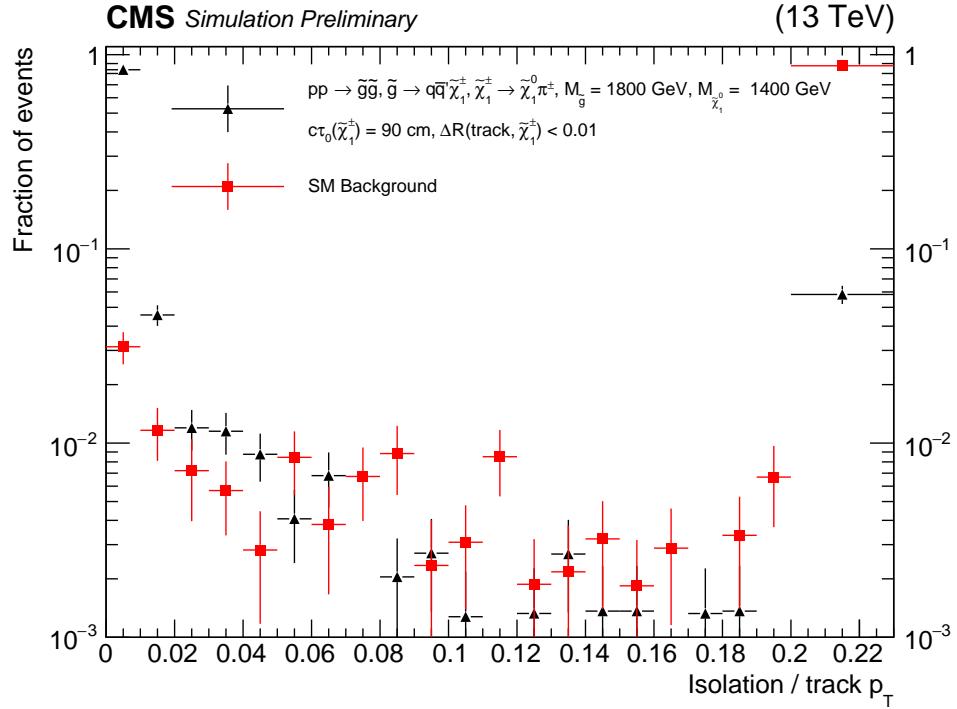


Figure 3.19: Disappearing tracks produced by $\tilde{\chi}^\pm$ decays (black triangles) tend to be much more isolated than disappearing tracks in background events (red squares). All selections listed in Table 3.6 are applied except the isolation selections. The last bin is an overflow bin, containing every track failing the relative isolation selection. The events containing those tracks will either populate control regions, or be rejected entirely, as described in Section 3.2.5. Additional material from [30].

The disappearing tracks produced by $\tilde{\chi}^\pm$ decays tend to be very high quality and highly isolated (shown in comparison with background in Figure 3.19), with negligible energy deposits in any of the calorimeter cells along the track's extrapolated path. The quality metrics in this context include the track's impact parameter with respect to the primary vertex, the size of the uncertainty in the track's assigned p_T , and the number of tracker layers in which the track failed to

leave a hit when one would be expected. Alongside isolation, shown in Figure 3.19, these quality metrics allow for signal-like disappearing tracks to be distinguished from disappearing tracks that are very likely to be due to a reconstruction error. This forms the basis of the background estimate described in Section 3.2.5.

3.2.4 Backgrounds

In contrast to the classic search, there are no irreducible disappearing track backgrounds produced by the Standard Model. However, there are a number of detector reconstruction issues that can lead to apparent disappearing tracks, many of which can be greatly or even entirely rejected by various cleaning selections.

Of the near-entirely reducible backgrounds, the largest before cleaning is that produced by electrons undergoing unusually strong bremsstrahlung inside the tracker, effectively converting the electron into a photon. While electrons are charged and therefore produce tracks, photons are neutral and do not. So, these tracks disappear, but can normally be rejected due to large photon energy deposits in the electromagnetic calorimeter cells to which the track extrapolates. However, some ECAL cells do not perform well for a variety of reasons and may fail to register the photon energy deposits, and so fail to veto the disappearing track. Fortunately, the majority of the ECAL performs well enough to make this background significant only for tracks extrapolating to defective cells.

Defective cells are mapped and vetoed using a tag-and-probe strategy, in which a PF electron “tag” identifies potential $Z \rightarrow e^+e^-$ events, and a disappearing track of opposite charge serves as the “probe.” If an event contains exactly one PF electron, the event is searched for a disappearing track of opposite charge, under the hypothesis that the event was in fact $Z \rightarrow e^+e^-$ and the second electron was not successfully reconstructed due to the conversion process described above. If such a track is found and the invariant mass of the electron-disappearing track pair is consistent with the mass of the Z boson, the region of the ECAL to which the disappearing track

extrapolated is deemed faulty. $Z \rightarrow e^+e^-$ events are sufficiently common that the map of faulty ECAL cells can be assembled as a function of time. As expected, the performance of the ECAL steadily degrades over time, so that the fraction of the ECAL that must be vetoed increases from 2016 to 2017 to 2018. This veto essentially entirely rejects the electron conversion background, but at a price. It is the single largest cause of signal inefficiency, as 15-20% of signal disappearing tracks point, by chance, into vetoed regions of the ECAL.

Another, much smaller reducible background is produced by the decay of strange baryons, including the Ξ^\pm , Σ^\pm , and Ω^\pm . These baryons have lifetimes in the proper range to produce tracks 10s of centimeters long, before decaying with a small mass splitting to a neutral baryon and a charged daughter that is sometimes low enough p_T to escape reconstruction. While these tracks disappear from the perspective of the tracker and the ECAL, the neutral baryon is still detectable in the hadronic calorimeter via its nuclear interactions, which allows for efficient rejection of this background. Additionally, strange baryons are almost always found inside jets, so the vast majority of their tracks can be vetoed on the basis of isolation. Even so, a region of the HCAL with known performance issues in 2018 showed significantly increased disappearing track counts that required it to be vetoed, likely in part because of this background.

The final two backgrounds are fake tracks and mesons, mostly pions, that undergo nuclear interactions in the tracker, neither of which can be entirely rejected.

Fake tracks are tracks comprised of hits that are not all associated to the same genuine particle. The probability that track reconstruction produces a fake track is small, and can be suppressed by requiring that tracks pass a sophisticated selection called high purity, as shown in Figure 2.5 for the 2016 tracker [33], which improved after the pixel tracker upgrade between 2016 and 2017 [72] as shown in Figure 2.8 (right). Furthermore, the probability that more than 3 or 4 unassociated hits will be strung together into a track is negligible, limiting the fake track background to very short tracks. It is also rare for unassociated hits to lie on a tight, relatively straight line; instead, they tend to be somewhat scattered. Therefore, fake tracks also tend to

assigned low p_T , since low energy track have greater curvature in the detector’s magnetic field. The disappearing tracks analysis considers only tracks with $p_T > 15$ GeV in part to avoid the majority of the fake track background, and divides tracks into low and high p_T bins partly in order to maintain a relatively fake-free high p_T region.

While the fake track background is limited to shorter tracks, the meson background is not and so constitutes the dominant longer disappearing track background. This background is generated when a meson undergoes a nuclear interaction inside the tracker, and showers in such a way that no single charged daughter is high enough energy to be reconstructed as a track, and any neutral products are sufficiently low energy or widely scattered that calorimeter deposits are insufficient to reject the track. While it is merely unusual for a meson to undergo a nuclear interaction in the tracker as shown in Figure 2.11, a shower leaving no detectable products is extraordinarily rare, occurring only a handful of times across three years of data taking. Mesons tend not to be isolated, allowing most of this background to be distinguished from signal-like tracks, but with one major exception. The τ lepton, almost always isolated when produced in the hard interaction, undergoes the decay $\tau^\pm \rightarrow h^\pm v$ for some meson h^\pm , almost always a π^\pm , with probability approximately 11.5% [73]. In this decay, the π^\pm effectively inherits the isolation of the τ . If this π^\pm undergoes the rare shower described above, it produces an isolated disappearing track of potentially significant length, scarcely distinguishable from a signal track. For this reason, observation of even a relatively long, isolated, maximally signal-like disappearing track is not itself sufficient to discover a signal. An estimate of the background’s rate is necessary, to allow a comparison to the observed frequency of disappearing tracks.

3.2.5 The Short Track Selection and Data-Driven Background Estimate

The two surviving disappearing track backgrounds, fake tracks and lost pions, are both produced by extremely rare failure modes of the detector. A data-driven background estimate is mandatory, as replicating such pathological edge cases is beyond the capabilities of any feasible

simulation. This requires definition of signal-depleted control regions that can be used to measure the rate at which known physics processes produce disappearing tracks. As discussed with respect to the classic search’s mismeasurement background in Sections 3.1.3 & 3.1.5, events with low M_{T2} are extremely background dominated. Low M_{T2} events compose one such control region. This dependence on M_{T2} to define a signal-depleted control region restricts the disappearing track search to multijet events, as no comparably powerful discriminant is known for monojet events. In high M_{T2} events, properties of the disappearing tracks themselves are used to divide events with signal-like disappearing tracks, called Short Tracks (ST), and background-like disappearing tracks, called Short Track Candidates (STCs), into ST signal regions and STC control regions. The ST and STC selections are summarized in Table 3.6. Signal disappearing tracks pass this selection in simulation with efficiency between 50% and 65%, while only around 1 in 1000 to 1 in 10,000 background events passing the baseline selection of the M_{T2} analysis possess a track passing the ST selection. The STC selection is not maximally background-like; rather, it selects tracks that nearly pass the ST selection, but are not STs. Instead, these tracks pass only a relaxed version of the ST selection. This choice ensures that STCs are as closely linked to STs produced by background as possible, without including a significant amount of true signal tracks. Only a few per cent of signal tracks are selected as STCs, usually due to an unlucky overlap with a jet causing a high isolation value as shown in Figure 3.19, but STCs are a few times more common in background events than STs.

The selections relaxed for STCs with respect to the ST definition can be grouped into two categories, those concerned with track quality, which are loosened by a factor of 3, and with isolation, which are loosened by a factor of 6. The specific track quality selections affected are the impact parameter, both along the beam axis and in the transverse plane, and the p_T error $\sigma(p_T)$. All isolation selections are loosened.

The efficiency for signal tracks to pass this selection is high, 50–65%, as shown in Figure 3.20.

Table 3.6: Selection requirements for STs and STCs. Some requirements differ for tracks of different lengths; the length categories are described in Section 3.2.5. For the subset of medium (M) length tracks that have just four tracking layers with a measurement, the minimum required number of layers of the pixel tracking detector with a measurement is three (\dagger). The selected tracks are required not to overlap with any identified electrons or muons, of any quality. The selected tracks are as well required to not be identified as PF candidates, and not to overlap with other tracks with $p_T > 15 \text{ GeV}$, even if those tracks are not associated to PF candidates. The table also reports the factor by which a selection is relaxed to define STCs, when applicable. If no factor is reported, the requirement is not relaxed for the selection of short track candidates.

Observable	Selection	Track length	STC factor
p_T [GeV]	> 15	All	
$ \eta $	< 2.4 and not $1.38 < \eta < 1.6$	All	
$\sigma(p_T) / p_T^2$ [GeV^{-1}]	$< 0.2; < 0.02; < 0.005$	P; M; L	$\times 3$
d_{xy} (from primary vertex) [cm]	< 0.02 (< 0.01)	P (M, L)	$\times 3$
d_z (from primary vertex) [cm]	< 0.05	All	$\times 3$
Neutral isolation ($\Delta R < 0.05$) [GeV]	< 10	All	$\times 6$
Neutral isolation / p_T	< 0.1	All	$\times 6$
Isolation ($\Delta R < 0.3$) [GeV]	< 10	All	$\times 6$
Isolation / p_T	< 0.2	All	$\times 6$
Number of pixel layers	≥ 3 (≥ 2)	P, M † (M, L)	
Number of tracker layers	$\geq 3; < 7; \geq 7$	P; M; L	
Number of lost inner hits	$= 0$	All	
Number of lost outer hits	≥ 2	M, L	
Is a PF candidate?	No	All	
PF lepton veto ($\Delta R < 0.1$)	Yes	All	
Lepton veto ($\Delta R < 0.2$)	Yes	All	
Track veto ($\Delta R < 0.1$)	Yes	All	
Bad calorimeter module veto	Yes	All	
M_T (track, \vec{E}_T) [GeV]	> 100 , if $p_T < 150 \text{ GeV}$	L	

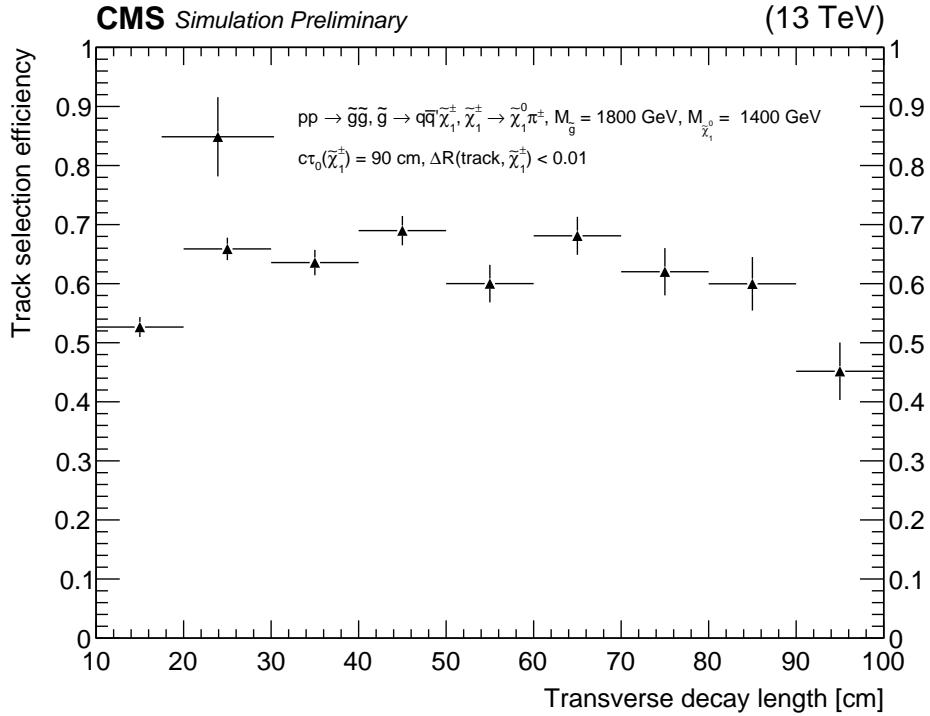


Figure 3.20: The signal efficiency with respect to the full short track selection after selecting events passing the baseline kinematic signal region selection. Additional material from [30].

Binning

The background for shorter disappearing tracks is dominated by fakes, while the background for longer disappearing tracks is dominated by lost pions. Additionally, signals with different LLP lifetimes produce tracks of different lengths. Thus, binning the signal region by track length allows backgrounds of different origins to be segregated in different bins, and maximizes sensitivity to a variety of signal lifetimes. The disappearing track search adopts 3 categories of tracks length.

- Pixel-only (P) tracks have hits only inside the pixel tracker. Fakes dominate the background.
In 2017 and 2018 data, these tracks are further subdivided into pixel tracks with hits in 3 distinct layers (P3) and pixel tracks with hits in 4 distinct layers (P4).
- Medium length (M) tracks have hits in fewer than 7 distinct layers, but have at least one hit

outside the pixel detector. Lost mesons dominate the background.

- Relatively long (L) tracks have hits in 7 or more distinct layers, but still disappear, having missing expected hits in at least the two outermost layers of the tracker. Lost mesons dominate the background, and the better opportunity to measure the track allows for an especially tight selection, allowing for extremely strong background suppression.

Additionally, P and M tracks are divided into tracks with $15 < p_T < 50$ GeV (“lo”) and tracks with $p_T \geq 50$ GeV (“hi”). Most background falls into the low p_T bins, and signals generically populate the high p_T bins, making the high p_T bins the drivers of sensitivity to most signals. This division is not adopted for L tracks, as statistics are too low to allow further division of that population of tracks.

Finally, events are binned based on the global event kinematics, using N_{jet} and H_T , again partly to segregate backgrounds produced by different signals and partly to enhance sensitivity to a variety of signal models. The N_{jet} bins are $2 < N_{\text{jet}} \leq 3$ (denoted by a trailing L, targeted at squark production), and $N_{\text{jet}} > 4$ (H, targeted at gluino production). The H_T bins are $250 < H_T < 575$ GeV (denoted by a leading “L”), $575 \leq H_T < 1200$ GeV (“M”), and $H_T \geq 1200$ GeV (“H”). For events with L length tracks, the L and M H_T bins are merged into “LM” to preserve statistics.

For example, the “L HLM” region is populated by events with L length tracks, with $N_{\text{jet}} \geq 4$ and $H_T \geq 1200$ GeV. The “M LM lo” region is populated by events with M length tracks of $15 < p_T < 50$ GeV, with $2N_{\text{jet}} \leq 3$ and $250 < H_T < 575$ GeV.

The STC control region binning mirrors the ST signal region, with each STC bin mapping to a corresponding ST bin.

Short Track Candidates and f_{short}

The background estimation procedure makes use of the ratio of the counts of STs and STCs, called f_{short} ,

$$f_{\text{short}} = \frac{N_{\text{ST}}}{N_{\text{STC}}}. \quad (3.5)$$

This ratio is measured in events with $60 < M_{\text{T2}} < 100$ GeV, as these events are so background-dominated that even STs are assuredly background-dominated. A separate f_{short} is measured for every bin, except that the measurement is inclusive in H_{T} , exploiting an empirical invariance of f_{short} with respect to H_{T} to improve the statistical precision of the measurement. The ratio is then applied at high M_{T2} to estimate the number of signal-like STs produced by background, N_{ST}^{B} , using the number of STC events observed in the corresponding control region, $N_{\text{STC}}^{\text{Obs}}$,

$$N_{\text{ST}}^{\text{B}} = f_{\text{short}} N_{\text{STC}}^{\text{Obs}}. \quad (3.6)$$

This estimate is correct up to statistical fluctuations, subject to the assumption that f_{short} is invariant with respect to M_{T2} , an assumption which is validated in data in the next subsection. Even without this explicit validation, there is no reason to expect that a ratio based on track-level observables, namely the impact parameter, p_{T} error, and isolation, ought to be sensitive to a global event kinematic variable like M_{T2} . By tethering the background estimate to the observed count of STCs in data, the analysis is able to remain agnostic of the details of background disappearing track production.

Validation

The f_{short} -based estimate of background ST production relies on the invariance of f_{short} with respect to M_{T2} . The f_{short} ratio is measured at $60 < M_{\text{T2}} < 100$ GeV, a background-dominated region, and applied in the $M_{\text{T2}} > 200$ GeV signal region. In Figure 3.21, the invariance of f_{short} with respect to M_{T2} is validated in the intermediate region, $100 < M_{\text{T2}} < 200$ GeV, in both 2016

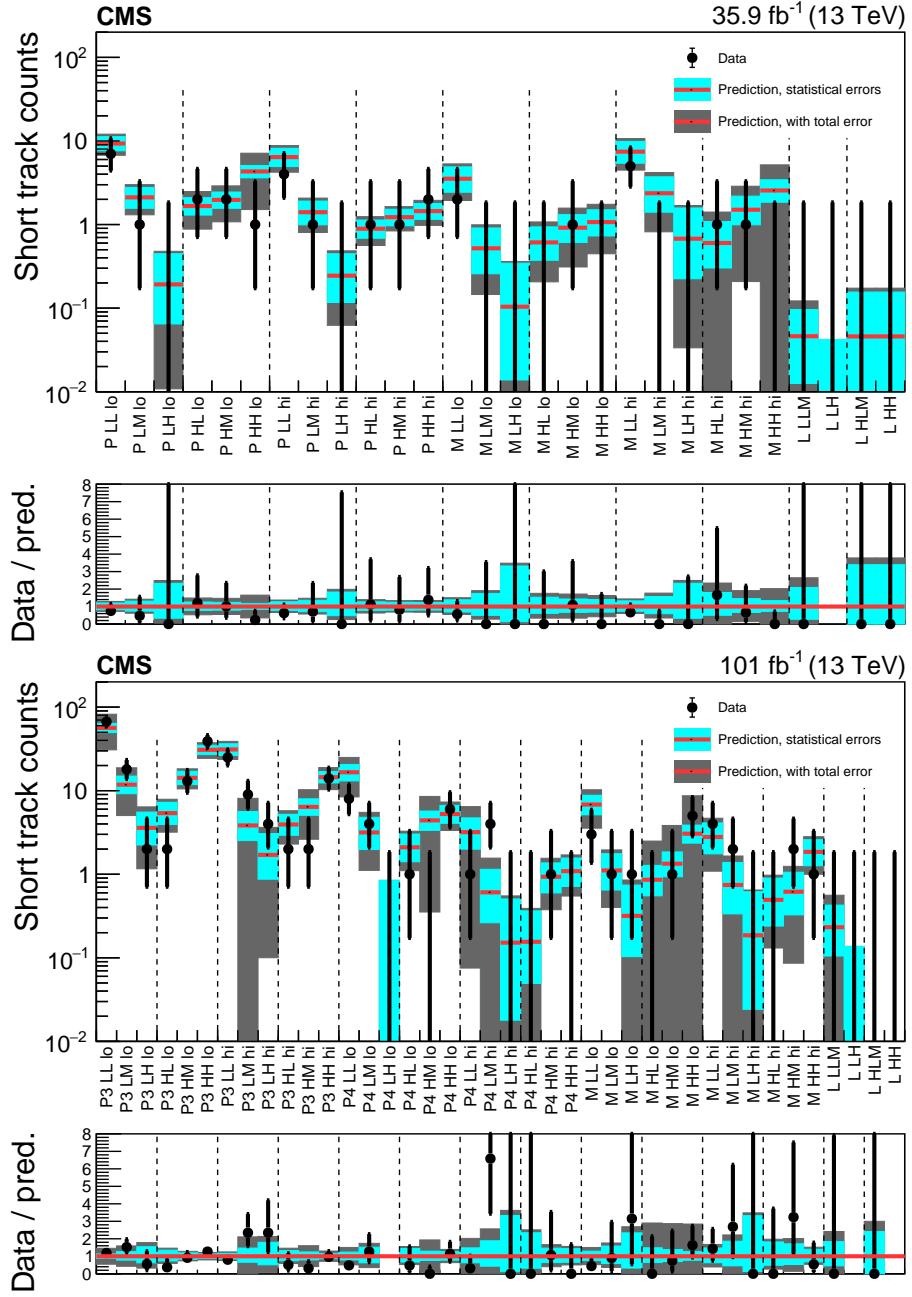


Figure 3.21: Validation of the disappearing tracks background estimate in (upper) 2016 and (lower) 2017–2018 data. There are no discrepancies inconsistent with statistical fluctuations. To be conservative, a systematic uncertainty (gray) is assessed region-by-region to account for any discrepancy greater than 1 statistical standard deviation (blue). Dotted vertical lines group bins that use the same measured value of f_{short} , because they differ only in H_T . Taken from [30].

and 2017–2018 data, by carrying out the background estimation procedure as used for the signal region in this M_{T2} sideband. There are no discrepancies inconsistent with statistical fluctuations,

so the estimate is considered validated. To be conservative, a systematic uncertainty is assessed bin-by-bin that is no less than the statistical uncertainty, to prevent any cases of accidental closure, and as large as is necessary to cover any observed discrepancy larger than 1 standard deviation of statistical uncertainty.

3.2.6 Systematics

Table 3.7: Summary of systematic uncertainties in the disappearing track background prediction, together with their typical size ranges across the search bins. The systematic uncertainties arising from the assumption of kinematic invariance of f_{short} and from the validation of the background prediction are always taken to be at least as large as the statistical uncertainties on the measured values of f_{short} and on the background prediction in the validation region, respectively. Taken from [30].

Source	Range [%]
Limited size of data control samples (high M_{T2} STC counts)	1–100
Limited size of data f_{short} measurement samples (low M_{T2} disappearing track counts)	5–45
Kinematic invariance of f_{short}	10–80
Validation of background prediction	25–75

The uncertainties uniquely affecting the disappearing tracks the background prediction, in addition to those obtained from the validation region, are summarized in Table 3.7 together with their typical size ranges across the search bins.

The first two uncertainties are purely statistical and dominate the total uncertainty.

The kinematic invariance of f_{short} is checked by varying the H_T and \cancel{E}_T definitions of the f_{short} measurement region. Any uncertainty greater than 1 statistical standard deviation is covered by applying a systematic, just as for the validation region systematic. Although the relative uncertainties thus assessed can be large, this is an artifact of extremely low statistics. Discrepancies easily consistent within 2 statistical standard deviations can nonetheless lead to assessment of a large relative systematic uncertainty. Still, these uncertainties are applied to ensure that the disappearing tracks search is conservative.

As for the classic search, the disappearing tracks signal simulation is subject to extra

uncertainties, all of which are shared except the uncertainties assessed to cover the reconstruction of signal disappearing tracks, produced by decaying $\tilde{\chi}^\pm$. There are three such uncertainties.

First, fast simulation is used for signals simulated in 2017 and 2018 conditions. The selection efficiency for signal tracks in fast simulation differs slightly from full simulation, and so a systematic is assessed to cover the difference.

A 10% uncertainty is assessed to cover any potential mismodeling of how efficiently the detector would reconstruct tracks from a long-lived $\tilde{\chi}^\pm$. This is half the difference between unity and the observed efficiency in simulation, for $\tilde{\chi}^\pm$ decays in the regions of the detector that ought to be able to reconstruct the tracks successfully.

Finally, an uncertainty of 1.5% is assessed to cover observed differences between full simulation and fast simulation in how often a $\tilde{\chi}^\pm$ that reaches the muon chamber is falsely reconstructed as a muon, triggering the lepton veto.

3.2.7 Results

Figure 3.22 shows the observed counts compared with the background estimate for both 2016 and 2017–2018 data, and corresponding tables are available in Appendix ???. It is worth noting how small the expected background is, especially in the high p_T M and L bins, in which only a few events are observed across the entire 13 TeV dataset recorded from 2016 to 2018, at an event rate on the order of 1 GHz. The background-only hypothesis is consistent with the data, and so these results are interpreted as exclusion limits on the signal models shown in Figure 3.18 with a variety of lifetimes. Since counts in the M and L bins are so small, signals populating the M and L bins need only to have a few expected events to be excluded at 95% CL by these results.

Signal Contamination

As with the contamination of lepton control regions in the classic search (Section 3.1.6), the disappearing tracks search must adjust its expected signal counts for contamination of its con-

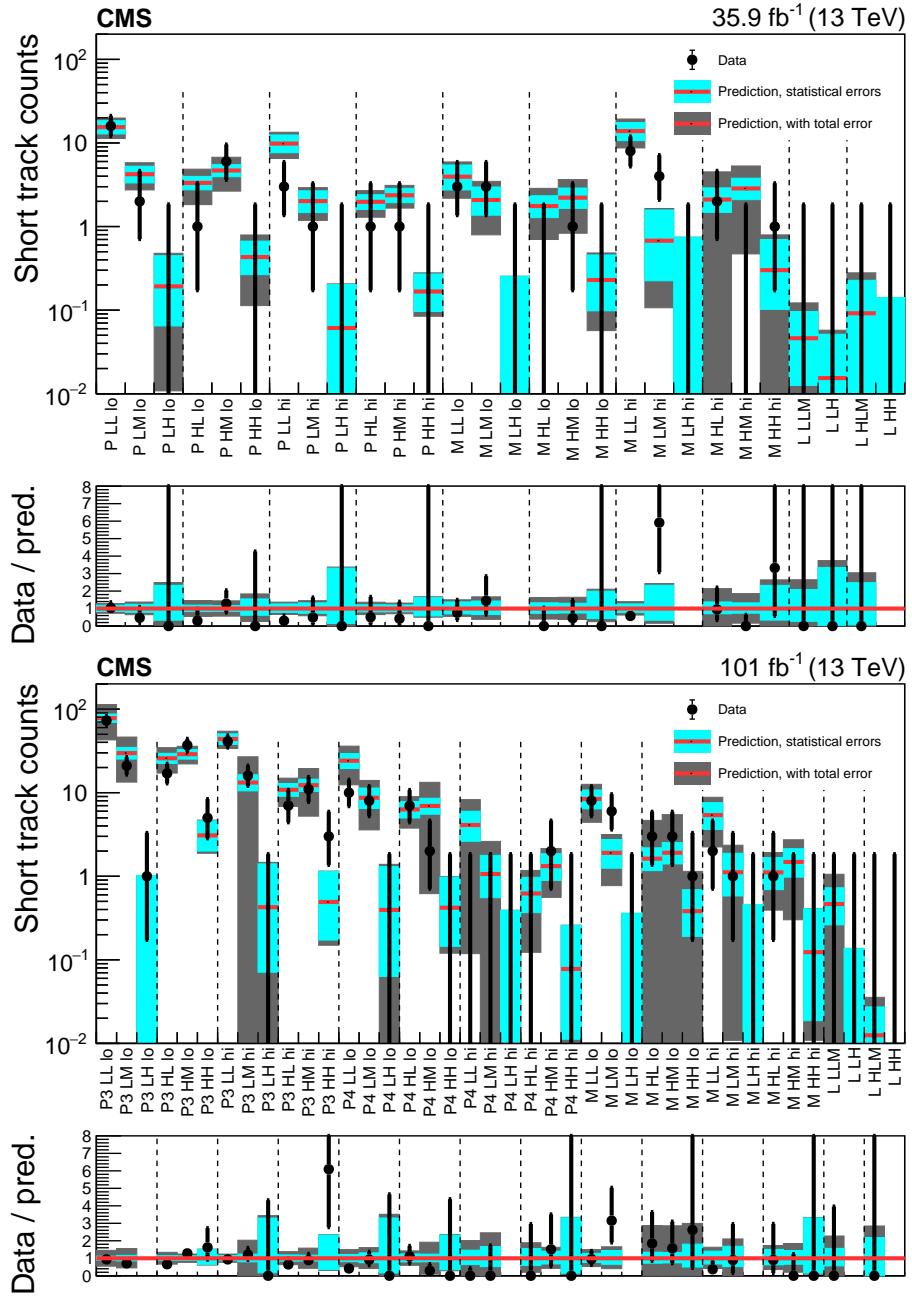


Figure 3.22: Taken from [30]

trol regions. In this case, there is some additional complexity because the signal can contaminate the background estimate by multiple routes. Any signal tracks in the f_{short} measurement region at low M_{T2} can bias the measurement of f_{short} , and signal STCs in the high M_{T2} STC control regions can bias the application of f_{short} . We can separate these two factors to express the effect

of contamination bias schematically as,

$$N_{\text{ST}}^{\text{Est}} + N_{\text{Signal}}^{\text{Bias}} = (f_{\text{short}} + \Delta f_{\text{short}})(N_{\text{STC}} + \Delta N_{\text{STC}}) = f_{\text{short}}N_{\text{STC}} + f_{\text{short}}\Delta N_{\text{STC}} + \Delta f_{\text{short}}N_{\text{STC}} + \Delta f_{\text{short}}\Delta N_{\text{STC}} \quad (3.7)$$

The bias is not linear in the signal strength, due both to the manifestly nonlinear final term, $\Delta f_{\text{short}}\Delta N_{\text{STC}}$, and to the hidden nonlinearity in Δf_{short} ,

$$f_{\text{short}} + \Delta f_{\text{short}} = \frac{N_{\text{ST}}^{\text{B}} + N_{\text{ST}}^{\text{S}}}{N_{\text{STC}}^{\text{B}} + N_{\text{STC}}^{\text{S}}} \quad (3.8)$$

(Recall that these quantities are measured in the $60 < M_{\text{T2}} < 100$ GeV sideband for the f_{short} measurement.)

If the overprediction of background due to contamination, $N_{\text{Signal}}^{\text{Bias}}$, is nonlinear, then the contamination factor C is a function of the signal strength μ , which means that the adjusted signal count must be recomputed for every trial signal strength during the statistical analysis. Alas, the effect of contamination is indeed nonlinear in μ ,

$$N_{\text{Signal}}^{\text{Adjusted}} = \mu \left(N_{\text{Signal}}^{\text{Raw}} - N_{\text{Signal}}^{\text{Bias}} \right) = \mu \left(1 - \frac{N_{\text{Signal}}^{\text{Bias}}}{N_{\text{Signal}}^{\text{Raw}}} \right) N_{\text{Signal}}^{\text{Raw}} = \mu C(\mu) N_{\text{Signal}}^{\text{Raw}}. \quad (3.9)$$

Recalculating the effect of signal contamination for every signal strength considered for every mass point of every model is utterly unfeasible. Fortunately, $N_{\text{Signal}}^{\text{Bias}}$ is *approximately* linear in μ as we will see, and the linearizing approximation's error is strictly conservative, so that adopting the approximation will never lead to a false discovery or false exclusion.

First, one may approximate Δf_{short} thus,

$$f_{\text{short}} + \Delta f_{\text{short}} = \frac{N_{\text{ST}}^{\text{B}} + N_{\text{ST}}^{\text{S}}}{N_{\text{STC}}^{\text{B}}}. \quad (3.10)$$

The omitted $N_{\text{STC}}^{\text{S}}$ is doubly small. First, signal is rare relative to background at low M_{T2} , and

second, only on the order of 1% of signal tracks fall into the STC category, so that $N_{\text{STC}}^{\text{S}} \ll N_{\text{STC}}^{\text{B}}$ in the f_{short} measurement region. This approximation linearizes Δf_{short} , and strictly makes Δf_{short} larger by reducing the denominator, tending to overestimate the bias and, therefore, produce conservative signal counts.

Next, adjusting the applied f_{short} to $f'_{\text{short}} = f_{\text{short}} - \Delta f_{\text{short}}$, Equation 3.7 becomes

$$N_{\text{ST}}^{\text{Est}} + N_{\text{Signal}}^{\text{Bias}} = f'_{\text{short}}(N_{\text{STC}} + \Delta N_{\text{STC}}) = f'_{\text{short}}N_{\text{STC}} + f'_{\text{short}}\Delta N_{\text{STC}} \quad (3.11)$$

The second component of $f'_{\text{short}}\Delta N_{\text{STC}} = f_{\text{short}}\Delta N_{\text{STC}} - \Delta f_{\text{short}}\Delta N_{\text{STC}}$ is nonlinear in the signal strength, but is again doubly small, for the same reasons as the first neglected term. Neglecting it again produces a strictly larger estimate of the contamination, making the approximation conservative once again.

Ultimately, then, there are two adjustments to the expected signal counts. First, the analysis subtracts the effect of signal contamination of low M_{T2} STs, which bias the f_{short} measurement. Then, the analysis subtracts the effect of signal contamination of high M_{T2} STCs, which bias the application of f_{short} . Both of these effects are small. Any terms that consider the combined effect of low M_{T2} signal and signal STCs are dropped, as signal is rare in both cases, and so doubly rare in the combined case. This procedure outputs an adjustment to expected signal counts that is linear in the signal strength, and conservative, allowing for a straightforward and safe statistical interpretation.

3.2.8 Limits

The statistical procedure for extracting limits in the disappearing tracks search is exactly identical to that used in the classic search, described in Section 3.1.8, with the addition of the LLP lifetime as a parameter of interest. Figures 3.23-3.25 show limits in the mass plane for gluinos decaying to light squarks, bottom squarks, and top squarks, respectively, with the $\tilde{\chi}^{\pm} c\tau$ set to

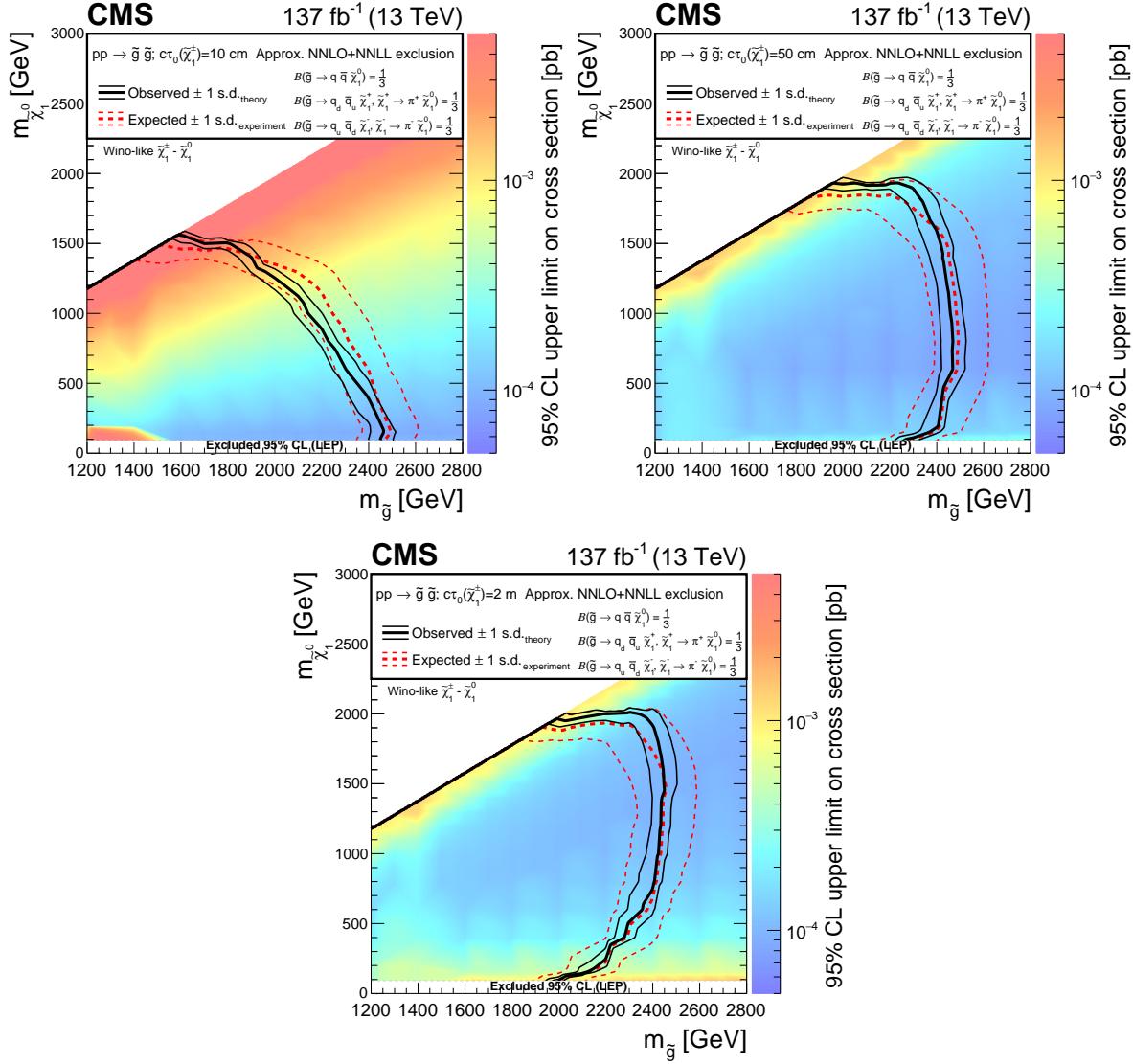


Figure 3.23: Exclusion limits at 95% CLs for direct gluino pair production where the gluinos decay to light-flavor quarks, with $c\tau_0(\tilde{\chi}^\pm) =$ (upper left) 10 cm, (upper right) 50 cm, and (lower) 200 cm. Taken from [30].

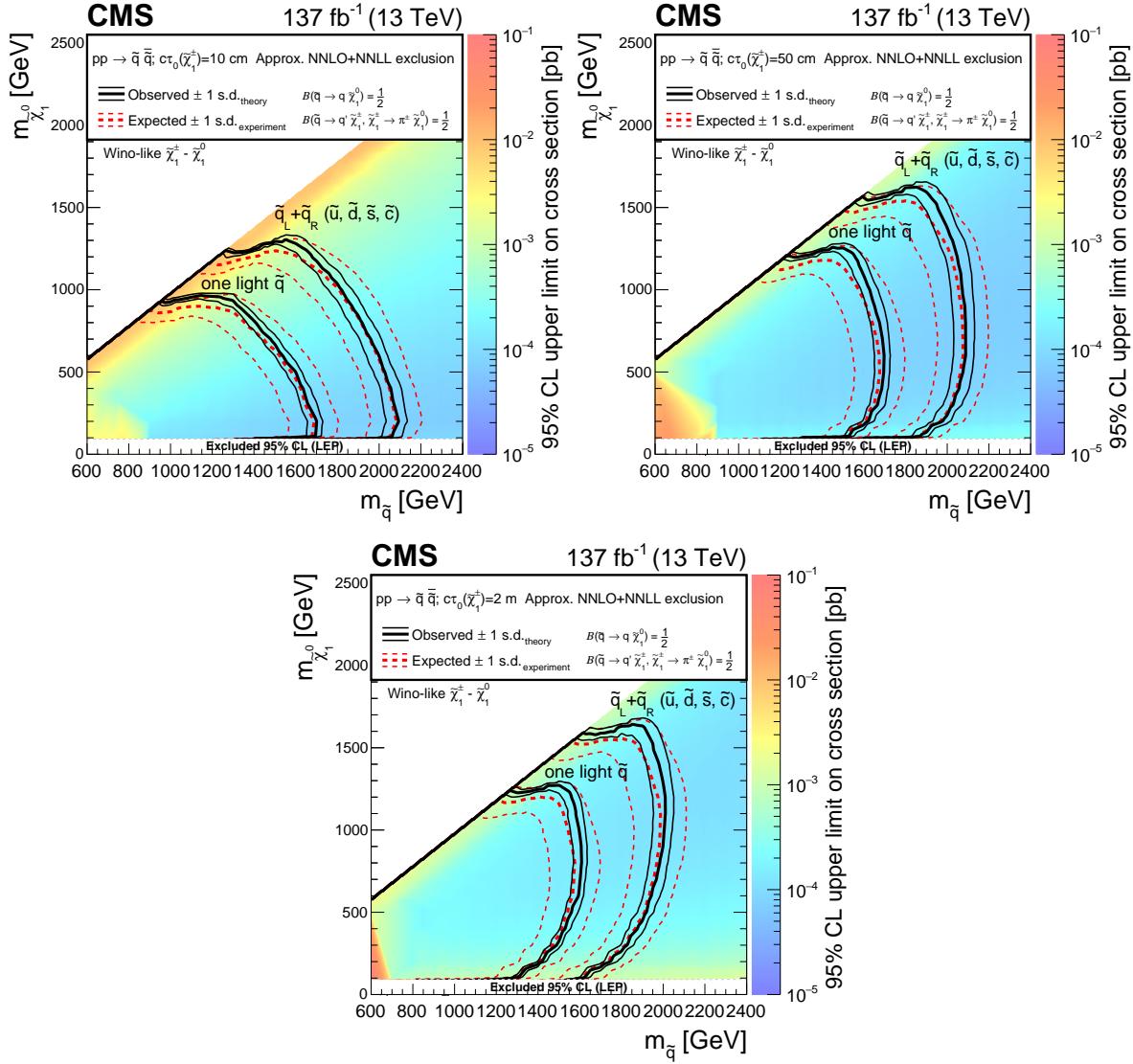


Figure 3.24: Exclusion limits at 95% CL_S for light squark pair production with $c\tau_0(\tilde{\chi}^\pm) =$ (upper left) 10 cm, (upper right) 50 cm, and (lower) 200 cm. Taken from [30].

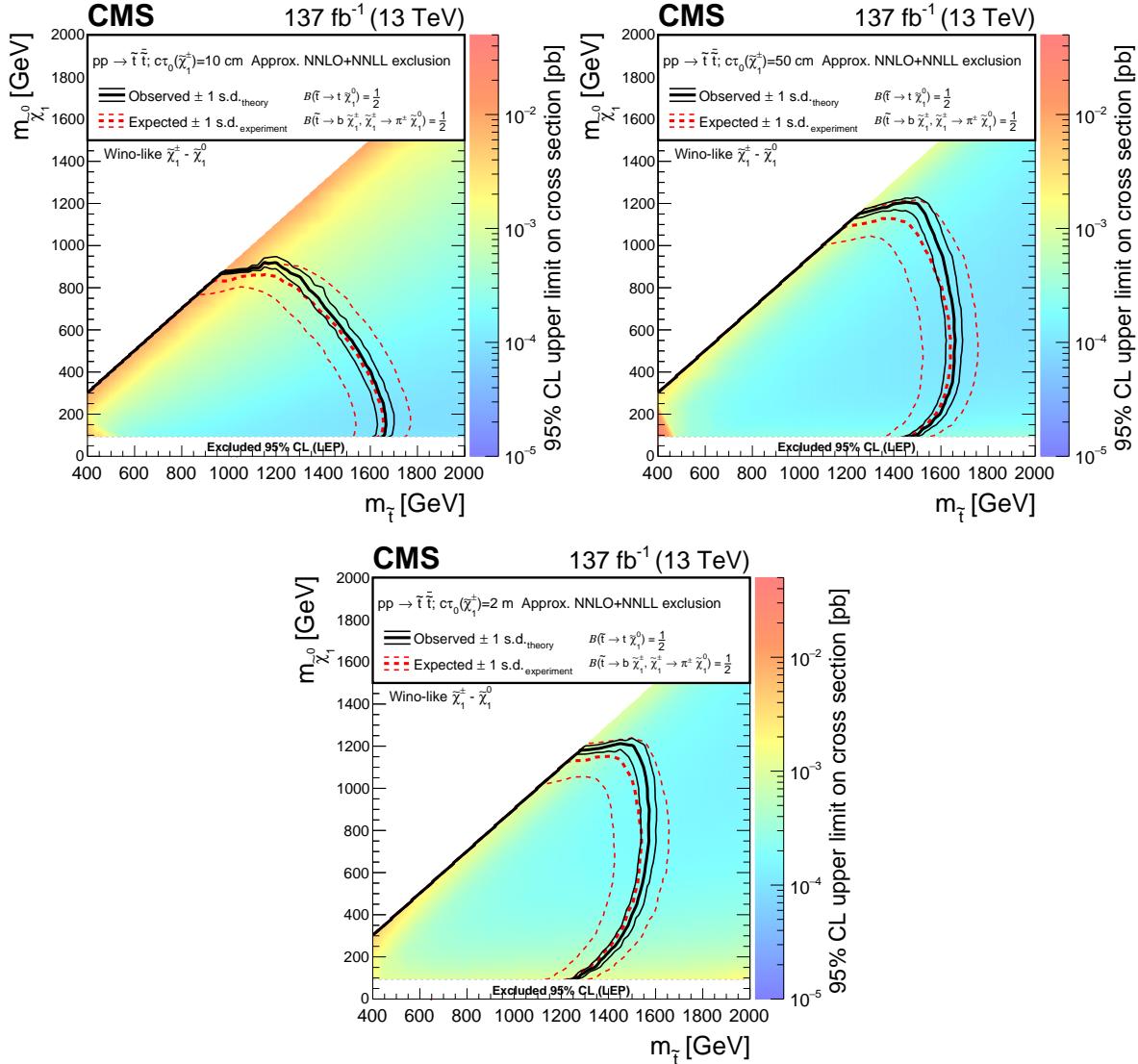


Figure 3.25: Exclusion limits at 95% CL_S for top squark pair production with $c\tau_0(\tilde{\chi}_1^\pm) =$ (upper left) 10cm (upper right) 50cm, and (lower) 200cm. Taken from [30].

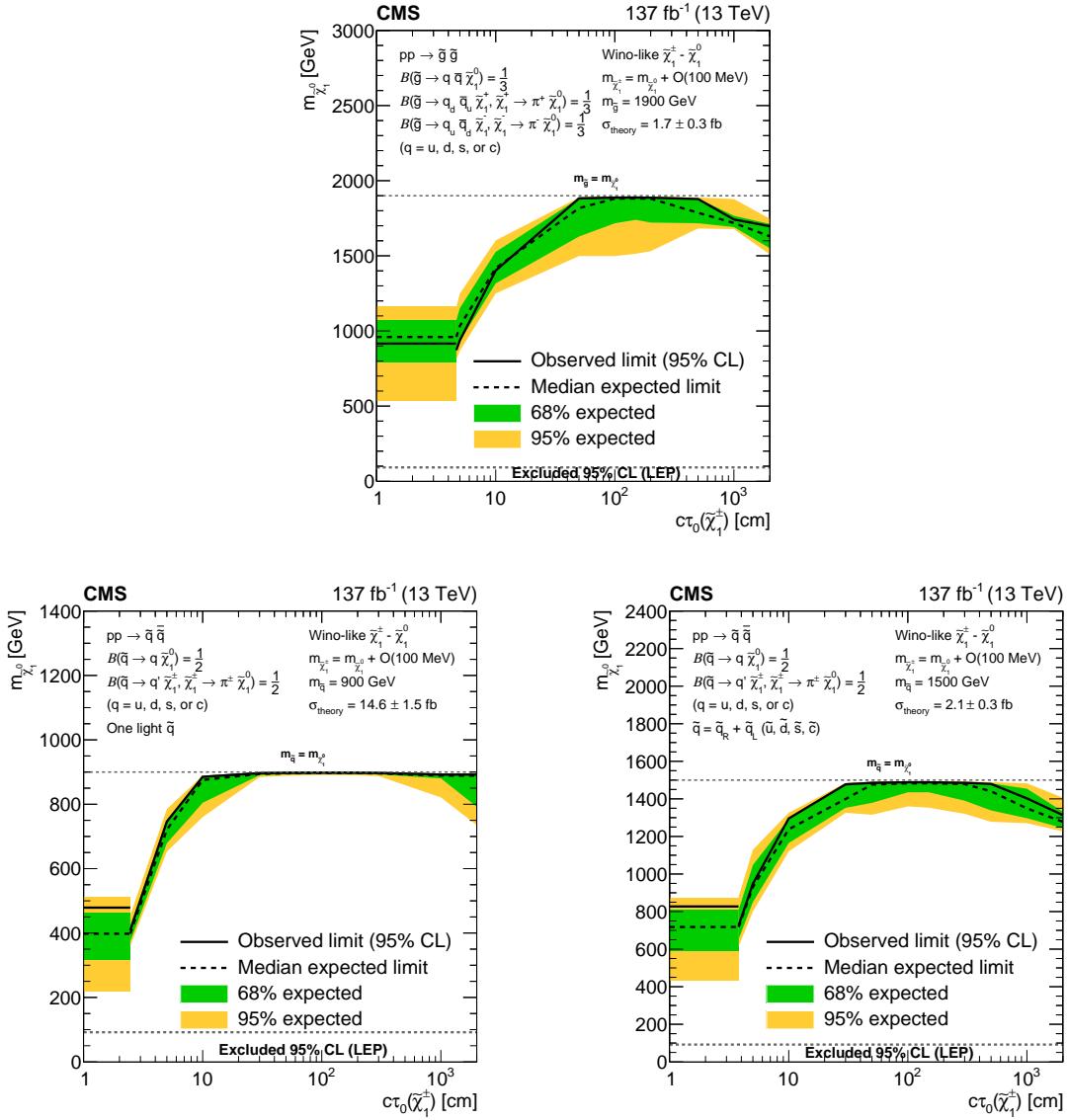


Figure 3.26: Exclusion limits at 95% CL_S on the $\tilde{\chi}_1^0$ mass, with $m_{\tilde{\chi}_1^\pm} = m_{\tilde{\chi}_1^0} + \mathcal{O}(100 \text{ MeV})$, as a function of the $\tilde{\chi}_1^\pm$ proper decay length, for (upper) direct gluino and (lower) direct light-flavor squark pair production, for representative gluino and squark masses. Taken from [30].

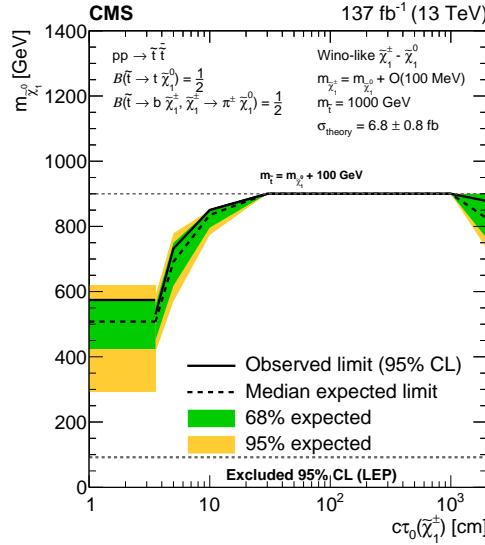


Figure 3.27: Exclusion limits at 95% CL_S on the $\tilde{\chi}_1^0$ mass, with $m_{\tilde{\chi}^\pm} = m_{\tilde{\chi}_1^0} + \mathcal{O}(100 \text{ MeV})$, as a function of the $\tilde{\chi}^\pm$ proper decay length, for direct top squark pair production, as obtained for a representative top squark mass. Taken from [30].

10 cm (upper left), 50 cm (upper right), and 200 cm (bottom). These limits can be compared directly to those obtained by the classic search, shown respectively in Figures 3.11 (upper), 3.13 (upper left), and 3.14 (upper right). In each case, the limits improve by hundreds of GeV along each axis, due entirely to the strong background suppression of the ST requirement.

An interesting feature of these exclusion curves is evident when comparing longer to shorter $\tilde{\chi}^\pm$ lifetimes. For $c\tau = 200$ cm, the limit curve abruptly turns inward at the bottom right of each plot, where the gluino or squark is much more massive than $\tilde{\chi}^\pm$ and $\tilde{\chi}_1^0$. At $c\tau = 50$ cm, this feature is reduced in prominence, and it is inverted when $c\tau = 10$ cm. This is caused by the large Lorentz boost and correspondingly extended lab frame lifetime of $\tilde{\chi}^\pm$ in these mass scenarios. At longer lifetimes, $\tilde{\chi}^\pm$ often lives long enough to escape the tracker, or at least make it near enough to the edge not to leave a disappearing track, defined as a track with at least two tracker layers missing expected hits. With no disappearing track, the event fails the ST selection, so the signal efficiency for long $\tilde{\chi}^\pm$ lifetime is reduced in the large mass splitting scenario, and limits are weakened. In contrast, shorter lifetimes benefit from the boost, as tracks more often extend

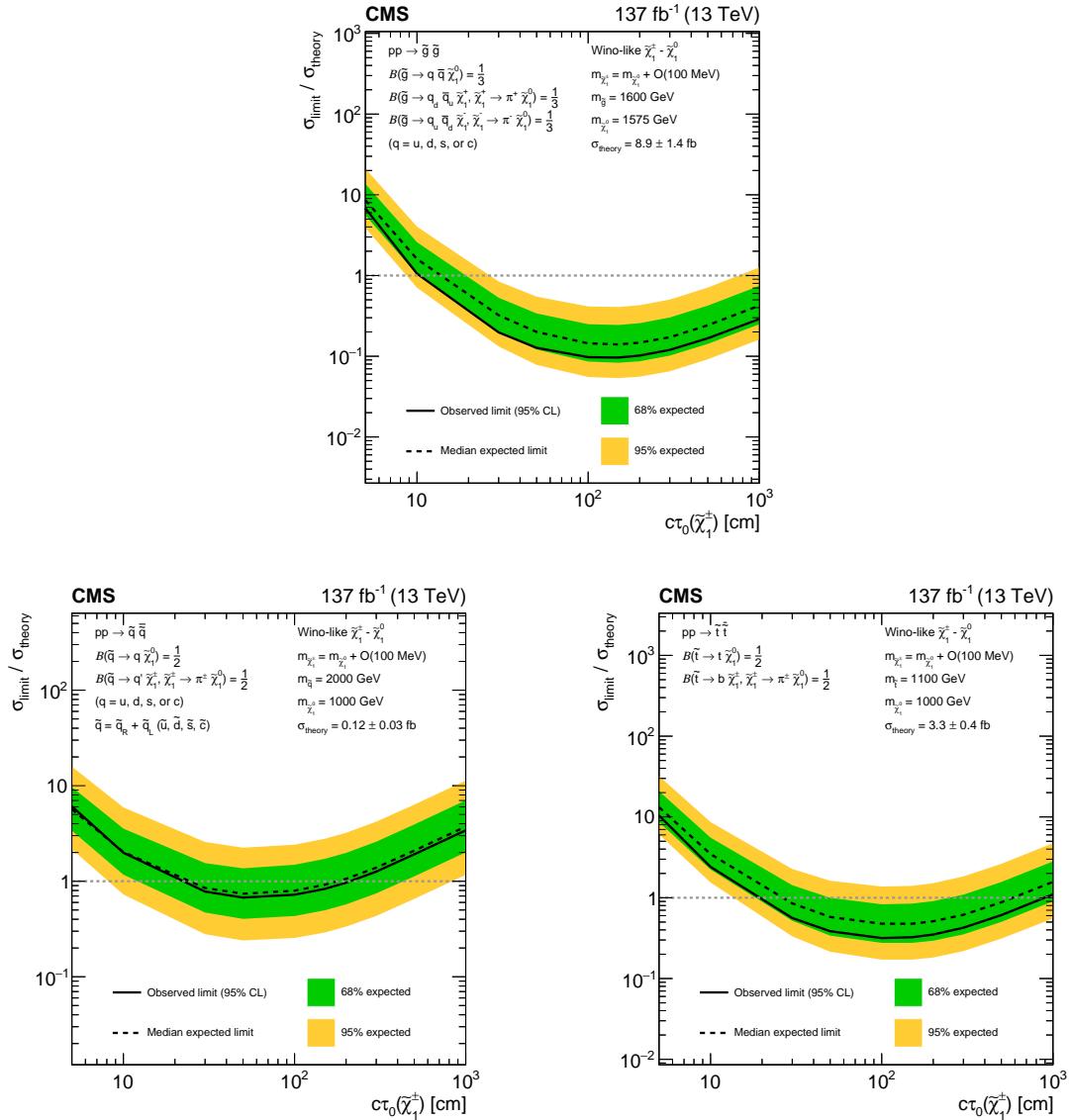


Figure 3.28: Exclusion limits at 95% CLs on $\sigma/\sigma_{\text{theory}}$ as a function of the $\tilde{\chi}^{\pm}$ decay length, for (upper) gluino pair production where the gluinos decay to light-flavor quarks, (lower left) light-flavor squark pair production, and (lower right) top squark pair production, as obtained from the search for disappearing tracks. Taken from [30].

into the more sensitive M and L bins. Masses of $\tilde{\chi}^\pm$ below 91.9 GeV are excluded by searches performed at CERN’s Large Electron Positron collider [51], the previous occupant of the LHC tunnel, mitigating the impact of this effect.

Figures 3.26 and 3.27 show limits at 95% CL on the $\tilde{\chi}_1^0$ mass as a function of the lifetime of $\tilde{\chi}^\pm$, at representative fixed masses of the considered gluino or squark, with the 68% and 95% uncertainty bands shown in green and yellow, respectively. Figure 3.26 (upper) considers gluino pair production at gluino mass 1900 GeV. Figure 3.26 (lower left) considers light squark pair production in the case of a single light squark of mass 900 GeV. Figure 3.26 (lower right) considers light squark pair production in the case of eight degenerate light squarks of mass 1500 GeV. Figure 3.27 considers top squark pair production and at top squark mass 1000 GeV.

There is a discontinuity at low $\tilde{\chi}^\pm$ lifetime in the exclusion curves of all four of these figures. At this discontinuity, the limits obtained for the classic search, applicable to any $\tilde{\chi}^\pm$ of lifetime short enough not to reach the calorimeters, become superior to those obtained by the disappearing track search. At these very short lifetimes, many of the charginos decay before reaching the tracker or soon after, and so do not produce tracks. As the disappearing track search only selects events with a disappearing track, it suffers from large selection inefficiencies at these very short lifetimes. The classic search does not require a track, and so avoids this selection inefficiency, at the cost of a larger background. For short enough $\tilde{\chi}^\pm$ lifetimes, to the left of the discontinuity in each figure, this trade-off becomes advantageous as measured by the expected sensitivity.

When the exclusion curves approach the kinematic limit, at which the mass of the decaying particle and the mass of the $\tilde{\chi}^\pm$ are nearly equal, the errors tend to become greatly compressed.

Finally, Figure 3.28 show the exclusion limits at 95% CL on the production cross section of gluinos and squarks as a function of $\tilde{\chi}^\pm$ lifetime for fixed gluino or squark mass and $\tilde{\chi}_1^0$ mass, for gluinos of mass 1600 GeV decaying to light quarks and $\tilde{\chi}_1^0$ of mass 1575 GeV (upper), light squarks of mass 2000 GeV and $\tilde{\chi}_1^0$ mass 1000 GeV (lower left), and top squarks of mass

1100 GeV and $\tilde{\chi}_1^0$ mass 1000 GeV (lower right). The limits are strongest for intermediate $\tilde{\chi}^\pm$ lifetimes for which the fraction of $\tilde{\chi}^\pm$ decays producing disappearing tracks tens of centimeters long is maximized. Sensitivity is lesser for longer and shorter lifetimes due to reduced efficiency to produce disappearing tracks, and consequently reduced signal selection efficiency.

All of these limits are the strongest produced to date by any analysis, and constitute stringent constraints on the entire class of signal models discussed in Section 3.2.3.

3.2.9 Future

The subject of LLPs in supersymmetric extensions of the Standard Model, and the potential sensitivity gains that LLPs can provide for new physics searches in general, will continue to inspire interest for the foreseeable future of the LHC’s physics program. The M_{T2} disappearing track extension is no exception. Furthermore, unlike the classic search, which has non-negligible backgrounds that will grow larger with increased integrated luminosity, the most sensitive bins of the disappearing tracks search are effectively entirely background-depleted. In fact, while the background should not meaningfully increase in the foreseeable future, the statistical precision of the background estimate *will* improve with the addition of more data. As a result, the sensitivity of the disappearing track search projects to improve roughly linearly with increasing integrated luminosity rather than as its square root. With continuing interest and better scaling of the sensitivity, this search is a potentially worthwhile candidate for continued updates as the CMS dataset expands.

Additionally, the disappearing tracks search in its current form targets only tracks that disappear inside the silicon tracker, and does not consider the case of tracks disappearing inside the muon system, as mentioned in Section 2.2.5. Indeed, it explicitly vetoes tracks potentially associated with hits in the muon system. However, sufficiently long-lived charginos would leave hits in the muon system, and so there exists a possibility to extend the lifetime sensitivity range of the disappearing tracks search by expanding the definition to include disappearing tracks in the

muon system, or even non-disappearing tracks in the muon system that appear to be produced by a particle too heavy to be consistent with a muon.

In any case, the disappearing tracks signature certainly has not been fully exploited, and will remain an interesting search strategy in future LHC datasets.

Chapter 4

Conclusions

The Standard Model of Particle Physics is robust at low energy physics, but is not entirely satisfactory. A variety of observations, both theoretical and experimental, indicate that many of the problems identified with the Standard Model suggest new physics at the TeV scale. Among these new physics models are those containing candidate dark matter particles, including supersymmetric extensions of the Standard Model. Recently, elementary particle physicists have gained access to powerful new tools at CERN’s Large Hadron Collider, including the CMS detector, that can use high energy proton collisions to probe the physics of the weak scale. A pair of searches targeting pair-produced particles decaying to dark matter have been performed at CMS using its 137 fb^{-1} dataset of 13 TeV proton-proton collisions recorded from 2016 to 2018. Both searches produce the most stringent constraints to date on a variety of models, especially supersymmetric models. The constraints of the first search apply generally to any model in which pair-produced colored states decay semi-invisibly. The second targets similar models that also include long-lived particles producing disappearing tracks, a challenging final state to study and one of increasing interest in recent years as null results in searches for supersymmetry have increased the minimum mass scale of supersymmetry. This second model in particular projects to benefit significantly from increased luminosity, and will be able to extend the sensitivity of CMS

to some supersymmetric models by as much as a TeV.

Appendix A

Full Classic Binning and Results

Table A.1: Predictions and observations for the 12 search regions with $N_{\text{jet}} = 1$. For each of the background predictions, the first uncertainty listed is statistical (from the limited size of data control samples and Monte Carlo samples), and the second is systematic. Reprinted from [30].

$N_{\text{jet}} = 1$						
$N_{\text{jet}}, N_{\text{b-tag}}$	$p_T^{\text{jet}1}$ [GeV]	Lost lepton	$Z \rightarrow \nu\nu$	Multijet	Total background	Data
1j, 0b	250-350	$70700 \pm 400 \pm 4100$	$167000 \pm 1000 \pm 11000$	$530 \pm 20 \pm 160$	238000 $\pm 1000 \pm 14000$	251941
	350-450	$13440 \pm 130 \pm 790$	$40100 \pm 500 \pm 3100$	$55 \pm 5 \pm 16$	53600 $\pm 500 \pm 3700$	54870
	450-575	$3050 \pm 50 \pm 180$	$10850^{+230}_{-220} \pm 690$	$5.6 \pm 1.1 \pm 1.6$	13910 $\pm 230 \pm 840$	14473
	575-700	$603^{+20}_{-19} \pm 38$	$2590^{+110}_{-100} \pm 160$	$0.38 \pm 0.06 \pm 0.11$	3200 $\pm 110 \pm 190$	3432
	700-1000	$220 \pm 13 \pm 16$	$1076^{+70}_{-66} \pm 66$	$0.12 \pm 0.03 \pm 0.03$	$1295^{+71}_{-67} \pm 79$	1304
	1000-1200	$11.7^{+4.1}_{-3.2} \pm 0.9$	$86^{+23}_{-19} \pm 6$	< 0.01	$98^{+24}_{-19} \pm 7$	98
	≥ 1200	$2.8^{+2.7}_{-1.5} \pm 0.6$	$23^{+12}_{-8} \pm 2$	< 0.01	$26^{+13}_{-9} \pm 2$	30
1j, $\geq 1b$	250-350	$4210 \pm 110 \pm 260$	$9030 \pm 230 \pm 630$	$58 \pm 10 \pm 17$	$13310^{+260}_{-250} \pm 820$	13549
	350-450	$878 \pm 38 \pm 56$	$2180^{+110}_{-100} \pm 170$	$4.6 \pm 0.4 \pm 1.3$	3060 $\pm 110 \pm 220$	3078
	450-575	$211^{+16}_{-15} \pm 13$	$651^{+57}_{-53} \pm 44$	$0.63 \pm 0.18 \pm 0.18$	$863^{+59}_{-55} \pm 53$	810
	575-700	$40.3^{+6.0}_{-5.5} \pm 2.5$	$164^{+30}_{-26} \pm 11$	$0.04 \pm 0.02 \pm 0.02$	$205^{+31}_{-26} \pm 13$	184
	≥ 700	$19.2^{+5.7}_{-4.6} \pm 1.3$	$74^{+21}_{-16} \pm 7$	< 0.01	$94^{+21}_{-17} \pm 7$	83

Table A.2: Predictions and observations for the 30 search regions with $250 \leq H_T < 450$ GeV. For each of the background predictions, the first uncertainty listed is statistical (from the limited size of data control samples and Monte Carlo samples), and the second is systematic. Reprinted from [30].

250 $\leq H_T < 450$ GeV						
$N_{\text{jet}}, N_{\text{b-tag}}$	$M_{\text{T2}} [\text{GeV}]$	Lost lepton	$Z \rightarrow \nu\nu$	Multijet	Total background	Data
2-3j, 0b	200-300	$73700 \pm 500 \pm 5000$	$156000 \pm 1000 \pm 12000$	$580 \pm 20 \pm 140$	$231000 \pm 1000 \pm 16000$	240867
	300-400	$12030 \pm 200 \pm 820$	$31300 \pm 200 \pm 2500$	$50 \pm 5 \pm 10$	$43400 \pm 300 \pm 3200$	44074
	≥ 400	$417^{+51}_{-47} \pm 28$	$1450 \pm 10 \pm 140$	$0.44 \pm 0.09 \pm 0.09$	$1870 \pm 50 \pm 160$	2022
2-3j, 1b	200-300	$12450 \pm 170 \pm 820$	$18700 \pm 300 \pm 1500$	$90 \pm 8 \pm 21$	$31300 \pm 300 \pm 2200$	32120
	300-400	$2380 \pm 80 \pm 160$	$3750 \pm 60 \pm 310$	$6.9 \pm 1.0 \pm 1.5$	$6130 \pm 100 \pm 430$	6258
	≥ 400	$97 \pm 8 \pm 39$	$174 \pm 3 \pm 17$	$0.01 \pm 0.01 \pm 0.00$	$271^{+9}_{-8} \pm 45$	275
2-3j, 2b	200-300	$2240 \pm 70 \pm 150$	$2340^{+110}_{-100} \pm 200$	$9.7 \pm 1.1 \pm 2.3$	$4600^{+130}_{-120} \pm 320$	4709
	300-400	$398^{+34}_{-32} \pm 27$	$469^{+21}_{-20} \pm 39$	$0.68 \pm 0.17 \pm 0.15$	$868^{+40}_{-38} \pm 61$	984
	≥ 400	$13.3 \pm 2.3 \pm 5.4$	$21.7^{+1.0}_{-0.9} \pm 2.2$	< 0.01	$35.0 \pm 2.5 \pm 6.0$	30
2-6j, $\geq 3b$	200-300	$507^{+32}_{-31} \pm 38$	$179^{+35}_{-30} \pm 27$	$1.77 \pm 0.46 \pm 0.46$	$688^{+47}_{-43} \pm 54$	699
	300-400	$69 \pm 6 \pm 15$	$40.0^{+7.8}_{-6.6} \pm 6.0$	$0.16 \pm 0.12 \pm 0.04$	$109^{+10}_{-9} \pm 16$	102
	≥ 400	$1.50 \pm 0.80 \pm 0.61$	$1.43^{+0.28}_{-0.24} \pm 0.25$	< 0.01	$2.92^{+0.85}_{-0.83} \pm 0.67$	0
4-6j, 0b	200-300	$12500 \pm 180 \pm 800$	$21600 \pm 300 \pm 1800$	$250 \pm 17 \pm 58$	$34400 \pm 400 \pm 2400$	35187
	300-400	$2070 \pm 80 \pm 130$	$4660 \pm 70 \pm 410$	$18.2 \pm 3.6 \pm 3.8$	$6750 \pm 110 \pm 510$	6725
	≥ 400	$42 \pm 5 \pm 17$	$155 \pm 2 \pm 64$	$0.06 \pm 0.03 \pm 0.01$	$197 \pm 5 \pm 67$	170
4-6j, 1b	200-300	$5750 \pm 100 \pm 380$	$4300 \pm 150 \pm 360$	$61 \pm 7 \pm 15$	$10120 \pm 180 \pm 680$	10564
	300-400	$784^{+43}_{-42} \pm 52$	$928^{+32}_{-31} \pm 84$	$2.07 \pm 0.29 \pm 0.45$	$1710 \pm 50 \pm 120$	1769
	≥ 400	$14.0 \pm 2.5 \pm 5.7$	$31 \pm 1 \pm 13$	$0.04 \pm 0.02 \pm 0.01$	$45 \pm 3 \pm 14$	40
4-6j, 2b	200-300	$2550^{+70}_{-60} \pm 170$	$921^{+68}_{-63} \pm 87$	$10.0 \pm 1.5 \pm 2.2$	$3480 \pm 90 \pm 230$	3621
	300-400	$220^{+23}_{-21} \pm 15$	$198^{+15}_{-14} \pm 20$	$0.47 \pm 0.15 \pm 0.11$	$419^{+27}_{-25} \pm 31$	496
	≥ 400	$3.2 \pm 0.8 \pm 1.3$	$6.6 \pm 0.5 \pm 2.7$	< 0.01	$9.8 \pm 0.9 \pm 3.1$	14
$\geq 7j, 0b$	200-300	$55^{+15}_{-13} \pm 4$	$61^{+23}_{-17} \pm 26$	$2.64 \pm 0.39 \pm 0.57$	$119^{+28}_{-22} \pm 27$	108
	300-500	$3.8^{+2.1}_{-2.0} \pm 0.8$	$8.1^{+3.1}_{-2.3} \pm 4.3$	$0.08 \pm 0.04 \pm 0.02$	$12.0^{+3.7}_{-3.1} \pm 4.4$	30
	≥ 500	$0.0^{+3.2}_{-0.0} \pm 0.0$	$0.0^{+1.2}_{-0.0} \pm 0.0$	< 0.01	$0.0^{+3.4}_{-0.0} \pm 0.0$	0
$\geq 7j, 1b$	200-300	$48.0^{+9.1}_{-8.2} \pm 3.5$	$19^{+19}_{-11} \pm 10$	$0.33 \pm 0.14 \pm 0.09$	$68^{+21}_{-13} \pm 11$	95
	≥ 300	$3.0 \pm 1.4 \pm 1.2$	$2.5^{+2.4}_{-1.3} \pm 1.7$	$0.03 \pm 0.02 \pm 0.01$	$5.6^{+2.8}_{-1.9} \pm 2.1$	12
$\geq 7j, 2b$	200-300	$41.3^{+7.7}_{-7.0} \pm 3.1$	$6.0^{+5.8}_{-3.2} \pm 3.7$	$0.29 \pm 0.14 \pm 0.06$	$47.6^{+9.7}_{-7.7} \pm 5.0$	30
	≥ 300	$2.15^{+0.78}_{-0.76} \pm 0.87$	$0.74^{+0.72}_{-0.40} \pm 0.57$	< 0.01	$2.9^{+1.1}_{-0.9} \pm 1.1$	1
$\geq 7j, \geq 3b$	200-300	$7.3^{+1.7}_{-1.5} \pm 0.9$	$1.0^{+1.0}_{-0.6} \pm 1.1$	$0.04 \pm 0.04 \pm 0.01$	$8.4^{+1.9}_{-1.6} \pm 1.5$	17
	≥ 300	$0.47 \pm 0.35 \pm 0.20$	$0.12^{+0.11}_{-0.06} \pm 0.14$	< 0.01	$0.59^{+0.37}_{-0.35} \pm 0.24$	0

Table A.3: Predictions and observations for the 40 search regions with $450 \leq H_T < 575$ GeV. For each of the background predictions, the first uncertainty listed is statistical (from the limited size of data control samples and Monte Carlo samples), and the second is systematic. Reprinted from [30].

$450 \leq H_T < 575$ GeV						
$N_{\text{jet}}, N_{\text{b-tag}}$	M_{T2} [GeV]	Lost lepton	$Z \rightarrow vv$	Multijet	Total background	Data
2-3j, 0b	200-300	$8860 \pm 110 \pm 640$	$20100 \pm 200 \pm 1300$	$69 \pm 13 \pm 16$	29100 $\pm 300 \pm 1900$	28956
	300-400	$4230 \pm 80 \pm 300$	$11770 \pm 140 \pm 790$	$10.6 \pm 0.8 \pm 2.4$	16000 $\pm 200 \pm 1000$	15876
	400-500	$1510 \pm 60 \pm 110$	$5020 \pm 60 \pm 360$	$2.86 \pm 0.62 \pm 0.60$	6540 $\pm 80 \pm 440$	6527
	≥ 500	$121^{+24}_{-21} \pm 9$	$580 \pm 7 \pm 63$	$0.07 \pm 0.03 \pm 0.02$	701^{+25}_{-22} ± 68	740
2-3j, 1b	200-300	$1326 \pm 43 \pm 88$	$2500 \pm 80 \pm 170$	$17.0 \pm 8.4 \pm 3.8$	3840^{+100}_{-90} ± 240	3859
	300-400	$737 \pm 35 \pm 49$	$1464^{+49}_{-48} \pm 99$	$1.62 \pm 0.20 \pm 0.43$	2200 $\pm 60 \pm 140$	2065
	400-500	$259^{+25}_{-23} \pm 19$	$626^{+21}_{-20} \pm 45$	$0.49 \pm 0.10 \pm 0.12$	885^{+32}_{-31} ± 58	907
	≥ 500	$19.1^{+2.8}_{-2.7} \pm 7.8$	$72.4 \pm 2.4 \pm 7.9$	$0.04 \pm 0.02 \pm 0.02$	92 $\pm 4 \pm 11$	79
2-3j, 2b	200-300	$201 \pm 15 \pm 13$	$322^{+31}_{-28} \pm 25$	$1.34 \pm 0.62 \pm 0.47$	524^{+35}_{-32} ± 35	463
	300-400	$83.8^{+9.6}_{-9.1} \pm 9.1$	$188^{+18}_{-17} \pm 15$	$0.26 \pm 0.07 \pm 0.07$	272^{+21}_{-19} ± 20	304
	400-500	$31.8^{+4.1}_{-4.0} \pm 6.7$	$80.4^{+7.7}_{-7.1} \pm 6.6$	$0.02 \pm 0.01 \pm 0.01$	112^{+9}_{-8} ± 10	120
	≥ 500	$2.16^{+0.67}_{-0.66} \pm 0.88$	$9.3^{+0.9}_{-0.8} \pm 1.1$	< 0.01	11.4 $\pm 1.1 \pm 1.4$	15
2-6j, $\geq 3b$	200-300	$232^{+17}_{-16} \pm 15$	$57^{+17}_{-13} \pm 7$	$2.20 \pm 0.70 \pm 0.80$	291^{+24}_{-21} ± 19	297
	300-400	$81^{+12}_{-11} \pm 6$	$33.6^{+9.9}_{-7.8} \pm 4.3$	$0.26 \pm 0.08 \pm 0.08$	115^{+16}_{-14} ± 8	76
	400-500	$10.7^{+2.1}_{-2.0} \pm 2.3$	$11.4^{+3.4}_{-2.7} \pm 1.5$	< 0.01	22.1^{+4.0}_{-3.4} ± 2.8	24
	≥ 500	$1.08 \pm 0.58 \pm 0.44$	$1.03^{+0.30}_{-0.24} \pm 0.17$	< 0.01	2.11^{+0.65}_{-0.62} ± 0.48	0
4-6j, 0b	200-300	$5660 \pm 90 \pm 370$	$8560 \pm 170 \pm 600$	$143 \pm 7 \pm 35$	14360 $\pm 190 \pm 890$	15047
	300-400	$2250 \pm 60 \pm 150$	$4790^{+100}_{-90} \pm 350$	$24.3 \pm 2.6 \pm 6.2$	7060 $\pm 110 \pm 460$	6939
	400-500	$428^{+32}_{-30} \pm 28$	$1220 \pm 20 \pm 110$	$1.42 \pm 0.21 \pm 0.52$	1650 $\pm 40 \pm 130$	1817
	≥ 500	$14.8 \pm 2.2 \pm 6.0$	$86 \pm 2 \pm 35$	$0.04 \pm 0.02 \pm 0.01$	101 $\pm 3 \pm 36$	104
4-6j, 1b	200-300	$2810 \pm 60 \pm 190$	$1880 \pm 80 \pm 130$	$63 \pm 15 \pm 19$	4750 $\pm 100 \pm 300$	4736
	300-400	$937 \pm 36 \pm 63$	$1054^{+45}_{-43} \pm 78$	$5.4 \pm 0.4 \pm 1.4$	2000 $\pm 60 \pm 130$	2039
	400-500	$138^{+17}_{-16} \pm 10$	$269 \pm 11 \pm 25$	$0.36 \pm 0.10 \pm 0.10$	407^{+20}_{-19} ± 31	403
	≥ 500	$7.5 \pm 2.2 \pm 3.0$	$19.1 \pm 0.8 \pm 7.9$	$0.01 \pm 0.01 \pm 0.00$	26.5 $\pm 2.3 \pm 8.5$	27
4-6j, 2b	200-300	$1343^{+38}_{-37} \pm 89$	$414^{+39}_{-35} \pm 33$	$11.5 \pm 1.0 \pm 3.3$	1770 $\pm 50 \pm 110$	1767
	300-400	$418^{+24}_{-23} \pm 29$	$232^{+22}_{-20} \pm 19$	$1.35 \pm 0.35 \pm 0.39$	651^{+32}_{-31} ± 43	636
	400-500	$45.6^{+3.9}_{-3.8} \pm 9.6$	$59.1^{+5.5}_{-5.1} \pm 5.9$	$0.03 \pm 0.02 \pm 0.01$	105^{+7}_{-6} ± 12	120
	≥ 500	$1.59 \pm 0.89 \pm 0.65$	$4.2 \pm 0.4 \pm 1.7$	< 0.01	5.8 $\pm 1.0 \pm 1.9$	7
$\geq 7j, 0b$	200-300	$149^{+17}_{-16} \pm 13$	$169^{+31}_{-27} \pm 34$	$11.5 \pm 0.8 \pm 3.0$	329^{+36}_{-31} ± 38	354
	300-400	$38.9^{+5.8}_{-5.6} \pm 8.2$	$64^{+12}_{-10} \pm 17$	$1.24 \pm 0.42 \pm 0.32$	104^{+13}_{-12} ± 20	110
	≥ 400	$1.28 \pm 0.82 \pm 0.52$	$8.8^{+1.6}_{-1.4} \pm 3.8$	$0.03 \pm 0.02 \pm 0.01$	10.1^{+1.8}_{-1.6} ± 3.8	10
$\geq 7j, 1b$	200-300	$191^{+13}_{-12} \pm 15$	$67^{+19}_{-15} \pm 15$	$4.4 \pm 0.5 \pm 1.2$	262^{+23}_{-19} ± 23	268
	300-400	$37.8^{+3.4}_{-3.3} \pm 8.0$	$25.3^{+7.2}_{-5.7} \pm 7.3$	$0.30 \pm 0.07 \pm 0.08$	63^{+8}_{-7} ± 11	65
	≥ 400	$2.31 \pm 0.69 \pm 0.94$	$3.5^{+1.0}_{-0.8} \pm 1.5$	$0.01 \pm 0.01 \pm 0.00$	5.8^{+1.2}_{-1.0} ± 1.8	3
$\geq 7j, 2b$	200-300	$173^{+12}_{-11} \pm 13$	$19.9^{+5.7}_{-4.5} \pm 5.2$	$1.24 \pm 0.18 \pm 0.33$	194^{+13}_{-12} ± 15	197
	300-400	$26.8 \pm 2.6 \pm 5.7$	$7.6^{+2.2}_{-1.7} \pm 2.4$	$0.09 \pm 0.04 \pm 0.03$	34.6^{+3.4}_{-3.1} ± 6.3	44
	≥ 400	$1.40 \pm 0.44 \pm 0.57$	$1.02^{+0.29}_{-0.23} \pm 0.46$	< 0.01	2.42^{+0.53}_{-0.49} ± 0.73	3
$\geq 7j, \geq 3b$	200-300	$55.4^{+4.8}_{-4.7} \pm 7.3$	$2.3^{+0.7}_{-0.5} \pm 1.1$	$0.15 \pm 0.06 \pm 0.06$	57.8^{+4.8}_{-4.7} ± 7.4	37
	300-400	$6.4 \pm 1.2 \pm 1.5$	$0.86^{+0.25}_{-0.20} \pm 0.46$	$0.01 \pm 0.01 \pm 0.00$	7.3 $\pm 1.2 \pm 1.6$	9
	≥ 400	$0.06 \pm 0.01 \pm 0.03$	$0.12 \pm 0.03 \pm 0.06$	< 0.01	0.18^{+0.04}_{-0.03} ± 0.07	0

Table A.4: Predictions and observations for the 47 search regions with $575 \leq H_T < 1200$ GeV, $N_{\text{jet}} < 7$. For each of the background predictions, the first uncertainty listed is statistical (from the limited size of data control samples and Monte Carlo samples), and the second is systematic. Reprinted from [30].

$575 \leq H_T < 1200$ GeV, $N_{\text{jet}} < 7$						
$N_{\text{jet}}, N_{\text{b-tag}}$	$M_{\text{T2}} [\text{GeV}]$	Lost lepton	$Z \rightarrow vv$	Multijet	Total background	Data
2-3j, 0b	200-300	$5270 \pm 60 \pm 370$	$11550 \pm 160 \pm 790$	$93 \pm 20 \pm 30$	16900 $\pm 200 \pm 1100$	17256
	300-400	$2560 \pm 50 \pm 180$	$7770^{+110}_{-100} \pm 540$	$11.9 \pm 1.3 \pm 4.4$	10340 $^{+120}_{-110} \pm 680$	10145
	400-500	$1101^{+32}_{-31} \pm 77$	$3900 \pm 50 \pm 280$	$1.33 \pm 0.24 \pm 0.41$	5000 $\pm 60 \pm 340$	5021
	500-600	$502^{+24}_{-23} \pm 35$	$2250 \pm 30 \pm 170$	$0.37 \pm 0.07 \pm 0.12$	2760 $\pm 40 \pm 200$	2706
	600-700	$180^{+16}_{-15} \pm 13$	$746 \pm 10 \pm 73$	$0.09 \pm 0.03 \pm 0.03$	926 $^{+19}_{-18} \pm 80$	1066
	700-800	$52.1^{+7.3}_{-6.5} \pm 5.5$	$256 \pm 3 \pm 36$	$0.01 \pm 0.01 \pm 0.00$	308 $^{+8}_{-7} \pm 38$	347
	800-900	$17.7^{+2.6}_{-2.3} \pm 2.2$	$107 \pm 1 \pm 20$	< 0.01	125 $\pm 3 \pm 21$	111
	900-1000	$6.0 \pm 0.9 \pm 1.3$	$39.4 \pm 0.5 \pm 8.5$	$0.01 \pm 0.01 \pm 0.00$	45.4 $^{+1.1}_{-1.0} \pm 8.7$	39
	1000-1100	$3.3^{+1.1}_{-1.0} \pm 1.0$	$13.3 \pm 0.2 \pm 3.9$	< 0.01	16.6 $\pm 1.1 \pm 4.1$	11
	≥ 1100	$0.31^{+0.09}_{-0.08} \pm 0.12$	$2.5 \pm 0.0 \pm 1.1$	< 0.01	2.8 $\pm 0.1 \pm 1.1$	2
2-3j, 1b	200-300	$826^{+27}_{-26} \pm 54$	$1480^{+60}_{-50} \pm 100$	$38 \pm 15 \pm 12$	2340 $\pm 60 \pm 140$	2499
	300-400	$426^{+21}_{-20} \pm 28$	$994^{+38}_{-37} \pm 69$	$2.33 \pm 0.26 \pm 0.84$	1422 $^{+43}_{-42} \pm 90$	1366
	400-600	$282^{+18}_{-17} \pm 20$	$788^{+30}_{-29} \pm 55$	$0.27 \pm 0.06 \pm 0.10$	1071 $^{+35}_{-34} \pm 69$	1057
	600-800	$43.5^{+3.2}_{-3.1} \pm 6.5$	$129 \pm 5 \pm 12$	< 0.01	172 $\pm 6 \pm 15$	225
	800-1000	$4.6 \pm 0.7 \pm 1.3$	$18.8 \pm 0.7 \pm 3.3$	< 0.01	23.4 $\pm 1.0 \pm 3.6$	22
	≥ 1000	$0.34 \pm 0.08 \pm 0.14$	$2.05 \pm 0.08 \pm 0.90$	< 0.01	2.38 $\pm 0.11 \pm 0.91$	1
2-3j, 2b	200-300	$105.1^{+9.2}_{-8.7} \pm 7.6$	$181^{+20}_{-18} \pm 15$	$3.8 \pm 0.5 \pm 1.3$	290 $^{+22}_{-20} \pm 20$	316
	300-400	$55.0^{+6.7}_{-6.3} \pm 7.5$	$122^{+14}_{-12} \pm 10$	$0.27 \pm 0.06 \pm 0.10$	177 $^{+15}_{-14} \pm 14$	159
	400-600	$36.5^{+4.6}_{-4.3} \pm 5.5$	$97^{+11}_{-10} \pm 8$	$0.08 \pm 0.03 \pm 0.03$	133 $^{+12}_{-11} \pm 11$	107
	600-800	$4.7 \pm 0.8 \pm 1.3$	$15.8^{+1.8}_{-1.6} \pm 1.6$	< 0.01	20.6 $^{+1.9}_{-1.8} \pm 2.2$	21
	≥ 800	$0.59 \pm 0.19 \pm 0.24$	$2.56^{+0.29}_{-0.26} \pm 0.45$	< 0.01	3.14 $^{+0.35}_{-0.32} \pm 0.52$	1
2-6j, $\geq 3b$	200-300	$299^{+17}_{-16} \pm 22$	$73^{+15}_{-13} \pm 10$	$6.2 \pm 0.4 \pm 2.1$	379 $^{+22}_{-21} \pm 28$	345
	300-400	$100 \pm 10 \pm 7$	$43.5^{+8.8}_{-7.4} \pm 6.2$	$0.68 \pm 0.09 \pm 0.24$	144 $^{+14}_{-12} \pm 11$	132
	400-600	$32.5^{+6.3}_{-5.6} \pm 2.5$	$31.2^{+6.3}_{-5.3} \pm 4.4$	$0.08 \pm 0.03 \pm 0.03$	63.8 $^{+8.9}_{-7.7} \pm 5.8$	48
	600-800	$3.16^{+0.95}_{-0.90} \pm 0.68$	$5.4^{+1.1}_{-0.9} \pm 0.8$	< 0.01	8.6 $^{+1.4}_{-1.3} \pm 1.1$	4
	≥ 800	$0.10 \pm 0.03 \pm 0.04$	$0.71^{+0.14}_{-0.12} \pm 0.15$	< 0.01	0.81 $^{+0.15}_{-0.12} \pm 0.16$	0
4-6j, 0b	200-300	$6280 \pm 70 \pm 420$	$9470 \pm 160 \pm 650$	$360 \pm 20 \pm 110$	16100 $\pm 180 \pm 1000$	16292
	300-400	$2700 \pm 50 \pm 180$	$5410 \pm 90 \pm 380$	$53 \pm 1 \pm 17$	8160 $\pm 100 \pm 520$	8330
	400-500	$927^{+28}_{-27} \pm 62$	$2420 \pm 40 \pm 180$	$7.7 \pm 0.4 \pm 2.4$	3350 $\pm 50 \pm 230$	3576
	500-600	$324^{+17}_{-16} \pm 22$	$1171^{+20}_{-19} \pm 100$	$1.46 \pm 0.12 \pm 0.46$	1500 $\pm 30 \pm 110$	1516
	600-700	$95.4^{+9.4}_{-8.7} \pm 6.4$	$413 \pm 7 \pm 47$	$0.33 \pm 0.06 \pm 0.10$	509 $^{+12}_{-11} \pm 50$	543
	700-800	$35.6^{+5.0}_{-4.5} \pm 3.6$	$171 \pm 3 \pm 27$	$0.03 \pm 0.02 \pm 0.01$	206 $^{+6}_{-5} \pm 27$	178
	800-900	$13.4^{+2.0}_{-1.8} \pm 1.6$	$64 \pm 1 \pm 11$	$0.02 \pm 0.01 \pm 0.01$	77 $\pm 2 \pm 11$	62
	900-1000	$4.39^{+0.78}_{-0.73} \pm 0.93$	$23.6 \pm 0.4 \pm 5.3$	< 0.01	28.0 $^{+0.9}_{-0.8} \pm 5.4$	20
	1000-1100	$0.64 \pm 0.16 \pm 0.20$	$6.3 \pm 0.1 \pm 2.0$	< 0.01	6.9 $\pm 0.2 \pm 2.0$	3
	≥ 1100	$0.78 \pm 0.58 \pm 0.32$	$0.89^{+0.02}_{-0.01} \pm 0.40$	< 0.01	1.68 $\pm 0.58 \pm 0.52$	1
4-6j, 1b	200-300	$2900 \pm 50 \pm 200$	$2220^{+80}_{-70} \pm 150$	$154 \pm 16 \pm 50$	5270 $\pm 90 \pm 330$	5335
	300-400	$1066 \pm 29 \pm 74$	$1267^{+44}_{-42} \pm 89$	$19.2 \pm 0.9 \pm 6.2$	2350 $\pm 50 \pm 150$	2547
	400-600	$504^{+22}_{-21} \pm 35$	$840^{+29}_{-28} \pm 61$	$2.98 \pm 0.21 \pm 0.93$	1347 $^{+36}_{-35} \pm 88$	1284
	600-800	$35.3^{+5.9}_{-5.2} \pm 2.6$	$138 \pm 5 \pm 14$	$0.09 \pm 0.03 \pm 0.03$	174 $^{+8}_{-7} \pm 16$	151
	800-1000	$3.89^{+0.83}_{-0.77} \pm 0.82$	$19.3^{+0.7}_{-0.6} \pm 4.3$	$0.01 \pm 0.01 \pm 0.00$	23.2 $^{+1.1}_{-1.0} \pm 4.5$	18
	≥ 1000	$0.18 \pm 0.07 \pm 0.07$	$1.57 \pm 0.05 \pm 0.65$	< 0.01	1.75 $\pm 0.09 \pm 0.65$	1
4-6j, 2b	200-300	$1500 \pm 30 \pm 100$	$473^{+36}_{-33} \pm 36$	$42 \pm 2 \pm 13$	2020 $\pm 50 \pm 130$	1968
	300-400	$508 \pm 20 \pm 35$	$270^{+20}_{-19} \pm 21$	$4.9 \pm 0.3 \pm 1.6$	783 $^{+29}_{-28} \pm 50$	788
	400-600	$167 \pm 12 \pm 12$	$179^{+14}_{-13} \pm 14$	$0.57 \pm 0.08 \pm 0.18$	346 $^{+18}_{-17} \pm 23$	354
	600-800	$11.9^{+1.3}_{-1.2} \pm 2.5$	$29.5^{+2.2}_{-2.1} \pm 3.5$	$0.02 \pm 0.01 \pm 0.01$	41.4 $^{+2.6}_{-2.4} \pm 4.6$	37
	≥ 800	$0.91 \pm 0.23 \pm 0.37$	$4.4 \pm 1.3 \pm 1.8$	< 0.01	5.4 $\pm 0.4 \pm 1.9$	7

Table A.5: Predictions and observations for the 34 search regions with $575 \leq H_T < 1200$ GeV, $N_{\text{jet}} \geq 7$. For each of the background predictions, the first uncertainty listed is statistical (from the limited size of data control samples and Monte Carlo samples), and the second is systematic. Reprinted from [30].

$575 \leq H_T < 1200$ GeV, $N_{\text{jet}} \geq 7$						
$N_{\text{jet}}, N_{\text{b-tag}}$	$M_{\text{T2}} [\text{GeV}]$	Lost lepton	$Z \rightarrow vv$	Multijet	Total background	Data
7-9j, 0b	200-300	$589^{+27}_{-26} \pm 39$	$573^{+47}_{-43} \pm 64$	$90 \pm 10 \pm 28$	$1252^{+55}_{-52} \pm 93$	1340
	300-400	$265^{+19}_{-18} \pm 18$	$279^{+23}_{-21} \pm 42$	$14.9 \pm 0.5 \pm 4.7$	$559^{+29}_{-28} \pm 51$	581
	400-600	$92^{+10}_{-9} \pm 6$	$159^{+13}_{-12} \pm 28$	$2.72 \pm 0.18 \pm 0.85$	$253^{+16}_{-15} \pm 30$	243
	600-800	$8.6 \pm 1.2 \pm 1.8$	$22.8^{+1.9}_{-1.7} \pm 6.4$	$0.10 \pm 0.03 \pm 0.03$	$31.6^{+2.2}_{-2.1} \pm 6.8$	32
	≥ 800	$0.51 \pm 0.16 \pm 0.21$	$3.0 \pm 0.2 \pm 1.3$	< 0.01	$3.5 \pm 0.3 \pm 1.3$	2
7-9j, 1b	200-300	$733 \pm 21 \pm 52$	$278^{+28}_{-25} \pm 33$	$48 \pm 3 \pm 16$	$1059^{+35}_{-33} \pm 73$	1052
	300-400	$252^{+13}_{-12} \pm 18$	$135^{+14}_{-12} \pm 21$	$7.7 \pm 0.4 \pm 2.5$	$395^{+19}_{-17} \pm 32$	387
	400-600	$71.3^{+6.9}_{-5.5} \pm 5.2$	$77^{+8}_{-7} \pm 14$	$1.36 \pm 0.13 \pm 0.45$	$150 \pm 10 \pm 16$	131
	600-800	$4.26^{+0.73}_{-0.71} \pm 0.90$	$11.0^{+1.1}_{-1.0} \pm 3.1$	$0.03 \pm 0.02 \pm 0.01$	$15.3^{+1.3}_{-1.2} \pm 3.3$	20
	≥ 800	$0.11 \pm 0.04 \pm 0.05$	$1.48^{+0.15}_{-0.13} \pm 0.63$	< 0.01	$1.60^{+0.15}_{-0.14} \pm 0.63$	1
7-9j, 2b	200-300	$675 \pm 20 \pm 51$	$82^{+8}_{-7} \pm 10$	$20.9 \pm 3.0 \pm 6.7$	$777^{+22}_{-21} \pm 56$	750
	300-400	$211 \pm 11 \pm 16$	$39.8^{+4.0}_{-3.6} \pm 6.4$	$2.42 \pm 0.19 \pm 0.79$	$253^{+12}_{-11} \pm 19$	259
	400-600	$55.4^{+5.5}_{-5.2} \pm 4.2$	$22.7^{+2.3}_{-2.1} \pm 4.2$	$0.50 \pm 0.07 \pm 0.16$	$78.6^{+5.9}_{-5.6} \pm 6.6$	72
	600-800	$3.00^{+0.63}_{-0.62} \pm 0.64$	$3.25^{+0.32}_{-0.30} \pm 0.93$	$0.01 \pm 0.01 \pm 0.01$	$6.3 \pm 0.7 \pm 1.2$	7
	≥ 800	$0.27 \pm 0.20 \pm 0.11$	$0.44 \pm 0.04 \pm 0.19$	< 0.01	$0.71 \pm 0.20 \pm 0.22$	1
7-9j, 3b	200-300	$185 \pm 8 \pm 18$	$11.3^{+1.1}_{-1.0} \pm 1.9$	$3.6 \pm 0.2 \pm 1.2$	$200 \pm 8 \pm 18$	184
	300-400	$52.0 \pm 3.8 \pm 5.0$	$5.5 \pm 0.5 \pm 1.2$	$0.72 \pm 0.12 \pm 0.26$	$58.3^{+3.9}_{-3.8} \pm 5.3$	59
	400-600	$13.6 \pm 1.8 \pm 1.3$	$3.13^{+0.31}_{-0.29} \pm 0.82$	$0.05 \pm 0.02 \pm 0.02$	$16.8 \pm 1.8 \pm 1.6$	14
	≥ 600	$0.49 \pm 0.21 \pm 0.20$	$0.51 \pm 0.05 \pm 0.21$	< 0.01	$1.00 \pm 0.21 \pm 0.29$	2
7-9j, $\geq 4b$	200-300	$38.8 \pm 3.1 \pm 7.4$	$2.01^{+0.20}_{-0.18} \pm 0.71$	$0.55 \pm 0.08 \pm 0.19$	$41.3^{+3.2}_{-3.1} \pm 7.4$	38
	300-400	$14.5^{+2.0}_{-1.9} \pm 2.8$	$0.98^{+0.10}_{-0.09} \pm 0.43$	$0.06 \pm 0.02 \pm 0.02$	$15.6^{+2.0}_{-1.9} \pm 2.8$	16
	≥ 400	$3.75^{+0.98}_{-0.97} \pm 0.70$	$0.65 \pm 0.06 \pm 0.35$	< 0.01	$4.40^{+0.98}_{-0.97} \pm 0.79$	3
$\geq 10j$, 0b	200-300	$11.5 \pm 1.6 \pm 1.0$	$4.4^{+0.4}_{-0.3} \pm 2.3$	$3.1 \pm 0.8 \pm 1.1$	$19.0 \pm 1.8 \pm 2.8$	27
	300-500	$5.6 \pm 1.0 \pm 0.5$	$3.0 \pm 0.2 \pm 1.7$	$0.55 \pm 0.08 \pm 0.20$	$9.1 \pm 1.0 \pm 1.8$	4
	≥ 500	$0.30 \pm 0.11 \pm 0.12$	$0.44^{+0.04}_{-0.03} \pm 0.24$	$0.02 \pm 0.01 \pm 0.01$	$0.76 \pm 0.11 \pm 0.27$	3
$\geq 10j$, 1b	200-300	$21.0 \pm 1.8 \pm 1.6$	$3.5 \pm 0.3 \pm 1.9$	$1.92 \pm 0.18 \pm 0.72$	$26.4 \pm 1.8 \pm 2.7$	32
	300-500	$7.7 \pm 1.0 \pm 0.6$	$2.4 \pm 0.2 \pm 1.4$	$0.45 \pm 0.07 \pm 0.17$	$10.5 \pm 1.1 \pm 1.6$	15
	≥ 500	$0.83^{+0.42}_{-0.41} \pm 0.07$	$0.36^{+0.04}_{-0.03} \pm 0.20$	$0.02 \pm 0.01 \pm 0.01$	$1.20^{+0.42}_{-0.41} \pm 0.22$	0
$\geq 10j$, 2b	200-300	$21.8 \pm 1.8 \pm 1.6$	$1.05 \pm 0.10 \pm 0.66$	$0.64 \pm 0.08 \pm 0.24$	$23.5 \pm 1.8 \pm 1.8$	26
	300-500	$8.8 \pm 1.2 \pm 0.6$	$0.69^{+0.07}_{-0.06} \pm 0.45$	$0.16 \pm 0.04 \pm 0.06$	$9.6^{+1.3}_{-1.2} \pm 0.8$	9
	≥ 500	$0.22 \pm 0.13 \pm 0.02$	$0.10 \pm 0.01 \pm 0.06$	< 0.01	$0.32 \pm 0.13 \pm 0.07$	0
$\geq 10j$, 3b	200-300	$9.9 \pm 1.3 \pm 1.2$	$0.25 \pm 0.02 \pm 0.20$	$0.29 \pm 0.05 \pm 0.12$	$10.4 \pm 1.3 \pm 1.2$	14
	≥ 300	$1.59 \pm 0.50 \pm 0.18$	$0.19 \pm 0.02 \pm 0.16$	$0.02 \pm 0.01 \pm 0.01$	$1.80 \pm 0.50 \pm 0.25$	2
$\geq 10j$, $\geq 4b$	≥ 200	$3.9 \pm 1.2 \pm 0.8$	$0.00^{+0.17}_{-0.00} \pm 0.00$	$0.05 \pm 0.02 \pm 0.02$	$4.0 \pm 1.2 \pm 0.8$	6

Table A.6: Predictions and observations for the 37 search regions with $1200 \leq H_T < 1500$ GeV, $N_{\text{jet}} < 7$. For each of the background predictions, the first uncertainty listed is statistical (from the limited size of data control samples and Monte Carlo samples), and the second is systematic. Reprinted from [30].

$1200 \leq H_T < 1500$ GeV, $N_{\text{jet}} < 7$						
$N_{\text{jet}}, N_{\text{b-tag}}$	$M_{\text{T2}} [\text{GeV}]$	Lost lepton	$Z \rightarrow vv$	Multijet	Total background	Data
2-3j, 0b	200-400	$315 \pm 15 \pm 21$	$656^{+51}_{-47} \pm 73$	$39 \pm 16 \pm 12$	$1009^{+55}_{-52} \pm 85$	1128
	400-600	$43.0^{+5.2}_{-4.7} \pm 4.9$	$185^{+14}_{-13} \pm 30$	$0.03 \pm 0.02 \pm 0.01$	$228^{+15}_{-14} \pm 31$	207
	600-800	$14.1^{+2.1}_{-2.0} \pm 1.7$	$64 \pm 5 \pm 17$	< 0.01	$78 \pm 5 \pm 17$	83
	800-1000	$6.4^{+1.1}_{-1.0} \pm 1.3$	$32.5^{+2.5}_{-2.3} \pm 7.6$	< 0.01	$38.9^{+2.7}_{-2.5} \pm 7.8$	36
	1000-1200	$3.23^{+0.61}_{-0.59} \pm 0.99$	$17.5 \pm 1.3 \pm 5.2$	< 0.01	$20.7^{+1.5}_{-1.4} \pm 5.3$	19
	≥ 1200	$0.87^{+0.14}_{-0.13} \pm 0.35$	$6.0^{+0.5}_{-0.4} \pm 2.6$	< 0.01	$6.9 \pm 0.5 \pm 2.6$	4
2-3j, 1b	200-400	$61.5^{+7.2}_{-6.5} \pm 4.2$	$78^{+19}_{-16} \pm 10$	$9.7 \pm 0.7 \pm 3.0$	$149^{+21}_{-17} \pm 12$	157
	400-600	$10.1 \pm 1.4 \pm 1.0$	$21.9^{+5.4}_{-4.4} \pm 3.8$	$0.03 \pm 0.02 \pm 0.01$	$32.0^{+5.6}_{-4.6} \pm 4.1$	27
	600-800	$2.36^{+0.36}_{-0.35} \pm 0.41$	$7.5^{+1.9}_{-1.5} \pm 2.0$	< 0.01	$9.8^{+1.9}_{-1.6} \pm 2.1$	9
	800-1000	$0.78^{+0.16}_{-0.15} \pm 0.19$	$3.84^{+0.95}_{-0.78} \pm 0.93$	< 0.01	$4.62^{+0.97}_{-0.79} \pm 0.96$	6
	1000-1200	$0.43 \pm 0.08 \pm 0.14$	$2.13^{+0.53}_{-0.43} \pm 0.64$	< 0.01	$2.56^{+0.54}_{-0.44} \pm 0.66$	2
	≥ 1200	$0.14^{+0.05}_{-0.04} \pm 0.06$	$0.71^{+0.18}_{-0.14} \pm 0.31$	< 0.01	$0.86^{+0.18}_{-0.15} \pm 0.31$	0
2-3j, 2b	200-400	$4.8^{+2.0}_{-1.6} \pm 0.3$	$11^{+11}_{-6} \pm 2$	$1.38 \pm 0.13 \pm 0.43$	$18^{+11}_{-6} \pm 2$	18
	400-600	$0.61^{+0.30}_{-0.25} \pm 0.07$	$3.2^{+3.1}_{-1.7} \pm 0.7$	< 0.01	$3.8^{+3.1}_{-1.8} \pm 0.7$	5
	600-800	$0.21^{+0.11}_{-0.09} \pm 0.04$	$1.1^{+1.1}_{-0.6} \pm 0.4$	< 0.01	$1.3^{+1.1}_{-0.6} \pm 0.4$	2
	800-1000	$0.07^{+0.04}_{-0.03} \pm 0.02$	$0.56^{+0.55}_{-0.31} \pm 0.18$	< 0.01	$0.63^{+0.55}_{-0.31} \pm 0.18$	1
	≥ 1000	$0.03 \pm 0.02 \pm 0.01$	$0.42^{+0.41}_{-0.23} \pm 0.18$	< 0.01	$0.46^{+0.41}_{-0.23} \pm 0.18$	1
2-6j, $\geq 3b$	200-400	$22.6^{+4.7}_{-4.2} \pm 1.8$	$0.0^{+6.6}_{-0.0} \pm 0.0$	$4.4 \pm 0.2 \pm 1.5$	$27.0^{+8.1}_{-4.2} \pm 2.4$	25
	400-600	$1.58^{+0.51}_{-0.48} \pm 0.34$	$0.0^{+1.6}_{-0.0} \pm 0.0$	$0.02 \pm 0.01 \pm 0.01$	$1.6^{+1.7}_{-0.5} \pm 0.3$	3
	≥ 600	$0.47^{+0.27}_{-0.26} \pm 0.19$	$0.00^{+0.94}_{-0.00} \pm 0.00$	< 0.01	$0.47^{+0.98}_{-0.26} \pm 0.19$	4
4-6j, 0b	200-400	$606^{+21}_{-20} \pm 41$	$909^{+63}_{-59} \pm 90$	$208 \pm 12 \pm 64$	$1720^{+70}_{-60} \pm 130$	1768
	400-600	$84.3^{+7.4}_{-6.9} \pm 5.8$	$234^{+16}_{-15} \pm 34$	$0.88 \pm 0.09 \pm 0.27$	$319^{+18}_{-17} \pm 36$	301
	600-800	$21.1^{+3.2}_{-2.9} \pm 2.3$	$75 \pm 5 \pm 17$	$0.06 \pm 0.02 \pm 0.02$	$96 \pm 6 \pm 17$	99
	800-1000	$7.6^{+1.2}_{-1.1} \pm 1.1$	$35.2^{+2.4}_{-2.3} \pm 8.0$	$0.01 \pm 0.01 \pm 0.00$	$42.7^{+2.7}_{-2.5} \pm 8.2$	41
	1000-1200	$2.23^{+0.36}_{-0.33} \pm 0.61$	$14.1^{+1.0}_{-0.9} \pm 4.2$	< 0.01	$16.3 \pm 1.0 \pm 4.2$	15
	≥ 1200	$0.47^{+0.10}_{-0.09} \pm 0.19$	$3.0 \pm 0.2 \pm 1.3$	< 0.01	$3.5 \pm 0.2 \pm 1.3$	5
4-6j, 1b	200-400	$278^{+15}_{-14} \pm 20$	$254^{+33}_{-30} \pm 28$	$97 \pm 2 \pm 30$	$629^{+36}_{-33} \pm 50$	579
	400-600	$30.3^{+4.0}_{-3.7} \pm 2.7$	$65^{+9}_{-8} \pm 10$	$0.33 \pm 0.06 \pm 0.10$	$96^{+9}_{-8} \pm 11$	79
	600-800	$8.2^{+1.4}_{-1.3} \pm 1.0$	$21.0^{+2.8}_{-2.5} \pm 4.8$	$0.02 \pm 0.01 \pm 0.01$	$29.2^{+3.1}_{-2.8} \pm 5.0$	16
	800-1000	$2.36^{+0.56}_{-0.54} \pm 0.50$	$9.8^{+1.3}_{-1.1} \pm 2.3$	$0.01 \pm 0.01 \pm 0.00$	$12.2^{+1.4}_{-1.3} \pm 2.4$	9
	1000-1200	$1.00 \pm 0.24 \pm 0.31$	$4.0 \pm 0.5 \pm 1.2$	< 0.01	$5.0^{+0.6}_{-0.5} \pm 1.2$	6
	≥ 1200	$0.07 \pm 0.02 \pm 0.03$	$0.86^{+0.11}_{-0.10} \pm 0.37$	< 0.01	$0.92^{+0.11}_{-0.10} \pm 0.37$	1
4-6j, 2b	200-400	$120.4^{+9.1}_{-8.7} \pm 9.8$	$45^{+18}_{-13} \pm 5$	$26.0 \pm 0.6 \pm 8.1$	$191^{+20}_{-16} \pm 15$	194
	400-600	$11.9 \pm 1.4 \pm 1.5$	$11.5^{+4.6}_{-3.4} \pm 1.8$	$0.11 \pm 0.03 \pm 0.04$	$23.4^{+4.8}_{-3.7} \pm 2.6$	27
	600-800	$3.49 \pm 0.83 \pm 0.75$	$3.7^{+1.5}_{-1.1} \pm 1.0$	< 0.01	$7.2^{+1.7}_{-1.4} \pm 1.3$	7
	800-1000	$0.66 \pm 0.16 \pm 0.20$	$1.73^{+0.69}_{-0.51} \pm 0.48$	< 0.01	$2.38^{+0.71}_{-0.54} \pm 0.53$	3
	≥ 1000	$0.15 \pm 0.04 \pm 0.06$	$0.84^{+0.34}_{-0.25} \pm 0.36$	< 0.01	$1.00^{+0.34}_{-0.25} \pm 0.36$	0

Table A.7: Predictions and observations for the 31 search regions with $1200 \leq H_T < 1500$ GeV, $N_{\text{jet}} \geq 7$. For each of the background predictions, the first uncertainty listed is statistical (from the limited size of data control samples and Monte Carlo samples), and the second is systematic. Reprinted from [30].

$1200 \leq H_T < 1500$ GeV, $N_{\text{jet}} \geq 7$						
$N_{\text{jet}}, N_{\text{b-tag}}$	M_{T2} [GeV]	Lost lepton	$Z \rightarrow vv$	Multijet	Total background	Data
7-9j, 0b	200-400	$120.4^{+9.8}_{-9.2} \pm 9.0$	$108^{+26}_{-21} \pm 21$	$91 \pm 3 \pm 29$	$\mathbf{319}^{+28}_{-24} \pm 38$	379
	400-600	$16.5^{+1.9}_{-1.8} \pm 2.0$	$25.8^{+6.3}_{-5.1} \pm 5.7$	$0.80 \pm 0.09 \pm 0.25$	$\mathbf{43.1}^{+6.5}_{-5.4} \pm 6.3$	45
	600-800	$2.94 \pm 0.42 \pm 0.63$	$8.6^{+2.1}_{-1.7} \pm 2.1$	$0.06 \pm 0.02 \pm 0.02$	$\mathbf{11.6}^{+2.1}_{-1.8} \pm 2.2$	17
	800-1000	$0.77^{+0.14}_{-0.13} \pm 0.24$	$2.90^{+0.70}_{-0.58} \pm 1.00$	$0.01 \pm 0.01 \pm 0.00$	$\mathbf{3.7}^{+0.7}_{-0.6} \pm 1.0$	3
	≥ 1000	$0.11 \pm 0.03 \pm 0.05$	$1.09^{+0.26}_{-0.22} \pm 0.50$	< 0.01	$\mathbf{1.21}^{+0.27}_{-0.22} \pm 0.50$	0
7-9j, 1b	200-400	$133.8^{+8.0}_{-7.7} \pm 9.8$	$36^{+13}_{-10} \pm 8$	$58 \pm 2 \pm 18$	$\mathbf{228}^{+15}_{-13} \pm 23$	247
	400-600	$16.6^{+2.9}_{-2.7} \pm 1.3$	$8.7^{+3.2}_{-2.4} \pm 2.1$	$0.46 \pm 0.07 \pm 0.14$	$\mathbf{25.8}^{+4.3}_{-3.6} \pm 2.7$	23
	600-800	$1.83^{+0.43}_{-0.41} \pm 0.28$	$2.9^{+1.1}_{-0.8} \pm 0.8$	$0.03 \pm 0.02 \pm 0.01$	$\mathbf{4.8}^{+1.1}_{-0.9} \pm 0.8$	7
	800-1000	$0.65^{+0.24}_{-0.23} \pm 0.18$	$0.95^{+0.34}_{-0.26} \pm 0.34$	$0.02 \pm 0.01 \pm 0.01$	$\mathbf{1.62}^{+0.42}_{-0.35} \pm 0.39$	2
	≥ 1000	$0.22 \pm 0.19 \pm 0.09$	$0.36^{+0.13}_{-0.10} \pm 0.17$	< 0.01	$\mathbf{0.58}^{+0.23}_{-0.21} \pm 0.19$	0
7-9j, 2b	200-400	$124.0^{+7.6}_{-7.4} \pm 9.1$	$9.9^{+3.6}_{-2.7} \pm 2.5$	$21.4 \pm 0.5 \pm 6.9$	$\mathbf{155} \pm 8 \pm 12$	162
	400-600	$15.0^{+2.8}_{-2.6} \pm 1.3$	$2.41^{+0.87}_{-0.66} \pm 0.67$	$0.12 \pm 0.03 \pm 0.04$	$\mathbf{17.5}^{+3.0}_{-2.7} \pm 1.5$	18
	600-800	$2.47^{+0.78}_{-0.76} \pm 0.53$	$0.81^{+0.29}_{-0.22} \pm 0.26$	$0.01 \pm 0.01 \pm 0.00$	$\mathbf{3.29}^{+0.83}_{-0.79} \pm 0.60$	1
	≥ 800	$0.24 \pm 0.11 \pm 0.10$	$0.36^{+0.13}_{-0.10} \pm 0.16$	< 0.01	$\mathbf{0.60}^{+0.17}_{-0.15} \pm 0.19$	1
7-9j, 3b	200-400	$30.0 \pm 2.6 \pm 3.2$	$1.89^{+0.68}_{-0.52} \pm 0.64$	$5.0 \pm 0.3 \pm 1.8$	$\mathbf{36.9}^{+2.7}_{-2.6} \pm 3.8$	46
	400-600	$4.1^{+1.1}_{-1.0} \pm 0.6$	$0.45^{+0.16}_{-0.12} \pm 0.18$	$0.02 \pm 0.01 \pm 0.01$	$\mathbf{4.6}^{+1.1}_{-1.0} \pm 0.6$	2
	≥ 600	$0.92^{+0.50}_{-0.49} \pm 0.38$	$0.23^{+0.08}_{-0.06} \pm 0.11$	< 0.01	$\mathbf{1.15} \pm 0.50 \pm 0.40$	1
7-9j, $\geq 4b$	200-400	$9.1 \pm 1.6 \pm 1.8$	$0.26^{+0.10}_{-0.07} \pm 0.23$	$0.88 \pm 0.10 \pm 0.32$	$\mathbf{10.3} \pm 1.6 \pm 1.9$	9
	≥ 400	$0.44^{+0.24}_{-0.23} \pm 0.08$	$0.10^{+0.04}_{-0.03} \pm 0.09$	< 0.01	$\mathbf{0.53} \pm 0.24 \pm 0.12$	0
$\geq 10j$, 0b	200-400	$7.7^{+1.2}_{-1.1} \pm 0.8$	$2.7^{+0.6}_{-0.5} \pm 2.8$	$8.3 \pm 0.9 \pm 3.0$	$\mathbf{18.7}^{+1.6}_{-1.5} \pm 4.1$	17
	400-600	$1.00 \pm 0.32 \pm 0.22$	$0.56^{+0.13}_{-0.11} \pm 0.62$	$0.11 \pm 0.03 \pm 0.04$	$\mathbf{1.66}^{+0.35}_{-0.34} \pm 0.66$	1
	≥ 600	$0.10^{+0.35}_{-0.04} \pm 0.04$	$0.14^{+0.08}_{-0.03} \pm 0.14$	$0.01 \pm 0.01 \pm 0.00$	$\mathbf{0.24}^{+0.36}_{-0.05} \pm 0.15$	0
$\geq 10j$, 1b	200-400	$15.2 \pm 1.8 \pm 1.4$	$1.1^{+0.4}_{-0.3} \pm 1.2$	$5.3 \pm 0.2 \pm 1.9$	$\mathbf{21.6}^{+1.9}_{-1.8} \pm 2.7$	22
	400-600	$1.27^{+0.38}_{-0.36} \pm 0.11$	$0.22^{+0.08}_{-0.06} \pm 0.26$	$0.05 \pm 0.02 \pm 0.02$	$\mathbf{1.55}^{+0.39}_{-0.37} \pm 0.29$	6
	≥ 600	$0.03 \pm 0.02 \pm 0.01$	$0.05^{+0.10}_{-0.01} \pm 0.05$	< 0.01	$\mathbf{0.07}^{+0.11}_{-0.02} \pm 0.05$	0
$\geq 10j$, 2b	200-400	$16.9 \pm 1.8 \pm 1.5$	$0.44^{+0.16}_{-0.12} \pm 0.50$	$2.7 \pm 0.2 \pm 1.0$	$\mathbf{20.1} \pm 1.8 \pm 1.9$	16
	400-600	$2.62^{+0.71}_{-0.68} \pm 0.30$	$0.09 \pm 0.03 \pm 0.11$	$0.01 \pm 0.01 \pm 0.00$	$\mathbf{2.73}^{+0.71}_{-0.68} \pm 0.32$	2
	≥ 600	$0.23 \pm 0.15 \pm 0.10$	$0.02^{+0.08}_{-0.01} \pm 0.02$	< 0.01	$\mathbf{0.25}^{+0.17}_{-0.15} \pm 0.10$	0
$\geq 10j$, 3b	200-400	$5.58^{+0.86}_{-0.85} \pm 0.61$	$0.12^{+0.11}_{-0.03} \pm 0.16$	$1.04 \pm 0.10 \pm 0.42$	$\mathbf{6.74}^{+0.87}_{-0.86} \pm 0.76$	6
	≥ 400	$0.51 \pm 0.22 \pm 0.06$	$0.03^{+0.11}_{-0.01} \pm 0.04$	< 0.01	$\mathbf{0.54}^{+0.25}_{-0.22} \pm 0.08$	0
$\geq 10j$, $\geq 4b$	≥ 200	$2.59 \pm 0.82 \pm 0.62$	$0.10^{+0.13}_{-0.03} \pm 0.13$	$0.31 \pm 0.06 \pm 0.13$	$\mathbf{3.00}^{+0.83}_{-0.82} \pm 0.65$	7

Table A.8: Predictions and observations for the 30 search regions with $H_T \geq 1500$ GeV, $N_{\text{jet}} < 7$. For each of the background predictions, the first uncertainty listed is statistical (from the limited size of data control samples and Monte Carlo samples), and the second is systematic. Reprinted from [30].

$H_T \geq 1500$ GeV, $N_{\text{jet}} < 7$						
$N_{\text{jet}}, N_{\text{b-tag}}$	$M_{\text{T2}} [\text{GeV}]$	Lost lepton	$Z \rightarrow vv$	Multijet	Total background	Data
2-3j, 0b	400-600	$27.2^{+4.4}_{-3.9} \pm 2.5$	$150^{+14}_{-13} \pm 19$	$0.16 \pm 0.04 \pm 0.05$	$\mathbf{177}^{+15}_{-13} \pm 20$	125
	600-800	$7.8^{+1.4}_{-1.2} \pm 0.8$	$38.7^{+3.6}_{-3.3} \pm 8.4$	< 0.01	$\mathbf{46.5}^{+3.9}_{-3.6} \pm 8.6$	37
	800-1000	$2.29^{+0.39}_{-0.34} \pm 0.35$	$17.2^{+1.6}_{-1.5} \pm 3.4$	< 0.01	$\mathbf{19.5}^{+1.7}_{-1.5} \pm 3.4$	19
	1000-1200	$1.20^{+0.21}_{-0.19} \pm 0.26$	$9.0 \pm 0.8 \pm 1.8$	< 0.01	$\mathbf{10.2}^{+0.9}_{-0.8} \pm 1.9$	14
	1200-1400	$0.80^{+0.16}_{-0.14} \pm 0.22$	$4.9^{+0.5}_{-0.4} \pm 1.3$	< 0.01	$\mathbf{5.7}^{+0.5}_{-0.4} \pm 1.4$	4
	1400-1800	$0.43^{+0.09}_{-0.08} \pm 0.15$	$2.80^{+0.26}_{-0.24} \pm 0.98$	< 0.01	$\mathbf{3.23}^{+0.28}_{-0.26} \pm 0.99$	3
	≥ 1800	$0.05 \pm 0.02 \pm 0.02$	$0.41^{+0.04}_{-0.03} \pm 0.19$	< 0.01	$\mathbf{0.46} \pm 0.04 \pm 0.19$	0
2-3j, 1b	400-600	$5.2^{+1.1}_{-1.0} \pm 0.6$	$13.4^{+4.9}_{-3.7} \pm 1.9$	$0.09 \pm 0.03 \pm 0.03$	$\mathbf{18.7}^{+5.0}_{-3.8} \pm 2.1$	23
	600-800	$1.52^{+0.43}_{-0.41} \pm 0.27$	$3.5^{+1.3}_{-1.0} \pm 1.0$	< 0.01	$\mathbf{5.0}^{+1.3}_{-1.0} \pm 1.0$	3
	800-1000	$0.38 \pm 0.09 \pm 0.10$	$1.53^{+0.55}_{-0.42} \pm 0.35$	< 0.01	$\mathbf{1.90}^{+0.56}_{-0.43} \pm 0.37$	3
	1000-1200	$0.10 \pm 0.03 \pm 0.03$	$0.81^{+0.29}_{-0.22} \pm 0.24$	< 0.01	$\mathbf{0.91}^{+0.29}_{-0.22} \pm 0.24$	4
	≥ 1200	$0.19 \pm 0.06 \pm 0.08$	$0.73^{+0.26}_{-0.20} \pm 0.31$	< 0.01	$\mathbf{0.92}^{+0.27}_{-0.21} \pm 0.32$	0
2-3j, 2b	≥ 400	$0.63^{+0.49}_{-0.36} \pm 0.26$	$0.0^{+3.0}_{-0.0} \pm 0.0$	< 0.01	$\mathbf{0.6}^{+3.0}_{-0.4} \pm 0.3$	2
2-6j, $\geq 3b$	400-600	$1.72^{+0.73}_{-0.68} \pm 0.42$	$1.1^{+2.4}_{-0.9} \pm 0.3$	$0.03 \pm 0.02 \pm 0.01$	$\mathbf{2.8}^{+2.5}_{-1.1} \pm 0.6$	1
	≥ 600	$0.37^{+0.19}_{-0.18} \pm 0.16$	$0.5^{+1.2}_{-0.4} \pm 0.2$	< 0.01	$\mathbf{0.9}^{+1.2}_{-0.5} \pm 0.2$	0
4-6j, 0b	400-600	$46.4^{+5.6}_{-5.1} \pm 3.6$	$176^{+15}_{-14} \pm 23$	$1.62 \pm 0.13 \pm 0.46$	$\mathbf{224}^{+16}_{-15} \pm 24$	207
	600-800	$10.6^{+2.3}_{-1.9} \pm 1.2$	$45.5^{+4.0}_{-3.7} \pm 9.9$	$0.07 \pm 0.03 \pm 0.02$	$\mathbf{56}^{+5}_{-4} \pm 10$	62
	800-1000	$4.5^{+1.1}_{-1.0} \pm 0.5$	$20.3^{+1.8}_{-1.6} \pm 3.9$	< 0.01	$\mathbf{24.8}^{+2.1}_{-1.9} \pm 4.1$	31
	1000-1200	$1.35^{+0.30}_{-0.26} \pm 0.24$	$10.6 \pm 0.9 \pm 2.1$	< 0.01	$\mathbf{11.9}^{+1.0}_{-0.9} \pm 2.2$	12
	1200-1400	$0.89^{+0.27}_{-0.25} \pm 0.23$	$5.7 \pm 0.5 \pm 1.5$	< 0.01	$\mathbf{6.6}^{+0.6}_{-0.5} \pm 1.6$	9
	1400-1600	$0.20 \pm 0.05 \pm 0.07$	$2.64^{+0.23}_{-0.21} \pm 0.92$	< 0.01	$\mathbf{2.84}^{+0.24}_{-0.22} \pm 0.92$	3
	≥ 1600	$0.09 \pm 0.03 \pm 0.04$	$1.18 \pm 0.10 \pm 0.51$	< 0.01	$\mathbf{1.27}^{+0.11}_{-0.10} \pm 0.51$	2
4-6j, 1b	400-600	$21.0^{+3.7}_{-3.3} \pm 2.0$	$32.6^{+7.0}_{-5.8} \pm 5.5$	$0.81 \pm 0.09 \pm 0.23$	$\mathbf{54.5}^{+7.9}_{-6.7} \pm 6.3$	72
	600-800	$4.79^{+0.91}_{-0.83} \pm 0.62$	$8.4^{+1.8}_{-1.5} \pm 2.3$	$0.02 \pm 0.01 \pm 0.01$	$\mathbf{13.2}^{+2.0}_{-1.7} \pm 2.5$	20
	800-1000	$1.27^{+0.26}_{-0.24} \pm 0.27$	$3.71^{+0.79}_{-0.66} \pm 0.92$	$0.03 \pm 0.02 \pm 0.01$	$\mathbf{5.01}^{+0.84}_{-0.71} \pm 0.97$	8
	1000-1400	$0.89^{+0.21}_{-0.20} \pm 0.28$	$3.00^{+0.64}_{-0.54} \pm 0.93$	< 0.01	$\mathbf{3.89}^{+0.68}_{-0.57} \pm 0.98$	6
	≥ 1400	$0.40^{+0.34}_{-0.33} \pm 0.16$	$0.72^{+0.15}_{-0.13} \pm 0.31$	< 0.01	$\mathbf{1.12}^{+0.37}_{-0.36} \pm 0.36$	3
4-6j, 2b	400-600	$7.2^{+1.2}_{-1.1} \pm 1.1$	$4.3^{+2.9}_{-1.9} \pm 1.4$	$0.17 \pm 0.04 \pm 0.05$	$\mathbf{11.7}^{+3.2}_{-2.2} \pm 1.9$	11
	600-800	$1.66^{+0.41}_{-0.40} \pm 0.46$	$1.12^{+0.76}_{-0.48} \pm 0.55$	$0.01 \pm 0.01 \pm 0.00$	$\mathbf{2.79}^{+0.86}_{-0.63} \pm 0.73$	3
	≥ 800	$0.32 \pm 0.13 \pm 0.13$	$0.99^{+0.67}_{-0.43} \pm 0.52$	< 0.01	$\mathbf{1.31}^{+0.68}_{-0.45} \pm 0.54$	4

Table A.9: Predictions and observations for the 21 search regions with $H_T \geq 1500$ GeV, $N_{\text{jet}} \geq 7$. For each of the background predictions, the first uncertainty listed is statistical (from the limited size of data control samples and Monte Carlo samples), and the second is systematic. Reprinted from [30].

$H_T \geq 1500$ GeV, $N_{\text{jet}} \geq 7$						
$N_{\text{jet}}, N_{\text{b-tag}}$	M_{T2} [GeV]	Lost lepton	$Z \rightarrow vv$	Multijet	Total background	Data
7-9j, 0b	400-600	$14.3^{+1.8}_{-1.7} \pm 1.7$	$32.3^{+7.5}_{-6.2} \pm 4.3$	$1.50 \pm 0.13 \pm 0.44$	$\mathbf{48.1}^{+7.7}_{-6.4} \pm 5.0$	36
	600-800	$3.77^{+0.56}_{-0.55} \pm 0.69$	$8.3^{+1.9}_{-1.6} \pm 2.2$	$0.18 \pm 0.04 \pm 0.05$	$\mathbf{12.3}^{+2.0}_{-1.7} \pm 2.3$	9
	800-1000	$1.16^{+0.18}_{-0.17} \pm 0.30$	$3.70^{+0.86}_{-0.71} \pm 0.83$	$0.01 \pm 0.01 \pm 0.00$	$\mathbf{4.86}^{+0.88}_{-0.73} \pm 0.90$	6
	1000-1400	$0.58 \pm 0.11 \pm 0.19$	$2.96^{+0.69}_{-0.57} \pm 0.86$	$0.01 \pm 0.01 \pm 0.00$	$\mathbf{3.55}^{+0.69}_{-0.58} \pm 0.89$	4
	≥ 1400	$0.05 \pm 0.01 \pm 0.02$	$0.71^{+0.17}_{-0.14} \pm 0.30$	< 0.01	$\mathbf{0.76}^{+0.17}_{-0.14} \pm 0.30$	2
7-9j, 1b	400-600	$12.8^{+2.5}_{-2.3} \pm 1.6$	$9.2^{+4.2}_{-3.0} \pm 1.4$	$0.82 \pm 0.09 \pm 0.24$	$\mathbf{22.9}^{+4.9}_{-3.8} \pm 2.3$	25
	600-800	$3.49^{+0.94}_{-0.89} \pm 0.76$	$2.4^{+1.1}_{-0.8} \pm 1.0$	$0.06 \pm 0.02 \pm 0.02$	$\mathbf{5.9}^{+1.4}_{-1.2} \pm 1.2$	7
	≥ 800	$1.09^{+0.34}_{-0.32} \pm 0.45$	$2.10^{+0.96}_{-0.69} \pm 0.93$	< 0.01	$\mathbf{3.2}^{+1.0}_{-0.8} \pm 1.0$	2
7-9j, 2b	400-600	$8.1^{+1.8}_{-1.6} \pm 1.0$	$2.4^{+1.1}_{-0.8} \pm 0.4$	$0.35 \pm 0.06 \pm 0.10$	$\mathbf{10.9}^{+2.1}_{-1.8} \pm 1.2$	10
	600-800	$1.78^{+0.54}_{-0.52} \pm 0.40$	$0.62^{+0.28}_{-0.20} \pm 0.25$	$0.02 \pm 0.01 \pm 0.01$	$\mathbf{2.41}^{+0.61}_{-0.56} \pm 0.49$	5
	≥ 800	$0.40^{+0.19}_{-0.18} \pm 0.17$	$0.55^{+0.25}_{-0.18} \pm 0.25$	$0.01 \pm 0.01 \pm 0.00$	$\mathbf{0.96}^{+0.31}_{-0.26} \pm 0.30$	0
7-9j, 3b	400-800	$2.40^{+0.74}_{-0.72} \pm 0.29$	$0.32^{+0.15}_{-0.10} \pm 0.12$	$0.10 \pm 0.03 \pm 0.03$	$\mathbf{2.82}^{+0.76}_{-0.72} \pm 0.32$	2
	≥ 800	$0.16 \pm 0.09 \pm 0.07$	$0.08^{+0.04}_{-0.03} \pm 0.04$	< 0.01	$\mathbf{0.24} \pm 0.09 \pm 0.08$	0
7-9j, $\geq 4b$	≥ 400	$0.52^{+0.23}_{-0.22} \pm 0.08$	$0.07^{+0.03}_{-0.02} \pm 0.06$	$0.02 \pm 0.01 \pm 0.01$	$\mathbf{0.61}^{+0.23}_{-0.22} \pm 0.10$	1
$\geq 10j$, 0b	400-800	$1.41 \pm 0.38 \pm 0.33$	$1.52^{+0.35}_{-0.29} \pm 0.34$	$0.23 \pm 0.05 \pm 0.08$	$\mathbf{3.17}^{+0.52}_{-0.48} \pm 0.49$	11
	≥ 800	$0.05 \pm 0.02 \pm 0.02$	$0.37^{+0.09}_{-0.07} \pm 0.17$	$0.01 \pm 0.01 \pm 0.00$	$\mathbf{0.43}^{+0.09}_{-0.08} \pm 0.17$	0
$\geq 10j$, 1b	400-800	$2.16^{+0.71}_{-0.69} \pm 0.25$	$0.56^{+0.25}_{-0.18} \pm 0.16$	$0.14 \pm 0.04 \pm 0.05$	$\mathbf{2.85}^{+0.76}_{-0.71} \pm 0.31$	3
	≥ 800	$0.55 \pm 0.30 \pm 0.22$	$0.13^{+0.06}_{-0.04} \pm 0.07$	< 0.01	$\mathbf{0.68}^{+0.31}_{-0.30} \pm 0.23$	0
$\geq 10j$, 2b	≥ 400	$1.98^{+0.69}_{-0.67} \pm 0.24$	$0.30^{+0.14}_{-0.10} \pm 0.12$	$0.05 \pm 0.02 \pm 0.02$	$\mathbf{2.33}^{+0.70}_{-0.68} \pm 0.28$	0
$\geq 10j$, 3b	≥ 400	$0.77 \pm 0.35 \pm 0.09$	$0.00^{+0.45}_{-0.00} \pm 0.00$	$0.05 \pm 0.03 \pm 0.02$	$\mathbf{0.82}^{+0.57}_{-0.35} \pm 0.09$	1
$\geq 10j$, $\geq 4b$	≥ 400	$0.09 \pm 0.05 \pm 0.01$	$0.00^{+0.45}_{-0.00} \pm 0.00$	< 0.01	$\mathbf{0.09}^{+0.45}_{-0.05} \pm 0.01$	0

Appendix B

Full Disappearing Tracks Binning and Results

The following tables provide predicted background event counts and observed yields for all 68 disappearing tracks search regions (28 for 2016 data, and 40 for 2017–2018 data).

Table B.1: Summary of the 28 signal regions of the search for disappearing tracks, for the 2016 data set, together with the corresponding background predictions and observations. For the background predictions, the first uncertainty listed is statistical (from the limited size of control samples), and the second is systematic. The systematic uncertainty is not shown when it is negligible. Taken from [30].

	Track length	N_{jet}	H_T range [GeV]	Track p_T [GeV]	Label	Background	Data
P	2–3	[250, 450)	[15, 50)	P LL lo	15.5 _{-2.7} ^{+3.0} ± 3.2	16	
			[50, ∞)	P LL hi	9.8 _{-2.2} ^{+2.6} ± 2.5	3	
			[15, 50)	P LM lo	4.2 _{-0.9} ^{+1.0} ± 1.2	2	
		[450, 1200)	[50, ∞)	P LM hi	2.02 _{-0.55} ^{+0.66} ± 0.63	1	
			[1200, ∞)	P LH lo	0.19 _{-0.13} ^{+0.26} ± 0.13	0	
			[50, ∞)	P LH hi	0.06 _{-0.05} ^{+0.14} ± 0.03	0	
	≥ 4	[250, 450)	[15, 50)	P HL lo	3.3 _{-0.6} ^{+0.7} ± 1.4	1	
			[50, ∞)	P HL hi	1.98 _{-0.38} ^{+0.43} ± 0.57	1	
			[450, 1200)	P HM lo	4.7 _{-0.7} ^{+0.8} ± 1.9	6	
		[1200, ∞)	[50, ∞)	P HM hi	2.37 _{-0.44} ^{+0.50} ± 0.55	1	
			[15, 50)	P HH lo	0.43 _{-0.17} ^{+0.24} ± 0.27	0	
			[50, ∞)	P HH hi	0.17 _{-0.07} ^{+0.10} ± 0.04	0	
M	2–3	[250, 450)	[15, 50)	M LL lo	3.9 _{-1.2} ^{+1.5} ± 1.3	3	
			[50, ∞)	M LL hi	14 _{-3.2} ^{+3.7} ± 4.0	8	
		[450, 1200)	[15, 50)	M LM lo	2.1 _{-0.71} ^{+0.89} ± 1.1	3	
			[50, ∞)	M LM hi	0.68 _{-0.45} ^{+0.90} ± 0.35	4	
		[1200, ∞)	[15, 50)	M LH lo	0.0 _{-0.0} ^{+0.25} ± 0.0	0	
			[50, ∞)	M LH hi	0.0 _{-0.0} ^{+0.7}	0	
	≥ 4	[250, 450)	[15, 50)	M HL lo	1.8 _{-0.5} ^{+0.6} ± 0.9	0	
			[50, ∞)	M HL hi	2.1 _{-0.6} ^{+0.8} _{-2.1} ^{+2.3}	2	
			[450, 1200)	M HM lo	2.2 _{-0.6} ^{+0.7} ± 1.3	1	
		[1200, ∞)	[50, ∞)	M HM hi	2.9 _{-0.8} ^{+0.9} ± 2.3	0	
			[15, 50)	M HH lo	0.23 _{-0.13} ^{+0.23} ± 0.11	0	
			[50, ∞)	M HH hi	0.30 _{-0.20} ^{+0.40} ± 0.29	1	
L	2–3	[250, 1200)	[15, ∞)	L LLM	0.046 _{-0.034} ^{+0.050} _{-0.046} ^{+0.057}	0	
		[1200, ∞)	[15, ∞)	L LH	0.015 _{-0.015} ^{+0.036} _{-0.015} ^{+0.022}	0	
	≥ 4	[250, 1200)	[15, ∞)	L HLM	0.092 _{-0.085} ^{+0.136} _{-0.092} ^{+0.130}	0	
		[1200, ∞)	[15, ∞)	L HH	0.0 _{-0.0} ^{+0.1}	0	

Table B.2: Summary of the 40 signal regions of the search for disappearing tracks, for the 2017–2018 data set, together with the corresponding background predictions and observations. For the background predictions, the first uncertainty listed is statistical (from the limited size of control samples), and the second is systematic. The systematic uncertainty is not shown when it is negligible. Taken from [30].

	Track length	N_{jet}	H_T range [GeV]	Track p_T [GeV]	Label	Background	Data
P3	2–3		[250, 450)	[15, 50)	P3 LL lo	78 ₋₉ ⁺⁹ ± 34	73
				[50, ∞)	P3 LL hi	43.9 _{-6.2} ^{+6.7} ± 8.1	41
			[450, 1200)	[15, 50)	P3 LM lo	30 ₋₅ ⁺⁵ ± 16	21
				[50, ∞)	P3 LM hi	13 ₋₃ ⁺³ ± 13	16
				[1200, ∞)	P3 LH lo	0.0 _{0.0} ^{+1.0}	1
	≥ 4		[250, 450)	[50, ∞)	P3 LH hi	0.43 _{-0.36} ^{+0.98} ± 0.34	0
				[15, 50)	P3 HL lo	25.8 _{-3.4} ^{+3.8} ± 7.9	17
			[450, 1200)	[50, ∞)	P3 HL hi	10.8 _{-1.8} ^{+2.1} ± 3.5	7
				[15, 50)	P3 HM lo	28.9 _{-3.7} ^{+4.0} ± 5.7	37
				[50, ∞)	P3 HM hi	12.3 _{-1.9} ^{+2.2} ± 6.8	11
P4	2–3		[250, 450)	[15, 50)	P4 HH lo	3.1 _{-1.1} ^{+1.5} ± 0.5	5
				[50, ∞)	P4 HH hi	0.49 _{-0.32} ^{+0.65} ± 0.12	3
			[450, 1200)	[15, 50)	P4 LL lo	24 ₋₅ ⁺⁵ ± 11	10
				[50, ∞)	P4 LL hi	4.1 _{-1.5} ^{+1.9} ± 3.7	0
				[15, 50)	P4 LM lo	8.7 _{-2.2} ^{+2.7} ± 4.6	8
	≥ 4		[1200, ∞)	[50, ∞)	P4 LM hi	1.1 _{-0.5} ^{+0.7} ± ^{+1.4} _{-1.1}	0
				[15, 50)	P4 LH lo	0.40 _{-0.33} ^{+0.91} ± 0.40	0
			[450, 1200)	[50, ∞)	P4 LH hi	0.0 _{-0.0} ^{+0.39}	0
				[15, 50)	P4 HL lo	6.3 _{-1.3} ^{+1.6} ± 2.2	7
				[50, ∞)	P4 HL hi	0.62 _{-0.25} ^{+0.35} ± 0.43	0
M	2–3		[250, 450)	[15, 50)	P4 HM lo	6.9 _{-1.4} ^{+1.6} ± 6.2	2
				[50, ∞)	P4 HM hi	1.32 _{-0.43} ^{+0.54} ± 0.63	2
			[450, 1200)	[15, 50)	P4 HH lo	0.42 _{-0.28} ^{+0.56} ± 0.12	0
				[50, ∞)	P4 HH hi	0.08 _{-0.07} ^{+0.18} ± 0.03	0
				[15, 50)	M LL lo	8.4 _{-2.0} ^{+2.4} ± 3.4	8
	≥ 4		[250, 450)	[50, ∞)	M LL hi	5.4 _{-1.8} ^{+2.2} ± 2.6	2
				[15, 50)	M LM lo	1.90 _{-0.66} ^{+0.85} ± 0.92	6
			[450, 1200)	[50, ∞)	M LM hi	1.12 _{-0.54} ^{+0.77} ± 0.97	1
				[15, 50)	M LH lo	0.00 ₋₀ ^{+0.36}	0
				[50, ∞)	M LH hi	0.00 ₋₀ ^{+0.46}	0
L	2–3		[250, 450)	[15, 50)	M HL lo	1.6 _{-0.5} ^{+0.6} ± ^{+3.0} _{-1.6}	3
				[50, ∞)	M HL hi	1.11 _{-0.42} ^{+0.57} ± 0.58	1
			[450, 1200)	[15, 50)	M HM lo	1.9 _{-0.5} ^{+0.6} ± ^{+3.5} _{-1.9}	3
				[50, ∞)	M HM hi	1.5 _{-0.5} ^{+0.7} ± 1.1	0
	≥ 4		[1200, ∞)	[15, 50)	M HH lo	0.38 _{-0.19} ^{+0.31} ± ^{+0.70} _{-0.38}	1
				[50, ∞)	M HH hi	0.12 _{-0.10} ^{+0.29} ± 0.04	0
			[250, 1200)	[15, ∞)	L LLM	0.46 _{-0.20} ^{+0.26} ± ^{+0.53} _{-0.46}	0
	2–3		[1200, ∞)	[15, ∞)	L LH	0.00 ₋₀ ^{+0.14}	0
			[250, 1200)	[15, ∞)	L HLM	0.013 _{-0.014} ^{+0.015} ± ^{+0.018} _{-0.013}	0
	≥ 4		[1200, ∞)	[15, ∞)	L HH	0.000 ₋₀ ^{+0.008}	0

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