Canonical Quantization Procedure and Time-Energy Uncertainty

A Toy Model of Quantum Gravity

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The Time-Energy Uncertainty Relation

First postulated by Heisenberg [1], the Time-Energy Uncertainty is:

$$\Delta E \cdot \Delta t \ge \hbar/2 \tag{1}$$
 Hamiltonian \hat{H} eigenvalue $\stackrel{\frown}{---}$?

Problem 1: No self-adjoint time operator

No axiomatic derivation similar $\Delta \hat{x} \cdot \Delta \hat{p} > \hbar/2$ based on their commutation relation $[\hat{x}, \hat{p}] = i\hbar$, since a time operator \hat{t} does not exist.

Problem 2: Different interpretations of time uncertainty

- Energy-dispersion (ΔE) of a state and its lifetime (Δt)
- 2 Energy exchange (ΔE) and exchange time-frame (Δt)
- **Solution** Energy measurement (ΔE) and measurement/preparation time (Δt)

Introduction

Notation and Klein-Gordon equation

Notation:

• four vectors: $x^{\mu} = (t, \vec{x})$ and $p^{\mu} = (E, \vec{p})$

2 Minkowski metric: $\eta_{\mu\nu} = (-, +, +, +)$

3 Einstein summation convention: $p_{\mu}p^{\mu} = p_{\mu}p_{\nu}\eta^{\mu\nu} = -p_0^2 + \vec{p}^2$

Def: A constraint on a system is a dynamical relation that the system must obey, of the form:

$$\phi(x^{\mu}, p_{\mu}) = 0 \tag{2}$$

e.g. the mass-shell constraint:

$$p_{\mu}p^{\mu} + m^2c^2 = 0 \tag{3}$$

For $\hat{p}_{\mu} = -i\hbar\partial_{\mu}$, quantizing the above yields:

$$(\hat{\rho}_{\mu}\hat{\rho}^{\mu} + m^2c^2)|\psi\rangle = (-\hbar^2\partial_{\mu}\partial^{\mu} + m^2c^2)|\psi\rangle = 0 \tag{4}$$

i.e. the Klein-Gordon equation.

Free Relativistic Bosonic Particle

The Hamiltonian after quantizing the simplest example is:

$$\hat{H}(x^{\mu}, p_{\mu}) |\psi\rangle = \frac{1}{2} (\hat{p}_{\mu}\hat{p}^{\mu} + m^{2}c^{2}) |\psi\rangle = 0$$
 (5)

This is called a Hamiltonian Constraint.

Problem of Time in Quantum Gravity

$$i\hbar \frac{\partial}{\partial t} |\psi\rangle = \hat{H} |\psi\rangle = 0 \tag{6}$$

Proposed solution: 4+1 formalism

Introduce an extra parameter time λ giving the system's evolution as:

$$i\hbar\frac{\partial}{\partial\lambda}|\psi\rangle = \hat{H}_0|\psi\rangle \tag{7}$$

World Line Parametrization

An observer moves in space-time as parametrized by the time evolution parameter λ . [2]

e.g. identify the parameter with proper time:

$$\lambda = \alpha \tau \tag{8}$$

Just a specific gauge choice, in general, $\lambda = f(\tau)$.

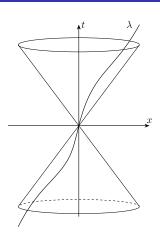


Figure 1: World line parametrization by λ

Free Relativistic Bosonic Particle (Again)

The Hamiltonian constraint after quantizing the simplest example is:

$$\hat{H} |\psi\rangle_{m} = \frac{1}{2} \left(\underbrace{\hat{p}_{\mu} \hat{p}^{\mu}}_{\text{operator}} + \underbrace{m^{2} c^{2}}_{\text{scalar value}} \right) |\psi\rangle_{m} = 0$$
 (9)

Reinterpret as parameter time independent equations:

$$\hat{H}_{0} |\psi\rangle_{m} = \hat{p}_{\mu} \hat{p}^{\mu} |\psi\rangle_{m} = i\hbar \frac{\partial}{\partial \lambda} |\psi\rangle_{m} \stackrel{!}{=} -m^{2} c^{2} |\psi\rangle_{m}$$
 (10)

solved by physical wave functions $|\psi\rangle_m \in \mathcal{H}_m \subset \mathcal{H}$.

But for arbitrary (not necessary on mass-shell) wave functions:

Parameter time evolution

$$\hat{H}_0 |\psi\rangle = \hat{\rho}_{\mu} \hat{\rho}^{\mu} |\psi\rangle = i\hbar \frac{\partial |\psi\rangle}{\partial \lambda} \tag{11}$$

Extension to Four Coordinate Dimensions

The wave functions $|\psi\rangle$ live in a four dimensional Hilbert Space, taken to be $\mathcal{H}=L^2(\mathbb{R}^4)$.

$$\langle \psi_1 | \psi_2 \rangle = \int_{\mathbb{R}^4} d^4 x \, \psi_1^*(x^\mu) \psi_2(x^\mu) \quad \text{with} \quad \|\psi\| = \sqrt{\langle \psi | \psi \rangle} \le \infty \quad (12)$$

Problem: How to normalize wave functions in time?

Take $U(t) = \exp\{-iEt\}$, then any wave function would be non-normalizable, as:

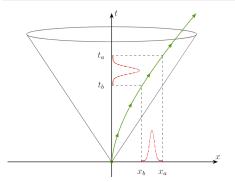
$$\langle \psi | \psi \rangle = \int_{\mathbb{R}} dt U^*(t) U(t) \int_{\mathbb{R}^3} d^3 x \, \psi^*(\vec{x}) \psi(\vec{x}) = \int_{\mathbb{R}} dt \to \infty$$
 (13)

Time Gaussian Wave Packets!

Time-localized Wave Functions

Particles localized in time

Wave packets evolve in parameter time λ and propagate in the space-time.



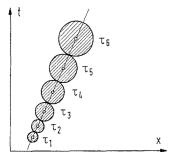


Figure 2: Space-time probabilities

Figure 3: Gaussian Spreading [2]

Probability Interpretation

The Born probability interpretation [3] holds! By construction, the probability density $\rho = \psi^* \psi$ is positive and moreover:

Continuity equation

$$\partial_{\lambda}\rho + \partial_{\mu}j^{\mu} = 0 \tag{14}$$

where:

$$j^{\mu} = -i(\psi^* \partial^{\mu} \psi - \psi \partial^{\mu} \psi^*) \tag{15}$$

Time-Energy Uncertainty

Uncertainty relations [4] for operators \hat{A} and \hat{B} :

$$(\Delta \hat{A})^2 (\Delta \hat{B})^2 \ge \frac{1}{2} \left| \langle [\hat{A}, \hat{B}] \rangle \right|^2 \tag{16}$$

From the position-momentum commutation relation $[\hat{x}^{\mu}, \hat{p}_{\nu}] = i\hbar \delta^{\mu}_{\nu}$:

$$\Delta \hat{x}^{\mu} \Delta \hat{\rho}_{\nu} \ge \frac{1}{2} \left| \langle \psi | \left[\hat{x}^{\mu}, \hat{\rho}_{\nu} \right] | \psi \rangle \right| = \frac{\hbar}{2} \delta_{\nu}^{\mu} \tag{17}$$

The zeroth component is:

Main result: The Time-Energy Uncertainty

$$\Delta x^0 \Delta p_0 = \Delta t \Delta E \ge \frac{\hbar}{2} \tag{18}$$

Summary

- ① Introduced another time evolution parameter λ to implement the Hamiltonian constraint as an eigenvalue problem.
- ② Promoted coordinate time $t = x^0$ to an operator.
- **②** Reformulated Relativistic Quantum Mechanics for the Bosonic Free Particle accordingly: $\psi \in L^2(\mathbb{R}^4)$, probability interpretation, etc.
- Axiomatically obtained the Time-Energy Uncertainty Relation via operator commutation relations.

References

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