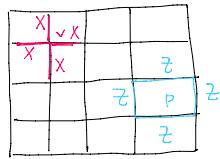


## Interpreting toric code as an instance of TV TQFT

Thursday, June 3, 2021 11:56 PM

Toric code:



$$\mathcal{H} = \text{Hilbert space} = \bigotimes_{\text{edges}} \mathbb{C}^2 \quad (= \bigotimes_{\text{edges}} \mathbb{C} G)$$

for  $G = \mathbb{Z}_2$

$$(\Sigma, \Delta) = (T^2, \text{N} \times \text{N} \text{ lattice})$$

surface      cell  
decomp

vertex operators:  $A_v$

$$\begin{array}{|c|} \hline X \\ \hline \end{array}$$

plaquette operators:  $B_p$

$$\begin{array}{|c|} \hline Z \\ \hline P \\ \hline Z \\ \hline \end{array}$$

(Note: more convenient to use  $\tilde{B}_p := \frac{1}{2}(I + B_p)$   
which is a projection, and  $\tilde{B}_p|\psi\rangle = |\psi\rangle$   
 $\downarrow$   
 $B_p|\psi\rangle = |\psi\rangle$ )

$$\text{Hamiltonian: } \sum_v I - A_v + \sum_p I - B_p = H$$

ground states: correspond to  $|\psi\rangle \in \mathcal{H}$  s.t.  $A_v|\psi\rangle = B_p|\psi\rangle = |\psi\rangle$

- dimension = 4 =  $\# H^1(T^2; \mathbb{Z}_2)$   
(in general =  $\# \left( \{ \text{hom } \pi_1(\Sigma) \rightarrow G \} / \text{conjugation by } G \right)$ )

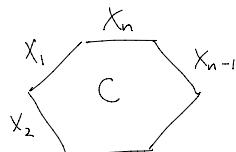
(in particular, the dim. of ground states is independent of  $N$ , actually a topological invariant of  $\Sigma$ ).

Turaev-Viro Theory for the lattice on a torus! (assigns a vector space  $\mathcal{Z}_{TV}(\Sigma)$  to the torus)

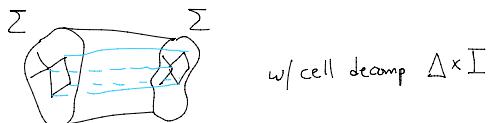
Given  $(\Sigma, \Delta)$  surface w/cell-decomp. and  $\mathcal{E} = \text{Rep } \mathbb{Z}_2$

Let  $\mathcal{H}^{TV}(\Sigma, \Delta) := \bigoplus_{\substack{\text{l simple} \\ \text{labeling of edges} \\ \text{by objects in } \text{Rep } \mathbb{Z}_2}} \bigotimes_{\text{face}} H(C, \mathcal{E})$

$$\text{where } H(C, \mathcal{E}) = \text{Hom}_\mathcal{E}(I, X_1 \otimes \dots \otimes X_n) \\ = \langle X_1, \dots, X_n \rangle$$



Consider  $\Sigma \times [0, 1]$  (cobordism  $\Sigma \Rightarrow \Sigma$ )



Then the TV theory assigns to this a linear map (using the defn given by Itai)

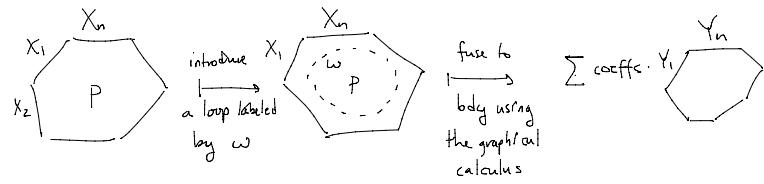
$$\mathcal{Z}^{TV}(\Sigma \times I, \Delta \times I) : H(\Sigma \times I) \rightarrow H^{TV}(\Sigma \times I) \quad (\text{in fact a projection})$$

and then  $\mathcal{Z}^{TV}(\Sigma) := \text{image of this projection.}$

Thm: this image is independent of the original cell-decomp  $\Delta$ , ie only depends on the topology of  $\Sigma$ .

Important technical result: This projection corresponds w/ the operator

$$\prod_P B_P^{TV} \quad \text{where} \quad B_P^{TV} \text{ acts locally at a plaquette by}$$



where  $\begin{array}{c} | \\ | \\ | \\ | \\ | \\ | \end{array}_w$  denotes an edge labeled by the object

$$c_w = \frac{1}{D^2} \sum_{X \in \text{Env}(e)} \dim_{\mathbb{C}}(X) X \quad \left( = \frac{1}{2} (\mathbb{I} \oplus \text{sign}) \text{ when } C = \text{Rep } \mathbb{Z}_2 \right)$$

think  $\mathbb{I} = |+\rangle\langle+|$   
 $\text{sign} = |- \rangle\langle -|$

Thm (Balsam-Kirillov, Thm 4.1)

$$\mathcal{Z}^{TV}(\Sigma) \cong \text{ground state space of Kitaev lattice model}$$

based on  $C = \text{Rep } \mathbb{Z}_2$       based on  $G = \mathbb{Z}_2$ .

Idea: The ground state of just the  $A_v$ 's reproduces  $H^{TV}(\Sigma, \Delta^*)$

Then  $B_P$  condition  $B_P |\psi\rangle = |\psi\rangle$  translates to  $(\prod_P B_P^{TV}) |\psi\rangle = |\psi\rangle$ .

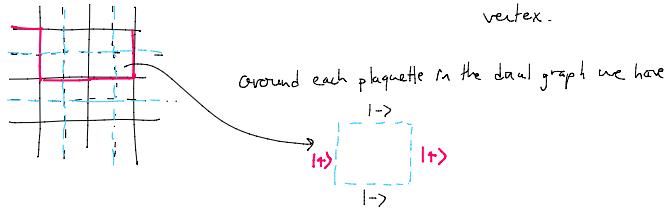
So the ground state for toric code ( $A_v$ 's &  $B_P$ 's) is the image of  $(\prod_P B_P^{TV}) \subseteq H^{TV}(\Sigma, \Delta^*)$ , which is  $\mathcal{Z}^{TV}(\Sigma)$ .

Toric code example: • The conditions  $A_v |\psi\rangle = |\psi\rangle$  mean  $|\psi\rangle$  is a linear combination of states w/  $|+\rangle$ 's and  $|-\rangle$ 's on every edge with an even # of  $|-\rangle$ 's around each vertex.

$$|\psi\rangle = \sum$$

in the dual graph

red =  $|-\rangle$   
black =  $|+\rangle$



simple labeling of edges s.t.

$$H^{TV}(C, \ell) = \langle X_1, X_2, X_3, X_4 \rangle \neq 0$$

$\Leftrightarrow X_1 \otimes X_2 \otimes X_3 \otimes X_4 = \mathbb{I}$   
 $\Leftrightarrow \text{even } \# \text{ of } |-\rangle$ 's).

$$\text{So } |\psi\rangle \in \bigoplus_{\substack{\ell \text{ simple} \\ \text{label}}} \bigotimes_{C \text{ in dual} \\ \text{graph}} H^{TV}(C, \ell)$$

• Now the condition  $B_P |\psi\rangle = |\psi\rangle$  for all  $|\psi\rangle$  corresponds to  $B_P^{TV} |\psi\rangle = |\psi\rangle \quad \forall P$ .

Examine  $B_P^{TV}$ :

$$X_1 \boxed{X_2 \boxed{X_3 \boxed{X_4 \dots}}} = \frac{1}{2} \left( X_1 \boxed{X_2 \boxed{X_3 \boxed{| \rangle}}} X_4 + X_1 \boxed{X_2 \boxed{X_3 \boxed{\text{sign}}} X_4} \right)$$

$$= \frac{1}{2} \left( \sum_{Y_1, Y_2} X_1 \boxed{X_2 \boxed{X_3 \boxed{Y_1 \dots}}} X_4 + \sum_{Y_1, Y_2} X_1 \boxed{X_2 \boxed{X_3 \boxed{Y_1 \dots}}} X_4 \right)$$

local relations from the :  $\begin{array}{c} | \\ | \\ | \\ | \\ | \\ | \end{array} = \sum \begin{array}{c} X \\ Y \\ \oplus \\ \ominus \end{array}$

local relations from the graphical calculus for  $\mathcal{C}$

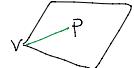
$$X \otimes Y = \sum Z \begin{array}{c} X \\ \otimes \\ Y \end{array}$$

$\uparrow$   
corresponds to  $Z \begin{array}{c} X \\ \otimes \\ Y \end{array}$   
 $Z \otimes Z$

usual basis code for magnetic operator.

Excited states:

In Kitaev code, excitations are localized at sites  $S = (v, p)$



$|\psi\rangle$  is an excitation localized at  $S = (v, p)$  when

$$A_w |\psi\rangle = |\psi\rangle \quad \forall w \neq v$$

$$B_q |\psi\rangle = |\psi\rangle \quad \forall q \neq p$$

At every site there is an algebra of local operators

$$\{I, A_v, B_p, A_v B_p\}$$

which generate an algebra which is a quotient of

$$D(\mathbb{C}\mathbb{Z}_2) = \mathbb{C}\mathbb{Z}_2 \otimes \text{Fun}(\mathbb{Z}_2, \mathbb{C})$$

$$\text{Rep}(D(\mathbb{C}\mathbb{Z}_2))$$

Dijkgraaf-Witten theory

The possible excitations at  $p$  are classified by irreps of  $D(\mathbb{C}\mathbb{Z}_2) = \text{objects of } \mathcal{Z}(\text{Rep } \mathbb{Z}_2)$

(In the case of  $D(\mathbb{C}\mathbb{Z}_2)$  there are 4 irreps

$$\begin{matrix} 1 & e & m & \psi \\ \text{trivial} & & & \end{matrix}, \quad \text{tensor product rules: } \begin{aligned} e \otimes e &= 1 \\ m \otimes m &= 1 \\ \psi \otimes \psi &= 1 \\ e \otimes m &= \psi \end{aligned} \quad \vdots$$

Turaev-Viro: excitations are described by the category  $\mathcal{Z}^{\text{TV}}(S^1)$

$$\mathcal{Z}^{\text{TV}}(S^1) = \text{Dijkgraaf-Witten theory of } \mathcal{C}.$$