

Supplementary Materials

The present study on the role of mixed fields for spin-flip loss evolved out of our continuing investigations into the collisional processes reported in references [1, 2]. It has become evident that a significant portion of the effects attributed to collisions could in fact be attributed to spin-flip losses. In what follows we provide the interested reader with our current best understanding of the situation.

We begin with reference [1] on electric field induced inelastic collisions. The authors identified and investigated the same single particle spin-flip loss enhancement process we discuss in the present work, and an attempt was made at deconvolution from the collisional effect. An appendix of [1] explains this well. Since that time, new observations prompted us to make an even more careful investigation, during which we discovered an important correction to the mathematics in the appendix of [1].

Relative to the approach taken in the appendix of [1], we make the same simplifying assumptions: loss only occurs in the $\vec{E} \perp \vec{B}$ plane, only the velocity orthogonal to this plane matters, and the population is a thermalized Maxwell-Boltzmann distribution. Our correction relates to the next step, where an integral calculation for the loss rate is performed. In [1] the integration spans the entire 3D spatial distribution, weighted by the frequency of crossing of the center plane and the chance of loss for each crossing:

$$\Gamma_{LZ} = \int_0^\infty 4\pi r^2 n(r) dr \int_0^\infty n(v_\vartheta) dv_\vartheta \left(\frac{v_\vartheta}{\pi r} P_{\text{hop}}(r, v_\vartheta) \right) \quad (1)$$

Here $n(r)$ is the radial distribution function, constrained to satisfy $\int_0^\infty 4\pi r^2 n(r) dr = 1$, and of the form $n(r) \propto e^{-\mu_B B^* r / kT}$. Likewise $n(v_\vartheta)$ is the usual normalized Maxwellian velocity distribution. Implicit in this integration is the simplifying assumption that molecules at a given radius r cross the center plane with a frequency of $v_\vartheta / \pi r$. Though not a bad place to start, this approximation is rather rough given that molecules are certainly not following circular orbits of constant v_ϑ but are in general following some unusual trap motion. Our correction is to perform an integration of flux through the loss plane directly:

$$\Gamma_{LZ} = \int_0^\infty 2\pi r n(r) dr \int_0^\infty n(v_z) dv_z (v_z P_{\text{hop}}(r, v_z)) \quad (2)$$

Here the same distributions are used, but the spatial distribution is integrated over the central plane only, hence the $2\pi r$ Jacobean, and the hopping probability is multiplied by the velocity v_z to give a flux. The coordinates z and ϑ are mathematically equivalent in the central plane,

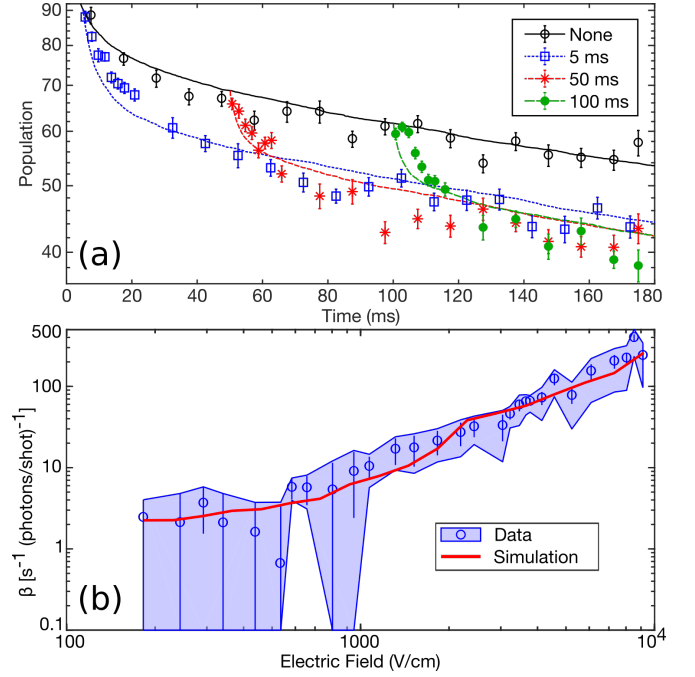


FIG. 1. Experimental electric field induced loss data with an attempted overlap to spin-flip loss simulations. (b) Two body fits from [1] with spin flip loss single particle simulation in red.

but z is more appropriate given that integration is no longer spanning the ϑ direction. This flux integral gives the desired loss rate without any approximations about molecule orbits or plane-crossing frequency. Although the two integrals differ significantly in their conception, mathematically the changes to the integrand and the Jacobean reduce to precisely an overall scaling factor of π .

The influence of this on the deconvolution procedure relates to the details of the two-body fitting routine. One plus two body fits were performed to various decay trap curves, with the one body component fixed to the value expected due to vacuum scattering and spin-flip loss. An example of such decay curves is shown in Fig. 1(a), where Electric field is turned on suddenly after various hold times. The problem is that with the stronger spin-flip loss, it seems no longer appropriate to assume this loss will be present in the data as a pure one-body decay. Rather, only those molecules whose orbits regularly intersect the loss region are lost, after which thermalization would be required to repopulate the loss prone regions of phase space. If thermalization is slow enough, spin-flip loss can actually masquerade as a two-body population curve, since it can lead to a loss rate that decreases over time.

We can perform purely single particle simulations of

spin-flip loss to investigate this, and we obtain curves such as shown above the time dependent experiment data points in Fig. 1(a). The simulations do predict a loss rate that turns off over time, and they do a decent job matching the magnitude of the loss induced by the electric field. We can even perform a two body fitting procedure like the one used in [1] to the single particle spin-flip loss simulation traces, see Fig. 1(b). This suggests that the effect attributed to two-body collisions could be largely explained by spin-flip losses. Still, there are notable differences, such as in the initial rate of the decays in fig. 1(a), which is faster in simulation than in experiment, and in the long time magnitude of the loss, which is maybe 20% larger in experiment than simulation. One avenue to try and be more quantitative would be to incorporate collisions in the simulation and see what collision rates yield the best agreement between simulation and experiment. Unfortunately there are many challenges in the quantitative application of simulations such as these, such as uncertainty regarding the initial distribution and the existence of various partially trapped substates. We think the best path forward is to perform future collisional studies with the single-particle effect removed, as described in the main text.

With regard to [2], the present study is one of a number of important modifications to our understanding that have come up in the past few years. It was originally thought that the measured spatial distribution in the trap was consistent with the buildup of a cold population in a region that was insensitive to our spectroscopy. This was supported by the direct optical detection of this population and by the goodness of fits performed under the hypothesis. Regarding the fits, an important correction has been discovered and is discussed in [3], where a thorough investigation of hyperfine shifts and external field effects was performed for OH. This correction results in a 15% increase in the magnetic dipole moment used to interpret spectroscopy, and thus to the magnetic field at a particular microwave frequency. The data show a sudden suppression below 480 G, consistent with the hypothesis of a suppressed population only if this is the location of a certain avoided crossing, but with the magnetic dipole moment correction the crossing is actually located closer to 400 G. Without this, the fits used to calculate temperature under the suppression hypothesis are no longer trustworthy. Without the suppression hypothesis, the temperatures fit in figure 3 of [2] are not reliable, and this is the same conclusion supported by the enhanced magnitude of spin-flip losses we are reporting. At the temperatures fitted to those spectra, the spin-flip loss caused by the electric field used during evaporation are large enough to significantly influence the population, see the table in the main text. Nonetheless, without fitting any temperatures, it is still possible to use normalized spectra to look for enhancements in density caused by the evaporation. The normalization simply rescales the

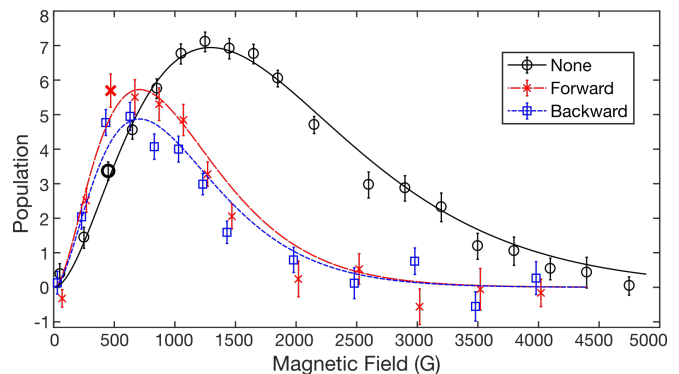


FIG. 2. Normalized spectra are performed under three conditions: one with no evaporation (black circles), one with an exponential RF evaporation ramp over 170 ms (red x's), and a third with this ramp applied in a time reversed manner (blue squares). Solid lines are fits to Maxwell-Boltzmann distributions with temperatures of 51 ± 2 in the case of no evaporation (black), and 30 ± 3 for both forward (red, long dashes) and backward (blue, short dashes) evaporations. The forward evaporation achieves a clear density enhancement in the vicinity of 500 G, where the line markers are bolded.

trace so that the area beneath adds up to the observed total population by laser induced fluorescence. Without it, spectroscopies of smaller populations yield distributions with enclosed areas that are decreased much more than the actual population. All of the evaporation traces in figure 3 of [2] show density enhancements in the region near 500 G, but a further correction changes this. By more carefully investigating the usage of microwave chirp rates and powers, we confirmed that what was previously interpreted as an offset of some kind was in fact an "oort cloud" persisting at larger fields [4]. After this, only figure 3c of [2] shows an enhancement, but it does so very clearly and repeatably as in fig. 2.

We have also developed a few more sensitive tools to look for collisional or evaporative effects, in addition to the microwave depletion during deceleration described and used in the main text. One is to compare the populations under two related conditions- the first a normal evaporation sequence and the second an evaporation with time-reversed microwave frequency. In other words, the cut goes backwards from deep to shallow. This comparison subjects all molecules to the same integrated microwave power, and thus the two conditions would be equivalent in a situation with only single particle effects. With evaporative effects, the normal condition ought to perform better. This is indeed what we consistently observe, at the 5% level, see fig. 2.

In conclusion, the roles of collisional effects in [1, 2] are reduced by spin-flip losses and other modifications to our understanding. Despite the effect of spin-flip losses and some other discoveries, spectroscopic comparisons and evaporation subtractions still point to an evaporative effect. The development of various more sensitive

tools has us poised to more unambiguously identify any future collisional effects in our next generation system described in the main text.

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 - [4] M. Inguscio, S. Stringari, C. E. Wieman, and S. italiana di Fisica., in *Varenna on Lake Como, Villa Monastero, 7-17 July 1998* (IOS Press, Amsterdam; Washington, DC, 1999).