

Aluminum cluster in the Fe-Al alloy to explain the anomalous behavior of the phase transition diagram.

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Abstract

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1. Introduction

The magnetic and structural properties of Iron aluminides has been widely studied during many years[1, 2], due its exhibit industrial applications properties such as oxidation and corrosion resistance, good ductility at room temperature, relatively low density, high magnetic permeability value, good vibration damping and insulation properties. Recent studies also have shown that alloys Fe-Al are promising candidates for data storage with magnetic patterns at room temperature [3, 4]. On the other hand, the magnetic properties of Fe-Al alloys are strongly linked with their chemical structure, the alloy can behave as ferromagnetic or paramagnetic at room temperature depending of composition and synthesis method [5, 6, 7, 8]. It is also possible to induce new types of critical behavior as superparamagnetism and reentrant spin glass phase [9, 10, 11, 12]. These characteristics have made this alloy one of the most studied binary systems in both experimental and theoretical research points of view.

The disordered Fe-Al alloys are arranged on body centered cubic (bcc) lattice structure, which may be represented by $\text{Fe}_{1-p}\text{Al}_p$, where p is the concentration of aluminum. In the range of $0 < p < 0.18$ the alloy is in a disordered ferromagnetic phase known as phase A2 [14, 15, 16], it is structurally disordered independent of the heat

treatment [2]. In this interval there is a small variation of the critical temperature as a function of the concentration q , and the temperatulanguage-latexre only decreases to $p > 0.1$ in contrast to the theoretical predictions, where for finite systems a finite inclination of the transition temperature at $q = 0$ [1]. For $p > 0.18$ obtaining the alloy in the disordered phase depends on the synthesis process. For $0.2 < p < 0.3$ the diagram shows a steep slope and vanishes in $p = 0.7$.

In order to try explain this anomalous behavior of phase transition diagram we propose an extended version of the site-diluted Ising model, following the same lines as the previous works on Fe-Al whose interaction of ferromagnetic exchange between Fe-Fe decreases linearly with the increase of p [? ? 17]. This is due to the expansion that is produced when Al atoms replace Fe atoms in a BCC network as a consequence of their larger atomic size. Also we consider in the model that Al atoms on small clusters can induce an extra super-exchange like ferromagnetic interaction between second-neighbors Fe atoms. Similar to the model used by Plascak and co-workers [17], differentiating only that the superexchange is constant and is not induced by Al atoms on large clusters.

We know from the advent of nanotechnology that the elements behave very differently when their dimensions are reduced to the point of being formed by a "countable" number of atoms [18]. Recently researchers have shown for example that the cluster of 102 gold atoms behaves as

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a giant molecule and not as metal. While clusters with 144 atoms exhibit a metallic behavior [19]. By analogy, this behavior depending only on the number of atoms in the cluster can be used to obtain the behavior of the phase diagram of the alloy of Fe-Al. For this we must visualize the system as being composed of different aluminum clusters dispersed in an iron lattice. In the proposed model we consider that some clusters of aluminum contribute to induce a superexchange interaction and others do not, depending of the clusters size.

In this work we report the calculated critical temperature e magnetics proprieties site bond diluted Ising model with superexchange interaction on BBC lattice. While the Ising model is not completely appropriate for the study of these alloys, there is some experimental evidence reinforcing the choice of such localized models instead of the itinerant ones, mainly due the constant average Fe magnetic moment as discoursed in Ref. [13].

2. Model and calculations

The system under investigation is given by a quenched site-diluted Ising model, Hamiltonian defined on a Body Centered Cubic lattice with N sites, that can be

$$\mathcal{H} = -J_1(p) \sum_{\langle NN \rangle} \epsilon_i \epsilon_j S_i S_j - J_2 \sum_{\langle NNN \rangle} \zeta_{ij} \epsilon_i \epsilon_j S_i S_j , \quad (1)$$

where the first sum is over nearest-neighbors (NN) pair and the second sum over next-nearest-neighbors (NNN) pairs, S_i represents a $1/2$ Ising spin. $J_1(p) = J(1-\alpha p)$ is the first-neighbor interaction, its a linear function of contraction p , and $J_2 = \beta J$ is the second-neighbor interaction, its a constant ferromagnetic interactions. In Eq.(1), ϵ_i are quenched, uncorrelated random variables, representing the existence of two kinds of particles in the system, namely the magnetic ones with $\epsilon_i = 1$, and non-magnetic ones with $\epsilon_i = 0$. The variable ϵ_i is chosen according to the probability

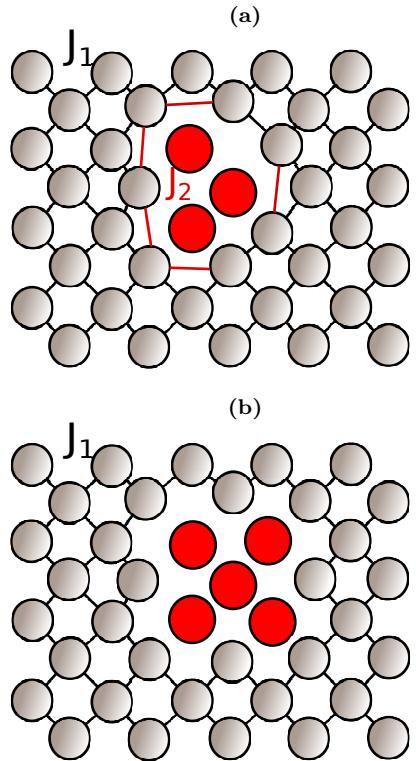


Figure 1: (Color online) A lattice sketch of the induced superexchange interaction J_2 by the cluster dopant (disconnected circles) for $\Gamma = 4$. In (a), tree dopant generates the NNN interaction J_2 . In (b), a cluster of five dopants does not induce J_2 .

distribution

$$P(\epsilon_i) = p\delta(\epsilon_i - 1) + (1-p)\delta(\epsilon_i), \quad (2)$$

where p is the concentration of no-magnetic sites, such that $p = 0$ corresponds to the pure case. In Eq.(1), $\zeta_{ij} = 1$ or 0 whether or not there is superexchange interaction. There are an Superexchange interaction between the next-nearest-neighbor pairs when an aluminum atom is present as NN , and if this atom is on a cluster with lesss equal of Γ aluminum atoms. In a quenched dilution, the increase in concentration p leads to the formation of Al clusters. Since the probability of finding a cluster with Γ Al atoms depends of p , the parameter Γ indirectly connects J_2 to the concentration p .

The above Hamiltonian has been studied through Monte Carlo simulations. First, we prepare a diluted sample, i.e., to a given value of p , the algorithm builds an random matrix $\{\epsilon\}$, another algorithm sweeps the $\{\epsilon\}$ measuring the Al

cluster size and using Γ parameter to builds the $\{\zeta\}$ matrix. In figure 1 a situation is represented for $\Gamma = 4$, for cluster with lessr than 5 impurities, Al atom induces superexchange interactions, otherwise it does not.

in this model, for every thermodynamic observable Q , we first calculate the thermal average $\langle Q_{\{\epsilon\}} \rangle$ for a given sample $\{\epsilon\}$ and the results for different samples are later average as $[\langle Q_{\{\epsilon\}} \rangle]_{\text{av}}$. In order to get the critical properties of the present model, for each sample of a given site-disorder configuration p , we used a hybrid Monte Carlo algorithm consisting of one single spin Metropolis algorithm [20], and one cluster wolff algorithm [21]. This hybrid Monte Carlo method [22, 23] has been implemented, and has been shown to reduce correlations between successive configurations in the simulation. Close to the transition temperature we have also resorted to single histogram techniques to get the corresponding thermodynamic quantities. We have first computed the magnetization, magnetic susceptibility, and the Binder cumulant given, respectively, by

$$m = \frac{1}{N} \sum_{i=1}^N S_i, \quad (3)$$

$$\chi = N \frac{\langle m^2 \rangle - \langle m \rangle^2}{T}, \quad (4)$$

$$u_4 = 1 - \frac{\langle m^4 \rangle}{3\langle m^2 \rangle^2}, \quad (5)$$

where $N = 2L^3$ and L is the linear size of the cubic lattice studied, T is the temperature given in units of J/k_B , k_B being the Boltzmann constant. We have also considered $J = 1$, in such a way that J_2 is measured in units of J . In addition, to reach the final results, for each dilution, temperature, and lattice size, the MC estimates $\langle Q_{\{\epsilon\}} \rangle$ of thermodynamic quantities, for a given random sample $\{\epsilon\}$ of diluted sites were averaged over different disorder realizations as

$$[\langle Q_{\{\epsilon\}} \rangle]_{\text{av}} = \frac{1}{\#\{\epsilon\}} \sum_{\{\epsilon\}} \langle Q_{\{\epsilon\}} \rangle, \quad (6)$$

with $\#\{\epsilon\}$ the number of total realizations considered. Now, regarding the simulational numbers,

for every sample the runs comprised 3×10^4 MCS per spin for equilibration and the measurements were made over 5×10^5 MCS. The lattice sizes raged from $L = 10, 20, 30$ the values being chosen so that $p \times L^3$ gives an integer number. We have used in the present paper 50 samples for all settings.

3. Results

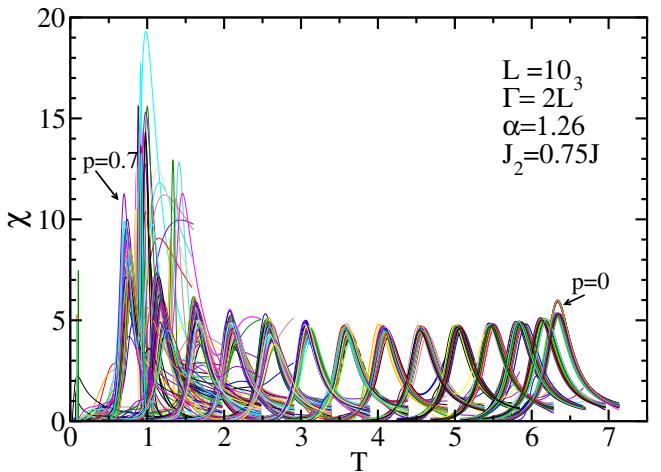


Figure 2: (color online) Magnetic susceptibility as a function of the temperature for different samples. In this example, to $L = 10$, $\Gamma = N = 2000$, $\alpha = 1.26$, $\beta = 0.75$. For a question of clarity, the Monte Carlo data have been omitted and only the lines are shown in this figure.

In Fig. 2 it is shown the behavior of χ as a function of temperature for fifty samples, for concentration ranging from $p = 0$ to $p = 0.7$ with steps of 0.05, lattce size $L = 10$, $\alpha = 1.26$, $\beta = 0.75$, and $\Gamma = N = 2000$, i.e., all dopants induce superexchange interaction J_2 . The temperature T here is measured in units of the Boltzmann constant. These curves were obtained using the simple histogram technique. For each sample, the system presented a temperature and susceptibility value that was slightly different. We can observe that for values of $p = 0$ up to $p = 0.5$ the susceptibility curves correspond to the same peak height. For $p > 0.5$ the peaks increase and a greater variation between the samples. The peak of each of these samples is used to estimate the transition temperature $T_c(p)$.

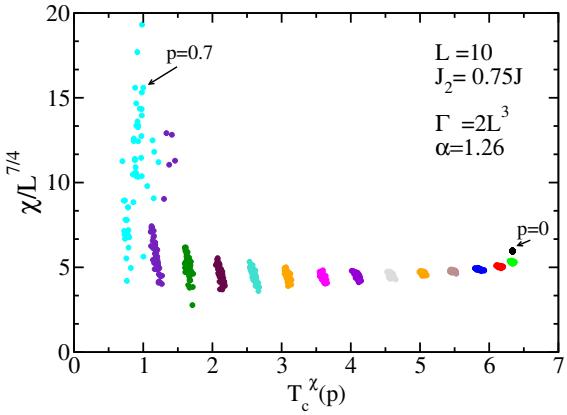


Figure 3: (Color online) Susceptibilidade vs critical temperature $T_c(p)$ for fifty samples and several values of p . With $\Gamma = 2L^3$, $\alpha = 1.26$ and $\beta = 0.75$ The simulational errors are smaller than the symbol sizes.

In Fig. 3 it is shown the peak of each of the simulated samples. We can observe that the points become more dispersed as the concentration of dopants increases. Taking the mean of the T_c of each sample and obtaining the temperature of the collection for a given configuration.

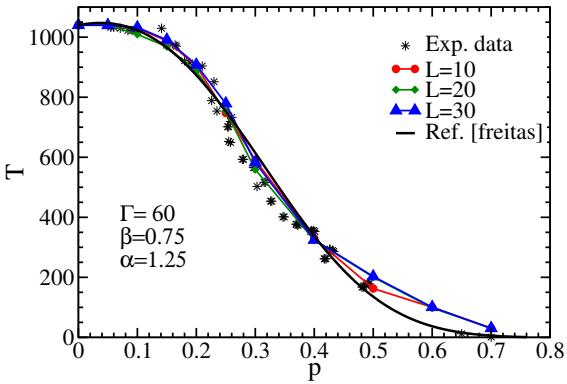


Figure 4: (Color online) Phase diagram in the Tq plane. Solid circles are the results of present model. Solid line is the previous fit using the Effective field theory Ref [24], star points are the experimental data. ??

Fig. 4 exhibits the transition temperature $T_c(p)$ as a function of p to the best result obtained with the adjustment of the α , β and Γ parameters in comparison with the experimental results and comparison with the results of Ref [24].

4. Conclusion

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