

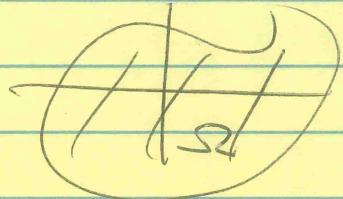
(1) 1D+2D

LECTURE 16

GREEN'S FUNCTIONS FOR ~~1D+2D~~ BVPs

We really want to solve for $u = u(x,y)$

$$\begin{cases} -Du = f & \text{in } \Omega \\ u = g & \text{on } \partial\Omega \end{cases}$$



For now this is too hard

So instead work in 1D on $[a,b]$, and solve ODE BVP

$$\begin{cases} -u'' = f & a < x < b \\ u(0) = 0 \\ u(1) = 0 \end{cases} \quad \Rightarrow \boxed{\text{LIKE } F=ma.}$$

FIRST SOLVE IN SPECIAL CASE $f = \delta_\xi$ AS UNIT IMPULSIVE FORCE

We call solution to

$$\begin{cases} -u'' = \delta_\xi & \text{on } [0,1] \text{ FOR } 0 < \xi < 1 \\ u(0) = 0 = u(1) \end{cases}$$

The Green's Function $u(\xi) = G_\xi(x) = f(\xi; \xi)$

$$\text{So } \begin{cases} -G_\xi'' = \delta_\xi \\ G_\xi(0) = 0 = G_\xi(1) \end{cases}$$

(2)

We saw before that $f_{\xi}'' = \delta_{\xi}$

where $f_{\xi}(x) = \begin{cases} 0 & \text{if } x < \xi \\ x-\xi & \text{if } x > \xi \end{cases}$ as RAMP FN.

Problem If $0 < \xi < 1$ Then

$$f_{\xi}(0) = 0 \quad \checkmark$$

$$\text{but } f_{\xi}(1) = 1 - \xi \neq 0.$$

So f_{ξ} does not satisfy BCs.

Solution Write $\ell_{\xi} = -f_{\xi} + z_{\xi} \quad z_{\xi} = z_{\xi}(x).$

and work out what ODE + VP z_{ξ} has to solve.

Well

$$z_{\xi} = \ell_{\xi} + f_{\xi}$$

$$-z_{\xi}'' = -\ell_{\xi}'' + f_{\xi}'' = \delta_{\xi} + -\delta_{\xi} = 0$$

and

$$z_{\xi}(0) = \ell_{\xi}(0) + f_{\xi}(0) = 0$$

$$z_{\xi}(1) = \ell_{\xi}(1) + f_{\xi}(1) = 0 + 1 - \xi = 1 - \xi$$

(3)

SOLUTION

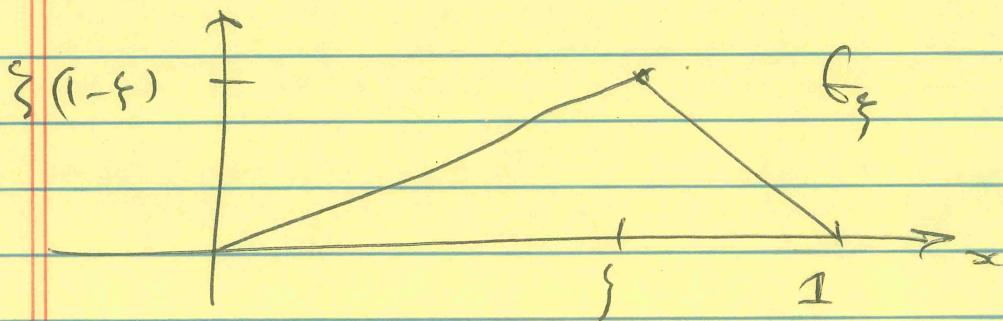
$$z_g'(x) = 0 \Rightarrow z_g(x) = ax + b.$$

$$0 = z_g(0) = b \Rightarrow z_g(x) = ax$$

$$1 - g = z_g(1) = a \Rightarrow z_g(x) = (1-g)x$$

So

$$G_g(x) = (1-g)x - p_g(x) = \begin{cases} (1-g)x & \text{if } x < \xi \\ g(1-x) & \text{if } x > \xi \end{cases}$$



SECOND Use superposition to solve for general forcing function f .

$$\begin{aligned} \text{INTUTION} \quad f(x) &= \int_x^\infty [f] = \int_0^1 \int_x^\infty g(\xi) f(\xi) d\xi \\ &= \int_0^1 \int_{x-\xi}^0 f(\xi) d\xi \\ &= \int_0^1 \int_{\xi-x}^\xi f(\xi) d\xi \end{aligned}$$

$$f(x) = \int_0^1 S_g(x) f(\xi) d\xi$$

(4)

Since

$$-\xi''(\alpha) = \delta_\xi(\alpha)$$

$$\xi(0) = 0 = \xi(1)$$

this suggests we define

$$u(\alpha) = \int_0^1 \xi(\alpha) f(s) ds$$

$$u(\alpha) = \int_0^1 \xi(\alpha; s) f(s) ds$$

INFORMALLY

$$-u''(\alpha) = -\int_0^1 \xi''(\alpha) f(s) ds$$

$$= \int_0^1 \delta_\xi(\alpha) f(s) ds = f(\alpha)$$

and $u(0) = \int_0^1 \xi(0) f(s) ds = \int_0^1 0 ds \rightarrow$
 $u(1) = 0$ ~~too~~ ✓

FORMALLY from formula for δ_ξ :

$$u(\alpha) = \int_0^\alpha \xi(1-s) f(s) ds + \int_\alpha^1 (1-s)\alpha f(s) ds$$

$$u(\alpha) = (1-\alpha) \int_0^\alpha \xi(s) f(s) ds + \alpha \int_\alpha^1 (1-s) f(s) ds$$

$$or u(\alpha) = (-\alpha) \int_0^\alpha \xi(s) f(s) ds - \alpha \int_1^\infty (1-s) f(s) ds \quad \textcircled{4}$$

(5)

CHECK $u(0) = 0 = u(1) \checkmark$

And

By FTC on \star

$$u'(x) = - \int_0^x f(s) ds + (L_x)_x f(x)$$

$$- \int_1^x (-s) f(s) ds - x(L_0) f(x)$$

$$u'(x) = - \int_0^x sf(s) ds - \int_1^x (-s) f(s) ds$$

So

$$u''(x) = - xf(x) = (1-x) f(x) = -f(x)$$

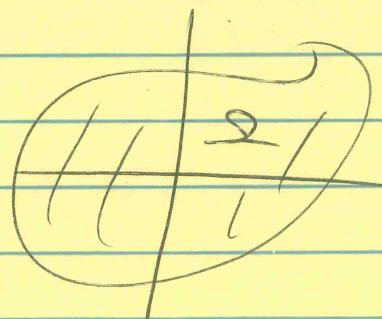
or $-u'' = f \checkmark$

 0

2 D CASE

FOR 2 SOLVE

$$\begin{cases} -\Delta u = f \text{ in } \Omega \\ u = h \text{ on } \partial\Omega \end{cases}$$



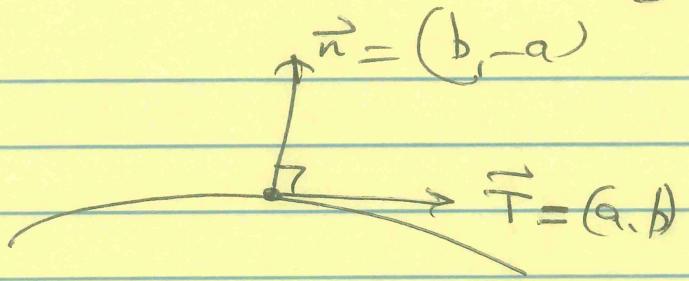
PRELIMINARIES: GREEN'S IDENTITIES (ANALOGUE OF INTEGR" BY PARTS)

GREEN'S THM Let $P = P(x, y)$, $Q = Q(x, y)$

$$\iint_{\Omega} \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dA = \int_{\partial\Omega} P dx + Q dy.$$

(6)

SET $\vec{F} = Q\vec{i} - P\vec{j}$.



Then

$$\nabla \cdot \vec{F} = \frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y}$$

and

$$\begin{aligned} \int_{\partial\Omega} Px + Qdy &= \int_{\partial\Omega} (P, Q) \cdot \vec{T} ds \\ &= \int_{\partial\Omega} (aP + bQ) ds \\ &= \int_{\partial\Omega} (\vec{F} \cdot \vec{n}) ds \end{aligned}$$

So get

[DIVERGENCE FORM OF GREEN'S THM]

$$\boxed{\int_{\Omega} \nabla \cdot \vec{F} dA = \int_{\partial\Omega} (\vec{F} \cdot \vec{n}) ds} \quad (1)$$

SET $\vec{F} = u\vec{v}$

$$u = u(x), \quad \vec{v} = \vec{v}(x)$$

FUNCTION VECTOR FIELD

Then

$$\nabla \cdot (u\vec{v}) = \nabla u \cdot \vec{v} + u \nabla \cdot \vec{v}$$

PRODUCT RULE (2)

PLUG into (1) :

$$\iint_{\Omega} (\nabla u \cdot \vec{v} + u \nabla \cdot \vec{v}) dA = \int_{\partial\Omega} u (\vec{v} \cdot \vec{n}) ds$$

INTEGRATION BY PARTS.

(7)

Set $\vec{v} = \nabla v$ where $v = v(x)$ is a function

and notice $\nabla \cdot \vec{v} = \nabla \cdot \nabla v = \Delta v$.

So get

$$\iint_{\Omega} (u \Delta v + \nabla u \cdot \nabla v) dA = \int_{\partial \Omega} u \frac{\partial v}{\partial n} ds \quad (3)$$

SWITCH u, v in (3) to get (3'). Subtract (3), (3')

to get

$$\boxed{\iint_{\Omega} (u \Delta v - v \Delta u) dA = \int_{\partial \Omega} (u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n}) ds} \quad (4)$$

- Green's 2nd Identity.

FUNDAMENTAL LAW TO LAPLACE'S EQUATION

Thm Let $F_{\vec{\xi}}(\vec{x}) = -\frac{1}{2\pi} \log |\vec{x} - \vec{\xi}|$

where $\vec{x} = (x, y)$, $\vec{\xi} = (\xi, \eta)$

Then

$$\boxed{-\Delta F_{\vec{\xi}} = f_{\vec{\xi}}} \quad$$

Antidote of
RAMP function

8

PF 10et

In polar coords (thinking \vec{z} is at origin)

$$\mathbf{F} = -\frac{1}{2\pi} \log(r) \hat{z}$$

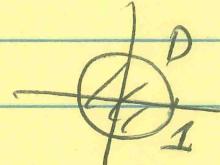
We saw by that

$$\Delta \log(r) = 0 \quad \text{for } r \neq 0.$$

So we guess $\exists C$:

$$+ \Delta \log(r) = C \oint_{\gamma}$$

By Property of \oint_{γ} :



$$C = \iint_D C \oint_{\gamma} dA$$

$$C = + \iint_D \Delta(\log(r)) dA$$

Apply Green's 2nd Identity ④ with
 $u = +1, v = \log(r)$ to get

$$C = + \iint_D \Delta(\log(r)) dA$$

$$= \iint_D (+1) \cdot \frac{\partial}{\partial n} (\log(r)) ds \quad \text{as } \Delta(+1) = 0$$

$$= + \iint_D \frac{\partial}{\partial r} (\log(r)) ds \quad \frac{\partial}{\partial n} (+1) = 0.$$

(9)

$$= + \int_{\partial D} \frac{1}{r} ds = + \int_{\partial D} \frac{1}{r} ds = 2\pi$$

$$\text{So } \Delta \left(\frac{1}{2\pi} \log r \right) = S_0.$$

Hence

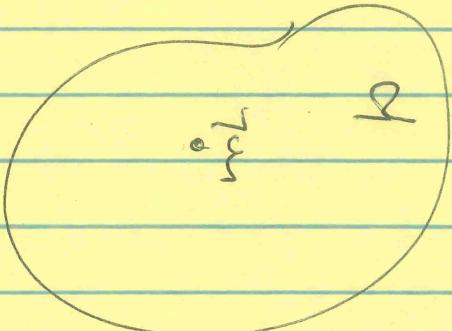
$$-\Delta \vec{F}_{\xi} = \vec{S}_{\xi}. \quad \checkmark$$

—————

Next fix $\vec{\xi} \in \Omega$.

SOLVE

$$\begin{cases} -\Delta u = \vec{S}_{\xi} & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega \end{cases}$$

CALL THE SOLN $u(\vec{x}) = \vec{G}_{\xi}(\vec{x}) = \vec{F}(\vec{x}; \vec{\xi})$.

As before we must adjust Fundamental Soln to get correct boundary values:

$$\text{Write } \vec{G}_{\xi} = \vec{F}_{\xi} + \vec{z}_{\xi}.$$

$$\text{Then } \Delta \vec{z}_{\xi} = \Delta \vec{F}_{\xi} - \Delta \vec{F}_{\xi}^* = -(\vec{F}_{\xi} - \vec{F}_{\xi}^*) = 0$$

$$\text{and on } \partial\Omega \quad \vec{z}_{\xi} = 0 - \vec{F}_{\xi}$$

(10)

SUPPOSE we can solve

$$\left\{ \begin{array}{l} \Delta z_{\xi} = 0 \quad \text{in } \Omega \\ z_{\xi}(\vec{x}) = \frac{1}{2\pi} \log |\vec{x} - \vec{\xi}| \quad \text{on } \partial \Omega \end{array} \right.$$
(5)

Then

$$G_{\xi} = F_{\xi} + z_{\xi} \quad \text{is our Green's F^n for } \Omega.$$

Finally to solve

$$\left\{ \begin{array}{l} -\Delta u = f \quad \text{in } \Omega \\ u = h \quad \text{on } \partial \Omega \end{array} \right.$$

Set $v = G_{\xi}$, $u = u$ in ④ to get

$$\iint_{\Omega} (u \Delta G_{\xi} - G_{\xi} \Delta u) dA = \int_{\partial \Omega} \left(u \frac{\partial G_{\xi}}{\partial \vec{n}} - G_{\xi} \frac{\partial u}{\partial \vec{n}} \right) ds$$

Since $\Delta G_{\xi} = -f_{\xi}$, $f_{\xi} = 0$ on $\partial \Omega$
 $\Delta u = f$ in Ω , $u = h$ on $\partial \Omega$

we get

~~$$u(\vec{x}) = \iint_{\Omega} S_{\xi}(\vec{x}) u(\vec{\xi}) dA = - \iint_{\Omega} u \Delta G_{\xi} dA$$~~

(11)

Suppose we can solve

$$\left\{ \begin{array}{l} \Delta \vec{z}_\xi = 0 \quad \text{in } \Omega \\ \vec{z}_\xi(\vec{z}) = \frac{1}{2\pi} \log |\vec{z} - \vec{\xi}| \quad \text{on } \partial\Omega \end{array} \right.$$

$\vec{\xi}$ on $\partial\Omega$ at $\vec{z} \neq \vec{\xi}$

* for ex If Ω = RECTANGLE or DISC can do this using Fourier Series.

Then Green's Function for D on Ω is

$$\vec{f}_\xi = \vec{f}_\xi + \vec{z}_\xi$$

$$G(\vec{z}; \vec{\xi}) = f(\vec{z}; \vec{\xi}) + z(\vec{z}; \vec{\xi})$$

Symmetry By construction

$$G(\vec{z}; \vec{\xi}) = G(\vec{\xi}; \vec{z})$$

So

$$\Delta_{\vec{\xi}} G(\vec{z}; \vec{\xi}) = \Delta_{\vec{\xi}} G(\vec{\xi}; \vec{z})$$

$$= \Delta_{\vec{z}} G(\vec{z}; \vec{\xi})$$

$$= -\delta_{\vec{\xi}}(\vec{z}) = -\delta(\vec{z} - \vec{\xi}) = -\delta_{\vec{z}}(\vec{\xi})$$

(12)

So

$$u(\vec{x}) = \iint_{\Omega} u(\vec{\xi}) S_{\vec{x}}(\vec{\xi}) dA(\xi)$$

$$= - \iint_{\Omega} u(\vec{\xi}) D_{\vec{\xi}} G(\vec{x}; \vec{\xi}) dA(\xi)$$

$$\stackrel{(4)}{=} - \iint_{\Omega} G(\vec{x}; \vec{\xi}) D^f u dA$$

$$- \int_{\partial\Omega} \left(u \frac{\partial G}{\partial n_{\xi}} - G \frac{\partial u}{\partial n_{\xi}} \right) ds$$

So Gauss's REPN THM

$$u(\vec{x}) = \iint_{\Omega} f(\vec{x}; \vec{\xi}) f(\vec{\xi}) dA - \int_{\partial\Omega} \frac{\partial G(\vec{x}; \vec{\xi})}{\partial n_{\xi}} h(\vec{\xi}) ds$$

Solves $\begin{cases} -\Delta u = f & \text{on } \Omega \\ u = h & \text{on } \partial\Omega \end{cases}$

IN SPECIAL CASE $\Omega = \text{DISC}$, $f = 0$
 recover POINCARÉ'S FORMULA.