



Rensselaer

## Decreasing Energy Aliasing in POD-ROMs

David Wells

Rensselaer Polytechnic Institute

In collaboration with:

T. Iliescu (VT), V. John (WIAS), S. Giere (WIAS)

June 28, 2017

Conference on Classical and Geophysical Fluid Dynamics

# Outline

- Governing Equations
- A Little on Streamline-Upwind Petrov-Galerkin (SUPG)
- The SUPG-ROM
- Numerical Results
- Summary

## Convection-Diffusion-Reaction Equation

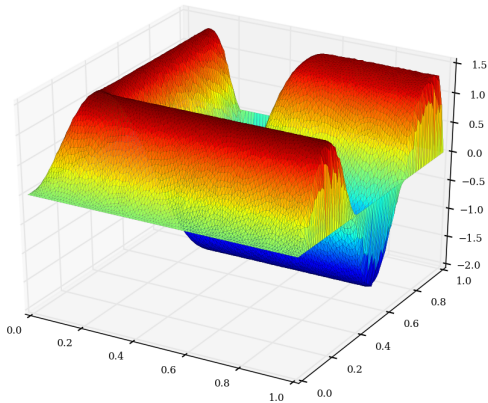
There are a few relevant model problems:

$$u_t - \varepsilon \Delta u + \vec{b} \cdot \nabla u + cu = f \quad (1)$$

$$-\varepsilon \Delta u + \vec{b} \cdot \nabla u + cu = f \quad (2)$$

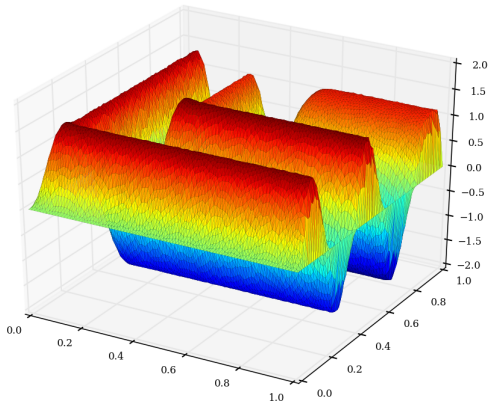
- $\varepsilon \ll 1$  and  $\varepsilon \ll h$
- Small diffusion implies wavelike behavior (*convection-dominated*): the relevant length scale is  $O(\sqrt{\varepsilon})$ , which may be too small to resolve.

## Convection-Diffusion-Reaction Equation



Third POD vector.

## Convection-Diffusion-Reaction Equation



Fifth POD vector.

## Stabilization by Streamline Upwinding

$$\varepsilon u_{xx} + u_x = 0, u(0) = 0, u(1) = 1 \quad (3)$$

$$\varepsilon \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1} - U_{i-1}}{2h} = 1 \quad (4)$$

known to be an unstable discretization for small  $\varepsilon$

$$\varepsilon \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1} - U_i}{h} = 1 \quad (5)$$

$$\varepsilon \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1}}{2h} + \frac{U_{i+1}}{2h} - \frac{2U_i}{2h} + \frac{1}{2h} (U_{i-1} - U_{i-1}) = 1 \quad (6)$$

$$\left( \varepsilon + \frac{h}{2} \right) \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1} - U_{i-1}}{2h} = 1 \quad (7)$$

Using a forward difference (*upwinding*) results in an order 1 method with a much larger diffusion length scale ( $O(\sqrt{h}) \gg O(\sqrt{\varepsilon})$ ).

## Stabilization by Streamline Upwinding

$$\varepsilon u_{xx} + u_x = 0, u(0) = 0, u(1) = 1 \quad (3)$$

$$\varepsilon \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1} - U_{i-1}}{2h} = 1 \quad (4)$$

known to be an unstable discretization for small  $\varepsilon$

$$\varepsilon \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1} - U_i}{h} = 1 \quad (5)$$

$$\varepsilon \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1}}{2h} + \frac{U_{i+1}}{2h} - \frac{2U_i}{2h} + \frac{1}{2h} (U_{i-1} - U_{i-1}) = 1 \quad (6)$$

$$\left( \varepsilon + \frac{h}{2} \right) \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1} - U_{i-1}}{2h} = 1 \quad (7)$$

Using a forward difference (*upwinding*) results in an order 1 method with a much larger diffusion length scale ( $O(\sqrt{h}) \gg O(\sqrt{\varepsilon})$ ).

## Stabilization by Streamline Upwinding

$$\varepsilon u_{xx} + u_x = 0, u(0) = 0, u(1) = 1 \quad (3)$$

$$\varepsilon \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1} - U_{i-1}}{2h} = 1 \quad (4)$$

known to be an unstable discretization for small  $\varepsilon$

$$\varepsilon \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1} - U_i}{h} = 1 \quad (5)$$

$$\varepsilon \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1}}{2h} + \frac{U_{i+1}}{2h} - \frac{2U_i}{2h} + \frac{1}{2h} (U_{i-1} - U_{i-1}) = 1 \quad (6)$$

$$\left( \varepsilon + \frac{h}{2} \right) \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1} - U_{i-1}}{2h} = 1 \quad (7)$$

Using a forward difference (*upwinding*) results in an order 1 method with a much larger diffusion length scale ( $O(\sqrt{h}) \gg O(\sqrt{\varepsilon})$ ).



## Stabilization by Streamline Upwinding

$$\varepsilon u_{xx} + u_x = 0, u(0) = 0, u(1) = 1 \quad (3)$$

$$\varepsilon \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1} - U_{i-1}}{2h} = 1 \quad (4)$$

known to be an unstable discretization for small  $\varepsilon$

$$\varepsilon \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1} - U_i}{h} = 1 \quad (5)$$

$$\varepsilon \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1}}{2h} + \frac{U_{i+1}}{2h} - \frac{2U_i}{2h} + \frac{1}{2h} (U_{i-1} - U_{i-1}) = 1 \quad (6)$$

$$\left( \varepsilon + \frac{h}{2} \right) \frac{U_{i+1} - 2U_i + U_{i-1}}{h^2} + \frac{U_{i+1} - U_{i-1}}{2h} = 1 \quad (7)$$

Using a forward difference (*upwinding*) results in an order 1 method with a much larger diffusion length scale ( $O(\sqrt{h}) \gg O(\sqrt{\varepsilon})$ ).

## Stabilization by Streamline Upwinding (Petrov-Galerkin)

We can get a similar result for Finite Elements with the Petrov-Galerkin method by choosing test functions  $\varphi + \delta\varphi_x$ :

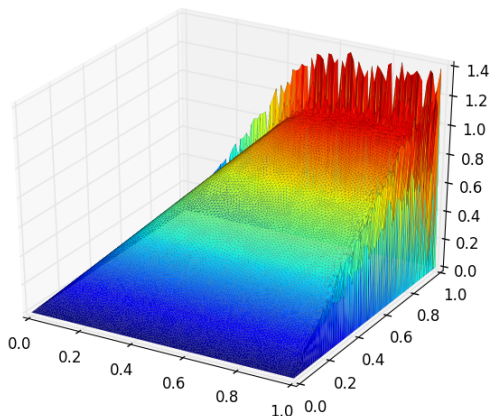
$$\varepsilon((\varphi + \delta\varphi_x)_x, u_x) + (\varphi + \delta\varphi_x, u_x) = (\varphi + \delta\varphi_x, f) \quad (8)$$

$$+(\varepsilon + \delta)(\varphi_x, u_x) + \delta(\varphi_{xx}, u_x) + (\varphi, u_x) = (\varphi + \delta\varphi_x, f) \quad (9)$$

$$(10)$$

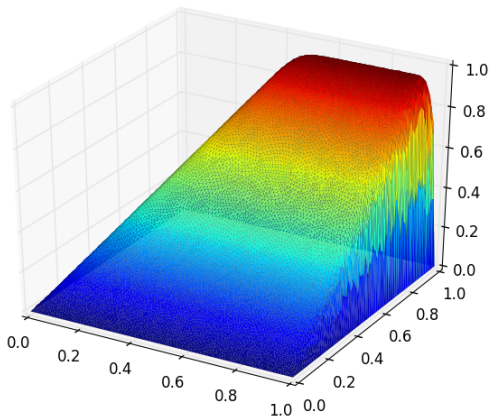
so the new viscous length scale is  $O(\sqrt{\delta})$ : if  $\delta = O(h)$  then this can be resolved.

## Stabilization by SUPG



$f = 1$ ,  $\varepsilon = 10^{-3}$ ,  $\vec{b} = (\cos(\pi/3), \sin(\pi/3))$ , and  $c = 10^{-3}$ : no stabilization.

## Stabilization by SUPG



$f = 1$ ,  $\varepsilon = 10^{-3}$ ,  $\vec{b} = (\cos(\pi/3), \sin(\pi/3))$ , and  $c = 10^{-3}$ , with SUPG,  $\delta = h$ .

## The SUPG-ROM

- SUPG helps us avoid instabilities arising from unresolved modes
- Unresolved modes  $\rightarrow$  aliasing onto the resolved parts (wiggles)
- ROMs, by definition, have lots of unresolved modes!

*Plan:* add SUPG to a POD-ROM for the convection-diffusion-reaction equation, with analysis done with both FEM and POD estimates.

## The SUPG-ROM

- SUPG helps us avoid instabilities arising from unresolved modes
- Unresolved modes  $\rightarrow$  aliasing onto the resolved parts (wiggles)
- ROMs, by definition, have lots of unresolved modes!

*Plan:* add SUPG to a POD-ROM for the convection-diffusion-reaction equation, with analysis done with both FEM and POD estimates.

## The SUPG-ROM

- SUPG helps us avoid instabilities arising from unresolved modes
- Unresolved modes  $\rightarrow$  aliasing onto the resolved parts (wiggles)
- ROMs, by definition, have lots of unresolved modes!

*Plan:* add SUPG to a POD-ROM for the convection-diffusion-reaction equation, with analysis done with both FEM and POD estimates.

## The SUPG-ROM

- SUPG helps us avoid instabilities arising from unresolved modes
- Unresolved modes  $\rightarrow$  aliasing onto the resolved parts (wiggles)
- ROMs, by definition, have lots of unresolved modes!

*Plan:* add SUPG to a POD-ROM for the convection-diffusion-reaction equation, with analysis done with both FEM and POD estimates.



## Error estimates

We assume a centered trajectory  $u_s$ , ROM approximating the fluctuations  $u_r$ , and projection onto the POD space  $P_r$ :

$$\begin{aligned} u - (u_s + u_r) &= (u - P_r(u_h)) + (P_r(u_h) - (u_s + u_r)) \\ &:= \eta - \phi_r \end{aligned} \tag{11}$$

Error estimate of the form (Giere, 2015 [1])

$$\begin{aligned} |||\phi_r|||_{\text{SUPG},r} \leq C &\left[ \left(1 + \frac{1}{\sqrt{\delta}} + \sqrt{\delta}\right) \|\eta\| + \sqrt{\varepsilon} \|\nabla \eta\| \right. \\ &\left. + \sqrt{\delta} \|\vec{b} \cdot \nabla \eta\| + \sqrt{\delta} \sqrt{\sum_T \|\Delta \eta\|_{0,T}^2} \right]. \end{aligned} \tag{12}$$

## Error estimates

We assume a centered trajectory  $u_s$ , ROM approximating the fluctuations  $u_r$ , and projection onto the POD space  $P_r$ :

$$\begin{aligned} u - (u_s + u_r) &= (u - P_r(u_h)) + (P_r(u_h) - (u_s + u_r)) \\ &:= \eta - \phi_r \end{aligned} \tag{11}$$

Error estimate of the form (Giere, 2015 [1])

$$\begin{aligned} |||\phi_r|||_{\text{SUPG},r} &\leq C \left[ \left( 1 + \frac{1}{\sqrt{\delta}} + \sqrt{\delta} \right) \|\eta\| + \sqrt{\varepsilon} \|\nabla \eta\| \right. \\ &\quad \left. + \sqrt{\delta} \|\vec{b} \cdot \nabla \eta\| + \sqrt{\delta} \sqrt{\sum_T \|\Delta \eta\|_{0,T}^2} \right]. \end{aligned} \tag{12}$$

## The POD Option

Seminorm bound on derivatives: [2, 4]

$$\sum_{n=0}^{R-1} \left\| u_n - \sum_{j=0}^{r-1} \langle u_n, \phi_j \rangle \phi_j \right\|_s^2 = \sum_{j=r}^{R-1} |\phi_j|_s^2 \sigma_j^2. \quad (13)$$

$$\|\eta\| \leq C_1 \sqrt{\sum_{j=r}^{R-1} \sigma_j^2} := C_1 \Lambda_0 \quad (14)$$

$$\|\nabla \eta\| \leq C_1 \sqrt{\sum_{j=r}^{R-1} \sigma_j^2 |\varphi_j|_1} := C_1 \Lambda_1 \quad (15)$$

$$\sqrt{\sum_{T \in \mathcal{T}_h} \|\delta \eta\|_T^2} \leq C_1 \sqrt{\sum_{j=r}^{R-1} \sigma_j^2 |\varphi_j|_2} := C_1 \Lambda_2. \quad (16)$$

$$\delta = \frac{\Lambda_0}{\Lambda_2 \varepsilon + \Lambda_0 + \Lambda_1}. \quad (17)$$

## The POD Option

Seminorm bound on derivatives: [2, 4]

$$\sum_{n=0}^{R-1} \left\| u_n - \sum_{j=0}^{r-1} \langle u_n, \phi_j \rangle \phi_j \right\|_s^2 = \sum_{j=r}^{R-1} |\phi_j|_s^2 \sigma_j^2. \quad (13)$$

$$\|\eta\| \leq C_1 \sqrt{\sum_{j=r}^{R-1} \sigma_j^2} := C_1 \Lambda_0 \quad (14)$$

$$\|\nabla \eta\| \leq C_1 \sqrt{\sum_{j=r}^{R-1} \sigma_j^2 |\varphi_j|_1} := C_1 \Lambda_1 \quad (15)$$

$$\sqrt{\sum_{T \in \mathcal{T}_h} \|\delta \eta\|_T^2} \leq C_1 \sqrt{\sum_{j=r}^{R-1} \sigma_j^2 |\varphi_j|_2} := C_1 \Lambda_2. \quad (16)$$

$$\delta = \frac{\Lambda_0}{\Lambda_2 \varepsilon + \Lambda_0 + \Lambda_1}. \quad (17)$$

## The Finite Element Option

Another option is to use finite element estimates [3] with order  $m$  polynomials and  $s$  derivatives:

$$|\eta|_s \leq Ch^{m+1/2-s} \quad (18)$$

Solve the optimization problem for  $\delta$ :

$$\delta = \frac{\Lambda_0 h^2 + h^{m+\frac{5}{2}}}{(h^2 + \varepsilon + h)\Lambda_0 + \varepsilon h^{m+\frac{1}{2}} + h^{m+\frac{5}{2}} + h^{m+\frac{3}{2}}} \quad (19)$$

If  $m \geq 2$ :

$$\delta \approx \frac{\Lambda_0 h^2}{h\Lambda_0} = h \quad (20)$$

We are back where we started! In practice the two stabilization parameters are within about 10%.

## The Finite Element Option

Another option is to use finite element estimates [3] with order  $m$  polynomials and  $s$  derivatives:

$$|\eta|_s \leq Ch^{m+1/2-s} \quad (18)$$

Solve the optimization problem for  $\delta$ :

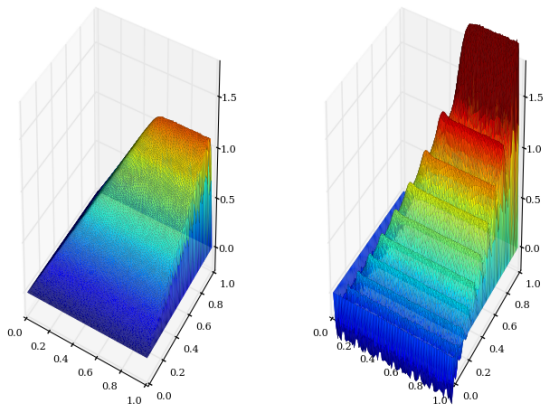
$$\delta = \frac{\Lambda_0 h^2 + h^{m+\frac{5}{2}}}{(h^2 + \varepsilon + h)\Lambda_0 + \varepsilon h^{m+\frac{1}{2}} + h^{m+\frac{5}{2}} + h^{m+\frac{3}{2}}} \quad (19)$$

If  $m \geq 2$ :

$$\delta \approx \frac{\Lambda_0 h^2}{h\Lambda_0} = h \quad (20)$$

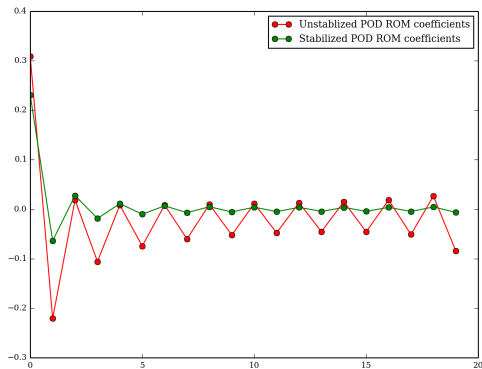
We are back where we started! In practice the two stabilization parameters are within about 10%.

## POD-ROM and SUPG-ROM



Stabilized (left) and unstabilized (right) POD-ROMs at  $t = 1.0$ .

## Comparison of POD coefficients



Coefficients at  $t = 1.0$  for the stabilized and unstabilized models.



## Summary

- ROMs for convection-dominated problems have aliasing issues when scales are unresolved.
- We can overcome this problem with SUPG; tentatively, this implies that the usual FEM stabilization methods may be useful.

Thank You!



S. Giere, T. Iliescu, V. John, and D. Wells.

SUPG reduced order models for convection-dominated convection-diffusion-reaction equations.  
*Comput. Methods Appl. Mech. Engrg.*, 289:454–474, 2015.



T. Iliescu and Z. Wang.

Variational multiscale proper orthogonal decomposition: Convection-dominated convection-diffusion-reaction equations.  
*Math. Comput.*, 82(283):1357–1378, 2013.



H. G. Roos, M. Stynes, and L. Tobiska.

*Robust Numerical Methods for Singularly Perturbed Differential Equations: Convection-Diffusion-Reaction and Flow Problems.*,  
volume 24 of *Springer Series in Computational Mathematics*.  
Springer, second edition, 2008.



J. R. Singler.

New POD error expressions, error bounds, and asymptotic results for reduced order models of parabolic PDEs.  
*SIAM J. Numer. Anal.*, 52(2):852–876, 2014.