

A CLASSIC THESIS STYLE

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An Homage to The Elements of Typographic Style

September 2015 – version 4.2

ABSTRACT

Short summary of the contents... a great guide by Kent Beck how to write good abstracts can be found here:

<https://plg.uwaterloo.ca/~migod/research/beck00PSLA.html>

*We have seen that computer programming is an art,
because it applies accumulated knowledge to the world,
because it requires skill and ingenuity, and especially
because it produces objects of beauty.*

— knuth:1974 [knuth:1974]

ACKNOWLEDGEMENTS

Put your acknowledgements here.

Many thanks to everybody who already sent me a postcard!

Regarding the typography and other help, many thanks go to Marco Kuhlmann, Philipp Lehman, Lothar Schlesier, Jim Young, Lorenzo Pantieri and Enrico Gregorio¹, Jörg Sommer, Joachim Köstler, Daniel Gottschlag, Denis Aydin, Paride Legovini, Steffen Prochnow, Nicolas Repp, Hinrich Harms, Roland Winkler, and the whole L^AT_EX-community for support, ideas and some great software.

Regarding L_YX: The L_YX port was initially done by *Nicholas Mariette* in March 2009 and continued by *Ivo Pletikosić* in 2011. Thank you very much for your work and the contributions to the original style.

¹ Members of GuIT (Gruppo Italiano Utilizzatori di T_EX e L^AT_EX)

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ACRONYMS

DRY Don't Repeat Yourself

API Application Programming Interface

UML Unified Modeling Language

INTRODUCTION

A major and succesful topic in astronomy has been the discovery and characterization of exoplanets. This intrest has sparked great progress in the field of high contrast imaging resulting in new imaging techniques and instuments. Though as of writing 3,949 planets [nasa] have been confirmed little is known about their formation. Certain is that a so called protoplanetary disk is an important stage in the formation of a planetary system. The disks can be observed best within the visible spectrum and infrared though ALMA has also had succes in observing disks in the radio spectrum.

It is however a great challenge to acquire observations of such a disk due to the high contrast with its star and the required angular resolution. ****explain why light from a star (dot) can hide planetary details in a normal telescope (is it only sys+atmosph psf for appature?)****. To observe with sufficient resolution adaptive optics are used. This negates most of the atmospheric seeing.

Reducing contrast is done with coronagraph, an instrument that blocks out the direct light from a star. The classic lyot coronagraph (Figure 1) blocks direct star light using two focusses. In the first focus an opaque mask blocks placed where the star is diffuses and absorbs the direct star light. Then between the focusses a ring shaped mask blocks most of the now diffused star light. Then at the second focus the image is made as usual.

However a classical coronagraph is not sufficient for observing disks. The vector Apodizing Phase Plate (vApp) [papervApp] is a different type of coronagraph placed at the pupil of a telescope. The vApp blocks starlight by modifying the phase of incoming light. This phase change flips the light in a D-shaped region around the star to the other side. Any faint object next to the star in the now dark re-

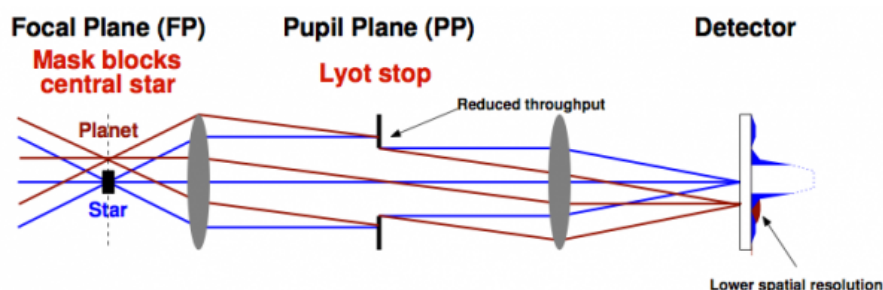


Figure 1: optical layout of a lyot coronagraph, by Matthew Kenworthy, from <https://home.strw.leidenuniv.nl/~kenworthy/app>

gion becomes detectable. Compared to the classic lyot coronagraph the vApp reduces the starlight to a greater degree, see Figure 2.

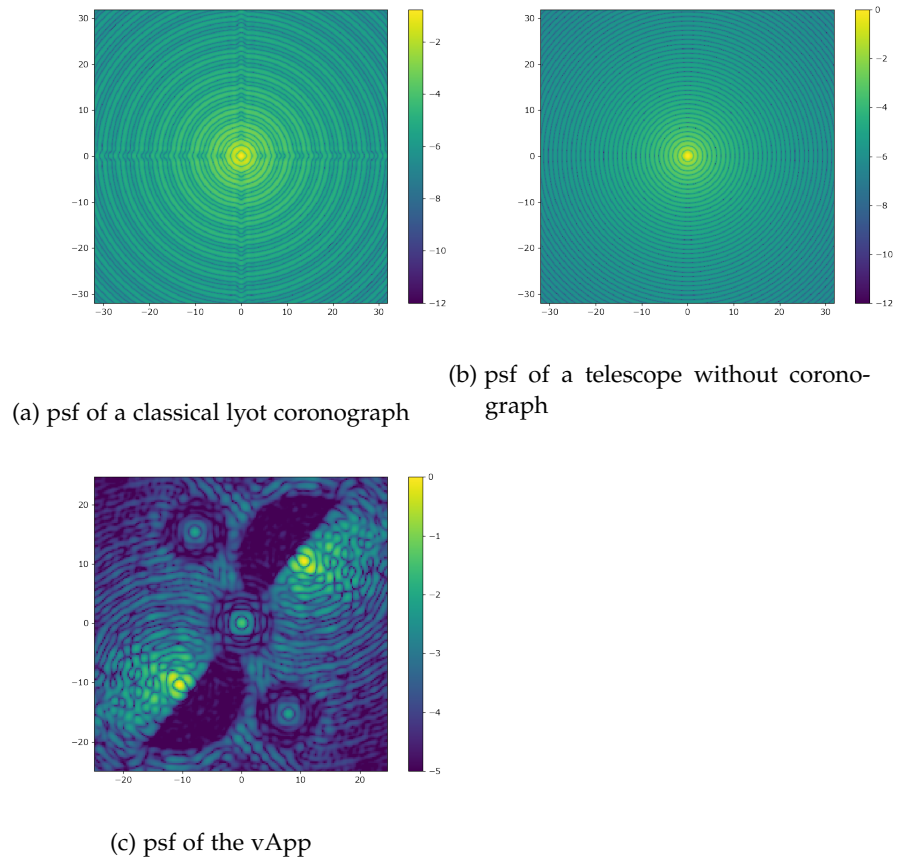


Figure 2: The point spread function (psf) is the image an instrument produces when looking at a point source such as a star. Here we see the psf of an instrument without a coronagraph, with the classical lyot coronagraph and one using the vApp

The vApp has been developed to detect rocky planets in the habitable zone of stars. However, it could also allow us to resolve disk features and study disks in greater detail directly. However, as the vApp changes the entire image, it is challenging to differentiate disk features from vApp artefacts.

Here we study what the effect is of both the vApp and reduction methods ADI and RDI on the apparent morphology of the disks. Here we try to characterise artificial disk observations with the vApp.

TODO overview thesis

<https://en.wikipedia.org/wiki/Coronagraph>

THEORY

Here we discuss what an psf is, how it helps us describe an optical system and how they naturally appear from basic optics. Then we look at the different types of disks that have been observed and what contrasts we expect.

2.1 PSF

We can describe what happens to light going through an optical system with its point spread function (PSF). It describes the light intensity on the focal plane (where the sensor is located) as a function of x and y when a single light ray is imaged on the center of the focal plane.

We derive a way to find the PSF from the Hygens-Fresnel Principle. It states any part of a wave can be described as a front of infinite point sources interfering with one another. An optical element can change these arrangement of these sources, for example an aperture allows only a small area to be filled with these point sources as illustrated in [Figure 3](#).

We find the electric field at a point P at distance R by summing up the fields of these infinite point sources taking into account the different distances to R . Writing the infinite sum as an integral we get [Equation 1](#) for the electric field at a point P some distance R from an aperture.

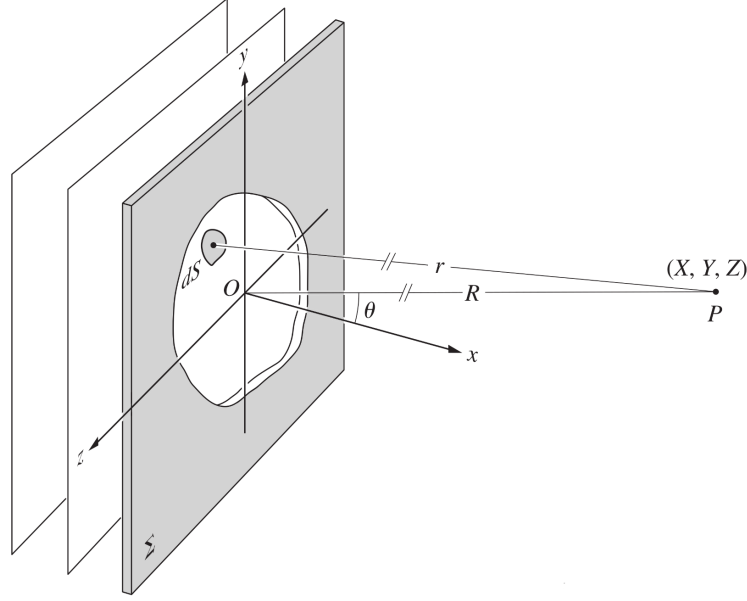
$$E = \frac{\epsilon_A e^{i(\omega t - kR)}}{R} \iint_{\text{Aperture}} e^{ik(Yy + Zz)/R} dS \quad (1a)$$

Here capital Y, Z describe the position in the imaging plane in which P lies as seen in [Figure 3](#). Small letters y and z are the position in the aperture plane. The integral is over the aperture, only integrating over the transparent parts.

To account for changes in phase and not only magnitude of the field caused by Optical instruments we including an aperture function instead of just integrating over the shape of the aperture. This results in:

$$\mathcal{A}(y, z) = \mathcal{A}_0(y, z) e^{i\phi(y, z)} \quad (2a)$$

Figure 3: Fraunhofer diffraction from an arbitrary aperture, r and R large compared to the size of the hole. Extracted from Optics 5th edition, by [4]



Here the amplitude of the apperture function comes from \mathcal{A}_0 and the phase from $e^{i\phi(y,z)}$.

$$E(Y, Z) = \iint_{-\infty}^{\infty} \mathcal{A}(y, z) e^{ik(Yy + Zz)/R} dy dz \quad (3a)$$

The expression for the E field at the point P (Equation 1) rewritten to make use of the apperture function.

We can rewrite this to get rid of the dependence on the distance by subsututing $K_y = kY/R$ and $K_z = kZ/R$ for Y and Z . This gives the final form:

$$E(K_Y, K_Z) = \iint_{-\infty}^{\infty} \mathcal{A}(y, z) e^{ik(K_Y y + K_Z z)/R} dy dz \quad (4a)$$

This is the 2 dimensional Foerier transformation of the aperture function. Thus "the field distribution in the fraunhofer diffraction pattern is the Fourier transform of the field distrubution across the aperture (e.i., the aperture function" [4]).

For the psf we are interested in the intesity which is not the electric field E but $|E|^2$. This means we can calculate the PSF of the vAPP by foerier transforming its apperture function and squaring the result. Note that the amplitude in the apperture function does not only have to depend on the shape of the apperture as there might be partially

transparent material. The aperture function does not even have to be an aperture.

With the PSF we know how a single point source would look when imaged by an optical system. If we assume the system is linear system (****isnt it always****?) we can use the superposition principle to image extended sources by convolving the psf with the extended source.

2.1.1 *Atmosphere*

The telescope or coronagraph are the only optical system at play. There are many differently moving layers of air between the telescope and space. These work as independent optical systems that change in time. Each layer moves in a different direction at a different speed as winds are different at various altitudes. This changes the phase of the light. The complete psf changes all the time. The changes are smaller at smaller timescales.

2.1.2 *Adaptive Optics and speckles*

Since the 1990 adaptive optics (AO) are used. These change shape to undo the phase change of the atmosphere. However as these effects are unpredictable they always lack behind slightly. Further more they have errors themselves, the phase is never completely corrected. Both these effects cause small distortions in the final image, these are known as speckles. Thus even with AO the total psf for the atmosphere, the AO and the instrument will keep changing in time, however the magnitude of the change is severely reduced. Adaptive optics don not correct the entire field, generally each system has a radius of correction in which the AO works.

-

2.2 VAPP

As mentioned in the introduction the vector apodising phase plate is not a normal coronagraph. The apodising phase plate (APP) is a coronagraph that changes amplitude of phase in the pupil plane to create destructive interference in an area of the psf. This creates a very dark area on the psf. Dim object imaged there can be resolved if the contrast between them and the star is smaller then the contrast between center peak of the psf and the dark zone. Because the coronagraph works in the pupil plane it is insensitive to the effects of speckles, further more unresolved stars do not limit how close to a star the coronagraph can function.

The APP does this by introduces differences in the optical path length thereby changing the phase. These changes are designed to

create the mentioned dark zone in the shape of a 180° half circle with a radius of 2 to 9 λ/D [1]. To do so optical the path length differences need to be different throughout the pupil, the design can be seen as a heightmap of path differences. The APP is manufactured by directly printing that heightmap using liquid-crystal technology.

The vAPP is an upgrade to the APP that applies a the path length

2.3 DISKS

A newly formed star is enveloped in a disk of gas and dust, a protoplanetary disk.

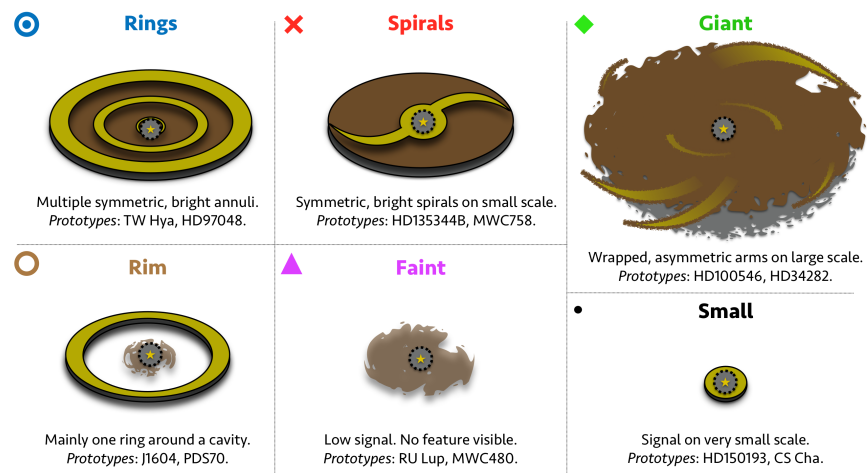
–something about the components/dust/gas refwilliams

From it planets can form. The process of planet formation is not yet understood. Direct imaging of newly formed planets is hard. Objects inside a disk disturb the shape and create features in the disk. By observing the features of a disk we can learn more about planet formation.

2.3.1 Categorisation

Roughly speaking disks seem to be ring or spiral shaped with some forms in between. [2] classifies them into 6 categories in Figure 4.

Figure 4: Sketch summarizing the different classifications of protoplanetary disks proposed by /ref



They then conclude that: -faint disks are young -spiral disks are almost starting their main sequence -ring disk have no outer stellar companion (??)

2.3.2 *Challenging to Observe*

It is challenging to observe a disk since its brightness is low compared to the star. To get an upper limit on the brightness of a disk assume all star light that hits the disk is reflected towards us. The light from a star drops quadratically as it gets farther away. Thus we will never have a disk brightness exceeding $1/R^2$. A disk at 1 AU from a sun like star will have a brightness 2.1×10^{-5} of the star. At 5 AU this drops to 8.6×10^{-7} .

DATA GENERATION

We choose not to use on sky data but create our own. This presents a number of advantages though the main reason is the lack of disk observations using the vApp. Generating the data from a model allows us to vary parameters as we wish. We can try a reduction method on a simple slightly inclined disk or a near face on disk with many rings. By reducing images of morphologically differing disks we can map out how well a reduction method works for each morphology. Furthermore we can clearly separate artifacts created during data reduction from disk features when we know what the disk looks like.

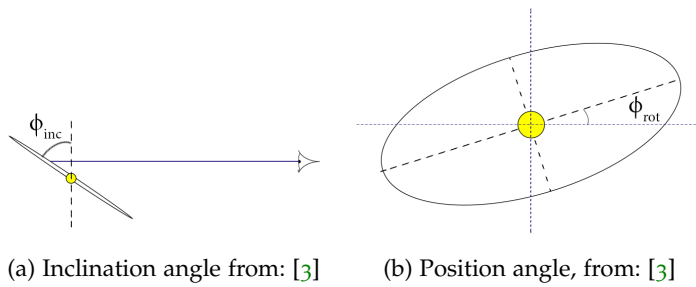
3.1 DISK MODEL

We use a 2-dimensional disk model based on the work by [3]. The model has 4 basic parameters:

INCLINATION The angle the disk is tilted towards the observer. A 0 degrees inclination gives a face on disk and 90 degrees a horizontal line being an edge on disk. As illustrated in ??

POSITION ANGLE After inclination the disk can be rotated around the line of sight from the observer, rotation to the left is positive. See ??

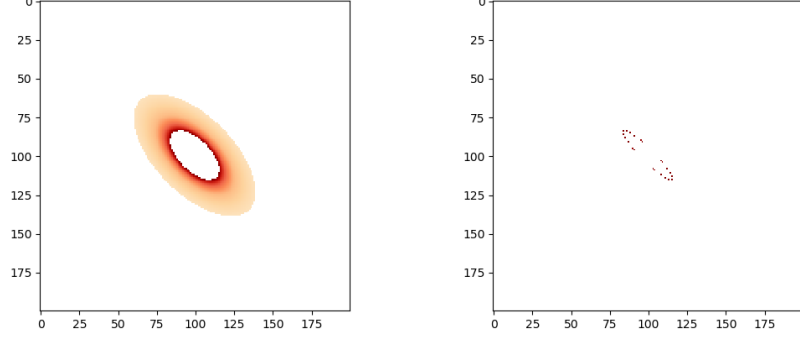
INNER AND OUTER RADIUS Many disks start and stop at some radius from star. The inner and outer radius are relative to the field size that defaults to 10. An Inner radius of 2 gives a hole in the disk with a diameter 20% of the image width.



The disk is modeled as optically thick and does not emit light on its own. At a certain radii between the inner and outer radii the disks brightness is given by:

$$B(r) = B_{\text{star}} \cdot \left(\frac{r}{R_{\text{star}}} \right)^2 \quad (5)$$

We rasterize the image onto a resolution of 200 by 200 pixels. This not only speeds up our calculations it also is the expected output of an actual observing session. The model behaves well for most disks however on this resolution features that are small in the observers plane are pixilated as we see in [Figure 6](#) below. This poses no problem as we do not expect to resolve such features.



(a) model output for an inclination of 60 deg, an inner radius of 2 and an outer radius of 5 (b) model output for an inclination of 80 deg, an inner radius of 2.2 and an outer radius of 2.3

Figure 6: two disks created by the model with different size and inclination both with a position angle of 45 degrees

3.2 PSF GENERATION

To simulate an observation do not use a single psf as given by `/cite-vApp` as during an observation the combined psf of the atmosphere and instrument changes. A complete simulation of these effects is out of the scope of this thesis. We instead decided to try and approach a similar morphology to the on sky psf.

First we tried modifying a give single `vApp` psf to get a set of disturbed psfs. To achieve this a pattern is added to the foerier transform of the psf before transforming it back from foerier space, see the equation below. The pattern is then shifted for every timestep.

$$\mathcal{F}_{2d}^{-1} \left(\text{intensity} \cdot \mathcal{F}_{2d}(\text{psf}) \cdot \text{pattern} + (1 - \text{intensity}) \cdot \mathcal{F}_{2d}(\text{psf}) \right) \quad (6)$$

The best results where achieved using a grid of blurred circles as pattern, see [Figure 7](#). Note the distortions to the psf are clustered around the center of its peaks but not randomly spread. Clearly this will not do.

Then we tried a simulation of the psfs using `HCIPy`, an open-source object-oriented framework written in Python for performing end-to-

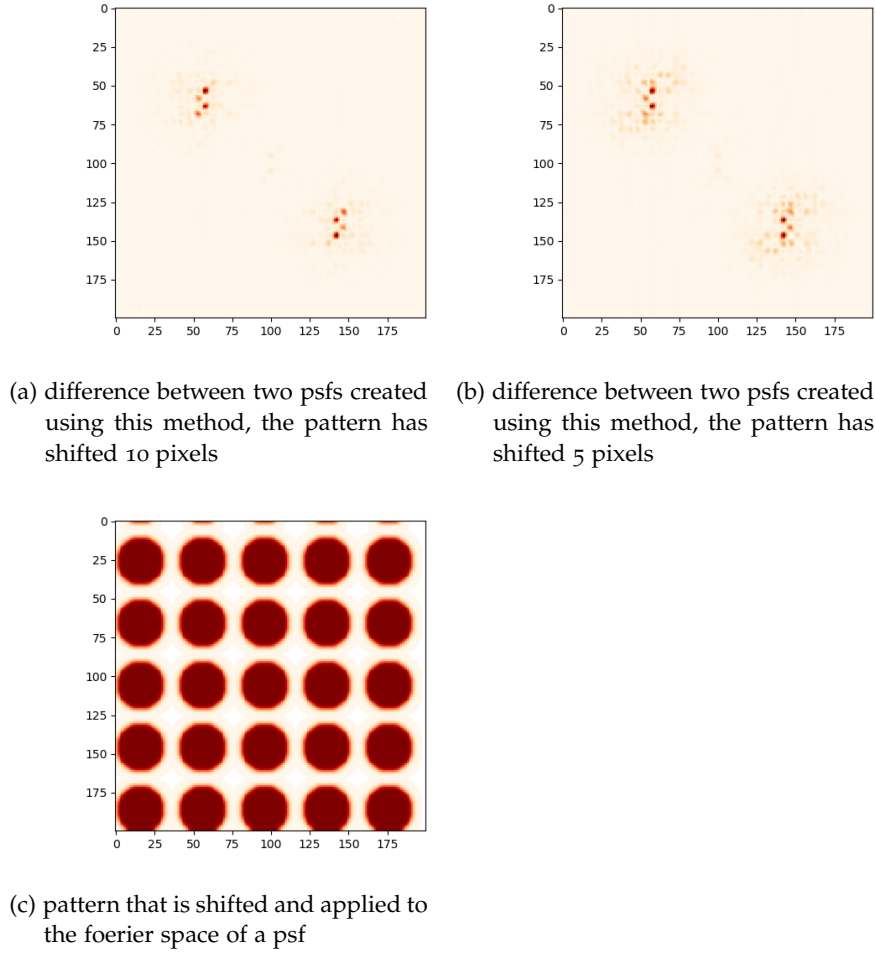


Figure 7: Pattern (a) and the diffrences between psfs ****distorted**** with this pattern using the method described above

end simulations of high-contrast imaging instruments [5]. The framework is used to do a very rough simulation. We use the provided methods to create a multi-layer atmospheric model that changes in time. Then we use that model, the vApp amplitude and vApp phase screen to generate a series of psfs. Instead of modelling an active optics system we modified the fried parameter for the atmosphere to get similar morphological changes in time to the available on sky images. This is sufficient as the disk will always lie within theradius of correction???..... which means the limiting effect of an AO system on the disk will only be a increase in resolution by reducing the seeing effect.

Wu use a telescope diamater of 8.2meters and a wavelength, of $1 \cdot 10^{-6}$ meters for the simulation. We got the right morphology with a fried paramater of 4m. To put this into perspective: to simulate excellent seeing conditions we would use 20cm. With these setting we use hcipy to create an orderd set of psfs changing through time

as the simulated atmosphere evolves. See [Figure 8](#) for the difference between the first psf in a set and later psfs in the same set. Note that as time evolves the differences grow.

3.3 COMBINING

To get a dataset that simulates an observation we create an ordered set of disk images each one rotated a bit to each other. This accounts for the field rotation caused by observing with an alt azimuth telescope. Now to get the simulated observation an ordered set is created by convolving each n-th image from the disk set with the n-th image of the psf set. These are the simulated observations through time.

Using variations on this method we create 3 different sets.

1. Place a single pixel with value one in the center of the disk before going through the above procedure. This is our observation set.
2. Swap the disk with the single pixel in the center representing the star. This set we use to check what artefacts are created by the reduction
3. Take the convolution of the n-th disk image with the first psf image of the psf set. This way we lose all atmospheric effects, since we haven't placed a star in the center we see only the effect of the vApp on the disk image.

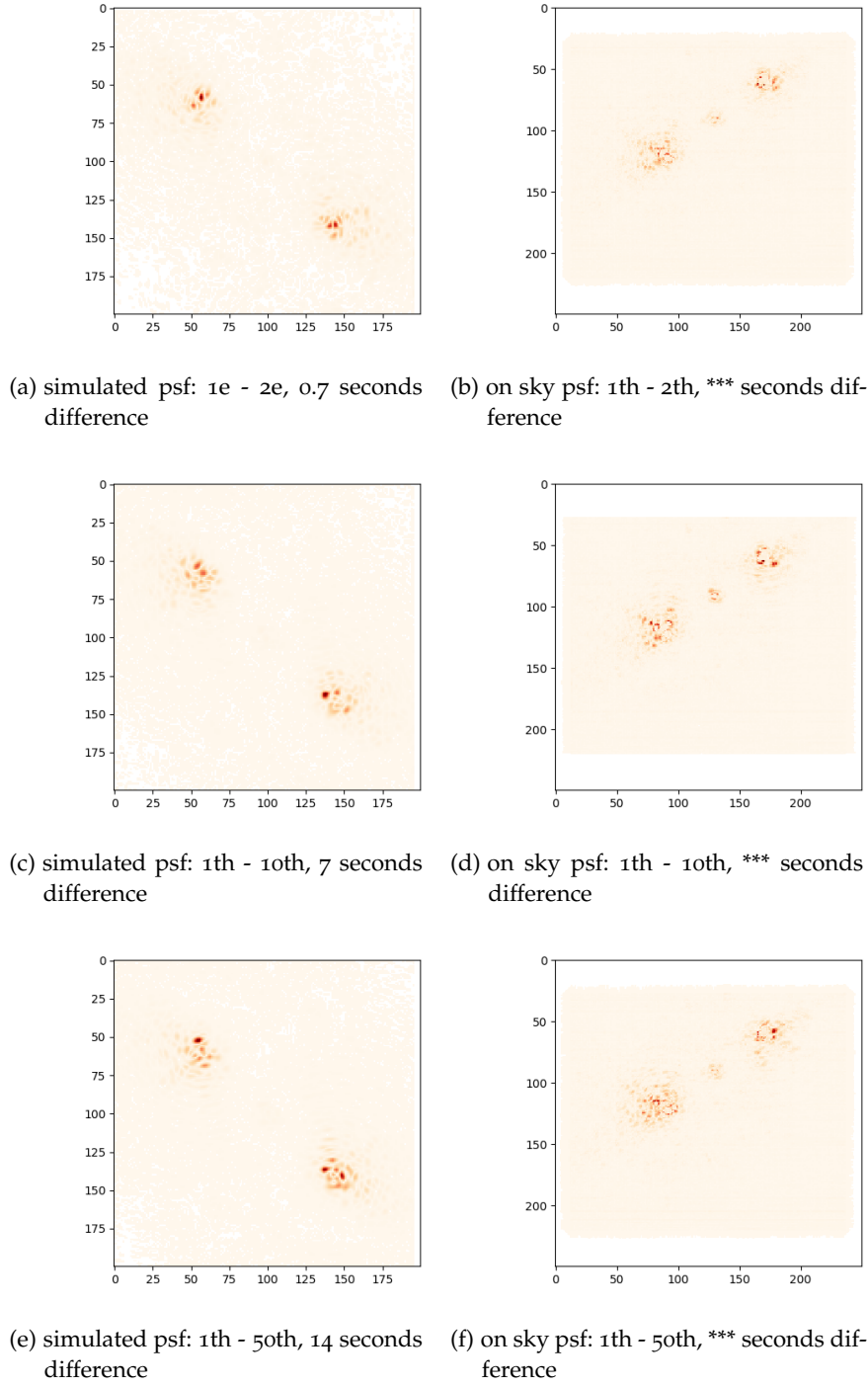


Figure 8: Differences between the first psf of an observation and later psfs. Compared between simulated observations on the left and on sky data on the right. The psfs were aligned then normalised on the maximum of the leakage term, finally the absolute value of the difference between the first and n-th psf was taken. The second, 10th and 50th simulated psf are 0.7, 7 and 14 seconds apart in the simulation

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