

Hierarchy of Pure States applied to Photosynthetic Complexes

Daniel Süß

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Light-Harvesting Complexes in Bacteria
Markovian Open Quantum Systems

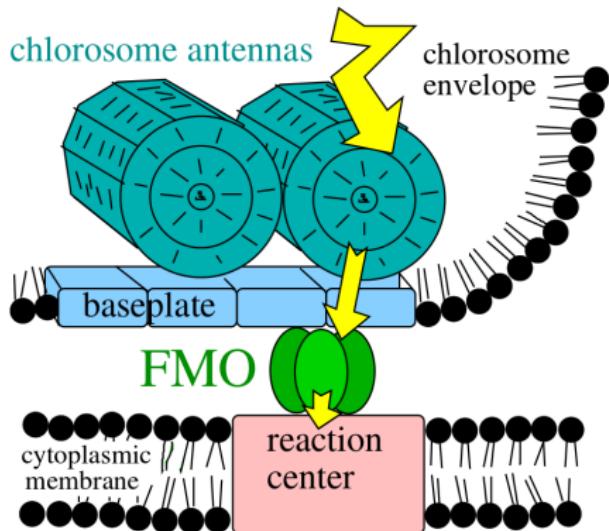
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Non-Markovian Stochastic Schrödinger Equation
Hierarchy of Pure States

③ HOPS applied to the FMO-Complex

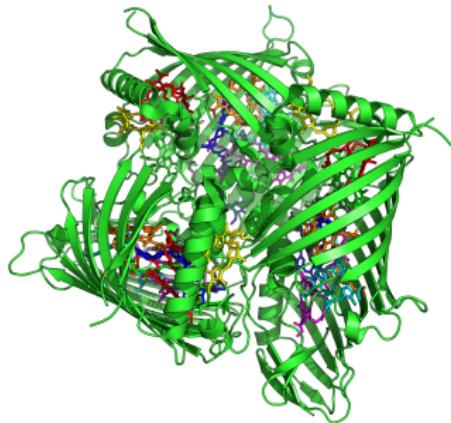
④ Conclusion & Outlook

Light-Harvesting Complexes



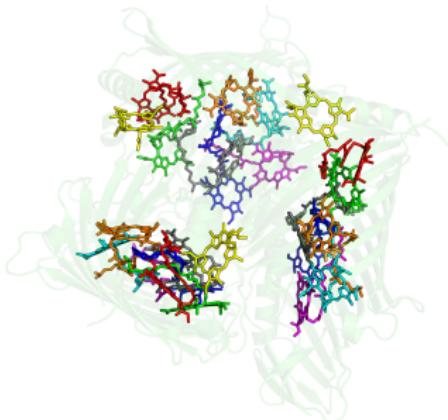
- ▶ light absorbed at chlorosome antennas; transport to reaction center
- ▶ thermal noise in biological environment
- ▶ lifetime $\sim 2 - 5$ ps

The Fenna-Matthews-Olson Complex



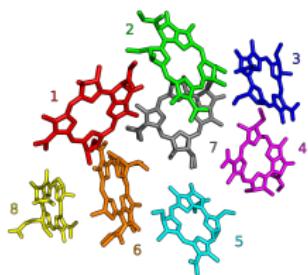
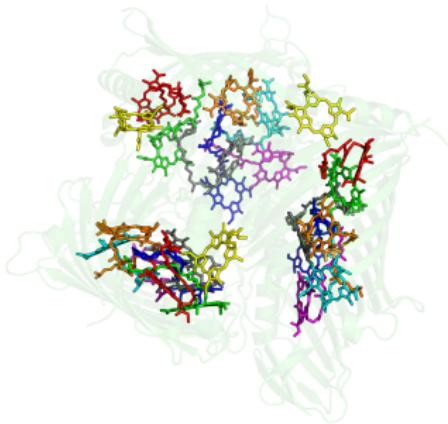
- ▶ found in green sulfur bacteria
- ▶ simplest pigment-protein complex
 - ▶ ideal test object
- ▶ quantum effects
- ▶ 3 identical, weakly coupled monomers + protein environment
- ▶ 8 Bacteriochlorophylls (BChls)
 - ▶ here: neglect BChl 8
- ▶ Energy received at BChls 1+6 transferred to sink at BChl 3

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Physical Model – Electronic System

- ▶ separation of electronic and vibrational degrees of freedom
 - ▶ electronic: valence electrons
 - ▶ vibrational: other electrons, nuclear DOF, environment
- ▶ one excited electronic state $|\pi_n\rangle$ per BChl (*exciton*)

$$H_{\text{el}} = \sum_n \epsilon_n |\pi_n\rangle\langle\pi_n| + \sum_{m,n} V_{mn} |\pi_m\rangle\langle\pi_n|$$

(site energy ϵ_n ; int. strength V_{mn})

- ▶ no effective transport \Rightarrow influence of environment important

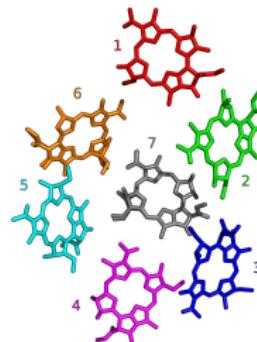
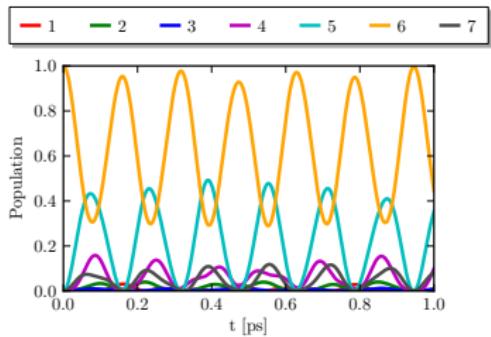
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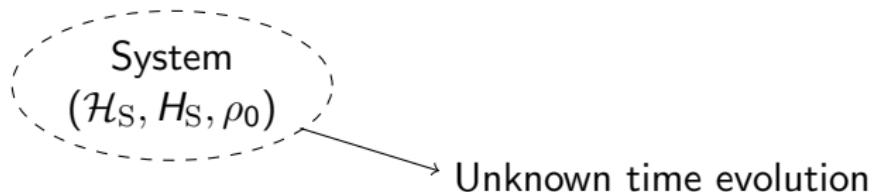
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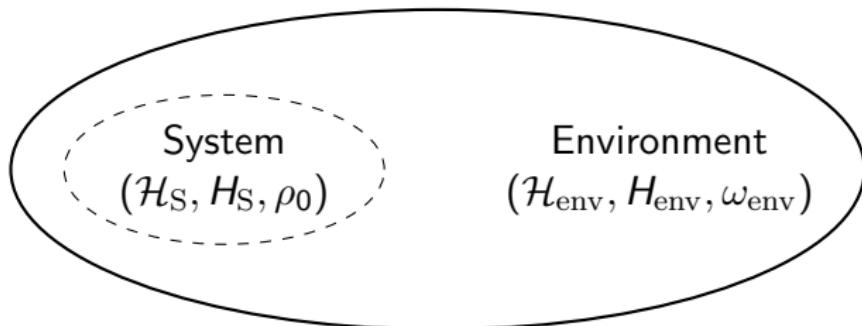
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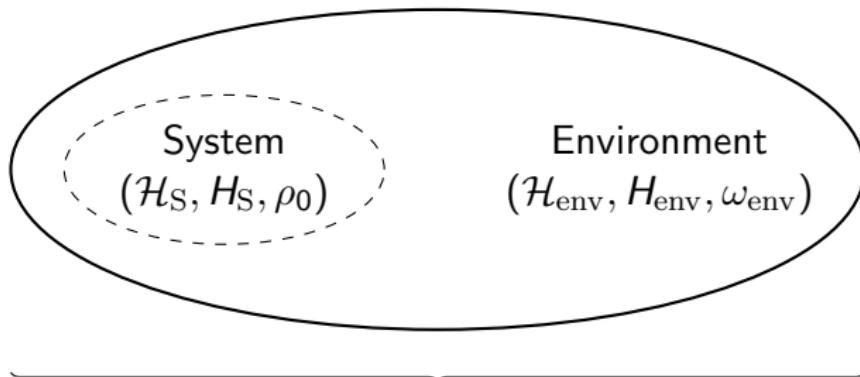
Open Quantum Systems



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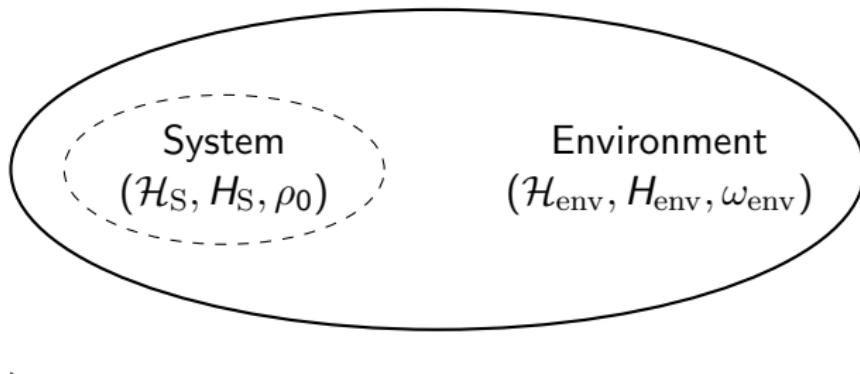


$$\mathcal{H}_S \otimes \mathcal{H}_{\text{env}}$$

$$H_S \otimes I + I \otimes H_{\text{env}} + H_{\text{int}}$$

$$\rho_0 \otimes \omega_{\text{env}}$$

Open Quantum Systems



$$\begin{gathered} \mathcal{H}_S \otimes \mathcal{H}_{\text{env}} \\ H_S \otimes I + I \otimes H_{\text{env}} + H_{\text{int}} \\ \rho_0 \otimes \omega_{\text{env}} \end{gathered}$$

$$\implies \rho_t = \Lambda_t \rho_0 = \text{Tr}_{\text{env}} \left(U_t \rho \otimes \omega_{\text{env}} U_t^\dagger \right)$$

→ Common approach: Derive evolution equation for ρ_t .

Markovian Open Quantum Systems

Born-Markov approximation,
Completely positive time evolution

Linblad Master Equation

Quantum Trajectories

Linblad Master Equation

$$\partial_t \rho_t = -i[H, \rho_t] + ([L \rho_t, L^\dagger] + [L, \rho_t L^\dagger])$$

- ▶ “free”, unitary time evolution + irreversible channel
- ▶ system of N^2 real valued ODEs

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Markovian Open Quantum Systems

Quantum Trajectories (diffusive, linear)

$$d|\psi_t\rangle = \left(-iH - \frac{1}{2}L^\dagger L \right) |\psi_t\rangle dt + L|\psi_t\rangle d\xi_t^*$$

- ▶ driven by complex white noise $\dot{\xi}_t$
- ▶ ρ_t recovered by “classical” average over independent realizations

$$\rho_t = \mathbb{E}(|\psi_t\rangle\langle\psi_t|) \approx \frac{1}{N} \sum_{n=1}^N |\psi_t^n\rangle\langle\psi_t^n|$$

- ▶ system of $2N$ real valued SDEs, trivial parallelization

→ failure of Markovian theory due to strong coupling, “memory” in time evolution (or initial entanglement)

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Standard Open System Model

- ▶ specific model for environment: harmonic oscillators, linearly coupled

$$H_{\text{tot}} = H_{\text{sys}} + \sum_{\lambda} \omega_{\lambda} a_{\lambda}^{\dagger} a_{\lambda} + \sum_{\lambda} \left(g_{\lambda}^* L a_{\lambda}^{\dagger} + g_{\lambda} L^{\dagger} a_{\lambda} \right)$$

- ▶ interaction picture with respect to H_{env}

$$H_{\text{tot}}(t) = H_{\text{sys}} + \sum_{\lambda} \left(g_{\lambda}^* L a_{\lambda}^{\dagger} e^{i\omega_{\lambda} t} + g_{\lambda} L^{\dagger} a_{\lambda} e^{-i\omega_{\lambda} t} \right)$$

- ▶ Schrödinger equation for system and environment

$$\partial_t |\Psi_t\rangle = -i H_{\text{tot}}(t) |\Psi_t\rangle, \quad |\Psi_0\rangle = |\psi_0\rangle \otimes |0\rangle$$

- ▶ coherent state representation for bath DOF $\psi_t(z^*) = \langle z|\Psi_t\rangle$

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Non-Markovian Stochastic Schrödinger Equation

$$\partial_t \psi_t(\mathbf{z}^*) = -i H_{\text{sys}} \psi_t(\mathbf{z}^*) + L Z_t^* \psi_t(\mathbf{z}^*) - L^\dagger \int_0^t \alpha(t-s) \frac{\delta \psi_t(\mathbf{z}^*)}{\delta Z_s^*} ds$$

- ▶ “stochastic process” $Z_t^* = -i \sum_{\lambda} g_{\lambda}^* z_{\lambda}^* e^{i \omega_{\lambda} t}$
- ▶ bath correlation function $\alpha(t) = \sum_{\lambda} |g_{\lambda}|^2 e^{-i \omega_{\lambda} t}$
- ▶ only dependence on z_{λ}^* through $Z_t^* \Rightarrow \psi_t(\mathbf{z}^*) \rightarrow \psi_t(Z^*)$
- ▶ stochastic Schrödinger equation for quantum trajectories $\psi_t(Z^*)$ with

$$\mathbb{E} Z_t = 0, \quad \mathbb{E}(Z_t Z_s) = 0, \quad \text{and} \quad \mathbb{E}(Z_t Z_s^*) = \alpha(t-s)$$

$$\rho_t = \text{Tr}_{\text{env}} |\Psi_t\rangle\langle\Psi_t| = \mathbb{E}(|\psi_t(Z^*)\rangle\langle\psi_t(Z^*)|)$$

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Hierarchy of Pure States (HOPS)

$$\partial_t \psi_t = -iH_{\text{sys}} \psi_t + LZ_t^* \psi_t - L^\dagger \underbrace{\int_0^t \alpha(t-s) \frac{\delta \psi_t}{\delta Z_s^*} ds}_{=: \psi_t^{(1)}}$$

- closed time evolution equation in case $\alpha(t) = g e^{-\gamma|t| - i\Omega t}$

$$\begin{aligned}
 \partial \psi_t^{(0)} &= (-iH_{\text{sys}} + LZ_t^*) \psi_t^{(0)} && - L^\dagger \psi_t^{(1)} \\
 && \uparrow & \downarrow \\
 \partial \psi_t^{(1)} &= (-iH_{\text{sys}} - (\gamma + i\Omega) + LZ_t^*) \psi_t^{(1)} && + g L \psi_t^{(0)} - L^\dagger \psi_t^{(2)} \\
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 &\vdots && \\
 \partial \psi_t^{(k)} &= (-iH_{\text{sys}} - k(\gamma + i\Omega) + LZ_t^*) \psi_t^{(k)} && + kg L \psi_t^{(k-1)} - L^\dagger \psi_t^{(k+1)} \\
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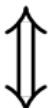
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Truncation of the Hierarchy

$$\partial_t \psi_t^{(k)} = (-iH_{\text{sys}} - k(\gamma + i\Omega) + LZ_t^*) \psi_t^{(k)} + kg L \psi_t^{(k-1)} - L^\dagger \psi_t^{(k+1)}$$



$$\psi_t^{(k)} = \int_0^t e^{-k(\gamma+i\Omega)(t-s)} \dots \left(kg L \psi_s^{(k-1)} - L^\dagger \psi_s^{(k+1)} \right) ds$$

- truncation at finite order D with *terminator* (here: $\Omega = 0$)

$$\partial_t \psi_t^{(D)} = \left(-iH_{\text{sys}} - D\gamma + LZ_t^* - \frac{g}{\gamma} L^\dagger L \right) \psi_t^{(D)} + Dg L \psi_t^{(D-1)}$$

- “Markovian” approximation for D^{th} order

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\Updownarrow

$$\psi_t^{(k)} \sim kg L \psi_t^{(k-1)} - L^\dagger \psi_t^{(k+1)}$$

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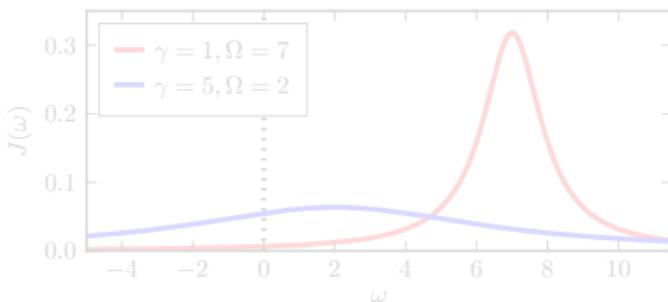
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Further Remarks

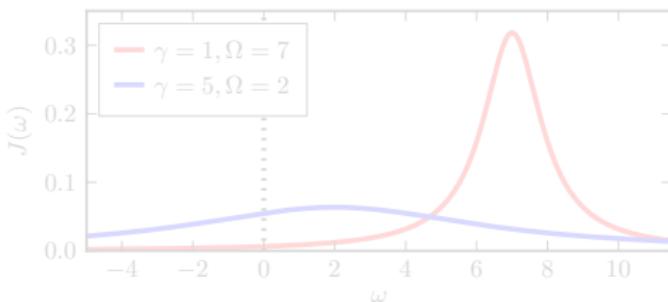
- ▶ no simple criterion to determine $D \rightarrow$ individual checking necessary
- ▶ rule of thumb: large truncation order D for
 - ▶ small memory time \iff large γ
 - ▶ strong coupling \iff large g
- ▶ nonlinear version \rightarrow proper importance sampling
- ▶ exponential BCF requires negative frequencies \rightarrow unphysical



- ▶ proper treatment for $T \neq 0$

Further Remarks

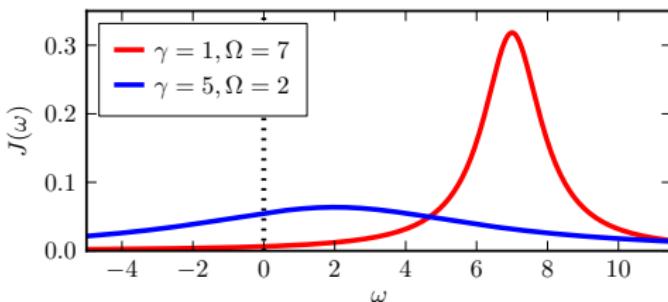
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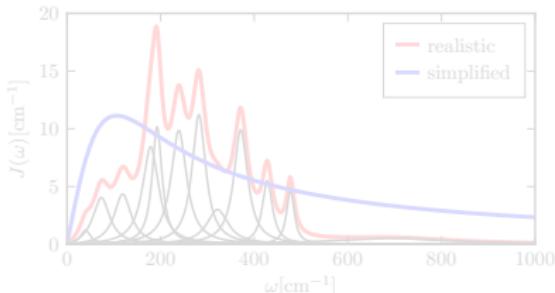
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The FMO-Complex – Full Model

- ▶ environment: independent, identical harmonic baths for each BChl
- ▶ all thermal effects included in *thermal correlation function*

$$\alpha(t) = \int_0^{\infty} J(\omega) \left(\coth\left(\frac{\omega}{2T}\right) \cos(\omega t) - i \sin(\omega t) \right) d\omega$$

- ▶ highly structured realistic spectral density
 - ▶ large number of exponential modes required
 - ▶ simplified spectral density (for $D = 2$: 120 vs. 15576 auxiliary states)

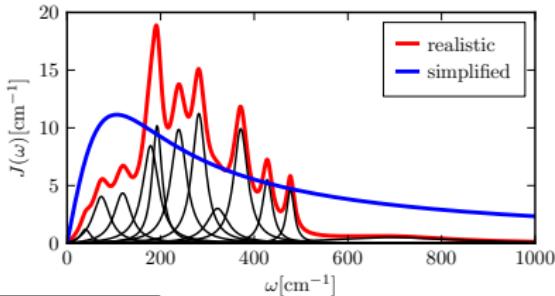


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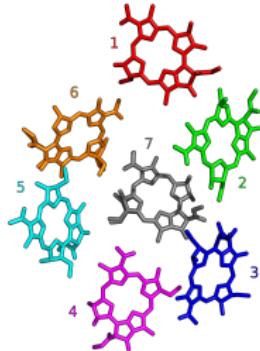
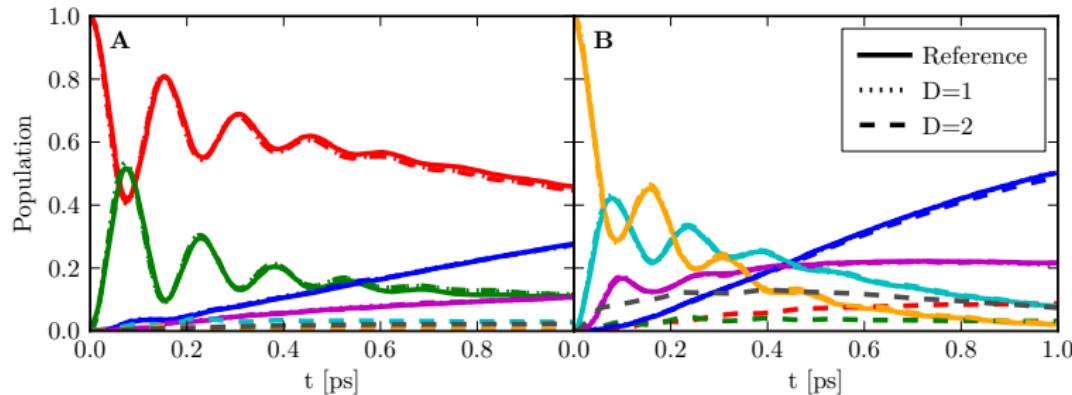
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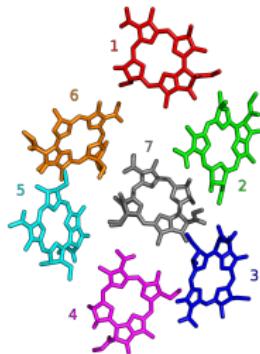
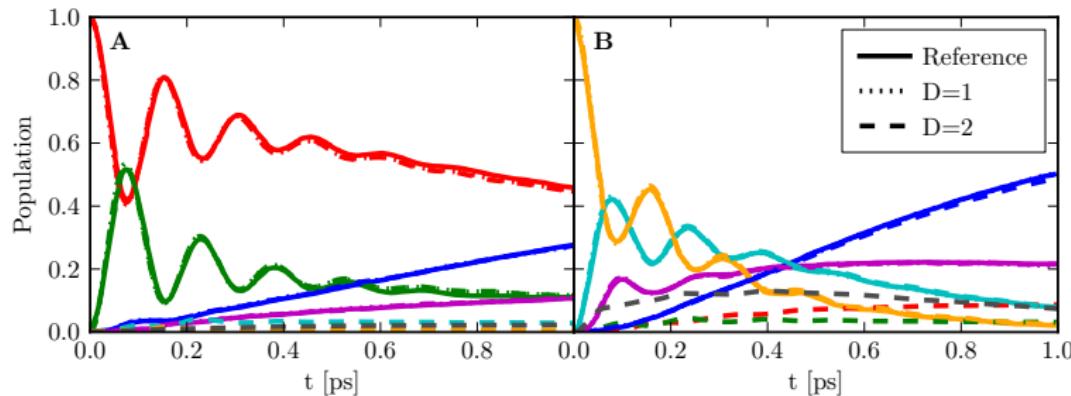


Exciton Transfer at Cryogenic Temperature



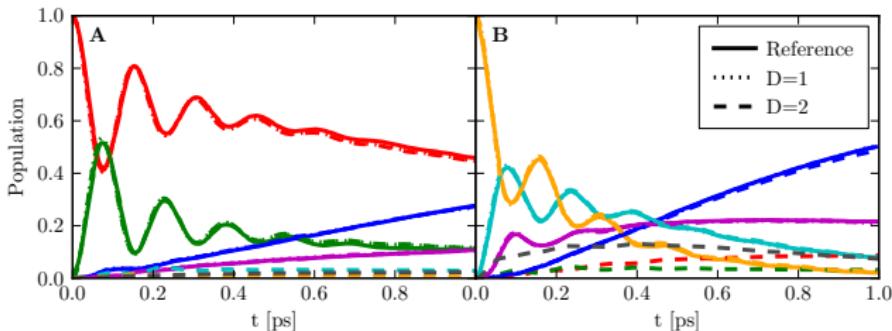
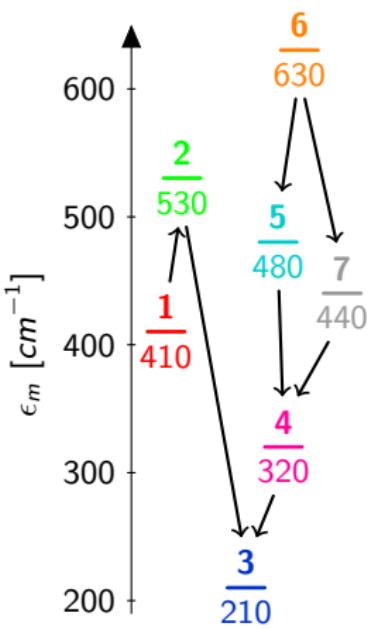
- ▶ directed energy transfer toward BChls 3 + 4
- ▶ yield of 40 % (resp. 75 %) on relevant BChls
- ▶ coherent, wavelike dynamics up to $t \approx 0.7$ ps

Exciton Transfer at Cryogenic Temperature



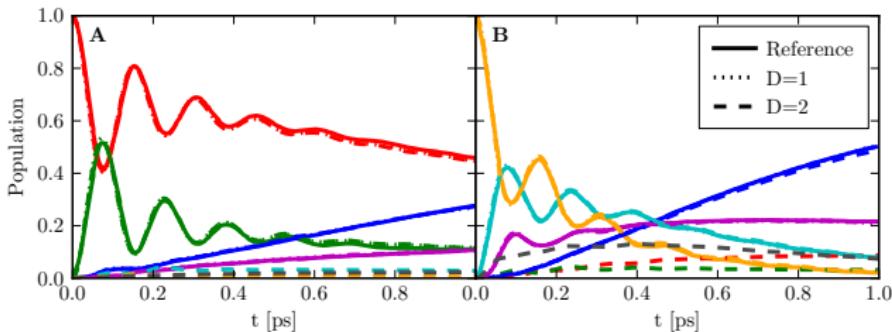
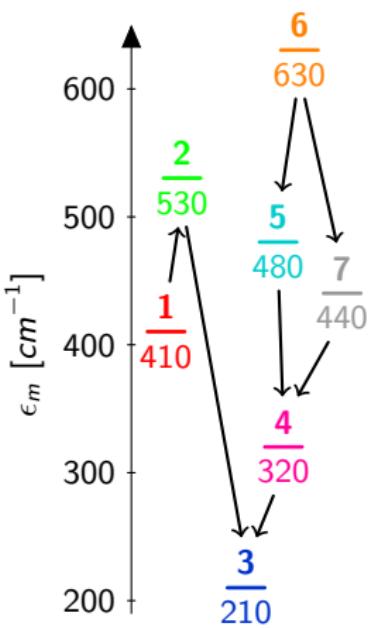
- ▶ directed energy transfer toward BChls 3 + 4
- ▶ yield of 40 % (resp. 75 %) on relevant BChls
- ▶ coherent, wavelike dynamics up to $t \approx 0.7$ ps

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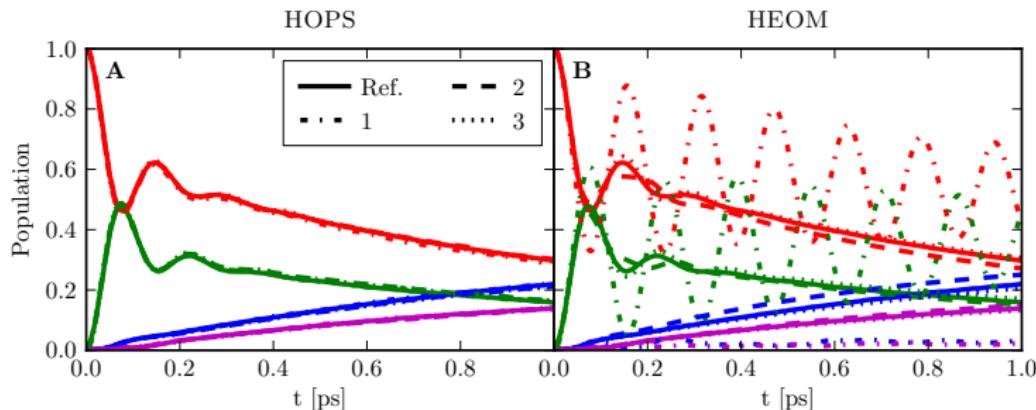
- ▶ two distinct transport channels
- ▶ potential step between BChls 1 & 2
 - ▶ higher yield with start on BChl 6
 - ▶ correction in full monomer
 - ▶ assumed to serve as barrier preventing depopulation of target
 - ▶ coherent tunneling may assist overcoming

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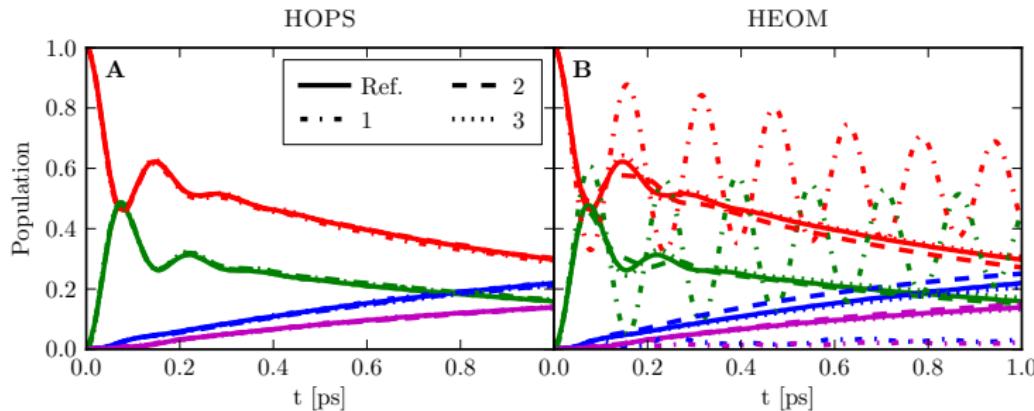
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Exciton Transfer at Physiological Temperature



- ▶ same qualitative result as for cryogenic temperature
- ▶ shorter timespan of wavelike dynamics (up to $t \approx 0.3$ ps)
- ▶ yield reduced by 25 %
- ▶ HEOM – hierarchical equations of motion
 - ▶ based on density matrix formalism
 - ▶ similar hierarchical structure

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Summary

- ▶ Hierarchy of pure states based on NMSSE
 - ▶ applicable to system coupled linearly to bosonic environment
 - ▶ numerically exact and valid for a large parameter regime
 - ▶ solves the “functional derivative problem” without any approximation
 - ▶ systematic check of convergence
 - ▶ higher precision at equal truncation order than HEOM-approach
- ▶ exciton-energy transfer in light-harvesting complexes
 - ▶ simplest example: FMO-complex in sulfur bacteria
 - ▶ reproduce established results using HOPS
 - ▶ remarkably high efficiency; possible explanation due to quantum effects
- ▶ We did not talk about...
 - ▶ nonlinear equations
 - ▶ thermal bath correlation function expansion
 - ▶ absorption spectra
 - ▶ ...

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- ▶ influence of spectral density

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Open Questions

- ▶ connection to HEOM
- ▶ non-exponential BCFs
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Thank you for your attention!

