

**ONE HOUR THIRTY MINUTES**

A list of constants is enclosed.

**UNIVERSITY OF MANCHESTER**

Advanced Quantum Mechanics

 2024,  - 

Answer **TWO** questions

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The use of calculators is permitted, as long as they cannot store text and have no graphing capability.

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The numbers are given as a guide to the relative weights of the different parts of each question.

You may assume the following formulae in any question *if proof is not explicitly requested*:

### Special integrals

$$\int_{-\infty}^{\infty} x^{2n} e^{-ax^2} dx = \frac{1 \cdot 3 \cdot 5 \cdots (2n-1)}{(2a)^n} \sqrt{\frac{\pi}{a}}, \quad \int_0^{\infty} x^n e^{-ax} dx = \frac{n!}{a^{n+1}}, \quad a > 0.$$

### 4-vectors

Use notation  $x^\mu = (x^0, \mathbf{r})$  for 4-position and  $p^\mu = (E/c, \mathbf{p})$  for 4-momentum. In electro-magnetism  $A^\mu = (\Phi/c, \mathbf{A})$  for 4-potential.

**Klein-Gordon equation** in scalar  $S$  and vector  $(V_0/c, \mathbf{V})$  potential field

$$[(\hat{\mathbf{p}} - \mathbf{V})^2 c^2 + (mc^2 + S)^2] \Psi = (i\hbar \partial_t - V_0)^2 \Psi.$$

**Dirac equation** in scalar  $S$  and vector  $(V_0/c, \mathbf{V})$  potential field

$$[\boldsymbol{\alpha} \cdot (\hat{\mathbf{p}} - \mathbf{V}) c + \beta(mc^2 + S)] \Psi = (i\hbar \partial_t - V_0) \Psi.$$

**Standard spinor solutions** to the free massive Dirac equation

$$u_s(E, \mathbf{p}) = \sqrt{E + mc^2} \begin{pmatrix} \phi_s \\ \frac{c\boldsymbol{\sigma} \cdot \mathbf{p}}{E + mc^2} \phi_s \end{pmatrix}, \quad \phi_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \text{ and } \phi_2 = \begin{pmatrix} 0 \\ 1 \end{pmatrix};$$

### Standard matrices

$$\begin{aligned} \sigma_x &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}; \\ \alpha_i &= \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix}, \quad \beta = \begin{pmatrix} I_2 & 0 \\ 0 & -I_2 \end{pmatrix}, \quad \Sigma_i = \begin{pmatrix} \sigma_i & 0 \\ 0 & \sigma_i \end{pmatrix}; \\ \gamma_0 &= \begin{pmatrix} I_2 & 0 \\ 0 & -I_2 \end{pmatrix}, \quad \gamma_1 = \begin{pmatrix} 0 & \sigma_x \\ -\sigma_x & 0 \end{pmatrix}, \quad \gamma_2 = \begin{pmatrix} 0 & \sigma_y \\ -\sigma_y & 0 \end{pmatrix}, \quad \gamma_3 = \begin{pmatrix} 0 & \sigma_z \\ -\sigma_z & 0 \end{pmatrix}. \end{aligned}$$

**Space-time metric** We use the metric  $\eta_{\mu\nu} = \text{diag}(1, -1, -1, -1)$ .

**Laplacian operator** in terms of angular momentum operator  $\hat{\mathbf{L}}$

$$\nabla^2 = \frac{1}{r} \frac{\partial^2}{\partial r^2} r - \frac{\hat{\mathbf{L}}^2}{\hbar^2 r^2}.$$

**Solution to Dirac equation** in spherical vector potential  $(V_0(r)/c, \vec{0})$

$$\begin{aligned} \psi_{jm}^\kappa(\mathbf{r}) &= \begin{pmatrix} f_j^\kappa(r) \mathcal{Y}_{jm}^\kappa(\theta, \phi) \\ ig_j^\kappa(r) \mathcal{Y}_{jm}^{-\kappa}(\theta, \phi) \end{pmatrix} \quad \text{with } F = rf \text{ and } G = rg, \\ \hbar c \left[ \frac{d}{dr} - \frac{\kappa}{r} \right] G_j^\kappa(r) &= -(E - V_0(r) - mc^2) F_j^\kappa(r), \\ \hbar c \left[ \frac{d}{dr} + \frac{\kappa}{r} \right] F_j^\kappa(r) &= (E - V_0(r) + mc^2) G_j^\kappa(r). \end{aligned}$$

1. a) Show that the angular momentum operator  $\hat{L}_x$  is the generator of infinitesimal rotations about the  $x$ -axis for a spinless particle. Derive an expression in terms of  $\hat{L}_x$  for the operator  $\hat{U}_\beta$  that represents a finite rotation of angle  $\beta$  about the  $x$ -axis. [8 marks]

- b) For a spinless particle with non-zero orbital angular momentum, write down the operator  $\hat{U}_{\alpha,\beta,\gamma}$  representing a rotation with Euler angles  $(\alpha, \beta, \gamma)$ . Within the space of states with angular momentum quantum number  $l = 1$ , the operators  $\hat{L}_x$ ,  $\hat{L}_y$ , and  $\hat{L}_z$  may be written in matrix representation as

$$\hat{L}_x = \frac{\hbar}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}, \quad \hat{L}_y = \frac{i\hbar}{\sqrt{2}} \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix}, \quad \hat{L}_z = \hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}.$$

Show that for the choice of Euler angles  $(0, 0, \gamma)$  the  $\hat{L}_x$  eigenstate

$$\frac{1}{\sqrt{2}} \begin{pmatrix} -1 \\ 0 \\ 1 \end{pmatrix},$$

can be transformed into an eigenstate of  $\hat{L}_y$  for an appropriate value of  $\gamma$ . Give the corresponding Wigner  $D$ -matrix. [7 marks]

- c) A charged quantum rotor of angular momentum quantum number  $l = 1$  is placed in a uniform magnetic field oriented along the  $z$ -axis. The Hamiltonian is given by

$$\hat{H}_0 = \alpha \hat{L}_z,$$

where  $\alpha$  is a positive constant. Initially the rotor is in the ground state of  $\hat{H}_0$ . At time  $t = 0$  a weak magnetic field is switched on in the  $x$ -direction, which subsequently decays in time. The Hamiltonian of the rotor is now

$$\hat{H}(t) = \hat{H}_0 + \beta \hat{L}_x \exp(-t/\tau),$$

where  $\beta$  and  $\tau$  are positive constants, with  $\beta \ll \alpha$ .

- i) Using first-order perturbation theory, calculate the transition probability  $P(t)$  to the first excited state of the rotor as a function of time. [7 marks]

- ii) Determine the long-time limit  $P(\infty)$  and find a condition on  $\tau$  to ensure its validity. [3 marks]

2. a) Consider a free relativistic spin-1/2 particle of mass  $m$ . Define the Dirac four-current  $j^\mu = (c\rho, \mathbf{j})$  and show that the Dirac equation implies the continuity equation,

$$\frac{\partial}{\partial t}\rho + \nabla \cdot \mathbf{j} = 0.$$

[4 marks]

- b) Now consider a relativistic electron moving in the  $x$ - $y$  plane with a uniform applied magnetic field in the  $z$ -direction, given by  $\mathbf{B} = (0, 0, B)$ .  
 i) Show that this magnetic field can be described by the vector potential  $\mathbf{A} = (0, xB, 0)$  and prove that the following identity holds for this potential:

$$[\boldsymbol{\sigma} \cdot (\hat{\mathbf{p}} + e\mathbf{A})]^2 = (\hat{\mathbf{p}} + e\mathbf{A})^2 + e\hbar\sigma_z B.$$

[4 marks]

- ii) The Hamiltonian for a non-relativistic electron in the same potential is

$$\hat{H}_{\text{NR}} = \frac{1}{2m}(\hat{\mathbf{p}} + e\mathbf{A})^2,$$

which has eigenvalues

$$E_n = \left(n + \frac{1}{2}\right) \hbar\omega_c.$$

Here,  $n = 0, 1, 2, \dots$  and  $\omega_c = eB/m$  is the cyclotron frequency.

Express the Dirac equation for this electron in terms of the upper components only. Find the energy eigenvalues by treating the upper components as simultaneous eigenstates of  $\hat{H}_{\text{NR}}$  and of the spin operator  $\hat{S}_z$ .

[8 marks]

- iii) For the weak-field case,  $\hbar\omega_c \ll mc^2$ , show that the energy levels of the relativistic electron are given by

$$E_n^\uparrow = mc^2 + (n+1)\hbar\omega_c, \quad E_n^\downarrow = mc^2 + n\hbar\omega_c, \quad n = 0, 1, 2, \dots$$

for spin-up and spin-down electrons, respectively.

[2 marks]

- c) A beam of massless relativistic particles with charge  $q = e$ , energy  $E_p$ , and in the spin-up eigenstate of  $\hat{S}_z$  (spinor  $\phi_1$ ) is moving in the positive  $z$ -direction and encounters an electrostatic step potential defined by

$$e\Phi(z) = \begin{cases} 0 & \text{for } z < 0, \\ V_0 & \text{for } z > 0. \end{cases}$$

Show that for  $E_p > V_0$  all particles are transmitted.

[7 marks]

3. a) Consider a non-relativistic spinless particle of charge  $q$  and mass  $m$  in the 4-potential  $A^\mu = (\Phi/c, \mathbf{A})$ .
- i) Write down the electromagnetic gauge transformation for the potentials  $\Phi$  and  $\mathbf{A}$  in terms of a differentiable function  $\lambda(\mathbf{r}, t)$ .

[2 marks]

- ii) The particle is placed in a homogenous oscillatory electric field of frequency  $\omega$  directed along the  $z$ -axis, with  $E_z = E_0 \cos(\omega t)$ . In two different gauges the electric field may be considered as arising purely from the potential  $\Phi$ , or purely from the potential  $\mathbf{A}$ . Write down the corresponding potentials and find a function  $\lambda(\mathbf{r}, t)$  that transforms from the first to the second. Given that  $|\psi_1\rangle$  is a solution to the time-dependent Schrödinger equation in the first case, show that

$$|\psi_2\rangle = \exp(iq\lambda(\mathbf{r}, t)/\hbar)|\psi_1\rangle,$$

is a solution in the second case.

[8 marks]

- b) Now consider a relativistic spinless particle of charge  $e$  that is bound to a nucleus by a Coulomb-like *scalar* potential  $S(r) = -Z\hbar c\alpha/r$ , where  $\alpha$  is the fine-structure constant.
- i) Write down the time-independent Klein-Gordon equation for the eigenstates of the particle with energy  $E$ .

[2 marks]

- ii) By considering an eigenstate of total angular momentum with quantum number  $l$ , show that the eigenvalue equation for the radial part  $\psi(r)$  of the wave function of the particle can be written as

$$\left[ 2mE' + \hbar^2 \left( \frac{1}{r} \frac{\partial^2}{\partial r^2} r - \frac{l'(l'+1)}{r^2} \right) + \frac{2m\hbar c Z \alpha}{r} \right] \psi(r) = 0. \quad (1)$$

Express  $E'$  and  $l'(l'+1)$  in terms of  $E$  and  $l$ .

[6 marks]

- iii) The eigenvalues of equation (1) can be written in the form

$$E' = -\frac{1}{2}mc^2 \left( \frac{Z\alpha}{n'} \right)^2,$$

where  $n'$  is related to  $l$  and the principal quantum number  $n = 1, 2, 3 \dots$  by

$$n' = n - (l + 1/2) + \sqrt{(l + 1/2)^2 + (Z\alpha)^2}.$$

Find the energy levels  $E_{n,l}$  of the particle and discuss whether or not there are any restrictions on the allowed values of  $Z\alpha$ . Considering the limit  $Z\alpha \ll 1$  find the first *relativistic* energy correction for the case  $l = 0$ .

[7 marks]