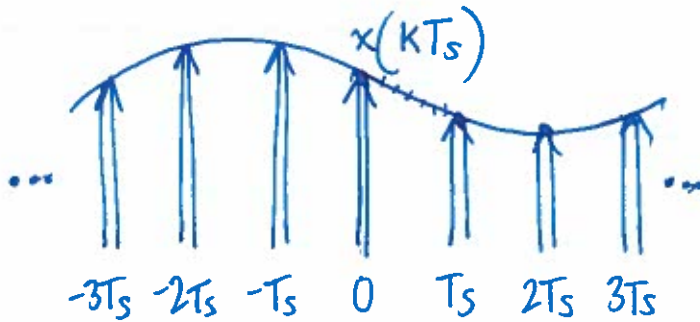


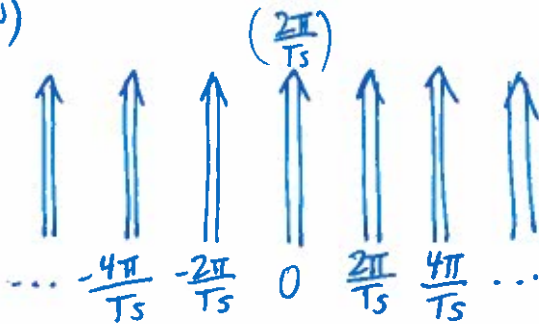
Signals and Systems PS08

Ryan Louie

1. a) $x_p(t)$. Sample $x(t)$ at each $p(t)$



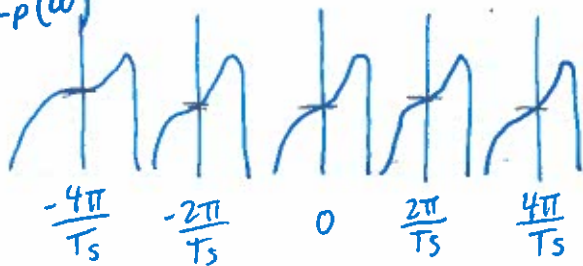
- b) $P(\omega)$



time	freq
$\delta(t - t_0)$	$e^{-j\omega t_0}$

$\omega_s = \frac{2\pi}{T_s}$?

- c) $X_p(\omega)$

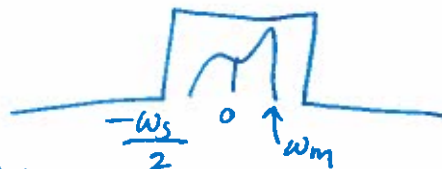


$$X(\omega) \cdot P(\omega) \cdot \frac{1}{2\pi}$$

- d) $\frac{2\pi}{T_s} - \omega_m > \omega_m$

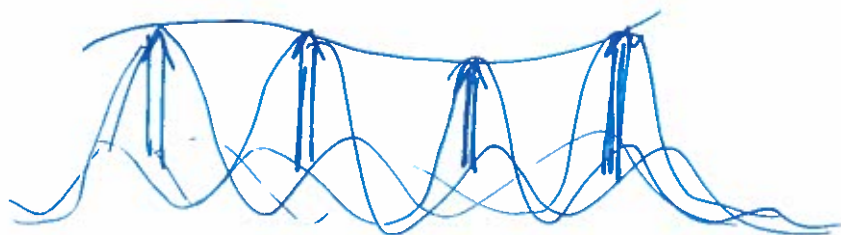
$$\frac{2\pi}{T_s} > 2\omega_m$$

$$2\omega_m < \omega_s$$



- e) Multiplying by an ideal low pass filter, corresponds to convolution of a Sinc function in the time domain.

Sinc functions ~~involved w/~~ ~~multiplied by~~ each impulse in the impulse train

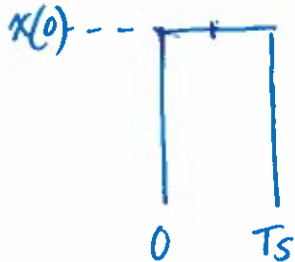
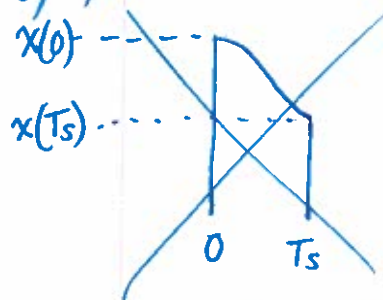


$$X(t) \cos(\omega_c t) \cos(\omega_c t) \rightarrow$$

$$\text{LPF} \rightarrow$$

$$\approx X(t)$$

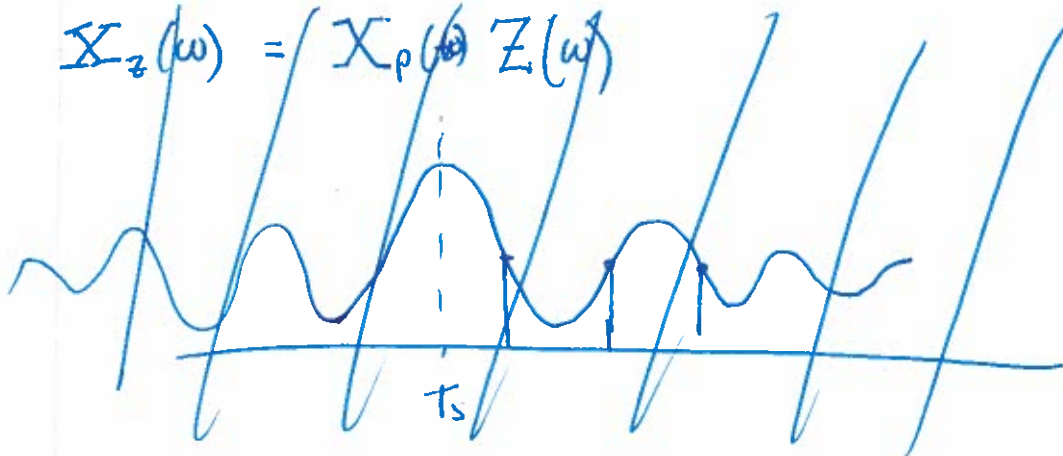
g) ~~g)~~ $X_z(t) = X_p * z(t)$



Zero order hold construction



h) $X_z(\omega) = X_p(\omega) Z(\omega)$

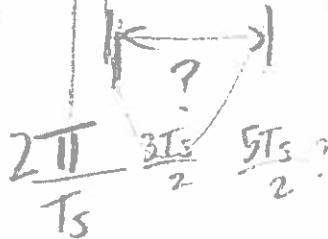


$$\begin{aligned}
 h) \quad X_z(\omega) &= \int_{-\infty}^{\infty} x(t) e^{-j\omega t} dt \\
 &= \int_{-T/2}^{T/2} x(t) e^{-j\omega t} dt
 \end{aligned}$$

$$\text{let } u = t + T/2, \quad t = u - T/2$$

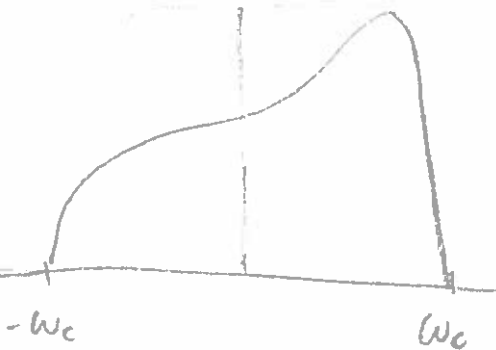
$$= \int_0^T x(u - T/2)$$

$$T_s/2$$



i)

$$\hat{X}(\omega) = X_p(\omega) H(\omega)$$

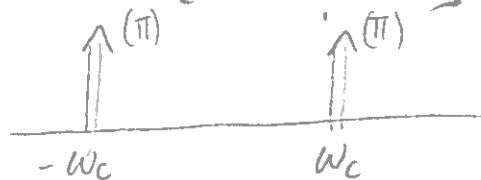


From the table of Canonical forms

2. a) $y(t) = x_1(t) \cos(\omega_1 t) + x_2(t) \cos(\omega_2 t)$

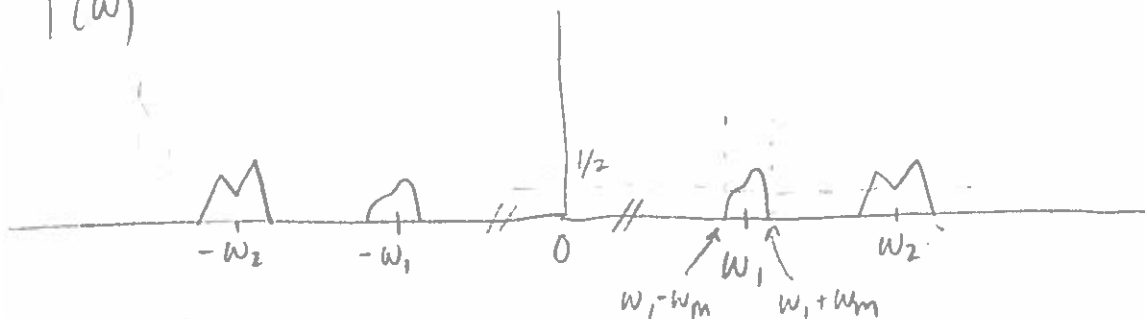
$$Y(\omega) = \frac{1}{2\pi} \left[X_1(\omega) \cdot \text{FT} \{ \cos(\omega_1 t) \} + X_2(\omega) \cdot \text{FT} \{ \cos(\omega_2 t) \} \right]$$

$$\text{FT} \{ \cos(\omega_c t) \}$$

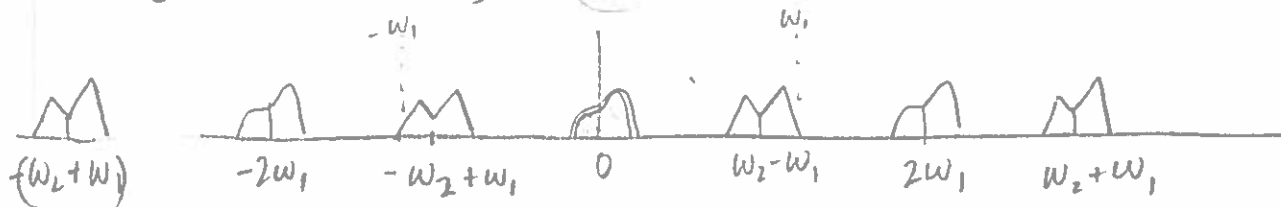


$$\begin{aligned} \omega_1 &\gg \omega_m \\ \omega_2 &\gg \omega_m \\ \omega_1 + 2\omega_m &< \omega_2 \end{aligned}$$

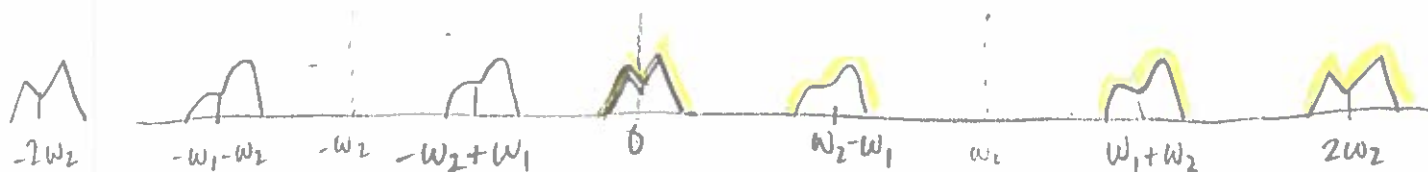
$$Y(\omega)$$



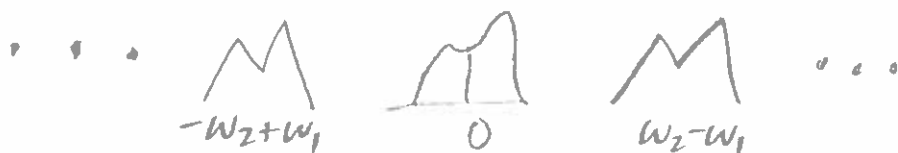
b) $\text{FT} \{ y(t) \cos(\omega_1 t) \} = Y(\omega) \cdot \text{FT} \{ \cos(\omega_1 t) \}$



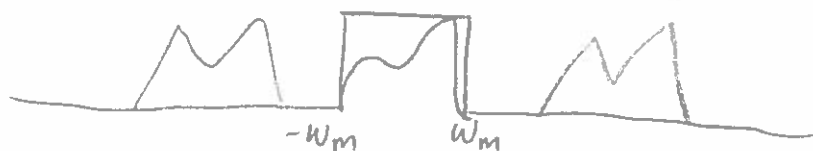
$$\text{FT} \{ y(t) \cos(\omega_2 t) \} = Y(\omega) \cdot \text{FT} \{ \cos(\omega_2 t) \}$$



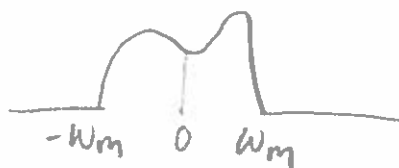
- c) To recover $x_1(t)$ from $y(t)$, we can convolve $y(t)$ w/ $\cos(\omega_1 t)$. This results in the wave which, in the frequency domain, looks like 2bI:



If we then apply an ideal low pass filter, with cutoff frequency ω_m ,

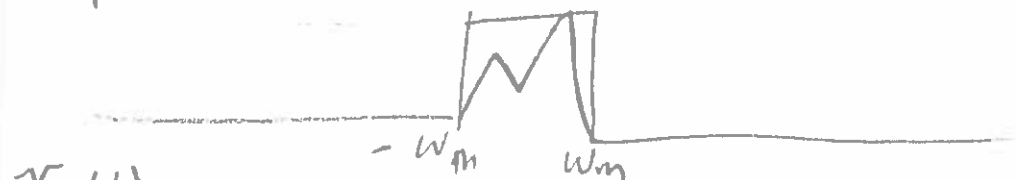
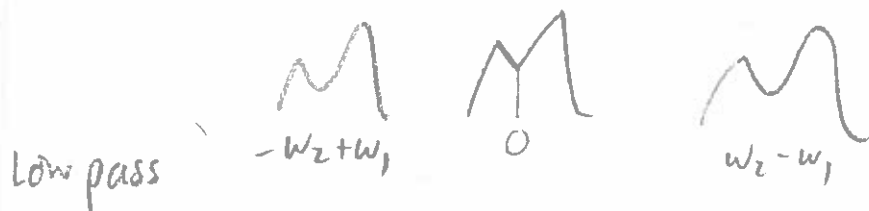


We are left w/ $X_1(t)$



which can easily be turned into $x_1(t)$.

Similarly, $y(t) \cdot \cos(\omega_2 t)$ looks like 2bII in frequency domain



$X_2(t)$



which can give us $x_2(t)$

3.

a)

 $V_{in}(t)$  $C = V_{out}$

$$i(t) = C \frac{d}{dt} V_{out}(t)$$

$$V_L = L \frac{d}{dt} i(t)$$

$$V_{in}(t) = V_R + V_L + V_{out}(t)$$

$$V_{in}(t) = R i(t) + L \frac{d}{dt} i(t) + V_{out}$$

$$V_{in}(t) = RC \frac{d}{dt} V_{out}(t) + LC \frac{d^2}{dt^2} V_{out}(t) + V_{out}(t)$$

b) $H(\omega) = \frac{V_{out}(t)}{V_{in}(t)}$ i.e. continuing w/ 3a...

$$V_{in}(t) = e^{j\omega t}$$

$$V_{out}(t) = H(\omega) e^{j\omega t}$$

$$e^{j\omega t} = \left[RC j\omega e^{j\omega t} + LC j^2 \omega^2 e^{j\omega t} + e^{j\omega t} \right] H(\omega)$$

$$H(\omega) = \frac{1}{RC j\omega - LC \omega^2 + 1}$$

$$c) \|H(\omega)\| = \sqrt{\frac{2}{C^2 L^2 \omega^4 + j(2CR\omega - 2C^2 LR\omega^3) - C^2 R^2 \omega^2 - 2CL\omega^2 + 1}}$$

thanks Wolfram!

$$= \sqrt{2} \cdot \frac{1}{-1 + C\omega(L\omega - jR)}$$

thanks Wolfram again!

thanks Wolfram

$$d) \frac{\partial}{\partial \omega} (|H(\omega)|) \stackrel{!}{=} \frac{-\sqrt{2}C(R+2iL\omega)}{[C\omega(R+iL\omega)-i][-1+C\omega(L\omega-iR)]} = 0$$

$$-\sqrt{2}C(R+2iL\omega) = 0$$

$$2iL\omega = -R$$

$$\omega = \frac{-R}{i2L}$$

uh oh, something went wrong Wolfram lied i
I will try again w/ C.

$$c) |H(\omega)| = \frac{1}{\sqrt{\omega^2 R^2 C^2 + (1 - \omega^2 LC)^2}}$$

d) minimize denom of $|H(\omega)|$ to maximize $|H(\omega)|$

$$0 = \frac{d}{d\omega} (\omega^2 R^2 C^2 + (1 - \omega^2 LC)^2)$$

$$0 = [2R^2 C^2]\omega + 2(1 - \omega^2 LC) \cdot (-2\omega)$$

$$0 = [2R^2 C^2]\omega - 4\omega + 4LC\omega^2$$

$$\omega = \frac{4 \pm \sqrt{16 - 4(4LC)(2R^2 C^2)}}{2(4LC)}$$

$$\omega = \frac{4 \pm \sqrt{16 - 16(2LR^2 C^3)}}{8LC}$$

cal pts.

$$\rightarrow \omega = \frac{1 \pm \sqrt{1 - 2LR^2 C^3}}{2LC}$$

when $\omega = \frac{1 + 4\sqrt{1 - 2LR^2C^3}}{2LC}$

denom of $H(\omega)$ is [thanks Wolfram]

Intermediate helps

$$\omega^2 = \frac{1 + 2[4\sqrt{1 - 2LR^2C^3}] + 16(1 - 2LR^2C^3)}{4L^2C^2}$$

$$\omega^2 = \frac{1 + 8\sqrt{1 - 2LR^2C^3} + 16(1 - 2LR^2C^3)}{4L^2C^2}$$

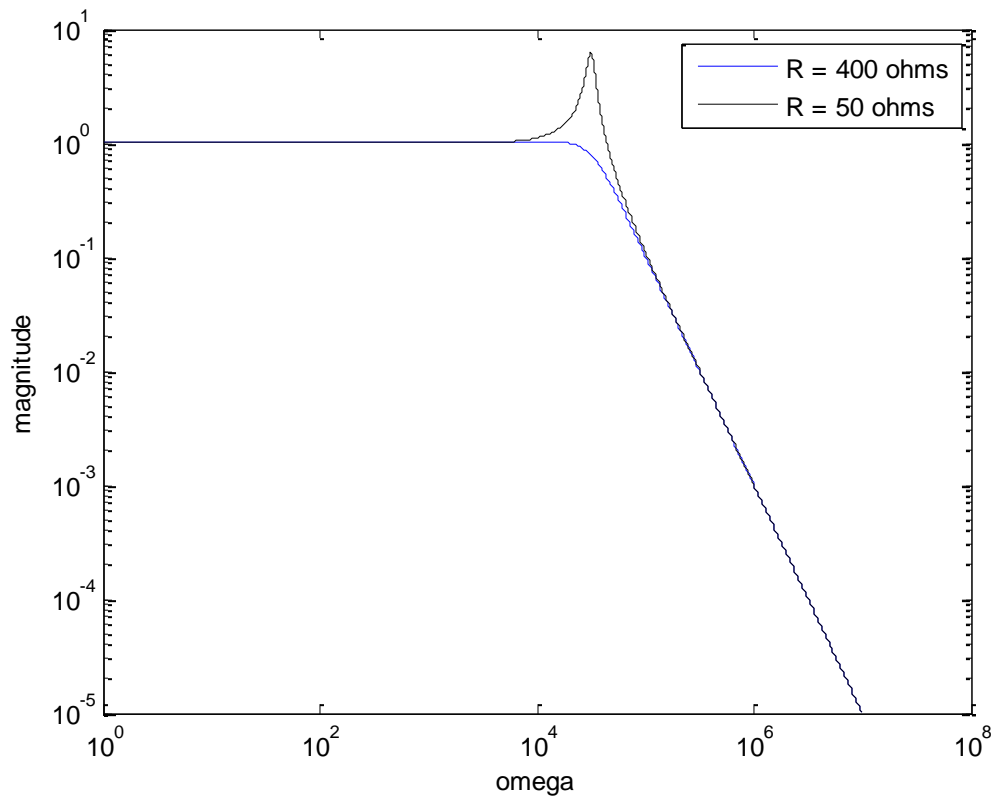
let $u = 1 - 2LR^2C^3$

$$\begin{aligned}\omega^4 &= 1 + 8\sqrt{u} + 16u + 64u + 256u\sqrt{u} + 16u + 16^2u^2 \\ &= 1 + 8\sqrt{u} + 96u + 256u\sqrt{u} + 256u^2\end{aligned}$$

These pluses are going to make denom large;
thus, its the other critical point that we want.

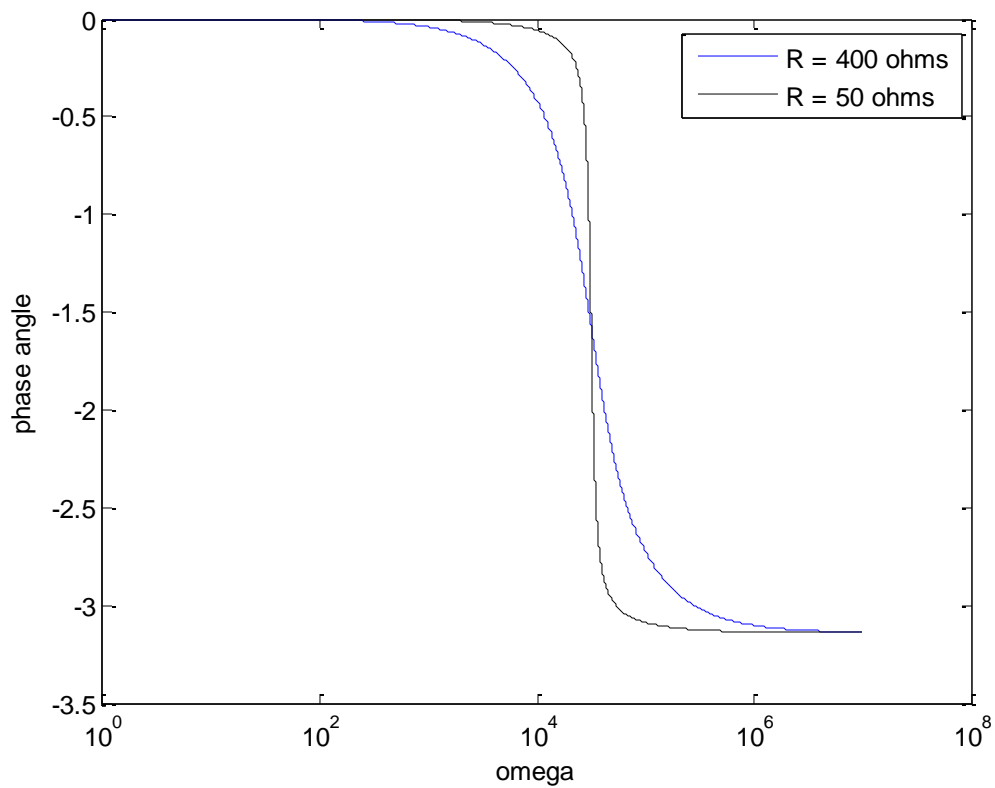
$$\boxed{\omega = \frac{1 - 4\sqrt{1 - 2LR^2C^3}}{2LC}}$$

NOTE: This was done as a team by Ryan and Jennifer



```
C = 10^-7; %F
L = 10^-2; %H
R1 = 400; %ohms
R2 = 50; %ohms
```

```
w = 1:1:10000000;
Mag1 = (1.0./ ( (w*R1*C).^2 + (1-(w.^2)*L*C).^2 ).^(0.5));
Mag2 = (1.0./ ( (w*R2*C).^2 + (1-(w.^2)*L*C).^2 ).^(0.5));
loglog(w,Mag1,'b',w,Mag2,'k');
legend('R = 400 ohms', 'R = 50 ohms');
xlabel('omega');
ylabel('magnitude');
```



```
figure;
Freq1 = 1.0./(1+(R1.*C.*j.*w)-((w.^2).*C.*L));
Freq2 = 1.0./(1+(R2.*C.*j.*w)-((w.^2).*C.*L));
phase1= angle(Freq1);
phase2 = angle(Freq2);
%phase1 = atan((j.*w.*R1.*C)./(-(w.^2).*L.*C).^2).*(360/2*pi);
%phase2 = atan((j.*w.*R2.*C)./(-(w.^2).*L.*C).^2).*(360/2*pi);
semilogx(w,phase1,'b',w,phase2,'k');
legend('R = 400 ohms', 'R = 50 ohms');
xlabel('omega');
ylabel('phase angle');
```