B03 Highly accurate simulation of electrokinetic transport processes in nanopores



Abstract

Presented is a multi-component mixture, single phase solver of ionically driven flow with the Finite Volume Method implemented in OpenFoam®. A flexible inheritance structure allows for further extensions of the Poisson-Nernst-Planck model, and quickly modifying case configurations based on dimensional and/or non-dimensional specifications at runtime. Implicit treatment of the electromigration flux provides numerical stability, while special boundary condition ensures better conservation of species flux and more accurate species concentration profiles at the boundary cells. The transport dynamics of this solver is preemptively verified with species transport between an instantaneously equilibrated reservoir and nanopore entrance.

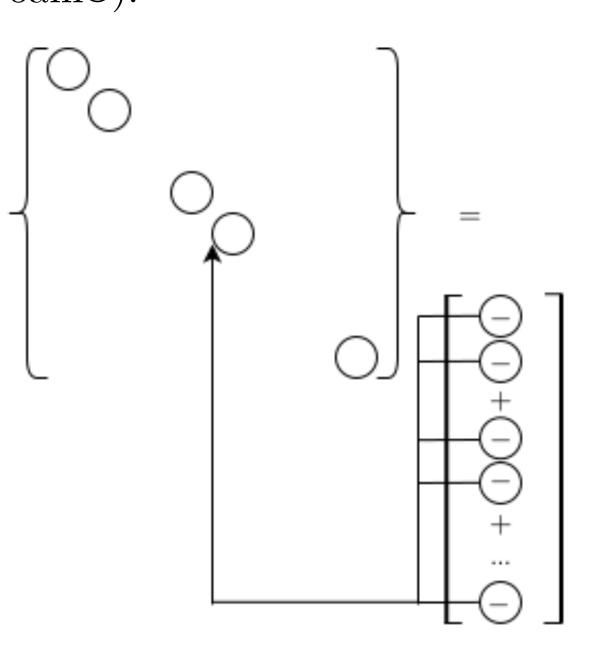
Numerical Implementation

Coupling Algorithm

Segregated solver stabilization with multiple nested iterations [1].

Electromigration Flux discretization

Implicit discretization of the negative terms in electromigration flux (supported natively within Open-Foam®).



Poisson solve iterations reached? Solve Poisson Equation Update electric potential field Solve species transport equation: Update concentration fields Solve pressure and concentration fields based on last time step No Poisson solve iterations reached? Yes Solve species transport equation: Update concentration fields Solve pressure field,

Boundary Conditions

Zero Ionic Flux

Neumann-type boundary condition ensuring no species flux on the domain wall boundaries:

$$\left[\nabla \cdot c_i\right]_f \cdot n_{\perp} \mathbf{A}_f = -\left[\frac{\mu_i}{D_i} \left(\nabla \Psi\right) c_i\right]_f \cdot \mathbf{A}_f$$

or

$$\nabla c_{i,P} + \frac{\mu_i}{D_i} \left(\nabla \Psi_f \right) c_{i,P} = 0$$

Subgrid Zero Ionic Flux

Dirichtlet-type boundary condition enforcing both no species flux on the domain walls, and exponential species distribution in the boundary cell. Proposed by [1] and first implemented in the rheoTool toolbox.

$$c_{i,P}^P = c_{i,f}^S \exp(\frac{\mu_i}{D_i} (\Psi_f - \Psi_P))$$

λ_D $c_{i,f}$ $c_{i,P}$ $c_{i} = f(x,z)$

velocity field

Outlook

- Verify late-time charging dynamics with Transmission Line Model. Doing so would require significant speedup of solver or adaptative time stepping.
- Verify early-times concentration and potential field solutions with analytical approximations provided by Henrique [5].
- Compare agreement of late-time charging dynamics with different modifications to the Poisson-Nernst-Planck equations: current discussions in literature regarding species transport dynamics lack verification with late-time charging.

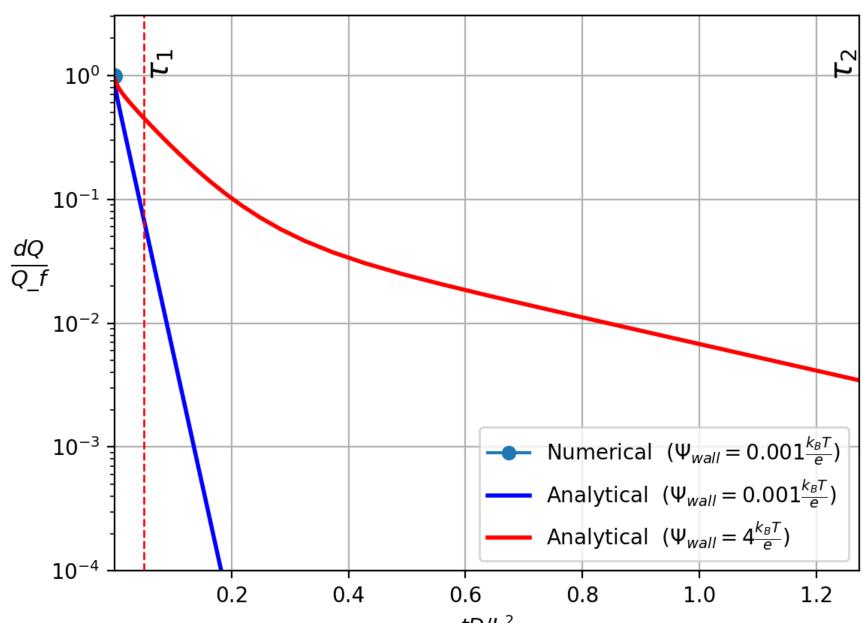


Fig.3 Analytical late-times charging rate, as compared with available early-times results.

Mathematical Modelling

Standard Poisson-Nernst-Planck equation, used in literature to derive the Transmission Line Model [2] and analytical solutions of nanopore charging [2,3], along with numerical simulations they are benchmarked against [4].

$$\int_{V(t)} \frac{\partial c_i}{\partial t} d\mathbf{x} = \int_{V(t)} \underbrace{\nabla \cdot (D_i \nabla c_i)}_{\text{Diffusive flux}} d\mathbf{x}$$

$$+ \int_{V(t)} \nabla \cdot \left[\underbrace{(D_i \frac{ez_i}{kT} \nabla \Psi) c_i}_{\mu_i} \right] d\mathbf{x}$$
Electromigration flux

$$\int_{V(t)} \nabla \cdot (\epsilon \nabla \Psi) d\mathbf{x} = \int_{V(t)} \underbrace{-\rho_E}_{-F \sum_{i=1}^{m} z_i c_i} d\mathbf{x}$$

Currently available extensions:

- Additional advection term: $\mathbf{u} \cdot \nabla c_i$
- Decomposed electric potential [4]
 - Poisson equations for intrisic and external contributions to the electric potential field:

$$\nabla \cdot (\epsilon \nabla \Psi) = \nabla \cdot (\epsilon \nabla \begin{bmatrix} \Psi_{Ext} \\ \Psi_{int} \end{bmatrix}) = -\rho_E$$

$$\longrightarrow \begin{cases} \nabla \cdot (\epsilon \nabla \Psi_{Ext}) = 0 \\ \nabla \cdot (\epsilon \nabla \Psi_{int}) = -\rho_E \end{cases}$$

- Electrical Force: $\mathbf{f}_E = -\rho_E \Psi_{Ext}$

Validation Test Cases

2D rectangular pore

Following Aslyamov's semi-infinite (close-ended) pore analysis [2].

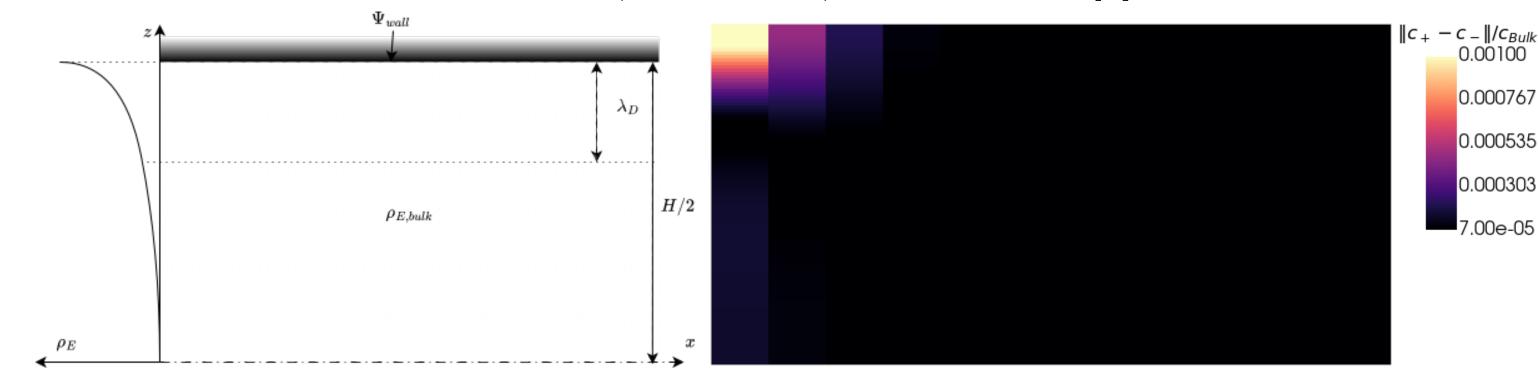


Fig.1 Sketch of case setup and simulation results for H = 3.77e - 06m, $D = 1e - 9m/s^2$, $\lambda_D = 2.97e - 07m$, $c_{Bulk} = 0.001M$, $\Psi_{wall} = 0.025 \frac{k_B T}{e}$

3D cylindrical pore

Simulation of a finite pore by Gupta [4]. Result analytically compared by Henrique [5], who also derived early-times solution of field quantities.

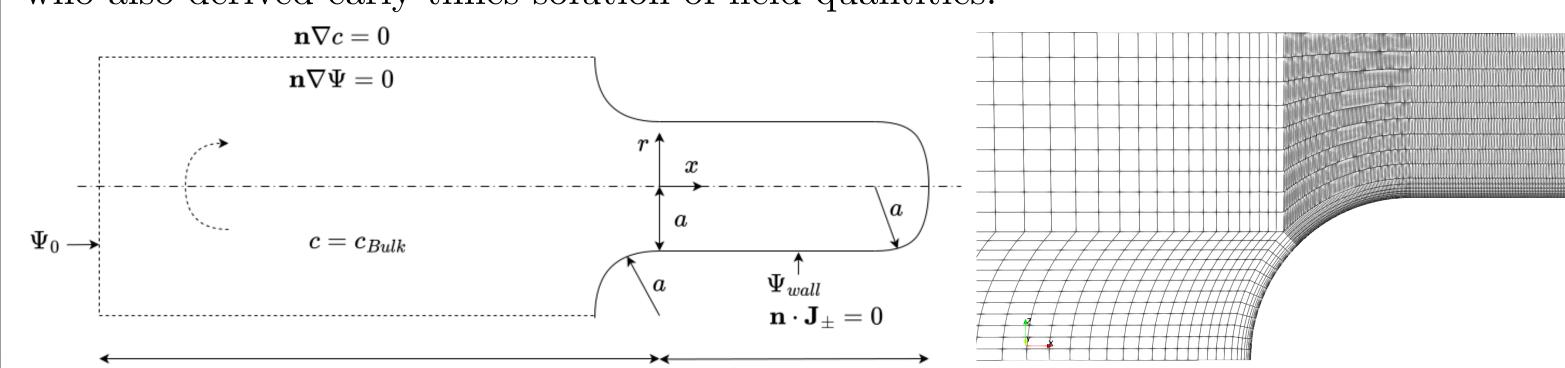


Fig.2 Sketch of case setup and mesh geometry details at pore opening (work in progress)

References

- [1] Francisco Pimenta and Manuel A. Alves. Numerical simulation of electrically-driven flows using openfoam, 2018.
- [2] Mathijs Janssen. Transmission line circuit and equation for an electrolyte- filled pore of finite length. Physical Review Letters, 126(13), March 2021.
- 3] Timur Aslyamov and Mathijs Janssen. Analytical solution to the poisson- nernst-planck equations for the charging of a long electrolyte-filled slit pore, 2022.
- Ankur Gupta, Pawel J. Zuk, and Howard A. Stone. Charging dynamics of overlapping double layers in a cylindrical nanopore. Physical Review Letters, 125(7), August 2020.
- [5] Filipe Henrique, Pawel J. Zuk, and Ankur Gupta. Charging dynamics of electrical double layers inside a cylindrical pore: predicting the effects of arbitrary pore size. Soft Matter, 18(1):198213, 2022.