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POD 20: Designing a Biodelector

- For large systems diffusion is sufficiently slow that we don't ever get to a point where the SL asymptote is reached!
- For small (microfluidic) systems this isn't true! This is important in things like biodelectors: ident. small conc. of molecules

Suppose we have a microchannel 100 μm deep by 200 μm wide

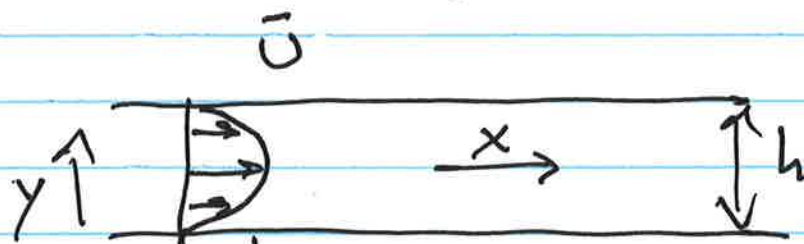
We have a flow rate of 0.2 $\mu\text{L}/\text{min}$

If our detector is 200 μm wide by 500 μm long, what is the max diffusivity that 50% of entering solute molecules would reach detector?

The easy way of doing this is determining how much remains at the end.

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That way we are dominated by the lead eigenfunction!



$$6 \bar{U} \frac{y}{h} \left(1 - \frac{y}{h}\right) \frac{\partial C}{\partial x} = D \frac{\partial^2 C}{\partial y^2}$$

$$C|_{x=0} = C_0 \quad C|_{y=0} = 0 \quad \frac{\partial C}{\partial y}|_{y=h} = 0$$

$$\text{so: } y^* = \frac{y}{h}, \quad C^* = \frac{C}{C_0}, \quad x^* = \frac{x}{x_c}$$

$$\therefore 6 y^* (1 - y^*) \frac{\partial C}{\partial x^*} \frac{\bar{U} C_0}{x_c} = D \frac{C_0}{h^2} \frac{\partial^2 C^*}{\partial y^{*2}}$$

Divide out:

$$6 y^* (1 - y^*) \frac{\partial C^*}{\partial x^*} = \left[\frac{D x_c}{\bar{U} h^2} \right] \frac{\partial^2 C^*}{\partial y^{*2}}$$

$$x_c = \frac{\bar{U} h^2}{D}$$

$$\text{Define } L^* \equiv \frac{L}{x_c} = \frac{L D}{\bar{U} h^2}$$

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We seek L^* s.t. the flow avg of c^* at $L^* = \frac{1}{2}$

So:

$$6y^*(1-y^*)\frac{\partial c^*}{\partial x^*} = \frac{\partial^2 c^*}{\partial y^{*2}}$$

$$c^*|_{x^*=0} = 1, \quad c^*|_{y^*=0} = 0, \quad \frac{\partial c^*}{\partial y^*}|_{y^*=1} = 0$$

We have already gotten homogeneous BC's!

$$\therefore c^* = G(x^*)F(y^*)$$

$$\frac{G'}{G} = \frac{F''}{6y^*(1-y^*)F} = -\sigma^2$$

$$\therefore G = e^{-\sigma^2 x^*}, \quad F'' + \sigma^2 6y^*(1-y^*)F = 0$$

$$F(0) = 0, \quad F'(1) = 0$$

To get an analytic estimate,

we replace parabolic velocity profile w/ average!

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$$\int_0^1 6y^*(1-y^*) dy^* = 1$$

so our approximate problem is:

$$F'' + \sigma^2 F = 0 \quad F(0) = 0 \quad F'(1) = 0$$

$$\therefore F = A \sin \sigma y^*, \quad \sigma = \left(n - \frac{1}{2}\right) \pi$$

$$A_n = \frac{\int_0^1 \sin \sigma y^* dy^*}{\int_0^1 \sin^2 \sigma y^* dy^*} = \frac{-\frac{1}{\sigma} \cos \sigma y^* \Big|_0^1}{\int_0^1 \sin^2 \sigma y^* dy^*}$$

$$= \frac{+2}{\left(n - \frac{1}{2}\right) \pi}$$

$$\therefore C_{app}^* = \sum_{n=1}^{\infty} \frac{2}{\left(n - \frac{1}{2}\right) \pi} e^{-\left(\left(n - \frac{1}{2}\right) \pi\right)^2 x^*} \sin \left(\left(n - \frac{1}{2}\right) \pi y^*\right)$$

Now we want $\frac{\int_0^h C|_{x=L} dy}{\int_0^h u dy} = \frac{C_0}{2}$

$$\text{so: } \int_0^1 C^* dy^* = \sum_{n=1}^{\infty} \frac{2}{\sigma_n^2} e^{-\sigma_n^2 x^*}$$

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We are dominated by the lead eigenvalue by the time $\tau^* \rightarrow \frac{1}{2}$

$$\text{So: } \frac{2}{\left(\frac{\pi^2}{4}\right)} e^{-\frac{\pi^2}{4} L^*} = \frac{1}{2} \quad (\text{target})$$

$$\therefore e^{-\frac{\pi^2}{4} L^*} = \frac{\pi^2}{16}$$

$$\text{or } L^* = \frac{4}{\pi^2} \ln\left(\frac{16}{\pi^2}\right) = 0.196 \approx \underline{0.2}$$

So :

$$0.2 = \frac{LD}{\bar{U} h^2}$$

$$\text{and } \bar{U} = \frac{Q}{Wh}$$

$$\therefore D_{min} = 0.2 \frac{\bar{U} h^2}{L} = 0.2 \frac{Qh}{WL}$$

$$\text{Now } Q = 0.2 \times 10^{-3} \text{ cm}^3/\text{min} = 3.3 \times 10^{-6} \text{ cm}^3/\text{s}$$

$$h = 0.01 \text{ cm}, \quad W = 0.02 \text{ cm}$$

$$\therefore \bar{U} = 0.0165 \text{ cm/s}$$

(for a 1 cm channel, this takes ~ 1 min to clear - a bit long)

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So we get

$$D_{\text{min}} = (0.2) \frac{(3.3 \times 10^{-6} \text{ cm}^3/\text{s})(10^{-2} \text{ cm})}{(2 \times 10^{-2} \text{ cm})(5 \times 10^{-2} \text{ cm})}$$
$$= 6.66 \times 10^{-6} \text{ cm}^2/\text{s}$$

This means it would only work (for these conditions) for a fairly small molecule.

How does D depend on the size of a molecule?

$$\text{Stokes-Einstein D.f} \equiv \frac{kT}{6\pi\eta a}$$

$$\text{Now } V = \frac{4}{3}\pi a^3 = \frac{MW}{\rho} \frac{1}{N_{\text{AV}}} \quad \text{where } N_{\text{AV}} \rightarrow 6.02 \times 10^{23}$$

$$\therefore MW = \frac{4}{3}\pi \left(\frac{1}{6\pi}\right)^3 \rho \left(\frac{kT}{\eta D}\right)^3 \times 6.02 \times 10^{23}$$
$$= \frac{1}{\pi^2} \frac{1}{27 \times 6} (1) \left(\frac{1.38 \times 10^{-16} \times 300}{(0.01)(6.7 \times 10^{-6})}\right)^3 \times 6.02 \times 10^{23}$$
$$= 89 \frac{\text{g}}{\text{mol}} \quad (\text{say, fairly small molecule - although the formula is off for this.})$$

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What if we need to catch a larger molecule? We have the scaling!

$$MW_{\max} \sim D_{\max}^{-3}$$

$$D_{\max} \sim \frac{Qh}{WL}$$

$$\therefore MW_{\max} \sim \left(\frac{WL}{Qh} \right)^3 !$$

If we increase L from $500\mu m$ to $5mm$ we would increase MW_{\max} by 10^3 !

This, for other parameters the same, yields $MW_{\max} \sim 90kDa$
this is comparable to, say, BSA.

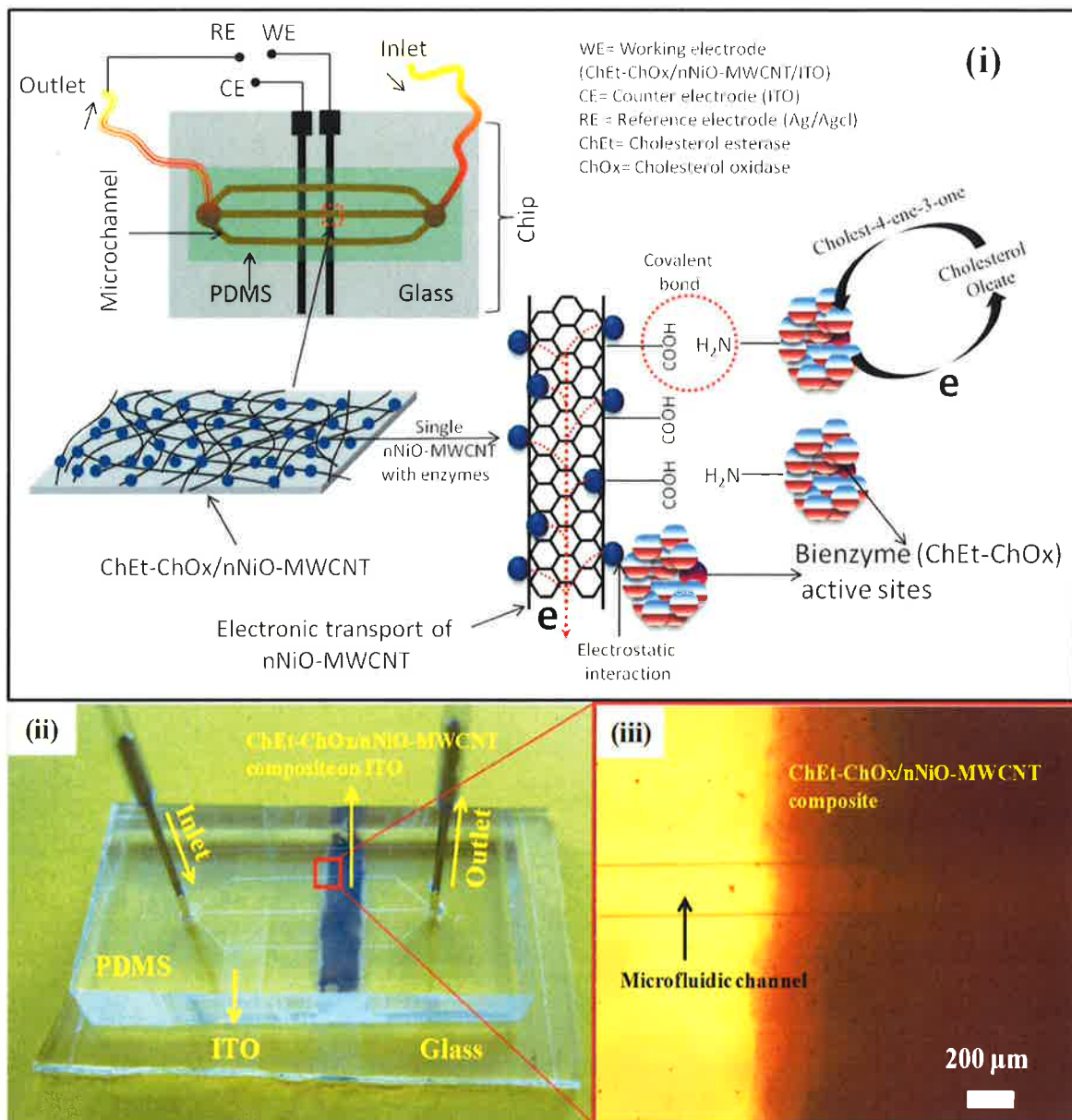


Figure 1 | (i) The schematic of the microfluidic biochip used for total cholesterol detection (the ordered arrangement of this microsystem is assumed). (ii) The photograph of real microfluidic biochip for cholesterol detection and (iii) the enlarged view of optical microscopic image of the microfluidic biochip.

open to reactions¹⁴. The MWCNTs are known to produce changes in energy bands close to the Fermi level^{15,16}. The exciting electronic properties and high electrochemical reactivity of MWCNTs suggest that fast electron transfer reaction occurs when they are used as the electrode in an electrochemical biochip^{15,16}. Lin et al. have developed a microfluidics electrochemical sensor for on-site, non-invasive monitoring of lead and chlorophenols¹⁷. Wisitsoraat et al. have developed an electrochemical biochip for cholesterol detection that has a sensitivity of 0.0512 nA/mg/dl, which is attributed to the direct growth of CNT on glass¹⁸. However, MWCNTs are known to agglomerate via *Van der Waals* interactions, resulting in poor film-forming ability. To overcome this problem, nanostructured metal oxides (NMOx) may be used to control the agglomeration of MWCNTs¹⁹. The covalent binding (or sidewall functionalization) of biomolecules (e.g., proteins, enzymes, and nucleic acids) to carboxyl-functionalized MWCNTs via diimide-activated amidation may provide improved stability and reproducibility^{20–24}. In such a case, the

large surface area of the MWCNTs and the presence of abundant functional groups may offer a suitable platform for biofunctionalization^{20–25}. Additionally, MWCNTs may facilitate continuous conducting pathways to transport the charge carriers, allowing for a higher sensitivity²⁵. Shim et al. have used functionalized CNT for biomolecular recognition in a streptavidin/biotin approach to investigate the adsorption of proteins on the sidewalls of carbon nanotubes²⁰.

A biosensor based on nanostructured nickel oxide (nNiO) has recently been explored to detect biomolecules such as DNA, antibody-antigen interactions, glucose, and cholesterol^{12b,27}. However, nNiO-based biosensors have limited applications due to the inherently poor electrical conductivity of nNiO²⁶. The non-covalent immobilization of enzymes onto nNiO-based biochip has recently been found to result in poor stability of the desired biomolecules^{26–28}. To improve the characteristics of a biosensing device, nNiO can be integrated with MWCNTs^{29,30}. Zhang et al. have used CNT-NMOx to develop solar cells and gas sensors²¹. The sp^2 hybridization and