## $\mathbf{QNC}_1 \subset \mathbf{noisy}\mathbf{-BQP}$

# Michael Benor David Ponarovsky September 18, 2024

## 1 Notations.

We denote by  $C_g$  the good qLDPC code [Din+22] [PK21] [LZ22b], and by  $C_{ft}$  the concatenation code presented at [AB99] (ft stands for fault tolerance). For a code  $C_y$ , we use  $\Phi_y$ ,  $E_y$ ,  $D_y$  to denote the channel maps circuits into the circuits computed in the code space, the encoder, and the decoder, respectively. We use  $\Phi_U$  to denote the 'Bell'-state storing the gate U. We say that a state  $|\psi\rangle$  is at a distance d from a quantum code C if there exists an operator U that sends  $|\psi\rangle$  into C such that U is spanned on Paulis with a degree of at most d. Sometimes, when the code being used is clear from the context, we will say that a block B of qubits has absorbed at most d noise if the state encoded on B is at a distance of at most d from that code.

## 2 The Noise Model

## 3 Fault Tolerance (With Resets gates) at Linear Depth.

**Claim 3.1.** There exists a value  $p_{th} \in (0,1)$  such that if  $p < p_{th}$ , then any quantum circuit C with a depth of D and a width of W can be computed by a p-noisy circuit C', which allows for resets. The depth of C' is at most  $\max \{O(D), O(\log(WD))\}$ .

#### 3.1 Initializing Magic for Teleportation gates and encodes ancillaries.

The Protocol:

- 1. Initialization of zeros: The qubits are divided into blocks of size |B|. Each block is encoded in  $C_g$  using  $D_{ft}\Phi_{ft}[E_g]|0^{|B|}\rangle$ .
- 2. Initialization of Magic for Teleportation gates: The gates in the original circuit are encoded in  $C_g$  using  $D_{ft}\Phi_{ft}[E_g]|\Phi_U\rangle$ .
- 3. Gate teleportation: Each gate in the original circuit is replaced by a gate teleportation.
- 4. Error reduction: After the initialization step, at each time tick, each block runs a single round of error reduction.

**Claim 3.2** (From [LZ22a]). Assuming that an error  $|e| = \gamma n$ , i.e e is supported on less than  $\gamma n$  bits, then a single correction round reduce e to an error e' such that  $|e'| < \nu |e|$ .

Claim 3.3. The gate  $D_{ft}\Phi_{ft}[E_g]$  initializes states encoded in  $C_g$  subject to 3p-noise channel.

Proof. Clearly  $\Phi_{ft}[E_g]$  success, with high probability, let's say  $1-\frac{1}{poly(n)}$ , to encode in to  $C_{ft}\circ C_g$ . Denote by  $E_i,D_i$  the encoder and the decoder at the ith level of the concatination construction. Recall that by definition  $D_iE_i=I$ , or in other words  $D_i=E_i^{\dagger}$ . Consider the decoder under  $\mathcal N$  action.  $P_2D_1P_2D_2,...,P_{i-1}D_iP_i$ , by the fault-tolerance construction a logical error happens at the ith stage occurs with probability  $p^{2^i}$ , therefore by the union bound the probability that in one of the steps the circuit absorbs

an error that is not corrected is less than  $p + p^2 + p^4 + ... < 2p$ . Hence any decoded qubit absorbs a noise with probability less than 2p.

Thus in overall we can bound the porobability a single qubit to be faulty by:

$$\begin{split} \mathbf{Pr}\left[\text{fault}\right] &= \mathbf{Pr}\left[\text{fault}|\Phi_{ft}[E_g]\right] \cdot \mathbf{Pr}\left[\Phi_{ft}[E_g]\right] + \mathbf{Pr}\left[\text{fault}|\overline{\Phi_{ft}[E_g]}\right] \cdot \mathbf{Pr}\left[\overline{\Phi_{ft}[E_g]}\right] \\ &\leq \mathbf{Pr}\left[\text{fault}|\Phi_{ft}[E_g]\right] + \mathbf{Pr}\left[\overline{\Phi_{ft}[E_g]}\right] \leq 2p + \frac{1}{poly(n)} \leq 3p \end{split}$$

**Remark 3.1.** In our construction we use the concatinate-code to encode  $\log(n)$ -length block, Thus any poly(n) in the above should be replaced by  $\log(n)$ . Yet it doesn't effect anything since the inequality dosn't depend on n.

**Claim 3.4.** With probability  $1 - \frac{WD}{|B|} \cdot D2e^{-2|B|(\beta-p)}$ , the total amount of noise been absorb in a block, in any time t, is less than  $\gamma n$ .

*Proof.* Consider the ith block, denoted by  $B_i$ . Using the Hoeffding's inequality we have that the probability that more than  $\beta|B|$  bits are flipped at time t is less than  $\leq 2e^{-2|B|(\beta-p)}$ . Using the union bounds over all the blocks at all the different time locations we get that with probability  $1 - \frac{WD}{|B|} \cdot D2e^{-2|B|(\beta-p)}$  the noise been absorbed in a block is less than  $|\beta|B$  for the whole computation.

Denote by  $X_t$  the support's size of the error over  $B_i$  at time t. Now using Claim 3.2, given that  $X_{t-1} \le \gamma n$ , it follows that the total amount of error absorbed by a block until time t can be bounded by:

$$X_t \le \nu \cdot (X_{t-1} + \beta |B|) \le \nu(\gamma + \beta)|B| \le \gamma |B|$$

### **Claim 3.5.** The total depth of the circuit is $O(\log n)$ .

*Proof.* The gate for encoding |B|-length blocks in  $C_g$ , is a Clifford and therefore can be computed in  $O(\log |B|)$  depth. The encoding of the magic/bell states, done by first compute them in the logical space (un-encoded qubits) and then by using the encoder. Hence it's foult-tolerence version of both initializing ancillaries and magic states /bell states cost  $O((\log |B|) \cdot \log^c(|B| \log |B|))^1$  depth [AB99]. Backing into  $C_g$  from  $C_{ft}$  by decoding the concatenation code takes exactly as the encoding namely.

Then using the bell measurements any of the logical gates takes O(1) depth and since we use perform only a single round of error correction we get that the reaming computation till the last decoding stage is a at most constant time of the original depth. Finally we pay  $O(\log |B|)$  for complete decoding. Summing all, we get:

$$O(\log |B| \cdot \log^c(|B| \log |B|)) + O(\text{original depth}) + O(\log |B|)$$

$$=O(\text{original depth}) + O(\log^c |B|)$$

Taking the block length to be  $|B|=\log((W\cdot D)^c)$  gives, by Claim 3.4, a linear fault tolerance construction that success with probability  $1-\frac{1}{\log^{c_2}(W\cdot D)}$ . Particularly, the fault tolerance version of circuits in  $\mathbf{QNC}_1$  has logarithmic depth. Then using the construction in [Aha+96] yields a polynomial fault tolerance circuit, in the only reversible gates setting.

 $<sup>^1 \</sup>text{The}$  width of the original circuit is  $|B|^2$  so the number of locations is  $|B|^2 \cdot \log |B|$ 

<sup>&</sup>lt;sup>2</sup>Assuming W is polynomial in D

## References

- [Aha+96] D. Aharonov et al. *Limitations of Noisy Reversible Computation*. 1996. arXiv: quant ph/9611028 [quant-ph]. url: https://arxiv.org/abs/quant-ph/9611028.
- [AB99] Dorit Aharonov and Michael Ben-Or. *Fault-Tolerant Quantum Computation With Constant Error Rate.* 1999. arXiv: quant-ph/9906129 [quant-ph].
- [PK21] Pavel Panteleev and Gleb Kalachev. Asymptotically Good Quantum and Locally Testable Classical LDPC Codes. 2021. DOI: 10.48550/ARXIV.2111.03654. URL: https://arxiv.org/abs/2111.03654.
- [Din+22] Irit Dinur et al. *Good Locally Testable Codes*. 2022. DOI: 10.48550/ARXIV.2207.11929. URL: https://arxiv.org/abs/2207.11929.
- [LZ22a] Anthony Leverrier and Gilles Zémor. Decoding quantum Tanner codes. 2022. arXiv: 2208. 05537 [quant-ph]. url: https://arxiv.org/abs/2208.05537.
- [LZ22b] Anthony Leverrier and Gilles Zémor. *Quantum Tanner codes.* 2022. arXiv: 2202. 13641 [quant-ph].