## $\mathbf{QNC}_1 \subset \mathbf{noisy}\mathbf{-BQP}$

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#### 1 Notations.

 $C_g$  - good qLDPC,  $C_{ft}$  - concatenation code (ft stands for fault tolerance). For a code  $C_y$  we use  $\Phi_y, E_y, D_y$  to denote the channel maps circuits into the circuits compute in the code space, the encoder, and the decoder. We use  $\Phi_U$  to denote the 'Bell'-state storing the gate U.

#### 2 The Noise Model

### 3 Fault Tolerance (With Resets gates) at Linear Depth.

**Claim 3.1.** There is  $p_{th} \in (0,1)$  such that if  $p < p_{th}$  then any quantum circuit C with depth D and width W can be computed by p-noisy, resets allowed, circuit C', with a depth at most  $\max\{D, \log(WD)\}$ .

#### 3.1 Initializing Magic for Teleportation gates and encodes ancillaries.

The Protocol:

- 1. Initializing zeros. Divide the qubits into |B|-size blocks. Encodes each block in  $C_g$  via  $D_{ft}\Phi_{ft}[E_g]|0^{|B|}\rangle$ .
- 2. Initializing Magic for Teleportation gates encoded in  $C_g$  via  $D_{ft}\Phi_{ft}[E_g]|\Phi_U\rangle$  for each gate U in the original circit .
- 3. Each gate is replaced by gate teleportation.
- 4. At any time tick, any block runs a single round of error reduction.

**Claim 3.2.** Assume that an error  $|e| = \gamma n$ , i.e e is supported on less than  $\gamma n$  bits, then a single correction round reduce e into an error e' such  $|e'| < \nu |e|$ .

**Claim 3.3.** The gate  $D_{ft}\Phi_{ft}[E_q]$  initializes states encoded in  $C_q$  subject to 3p-noise channel.

Proof. Clearly  $\Phi_{ft}[E_g]$  success, with high probability, let's say  $1-\frac{1}{poly(n)}$ , to encode in to  $C_{ft}\circ C_g$ . Denote by  $E_i,D_i$  the encoder and the decoder at the ith level of the concatination construction. Recall that by definition  $D_iE_i=I$ , or in other words  $D_i=E_i^{\dagger}$ . Consider the decoder under  $\mathcal N$  action.  $P_2D_1P_2D_2,...,P_{i-1}D_iP_i$ , by the fault-tolerance construction a logical error happens at the ith stage occurs with probability  $p^{2^i}$ , therefore by the union bound the probability that in one of the steps the circuit absorbs an error that is not corrected is less than  $p+p^2+p^4+...<2p$ . Hence any decoded qubit absorbs a noise with probability less than 2p.

Thus in overall we can bound the porobability a single qubit to be faulty by:

$$\begin{split} \mathbf{Pr}\left[\text{fault}\right] &= \mathbf{Pr}\left[\text{fault}|\Phi_{ft}[E_g]\right] \cdot \mathbf{Pr}\left[\Phi_{ft}[E_g]\right] + \mathbf{Pr}\left[\text{fault}|\overline{\Phi_{ft}[E_g]}\right] \cdot \mathbf{Pr}\left[\overline{\Phi_{ft}[E_g]}\right] \\ &\leq \mathbf{Pr}\left[\text{fault}|\Phi_{ft}[E_g]\right] + \mathbf{Pr}\left[\overline{\Phi_{ft}[E_g]}\right] \leq 2p + \frac{1}{poly(n)} \leq 3p \end{split}$$

**Remark 3.1.** In our construction we use the concatinate-code to encode  $\log(n)$ -length block, Thus any poly(n) in the above should be replaced by  $\log(n)$ . Yet it doesn't effect anything since the inequality dosn't depend on n.

**Claim 3.4.** With probability  $1 - \frac{WD}{|B|} \cdot D2e^{-2|B|(\beta-p)}$ , the total amount of noise been absorb in a block, in any time t, is less than  $\gamma n$ .

*Proof.* Consider the ith block, denoted by  $B_i$ . Using the Hoeffding's inequality we have that the probability that more than  $\beta|B|$  bits are flipped at time t is less than  $\leq 2e^{-2|B|(\beta-p)}$ . Using the union bounds over all the blocks at all the different time location we get that with probability  $1 - \frac{WD}{|B|} \cdot D2e^{-2|B|(\beta-p)}$ . Denote by  $X_t$  the support's size of the error over  $B_i$  at time t. Now using Claim 3.2, given that  $X_{t-1} \leq \gamma n$  it follows that total amount of error absorbed by a block until time t can be bounded by:

$$X_t \le \nu \cdot (X_{t-1} + \beta |B|) \le \nu(\gamma + \beta)|B| \le \gamma |B|$$

**Claim 3.5.** *The total depth of the circuit is*  $O(\log n)$ .

*Proof.* The gate for encoding |B|-length blocks in  $C_g$ , is a Clifford and therefore can be computed in  $O(\log |B|)$  depth. The encoding of the magic/bell states, done by first compute them in the logical space (un-encoded qubits) and then by using the encoder. Hence it's foult-tolerence version of both initializing ancillaries and magic states /bell states cost  $O((\log |B|) \cdot \log^c(|B| \log |B|))^1$  depth [AB99]. Backing into  $C_g$  from  $C_{ft}$  by decoding the concatenation code takes exactly as the encoding namely.

Then using the bell measurements any of the logical gates takes O(1) depth and since we use perform only a single round of error correction we get that the reaming computation till the last decoding stage is a at most constant time of the original depth. Finally we pay  $O(\log |B|)$  for complete decoding. Summing all, we get:

$$\begin{aligned} &O(\log|B| \cdot \log^c(|B|\log|B|)) + O(\text{original depth}) + O(\log|B|) \\ = &O(\text{original depth}) + O(\log^c|B|) \end{aligned}$$

Taking the block length to be  $|B| = \log((W \cdot D)^c)$  gives, by Claim 3.4, a linear<sup>2</sup> fault tolerance construction that success with probability  $1 - \frac{1}{\log^{c_2}(W \cdot D)}$ . Particularly, the fault tolerance version of circuits in QNC<sub>1</sub> has logarithmic depth. Then using the construction in [Aha+96] yields a polynomial fault tolerance circuit, in the only reversible gates setting.

#### References

[Aha+96] D. Aharonov et al. Limitations of Noisy Reversible Computation. 1996. arXiv: quant - ph/9611028 [quant-ph]. url: https://arxiv.org/abs/quant-ph/9611028.

[AB99] Dorit Aharonov and Michael Ben-Or. *Fault-Tolerant Quantum Computation With Constant Error Rate.* 1999. arXiv: quant-ph/9906129 [quant-ph].

<sup>&</sup>lt;sup>1</sup>The width of the original circuit is  $|B|^2$  so the number of locations is  $|B|^2 \cdot \log |B|$ 

 $<sup>^2</sup>$ Assuming W is polynomial in D